

ATF Program Advisory Committee & ATF Users' Meeting

April 2-3, 2009 - Brookhaven National Laboratory

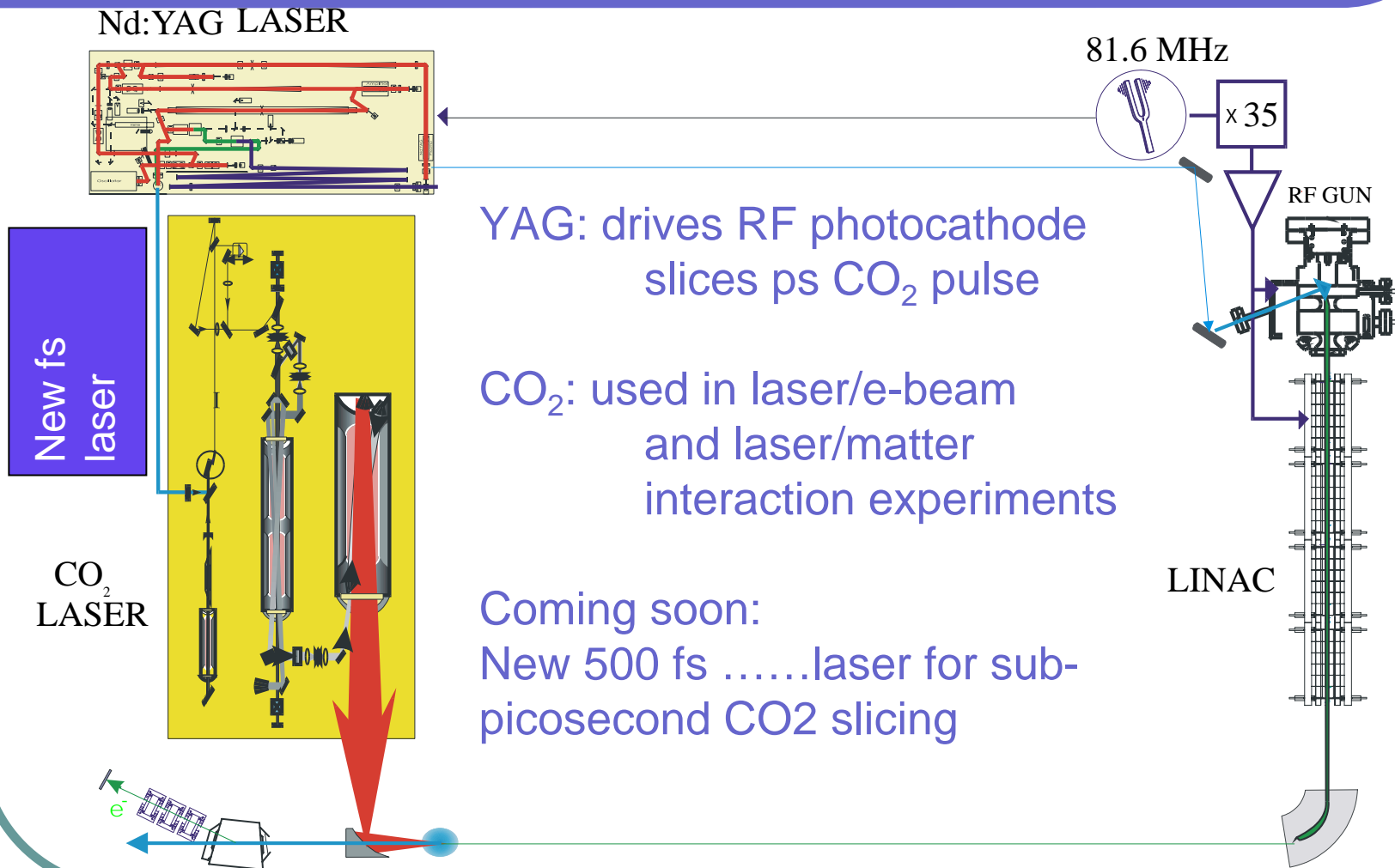
ATF CO₂ LASER

back to basics
present status
research highlights

Igor Pogorelsky



ATF block diagram

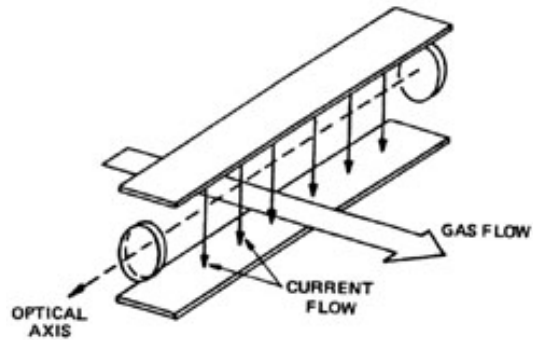


ATF laser personnel

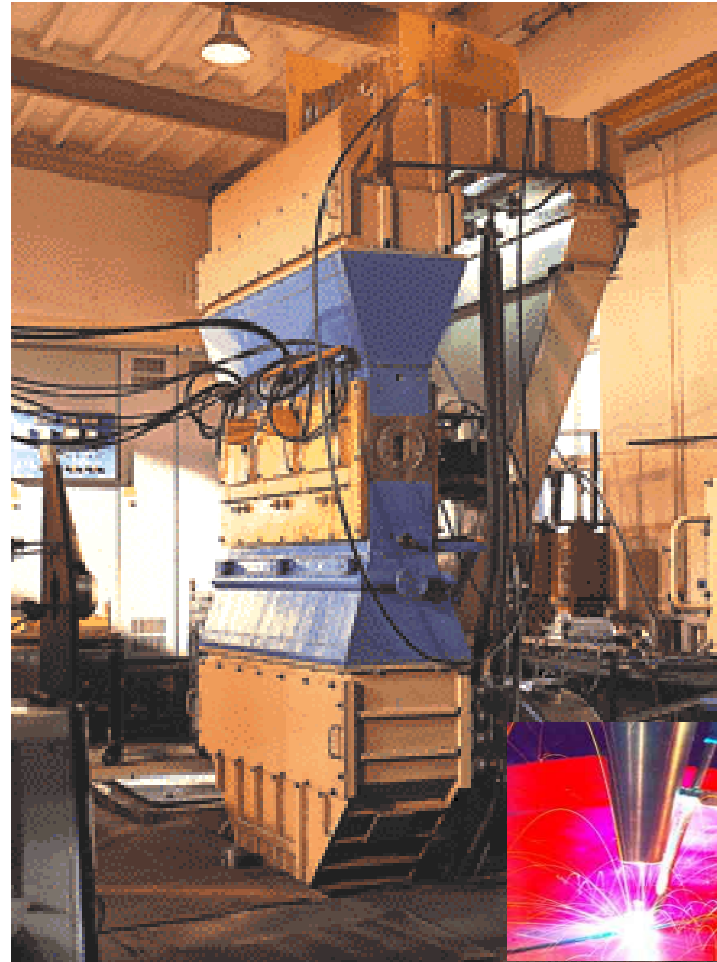


- Igor Pogorelsky CO₂ laser and experiments support
 - Mikhail Poliansky LDRD Postdoc, simulations, CO₂ diagnostics
 - Marcus Babzien YAG laser, general optical diagnostics
 - Daniil Stoliarov Post-doc, fs solid-state laser
 - Karl Kusche laser safety, computer controls
 - Vitaly Yakimenko global laser strategy
- + ATF computer engineer, electronic engineer, designers, technicians

Industrial CO₂ lasers



- Up to 100 kW average power
- Operate at low pressure $\ll 1$ atm
- Bandwidth $\sim P$ (10 atm supports a picosecond pulse)



Ultrafast gas lasers require high pressure

Inverse Fourier Transform for discrete spectrum results in a train of discrete pulses

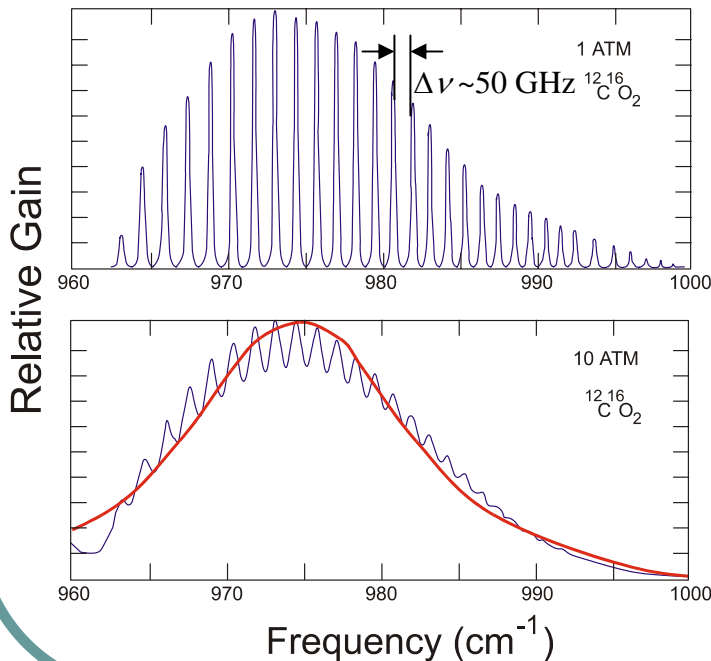
$$f(t) = \sum_{j=0}^n F_j(t) \exp[i2\pi(\nu_0 + \Delta\nu \times j)t]$$

where $n \approx 1/\tau_0 \Delta\nu$, τ_0 is the initial pulse width.

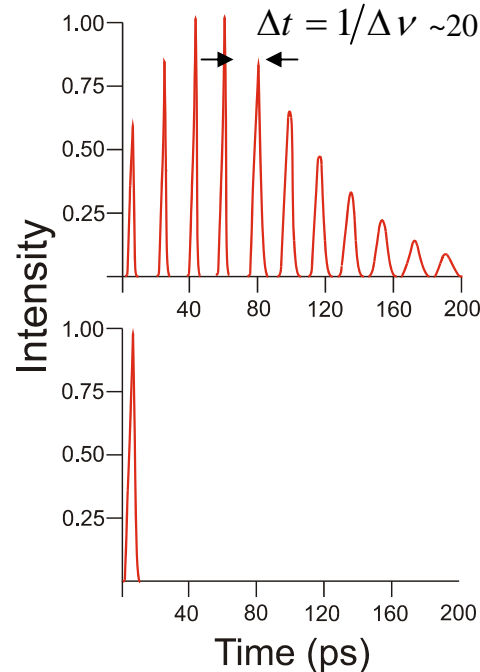
$$f(t) = \sum_{k=0}^m f_k(t_0 + \Delta t \times k)$$

where $\tau_k \approx \tau_0$, $m \approx 1/\tau_0 \delta\nu$, $\delta\nu$ - rotational linewidth defined by pressure broadening.

Gain Spectrum



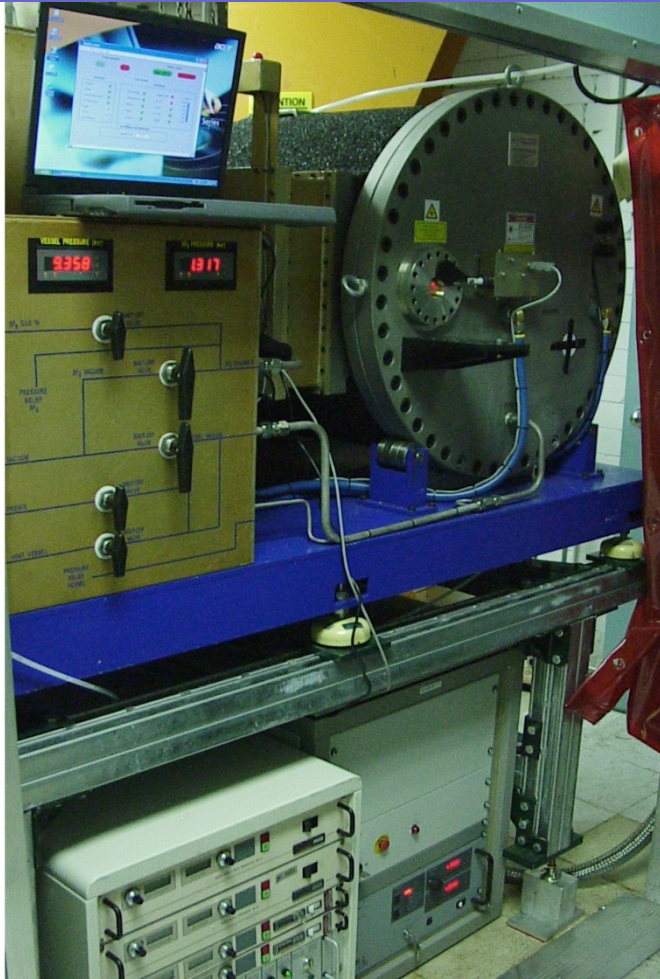
Amplified Picosecond Pulse



Strongly modulated rotational line structure of the CO_2 gain spectrum modifies the frequency content of picosecond pulses, changing their temporal structure.

At 10 atmospheres, collisional broadening produces overlap of the rotational lines into the 1 THz wide quasi-continuous gain spectrum, and pulses as short as 1 ps can be amplified without distortion.

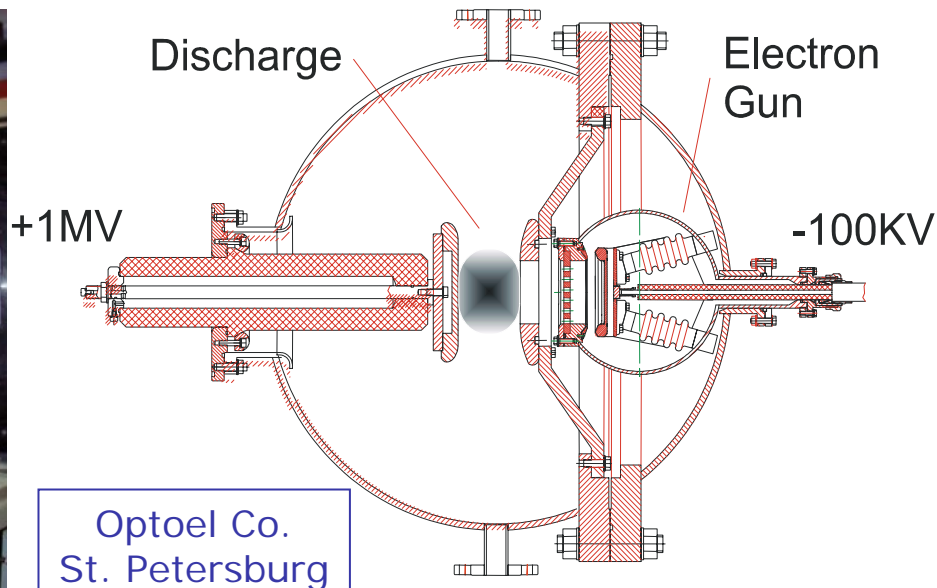
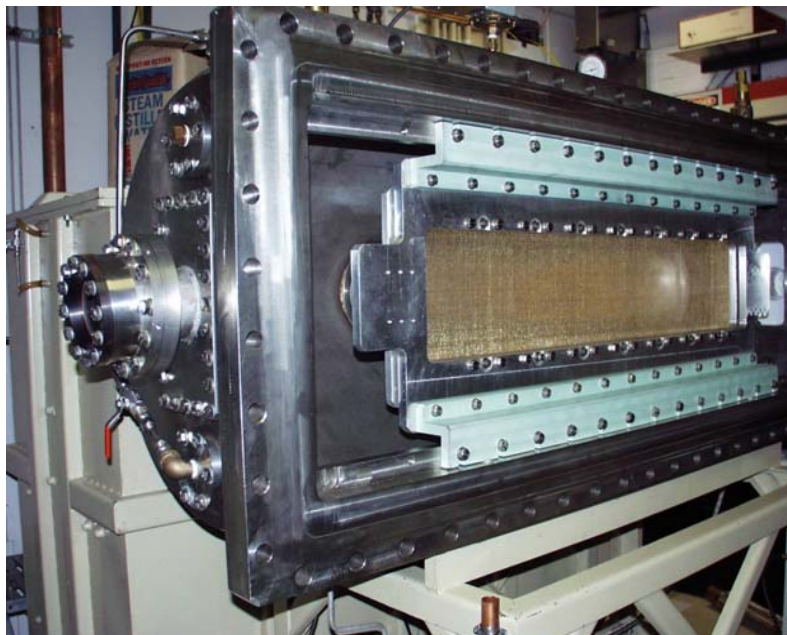
Commercially Available High-Pressure CO₂ Laser with UV-ionized discharge



UV-preionized	
Pressure	10 atm
Beam Size	13 x 13 mm ²
Repetition Rate	20 -500 Hz
Pulse Energy	1.5 J
Average Power	750 W



Custom high-pressure x-ray ionized CO₂ amplifier operated at BNL/ATF

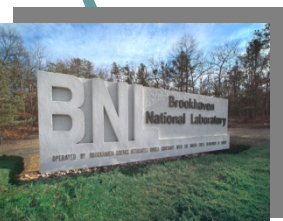


Amplifier cross-section diagram

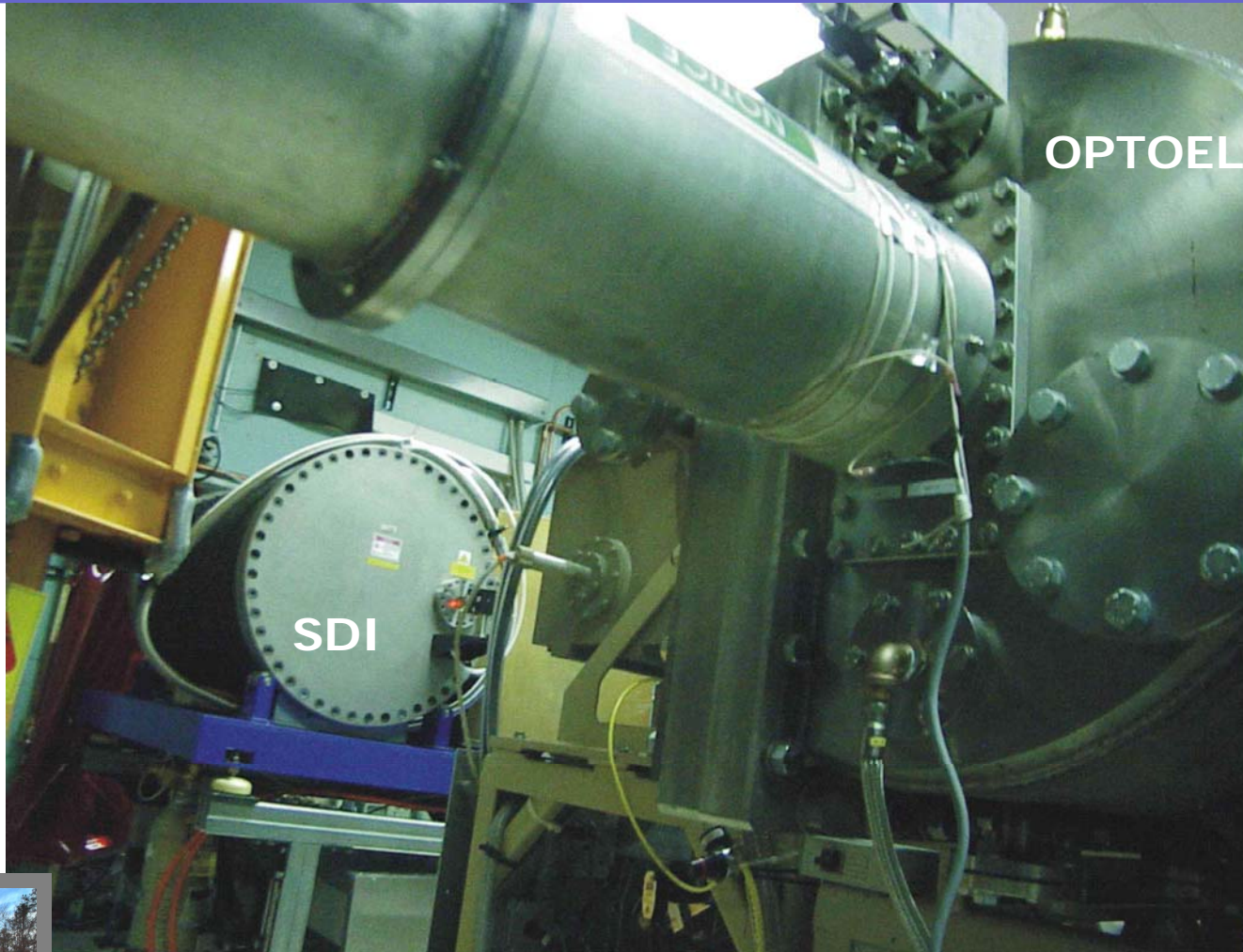
Discharge 10x10x100 cm³
 Pressure 10 atm
 Voltage 1 MV

Stored energy $h\nu \times \Delta N = 10 \text{ mJ/cm}^3$
 Small-signal gain $\sigma \times \Delta N = 2\%/cm$
 Saturation fluence $h\nu/\sigma = 0.5 \text{ J/cm}^2$

$$h\nu = 0.1 \text{ eV} = 1.6 \times 10^{-20} \text{ J}, \quad \sigma = 3 \times 10^{-20} \text{ cm}^2, \quad \Delta N = 10^{18} \text{ cm}^{-3}$$



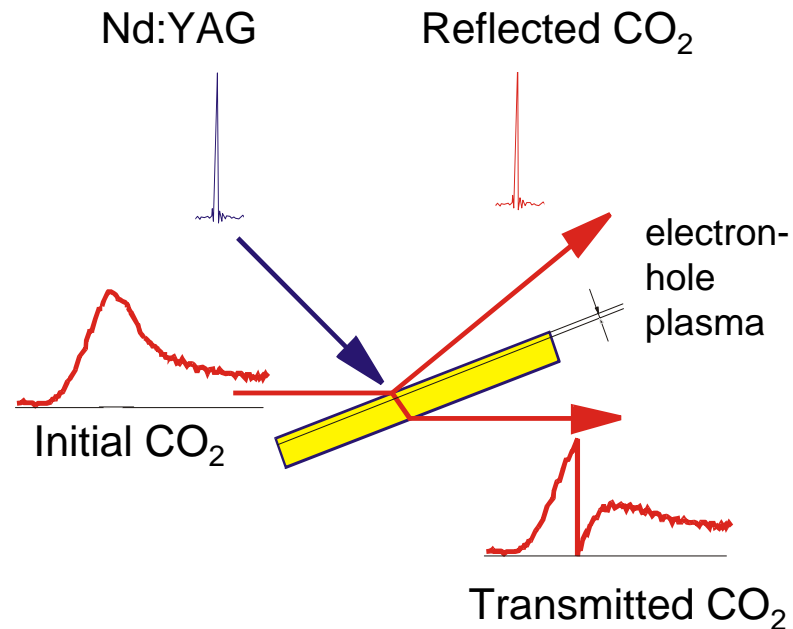
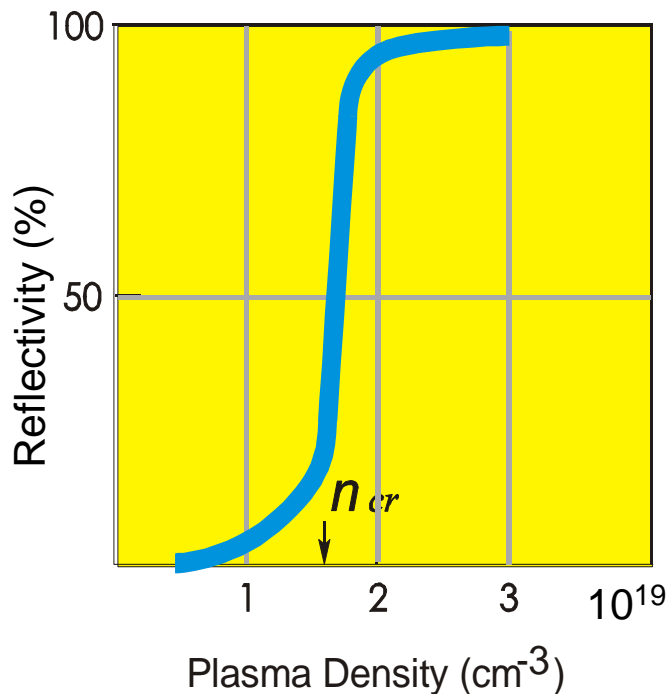
High-pressure CO₂ amplifiers operated at BNL/ATF



Methods of short pulse generation

Semiconductor optical switch

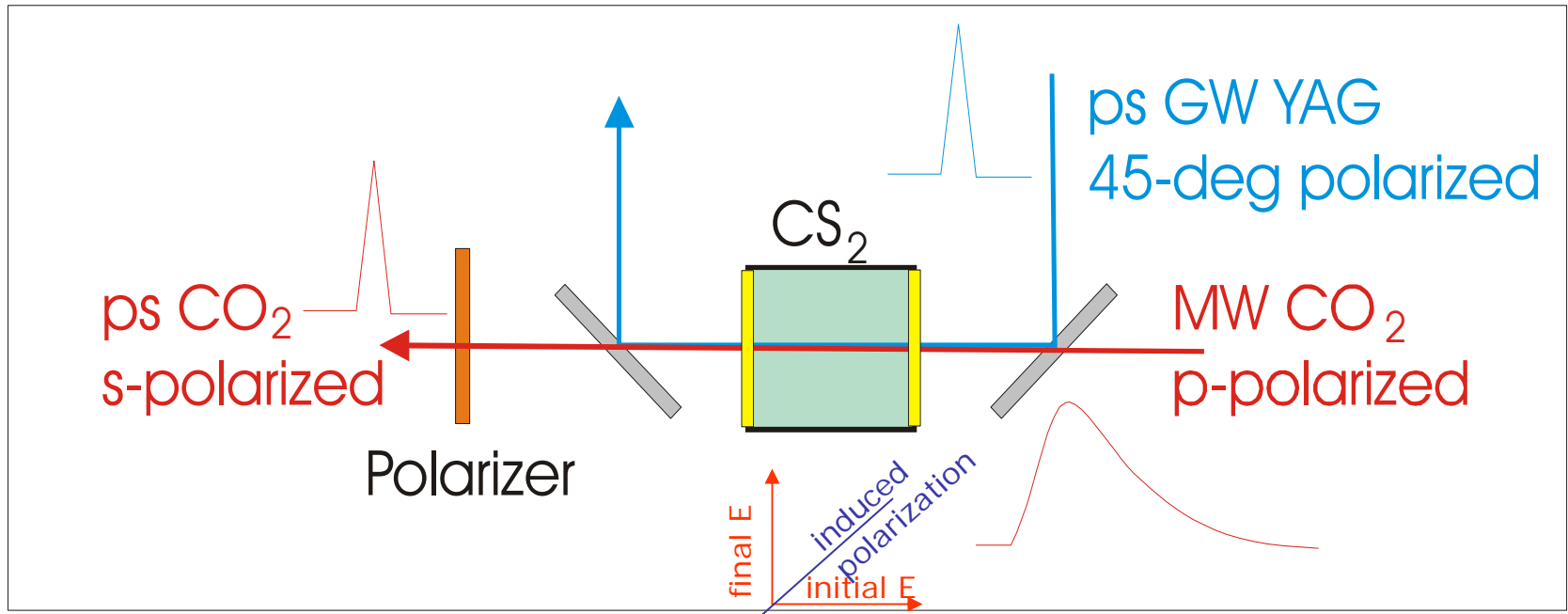
- Reflection by electron-hole super-critical plasma controlled by a short wavelength laser.



- Reflection lasts for 200 ps defined by relaxation of free carriers.
- Two switches allow to slice CO₂ pulse to the duration of a control pulse.

Methods of short pulse generation

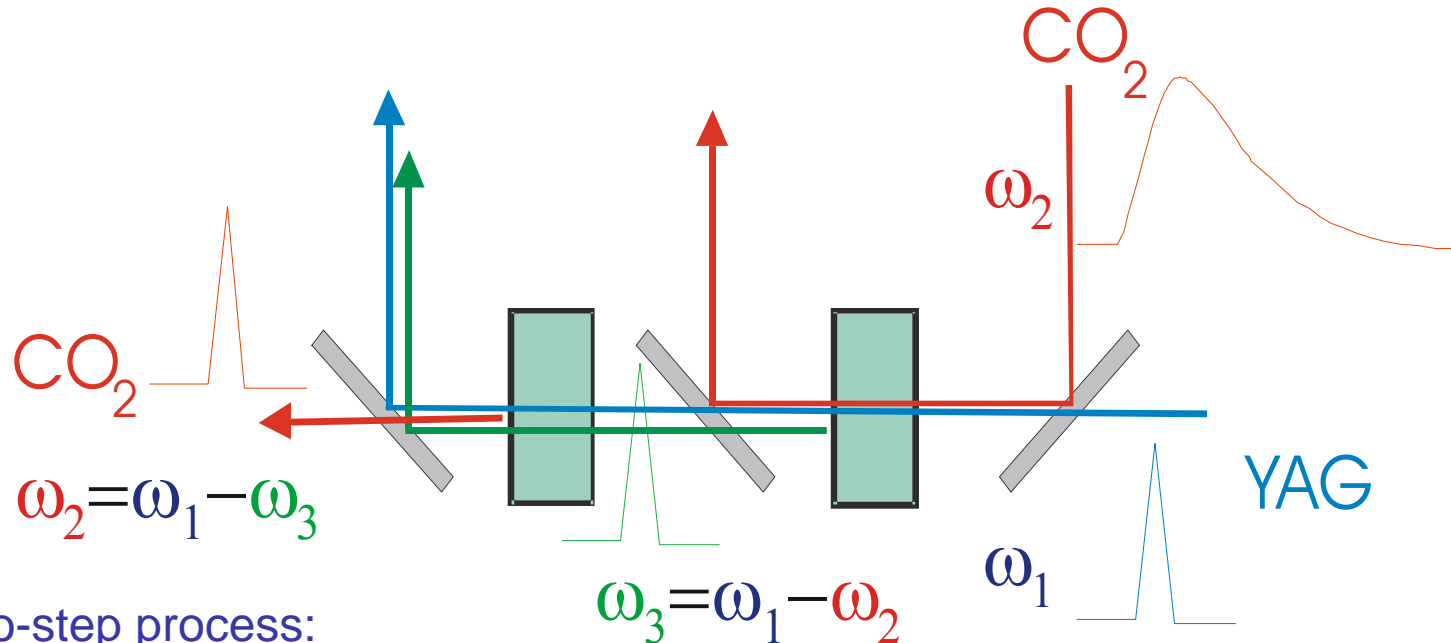
Kerr effect



- A higher-power laser induces polarization in optically active CS₂ liquid.
- This rotates polarization of a portion of a CO₂ pulse overlapped with YAG.
- Relaxation of CS₂ molecules in 1-2 ps.

Methods of short pulse generation

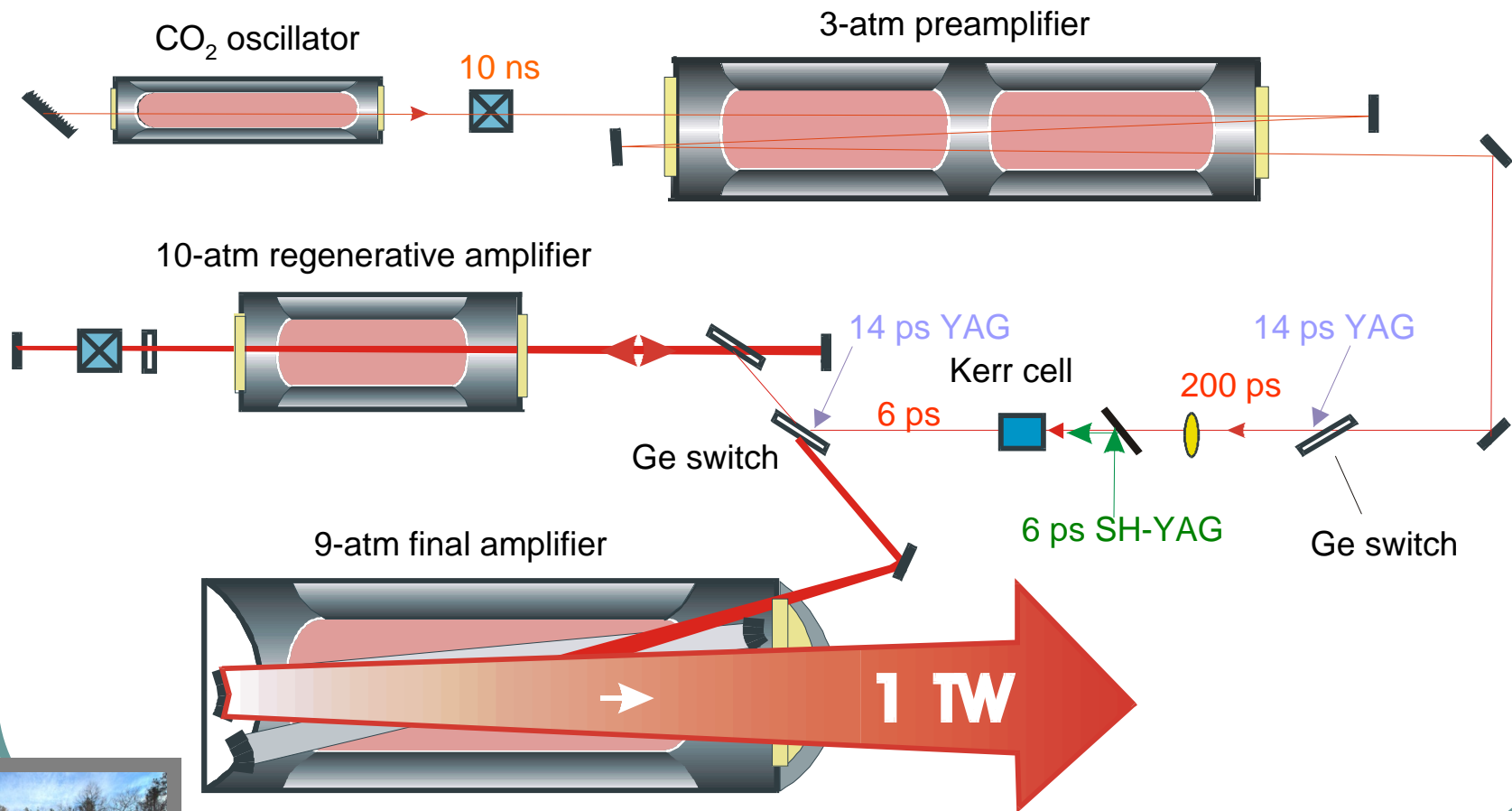
Parametric generator



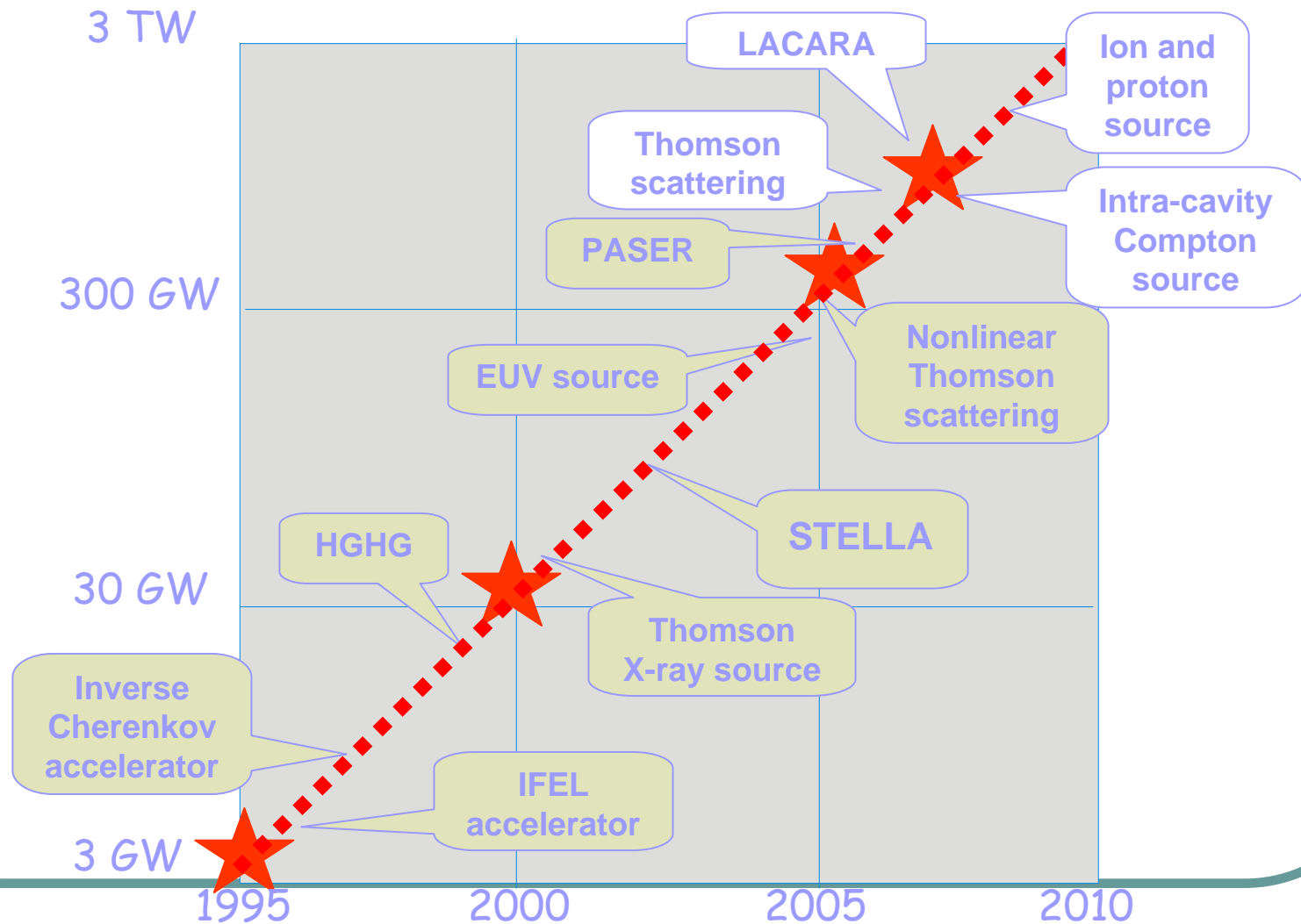
Two-step process:

- Mixing two laser beams in a nonlinear crystal with production of differential frequency.
- Mixing ω_3 and ω_1 with production of a CO₂ pulse at the duration of YAG or shorter.
- Insignificant stretch in thin crystals, 1 ps is possible.

CO₂ laser system delivers 0.5 TW, 6 ps pulses



ATF success story



Benefits from using CO₂ laser

ATF pioneers a *picosecond* CO₂ gas laser for strong-field physics applications.

This provides a new platform for exploring novel methods of particle acceleration and radiation sources.

CO₂ ($\lambda=10 \mu\text{m}$) advantages

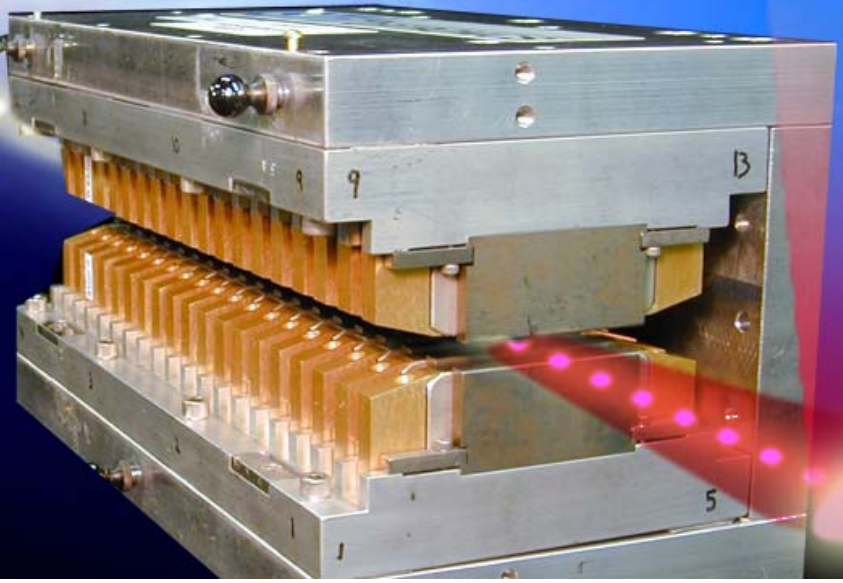
as compared to solid-state lasers ($\lambda\sim 1 \mu\text{m}$) :

- 10 times more photons per Joule
- favorable scaling of accelerating structures, better electron phasing into the field
- stronger ponderomotive effects at the same laser intensity

Long-wavelength benefits: better electron phasing into the field

STELLA

Staged Electron Laser Acceleration



**First Monoenergetic
Laser Acceleration**

Long-wavelength benefits: stronger ponderomotive effects

Energy of the electron quiver motion in laser field E

$$\Phi = \frac{mv^2}{2}$$

\Rightarrow

$$v \sim \frac{\dot{v}}{\omega}$$

\Rightarrow

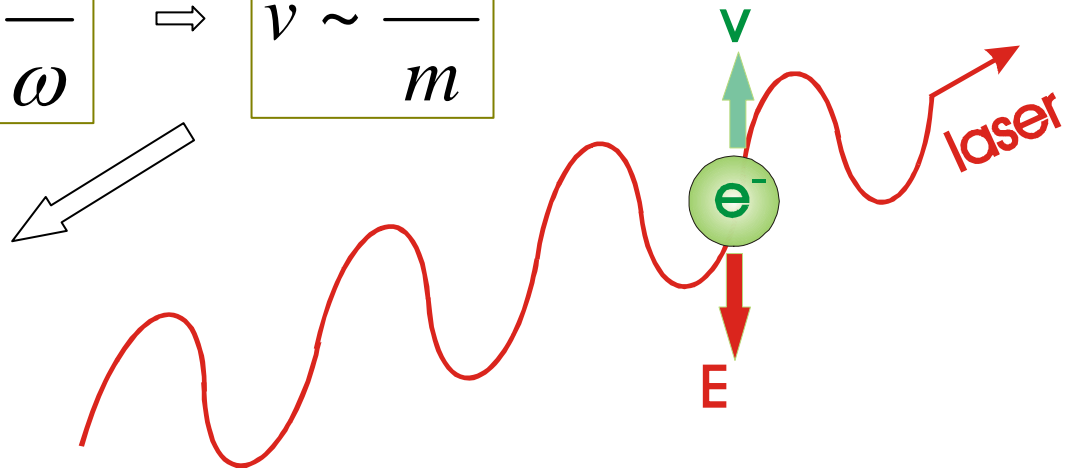
$$\dot{v} \sim \frac{eE}{m}$$

$$\Phi \sim \frac{e^2 E^2}{2m\omega^2}$$

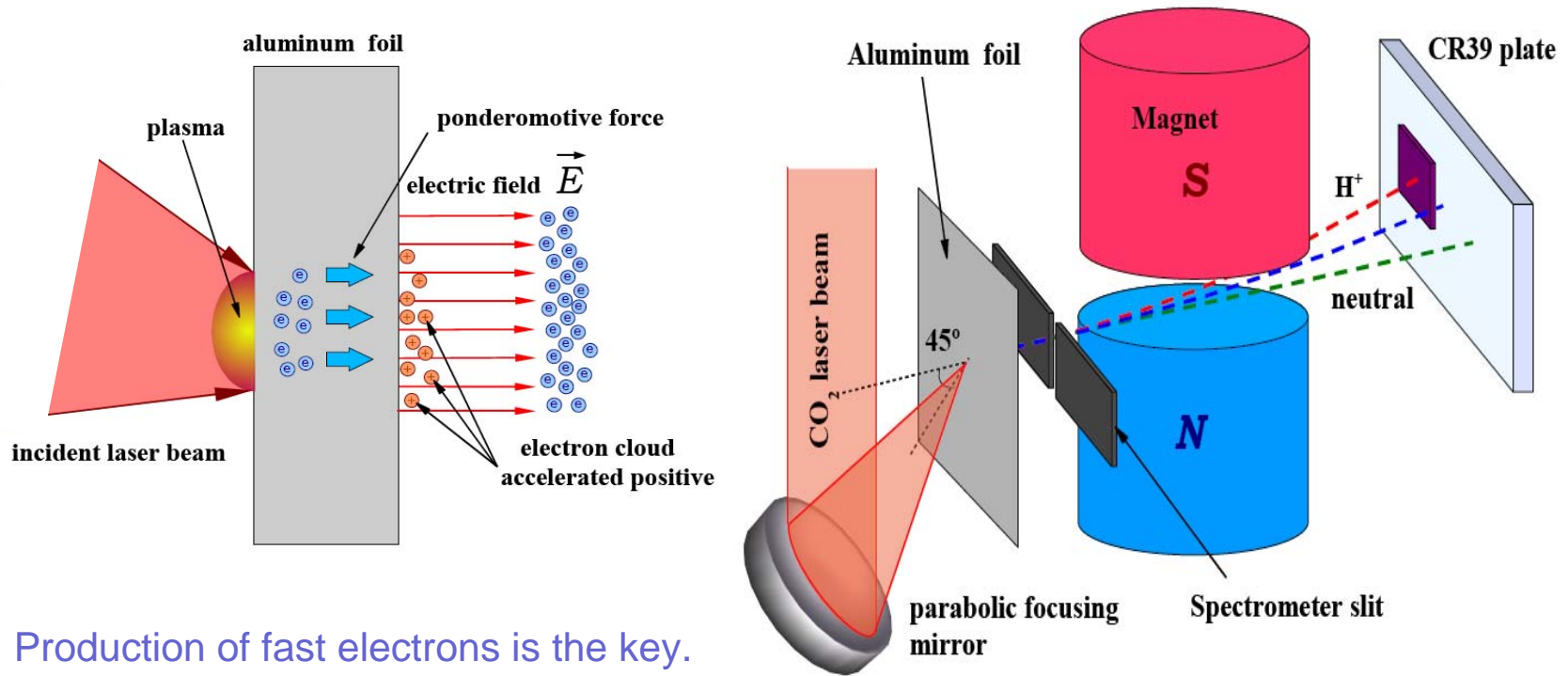
\Downarrow

$$\Phi \sim \frac{I}{\omega^2}$$

note the ω^2 dependence



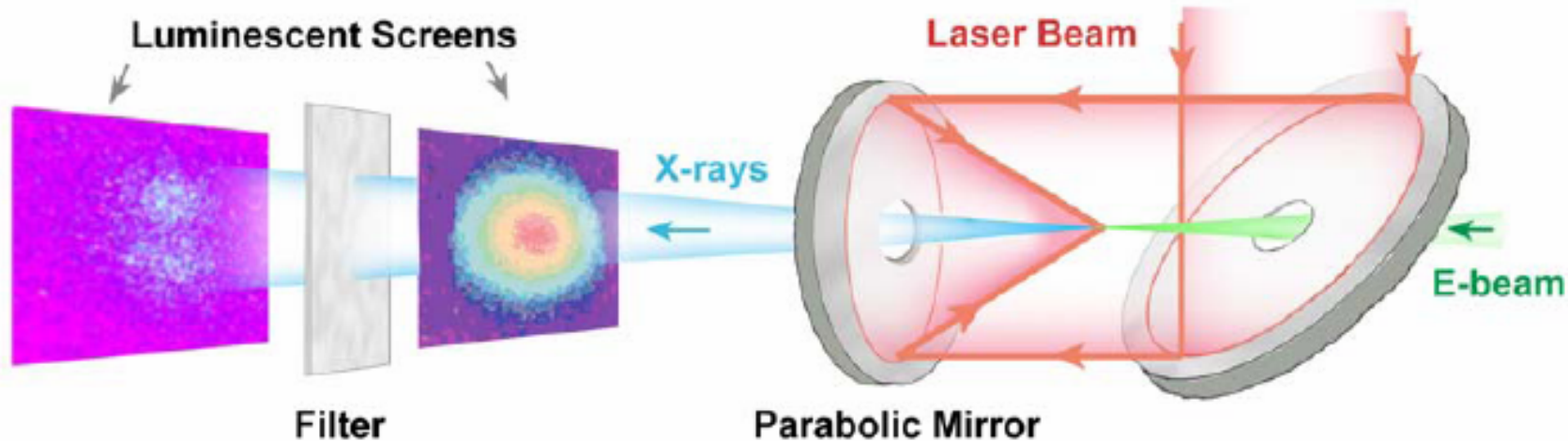
Proton accelerator



Production of fast electrons is the key.

- At the same power and energy, CO₂ laser will provide the same ponderomotive action within $\sim \lambda^2$ (100 times) bigger area or $\sim \lambda^3$ (1000 times) bigger volume.
- Accordingly, we expect that the number of accelerated ions would grow with λ .

Demonstration of record yield and 2nd harmonic in relativistic Thomson scattering



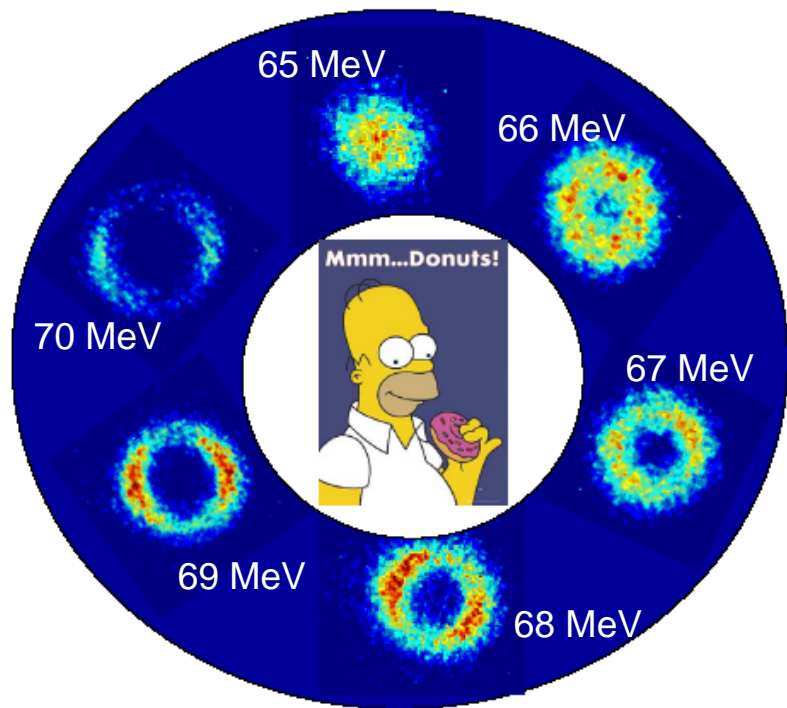
With 5 J CO₂ laser focused to $\sigma = 35 \mu\text{m}$, we demonstrated record x-ray yield $N_x/N_{e^-} \sim 1$ And 2nd harmonic in relativistic Thomson scattering

note that 10- μm laser produces 10 times more photons per Joule than 1-mm laser

$$\frac{N_x}{N_e} = \frac{N_L}{\pi w_0^2} \sigma_T$$

Spin-offs from earlier Thomson scattering experiments

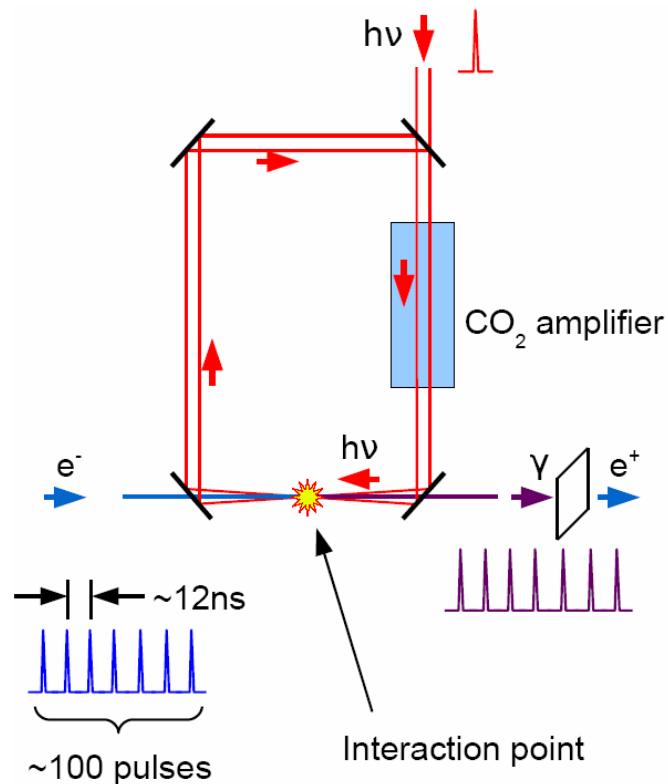
Parametric studies



credits to 

UCLA

Intra-cavity gamma-source



Next laser talks

- Marcus Babzien “*Experimental support and instrumentation*”
- Mikhail Polyanskiy ”*Simulation and diagnostic tools for better understanding and upgrade of the CO₂ laser*”
- Daniil Stolyarov “*A femtosecond CO₂ front end based on frequency mixing with an Ytterbium laser*”
- Igor Pogorelsky “*CO₂ laser: Near-term plans*”