

Exotic Concepts and HED: Working Group Summary

By R. Noble, T. Katsouleas

- Creative, paradigm changing ideas
- Confluence of HED, AAC and Applications
- Charge -- design proton injector

Exotic Concepts WG

- J. Alessi
- D. Umstadter
- K. Mima
- A. Ogata
- J. Fuchs
- P. Awtia
- P. Bolton
- J. Lewellen
- J. Kim
- G. Shvets
- I. Pogorelsky
- B. Bowes
- I. Pavlishin
- J. Smedley
- J. Wurtele
- P. Stoltz
- C. Toth
- T. Cowan
- R. Noble
- T. Katsouleas

From attendee list Day 1

We heard some nice talks...

- Proton/ion acceleration from laser-driven foils
 - Experiments: Lin, Bolton, Ogata, Fuchs, Cowan, Mima
 - Simulations: Silva, Mima, Messmer
 - Applications -- *the charge*: Injectors (Alessi), HIF/HED (Cowan, Wurtele, Bowes)

We heard some nice talks... (II)

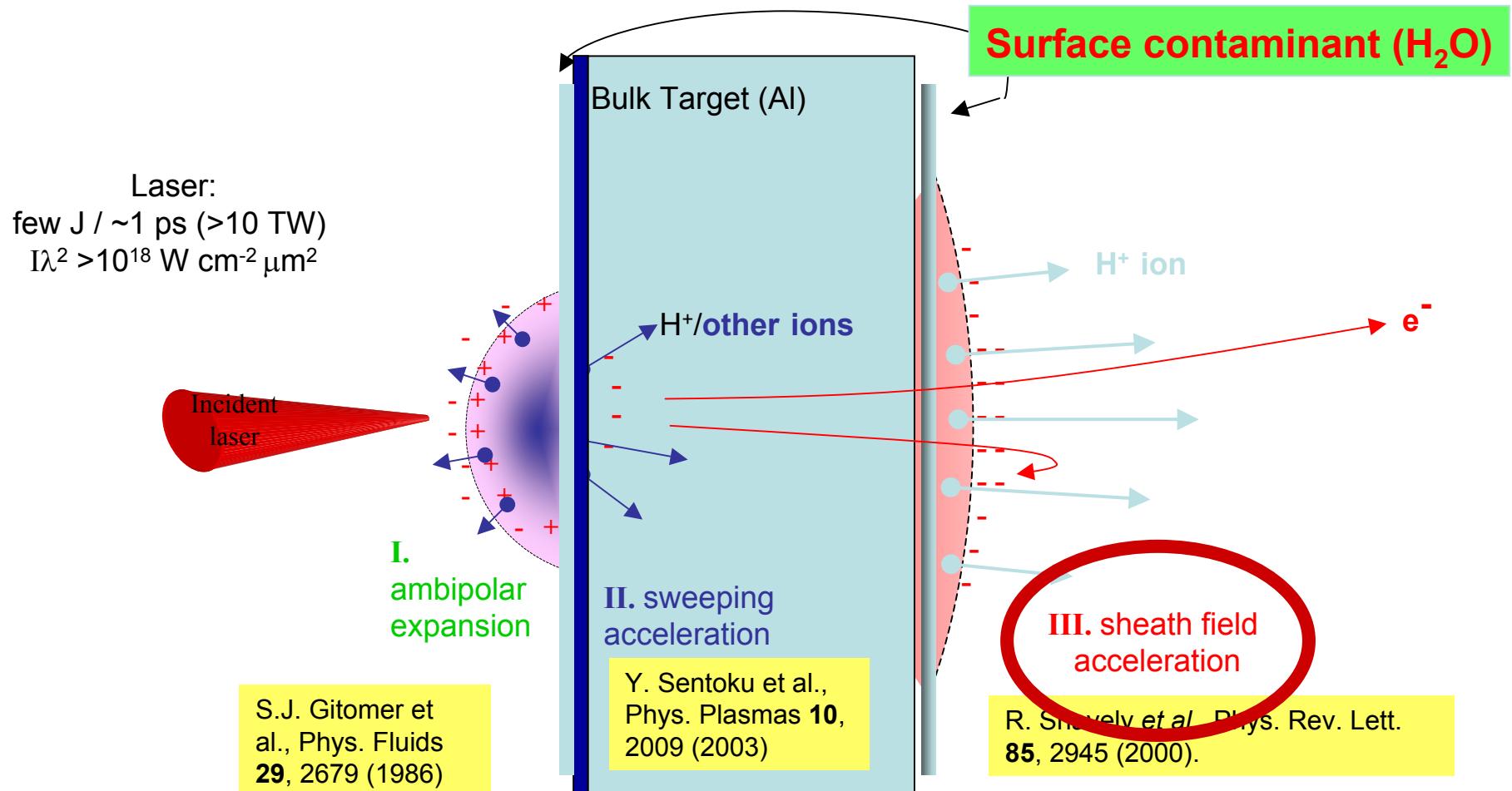
- Vacuum Acceleration (Umstadter, Nakajima)
- Active Medium Acceleration (Shachter)
- Ferroelectrics (Kanareykin)
- Beam conditioning for FELs (Wurtele, Esarey, Schroeder)

Addressing the Charge:

*“Design a laser-ion source
that exceeds existing ion sources in some key parameter
for an application”*

- Injector for an Ion Accelerator
 - for HEP, medicine, or fast physics
- Heavy Ion Fusion/High Energy Density Physics Driver

Mechanisms of laser-acceleration of *ions*

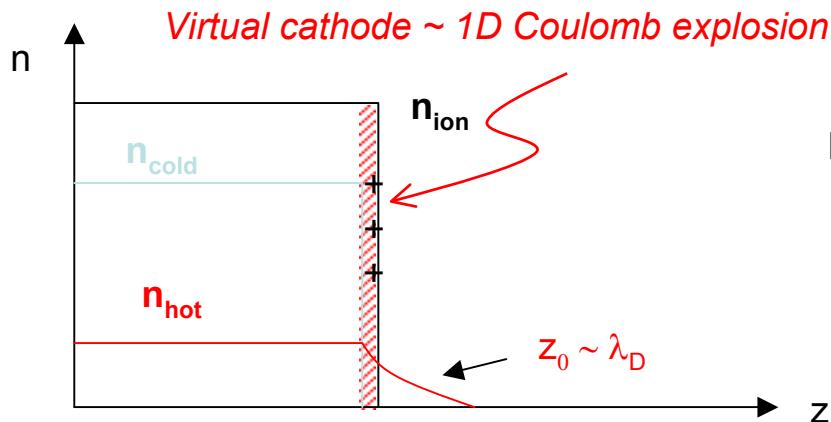


if target is heated → efficient acceleration of heavy ions

[M. Hegelich et al., Phys. Rev. Lett. **89**, 085002 (2002).]

Summary of rear-surface acceleration of protons

①

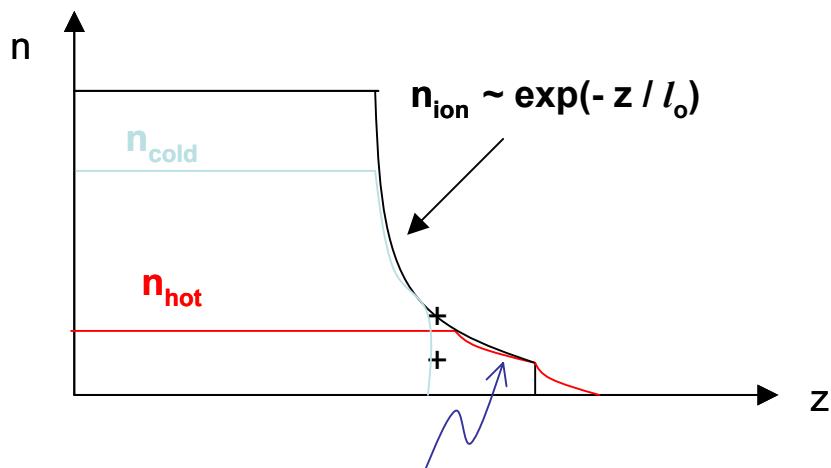


Initial Electric Field:

$$E_z = kT_{\text{hot}}/e\lambda_D$$

$$\sim 10^{12} \text{ V / m}$$

②



Quasineutral, ambipolar expansion

In expanding sheath:

$$E_z(t) \sim kT_{\text{hot}}(t)/e l_o(t)$$

Self-similar expansion of isothermal plasma:

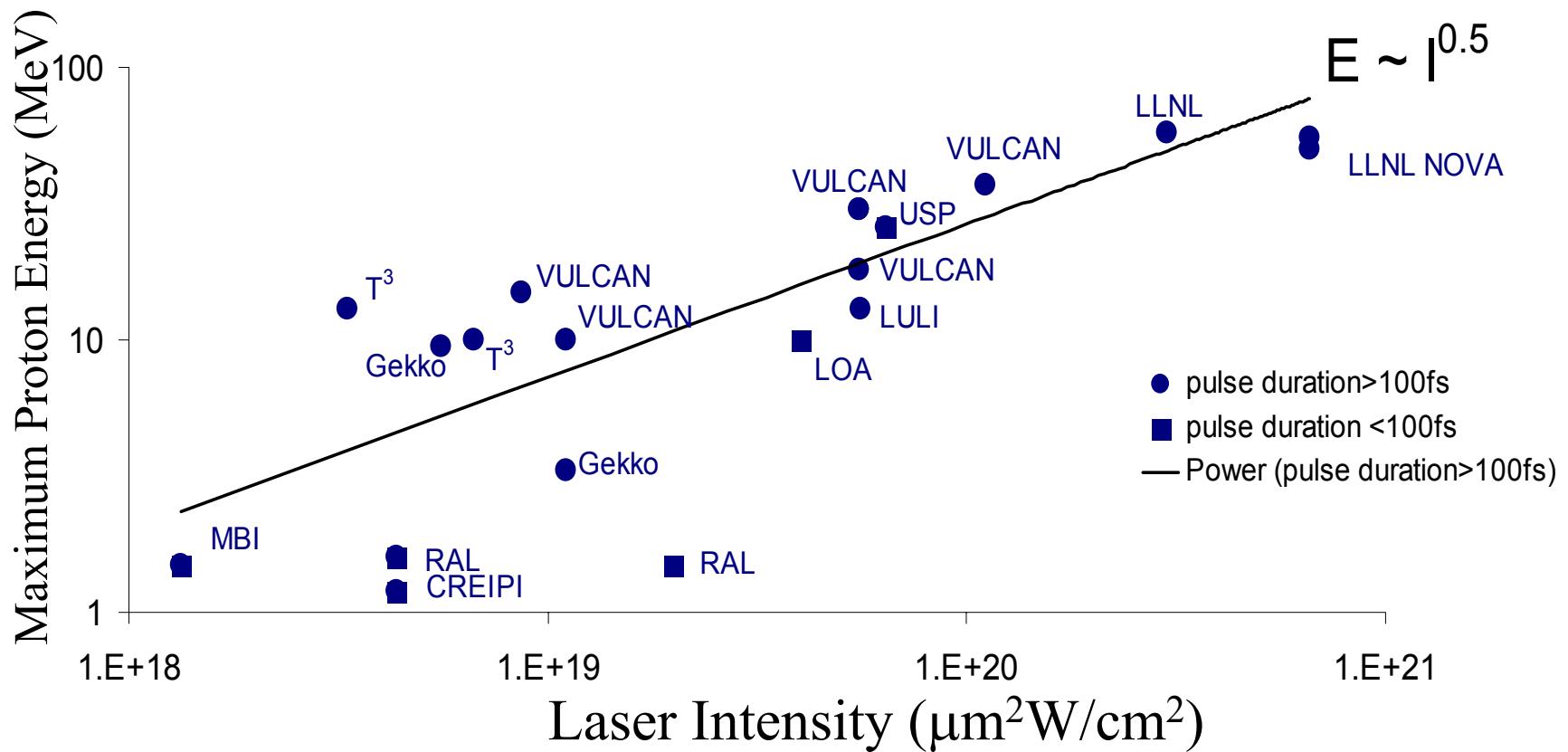
$$l_o(t) = c_s t$$

Then adiabatic:

E_z decreases from expansion, l_o , and cooling of T_{hot} by transfer to ions.



Proton Energy Scaling (T. Lin)



We did some work...

Working Group Result:

$$T_h / mc^2 \sim .8a_o \quad \text{Forslund, 70's}$$

$$V_{\max} = 2c_s \ln(2\omega_{pi} \tau_p) \quad \text{Mora, '03}$$

$$c_s = \sqrt{\frac{T_h}{M}}$$

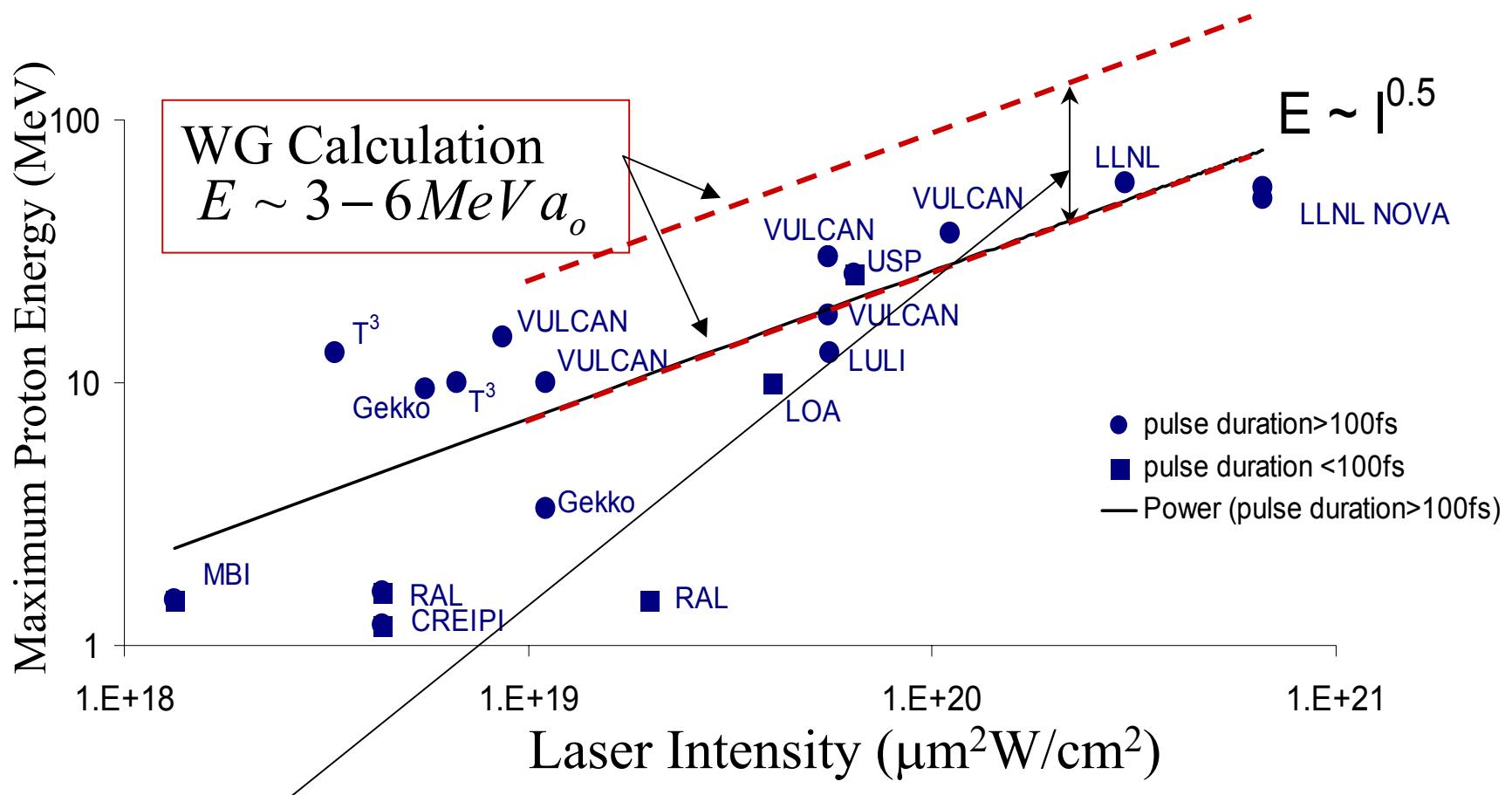
$$E = \frac{1}{2} M V_{\max}^2$$

$$\Rightarrow E \sim 3 - 6 MeV a_o \quad a_o \gg 1$$

Note: this is about 10x $\phi_{\text{sheath}} \sim T_h$



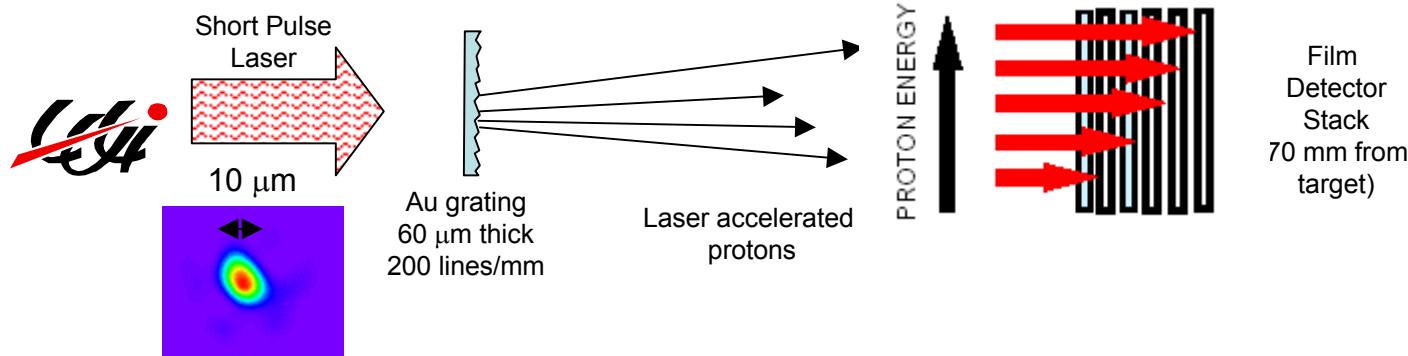
Proton Energy Scaling (T. Lin)



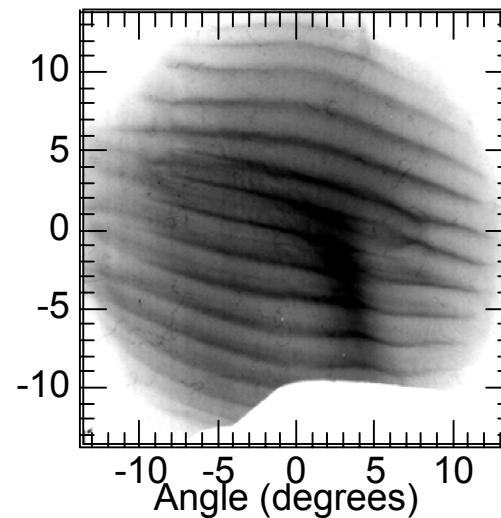
Pulse length dependence in theory
($\ln 2 \omega_{\text{pi}} \tau_p$); Mora, PRL (2003)

Recent measurements show normalized $\varepsilon < .004$ mm-mrad!

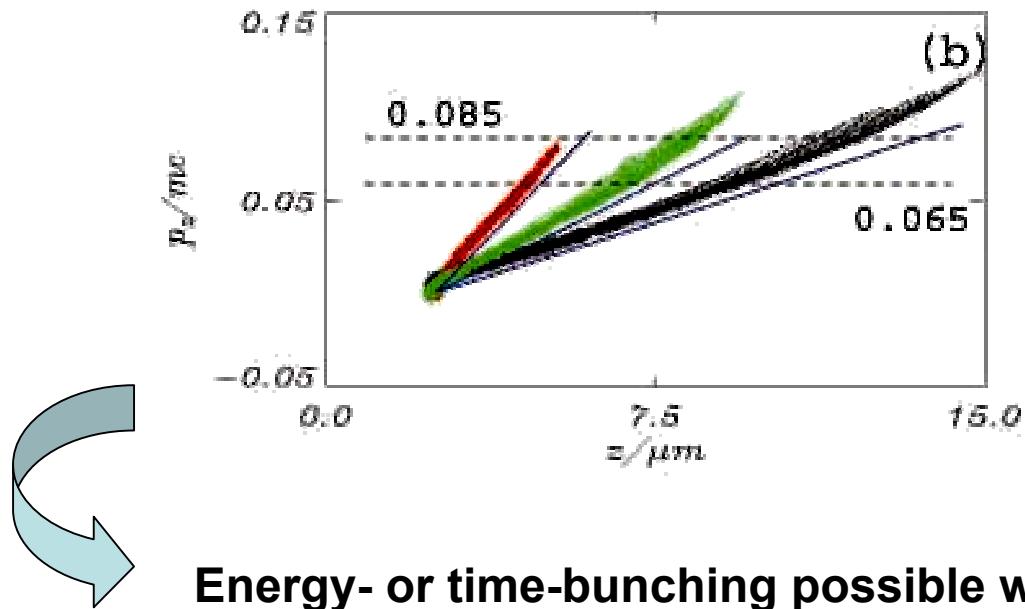
T. Cowan, J. Fuchs, H. Ruhl *et al.*, Phys. Rev. Lett. **92**, 204801 (2004).



8 MeV layer

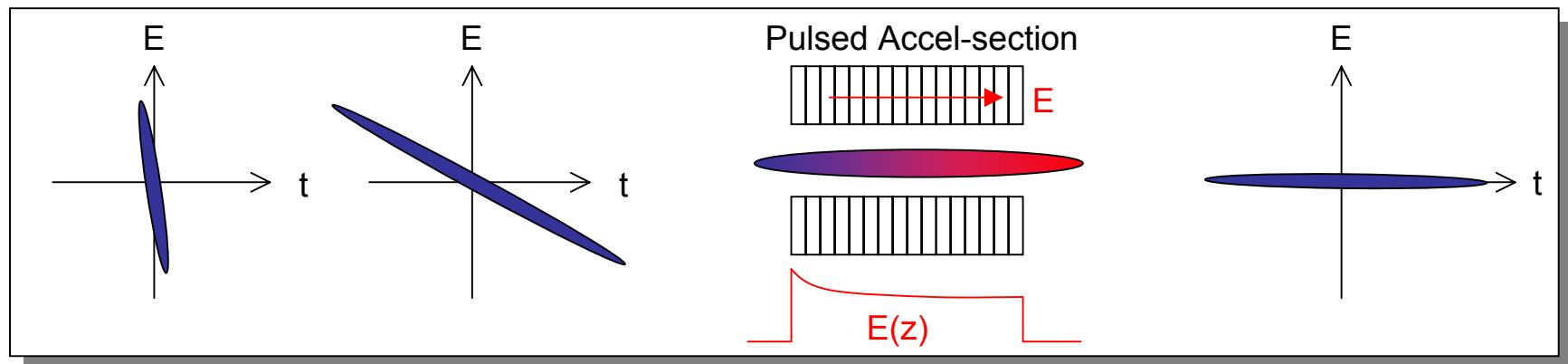


PIC simulations also show an excellent longitudinal emittance



Rapid acceleration produces strong ΔE - Δt correlation

Energy- or time-bunching possible with post-acceleration



Laser-accelerated ions as sources for conventional ion injectors ?

Noble, Alessi, Fuchs

Typical HE linac requirements (proton)

- 10^{13} - 10^{14} p/pulse $\sim 10^4$ nC
- pulse duration: 5-500 μ s
(bunch: <1 ns)
- $\Delta E/E < 10^{-2}$ at injection
- freq: 5-50 Hz
- Injection E: 50 keV
- emittance N: 0.05-0.5 mm.mrad at injection

Laser-accelerated protons

- 10^{10} - 10^{13} p/pulse [TW to PW]
- bunch duration: 1 ps at source debunches to 100 ps in cm
- $\Delta E/E \sim 1$ ($E_{\max} \sim 10$ -50 MeV)
- 10 Hz - hour [TW to PW]
- 10 MeV
- emittance N: 0.005

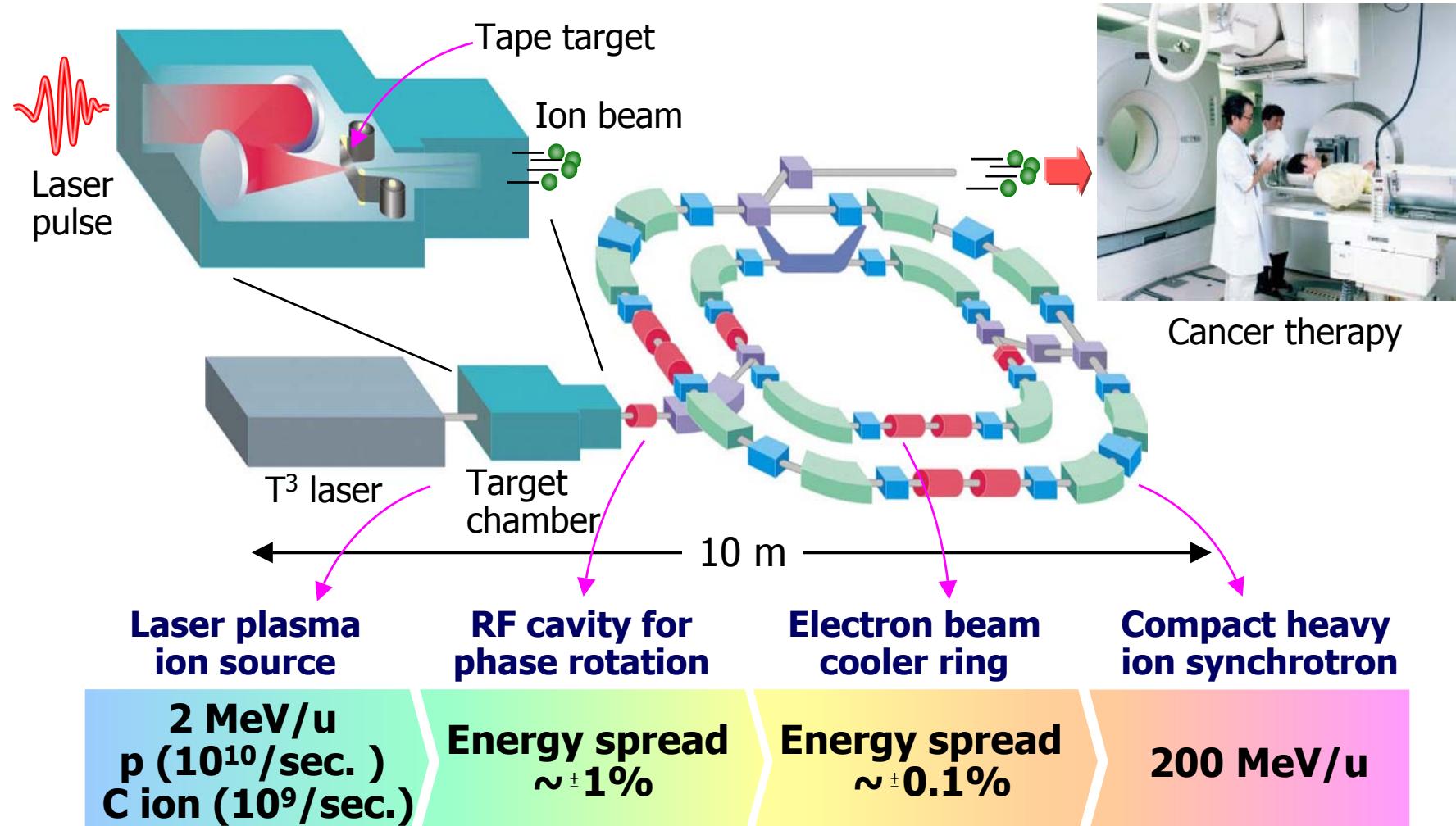
most Synchr. injectors use H⁻

Laser-accelerated ions as sources for conventional ion injectors ?

- How can we get the needed charge?
accelerator- 10^4 nC/pulse vs. TW laser-1 nC/bunch
- Laser average power needs to increase
 - 10 Hz TW laser now ~ 1 W average power
 - Need 1-10 kW (semi-conductor lithog. Lasers)
- Phase-space rotation needed to reduce ΔE once naturally debunched
- Higher injection energy + smaller emittance
could lead to higher brightness out of linac
- Other Applications:
 - low-av. current Carbon (others) medical use
 - short-duration source for pump-probe time-resolved exp.
 - time-resolved radiography (brightness but low number of charges)

Development of Compact Ion Accelerator

National Institute of Radiological Sciences, JAPAN



JAERI, Univ. of Tokyo,
Hiroshima Univ.

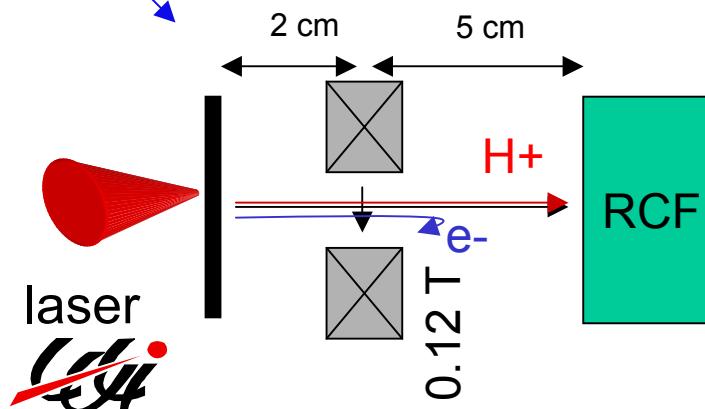
Kyoto Univ.,
JAERI

Kyoto Univ.

KEK
Fukumi, Ogata, ...

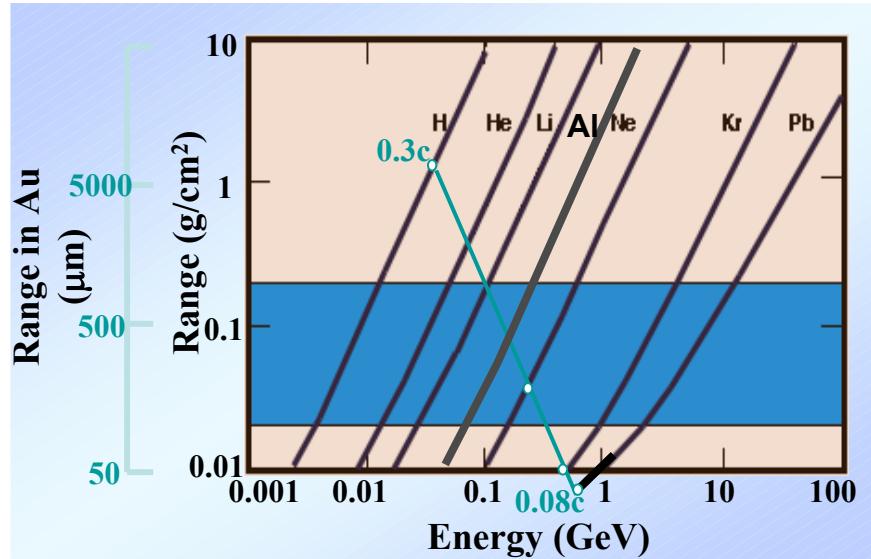
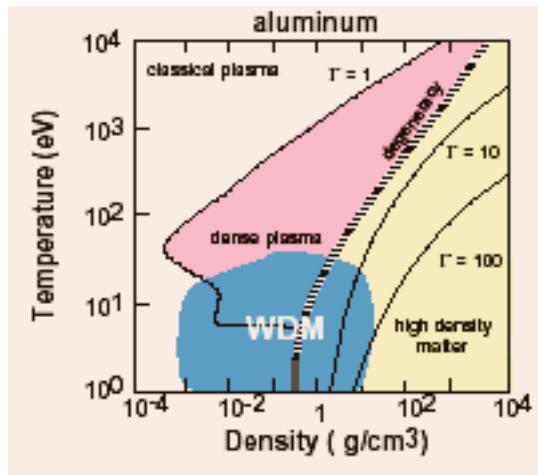
We identified some open questions...

- Finer dependencies of E
 - (spot size, thickness, λ -- C0₂?)
- Can we extend the E ~ I scaling (at low a_o)?
 - Large pre-plasma (1-D sims ok, 2-D?)
- Can we tailor/control acc/energy spectra?
 - Multiple foils, droplets?
 - Tailored plasma density -- near critical density
- Key Milestone -- **extract & post-accelerate** to rotate p_z-z

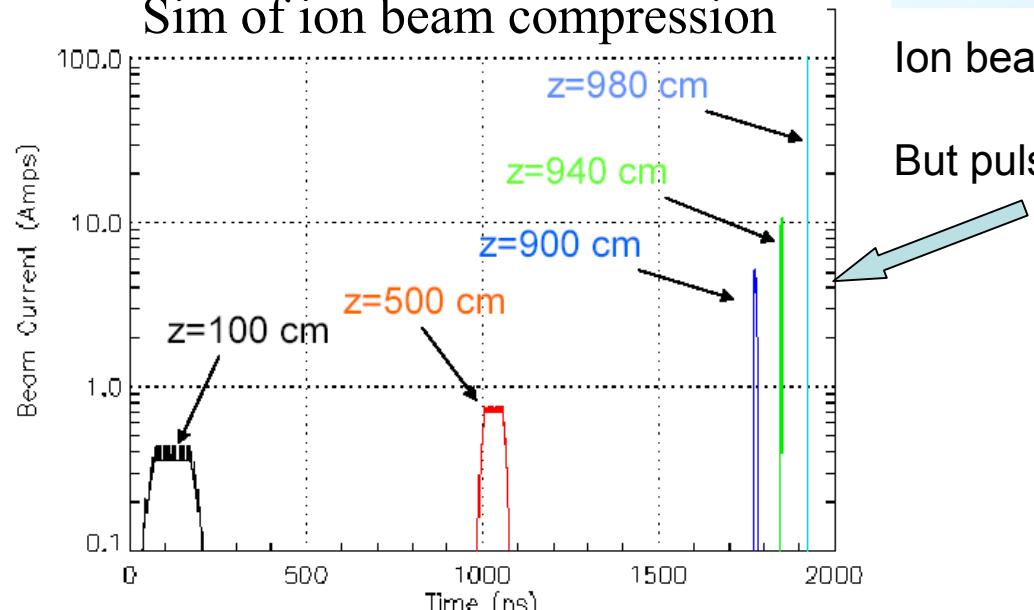


Heavy Ion High Energy Density Physics

What is the nature of “warm dense matter”?



Sim of ion beam compression



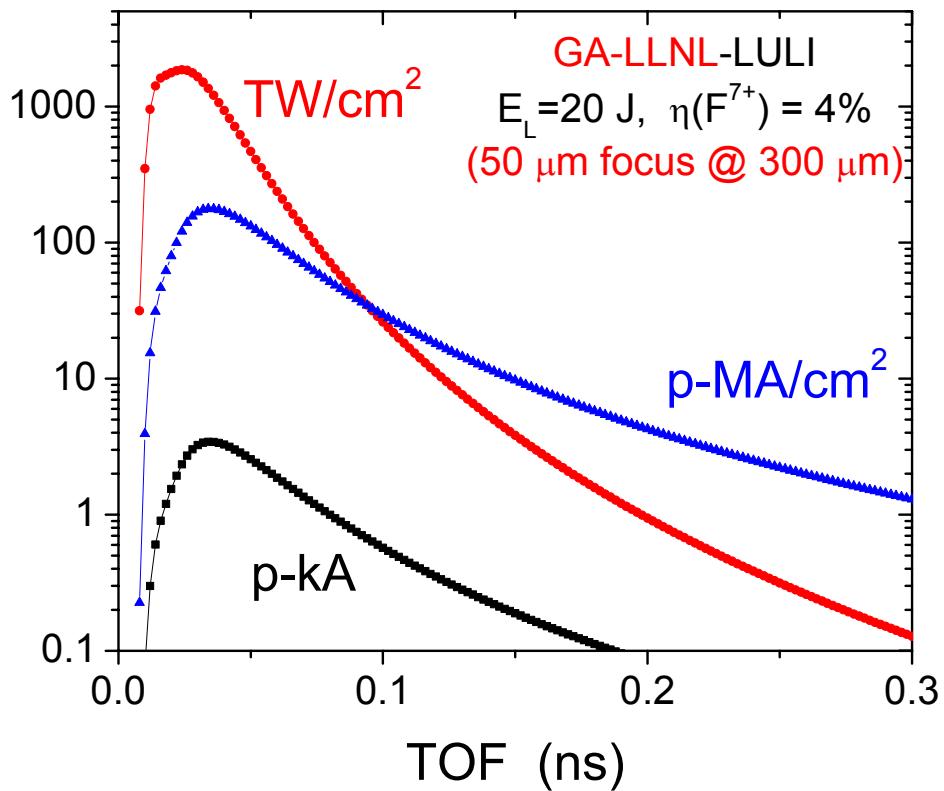
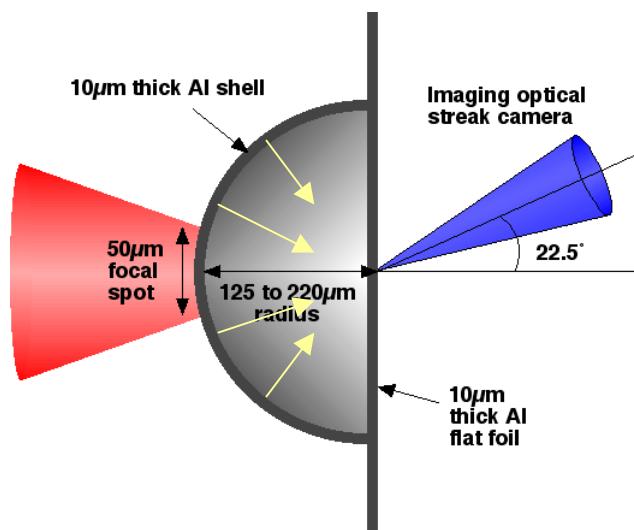
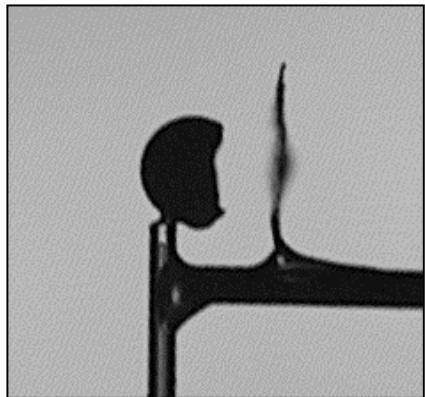
Ion beam range ideal for heating solid matter

But pulse must be short to beat expansion
requires neutralized transport

5 yr Goal: (10^{11} J/m^3)
 100 p-MA/cm^2 !
10 MeV
1 ns

Can laser-irradiated foils do the job? (T. Cowan)

$>10^8$ p-A/cm², and > 1000 TW/cm² on target



HED w/ Ion Beams Summary

5 yr Goal: (10^{11} J/m³)

100 p-MA/cm²

10 MeV

1 ns

Now w/ 20J lasers

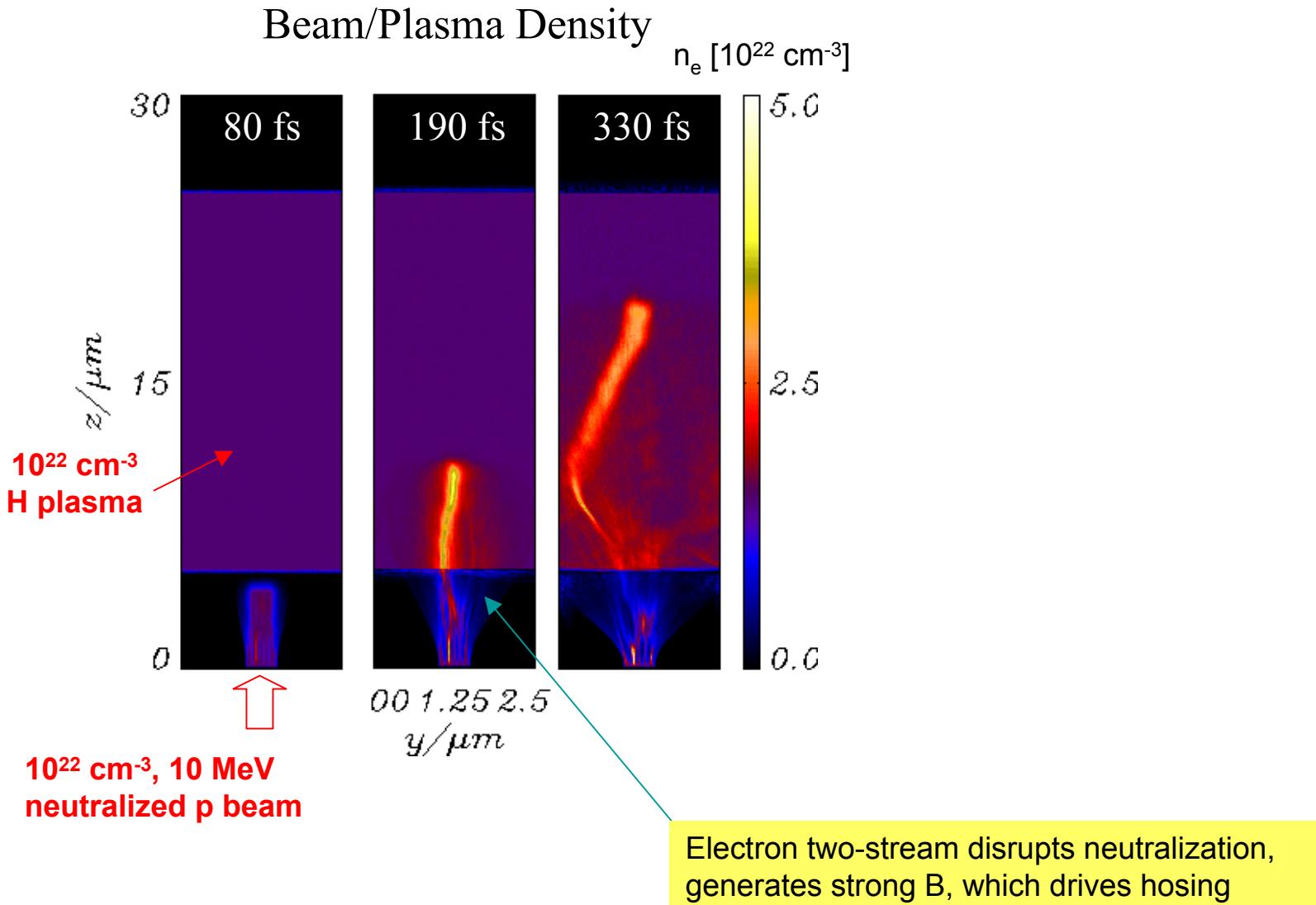
100 p-MA/cm²

10 MeV

10 ps

New 2kJ PW lasers (e.g., Sandia ZBL PW) should enable ion current densities relevant to HIF beams (Cowan)

Beam-plasma instabilities affecting neutralized beam transport might be studied with laser-foil beams



Highlights of Some Talks



Ultrafast 2-D radiative transport in a micron-scale aluminum plasma excited at relativistic intensity

Ben Bowes and M. C. Downer

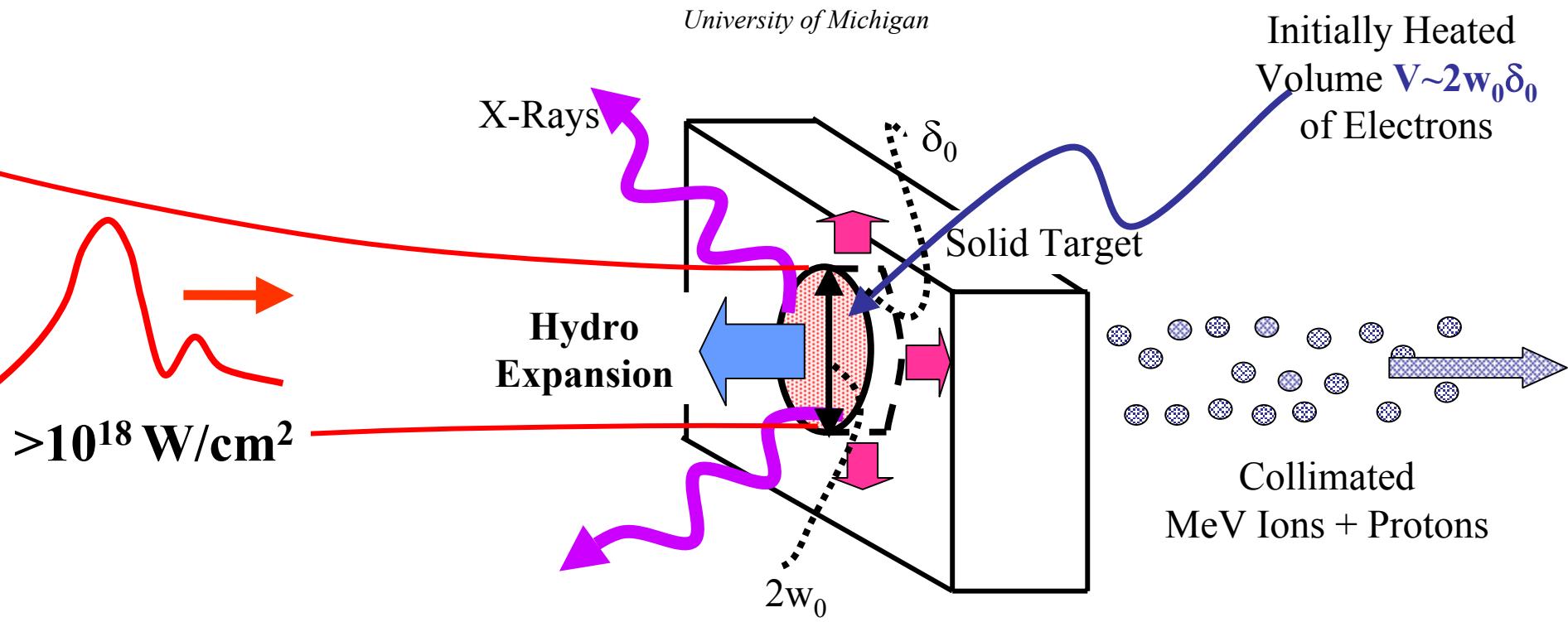
University of Texas at Austin

H. Langhoff

Physikalisches Institut der Universität Würzberg

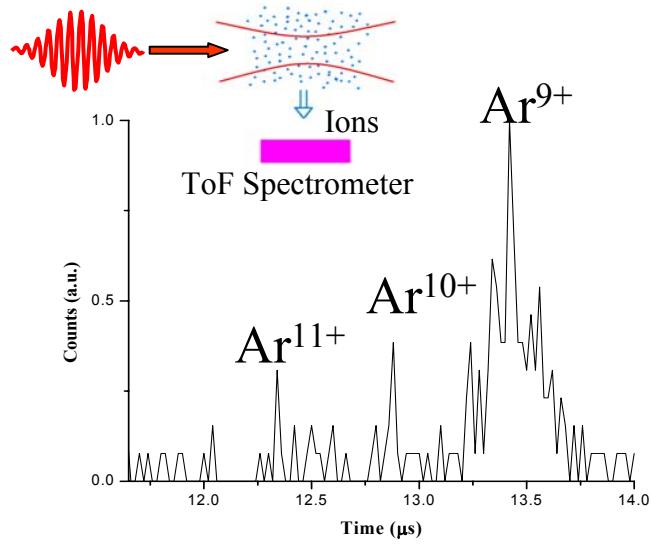
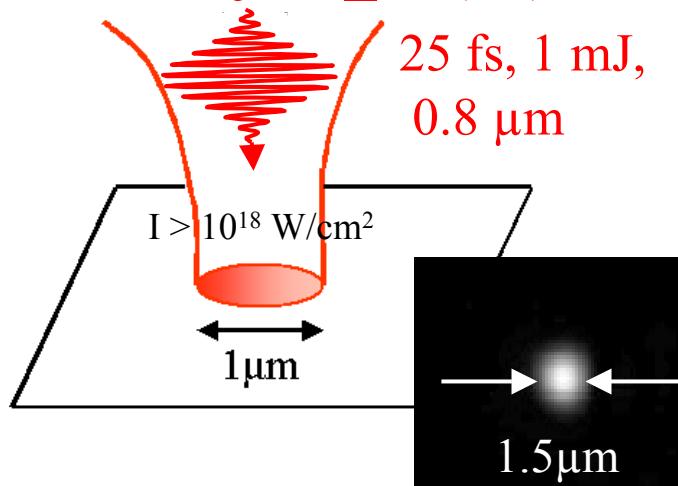
M. Wilcox, B. Hou, J. Nees and G. Mourou

University of Michigan



λ^3 laser system: relativistic intensity at 1 kHz

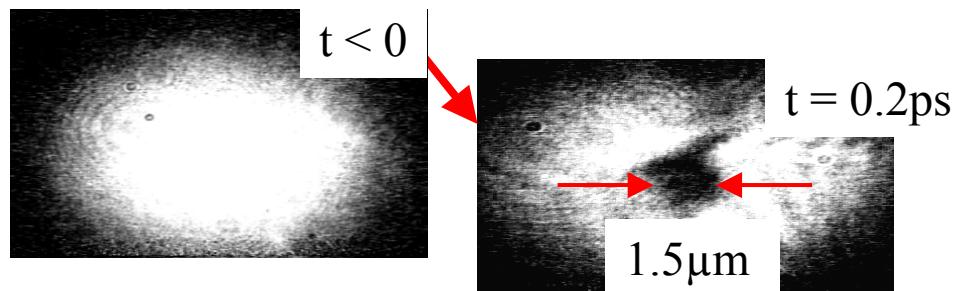
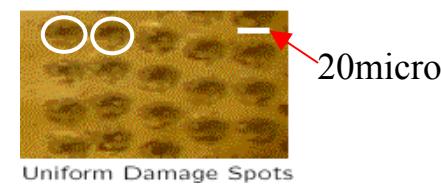
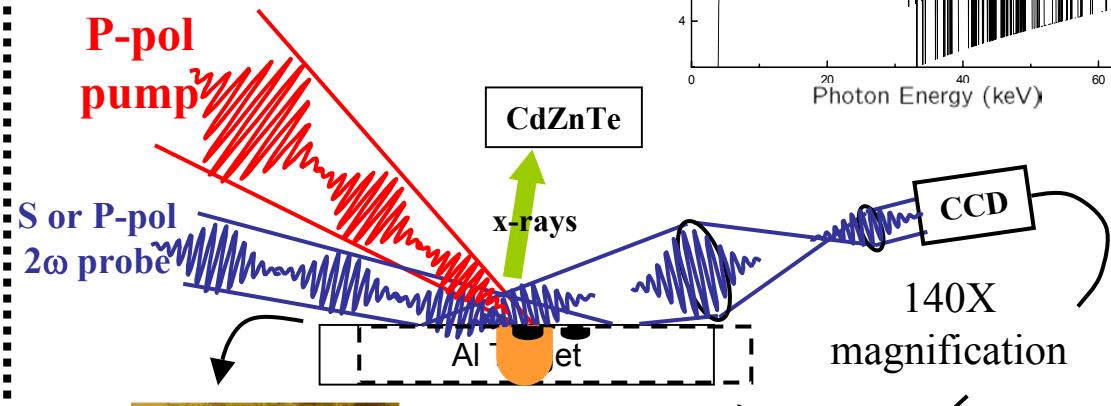
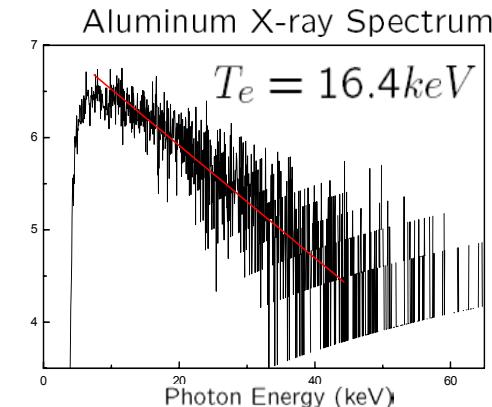
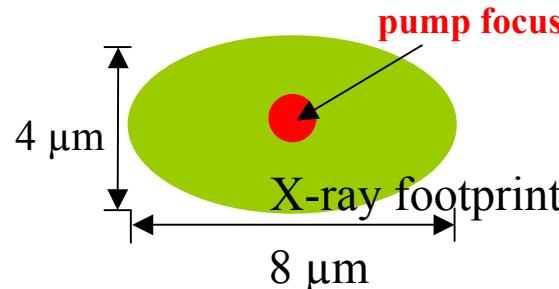
Albert, O.; Opt. Lett. **25** 1125 (2000)



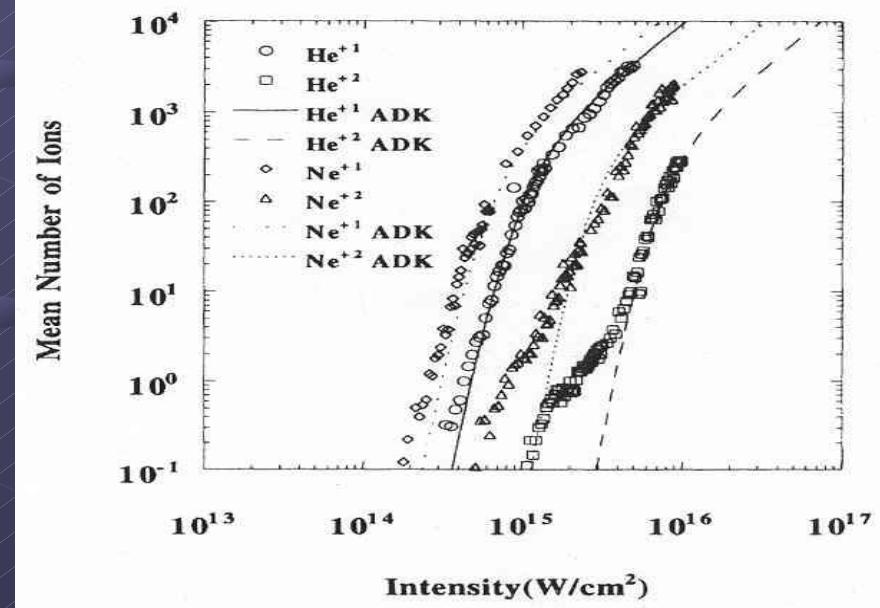
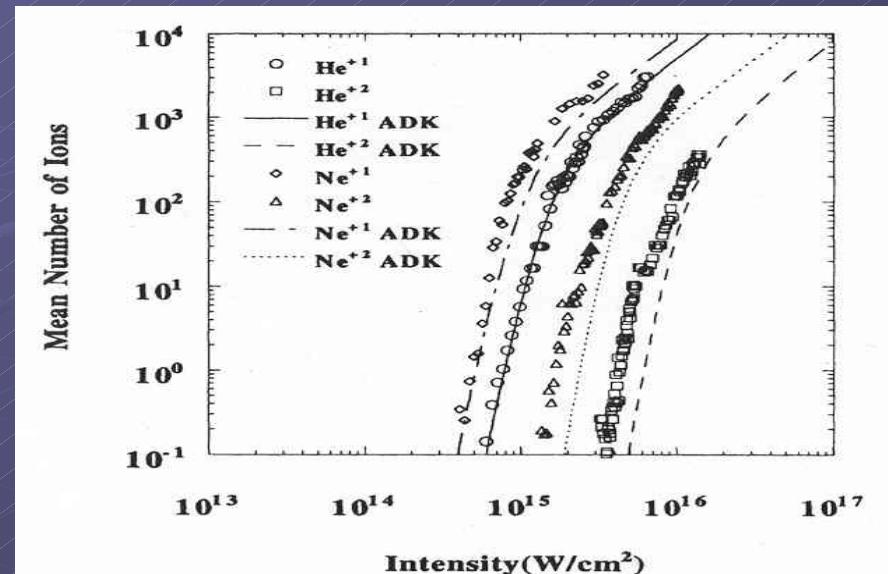
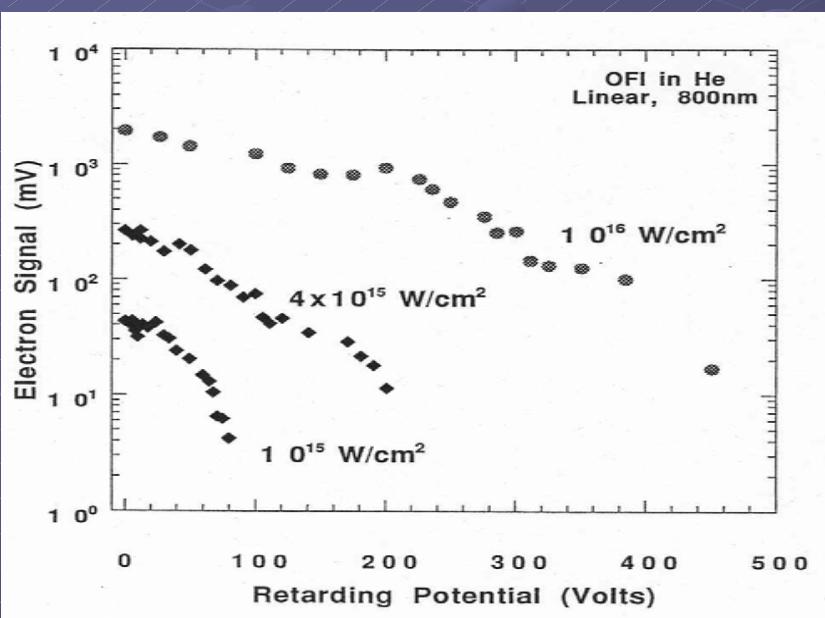
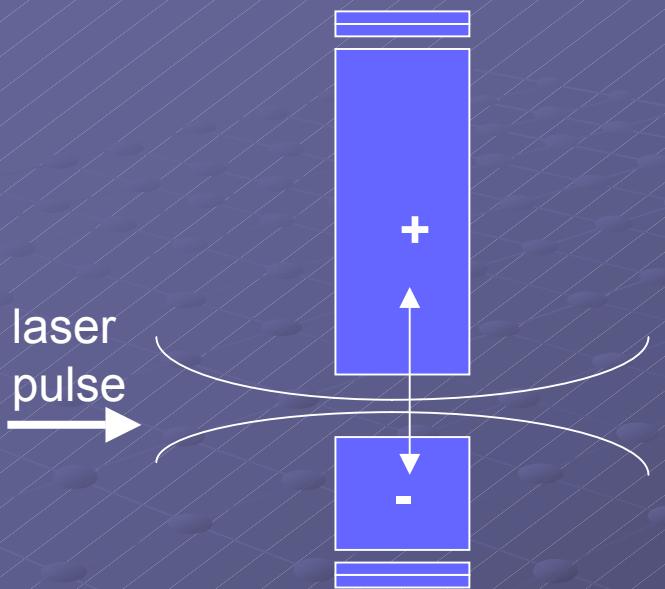
$$0.1\% \text{ Ionization Rate for } \text{Ar}^{11+} \\ \Rightarrow I = 3 \times 10^{18} \text{ W/cm}^2$$

S. Augst; JOSA-B **8** 858 (1991)

Experimental set-up for fs reflective imaging



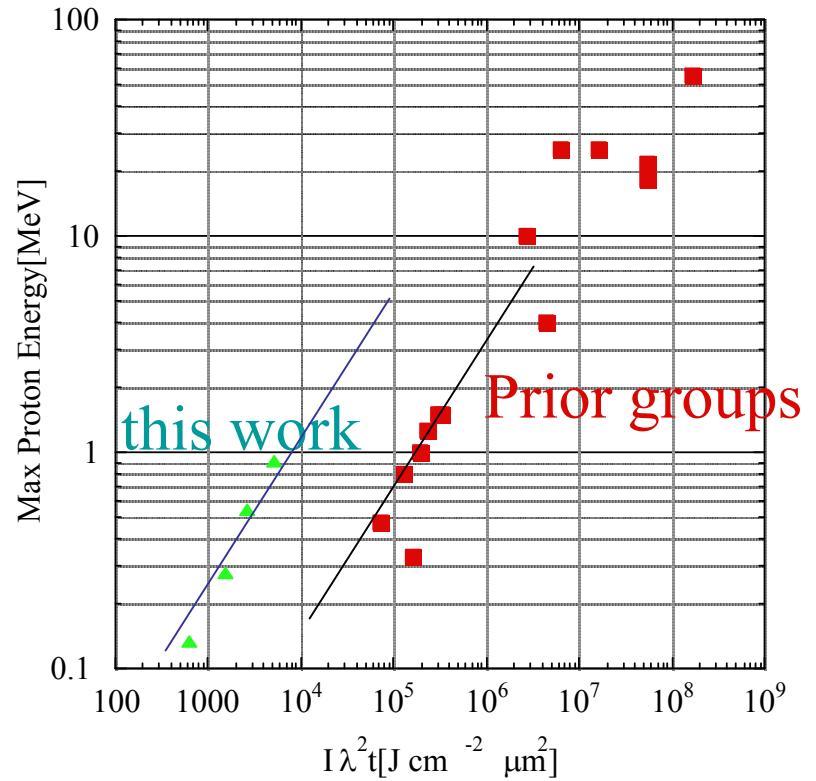
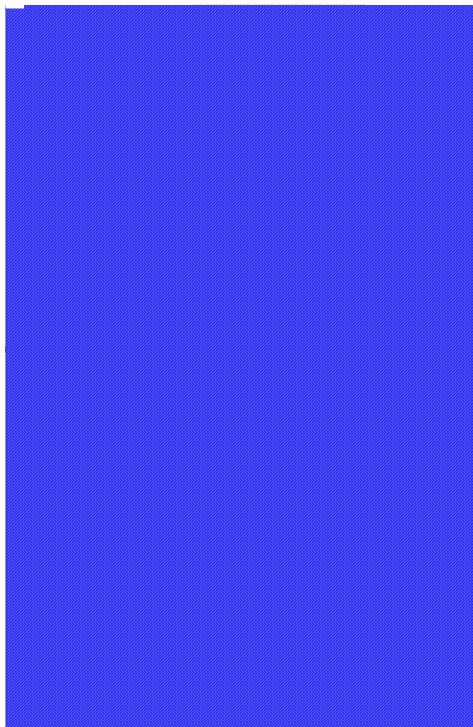
Multiple charge ions – P. Bolton



A. Ogata

Getting more acceleration with less laser

T³ laser with a 10⁻³ prepulse



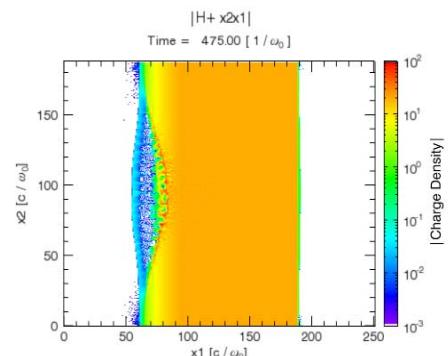
Shocks in laser-plasma interactions

Luís O. Silva

J. Davies, R. A. Fonseca, M. Marti, F. Peano
(IST/Portugal)

J. Fahlen, C. Ren, F. Tsung, W. B. Mori
(UCLA)

And **P. Messmer**, D. Bruhwiler, Tech-X
Shocks and THz generation



Acceleration mechanisms

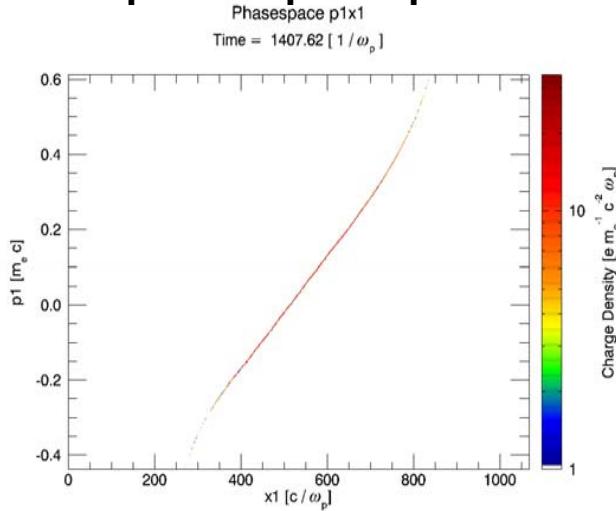
No shock

Shock with higher
energy than rear

Shock with lower
energy than rear

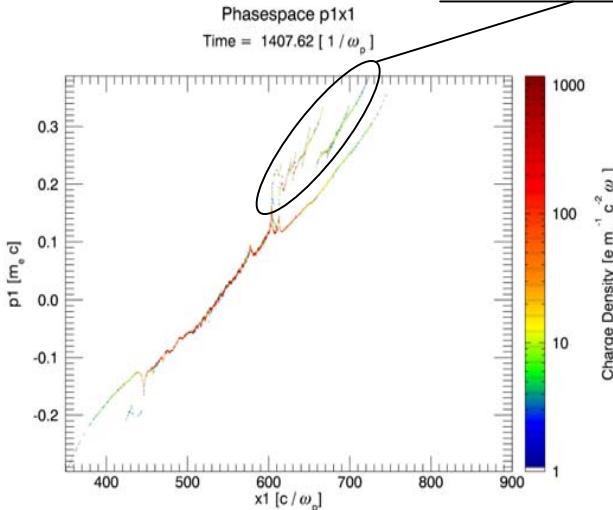
Diffusive/sheath acceleration

Ion phase space p1x1

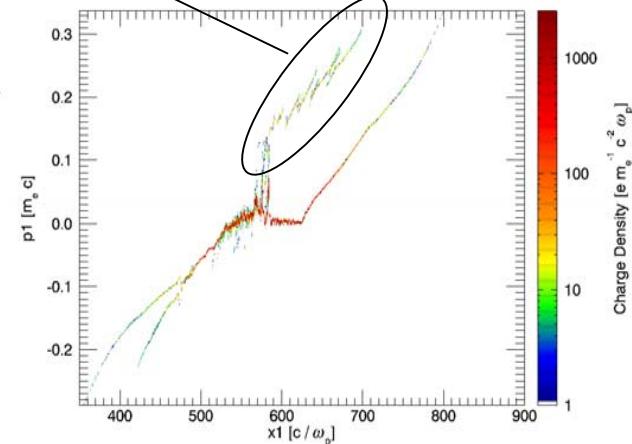


shock acceleration

Phasespace p1x1
Time = 1407.62 [1 / ω_p]



Phasespace p1x1
Time = 1407.62 [1 / ω_p]



$$a_0 = 16$$

$$\Delta = 1 \mu\text{m}$$

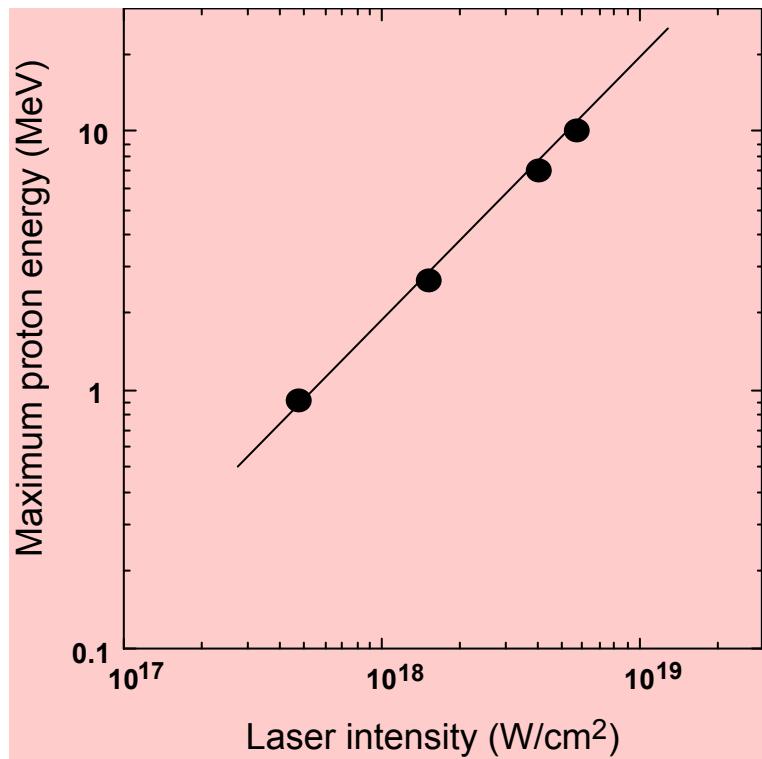
$$a_0 = 16$$

$$\Delta = 6.2 \mu\text{m}$$

$$a_0 = 16$$

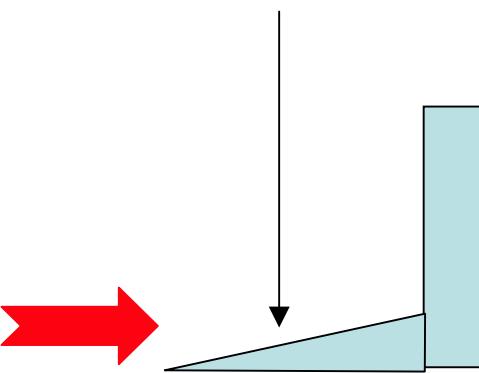
$$\Delta = 8.8 \mu\text{m}$$

CUOS Experiments Shows Maximum Proton Energy is Enhanced and Proportional to the Laser Intensity [3]



(1 ~10) TW Laser with intensity contrast ratio; $5 \cdot 10^{-5}$

Estimated preplasma scale is ≤ 10 microns.

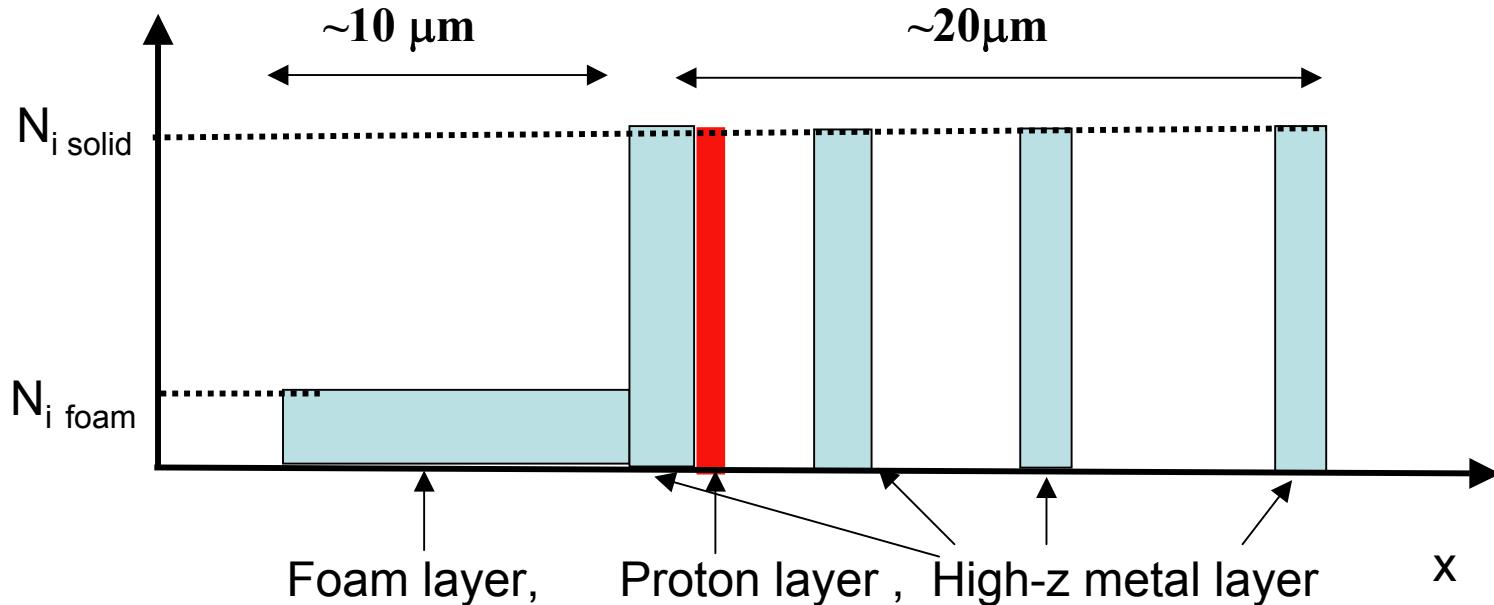


For prepulse level is 10^{-7} proton energy was much less.

[3] K. Nemoto, A. Maksimchuk, S. Banerjee *et al.*, Appl. Phys. Lett. **78**, 595 (2001).

[Concluding Remark]

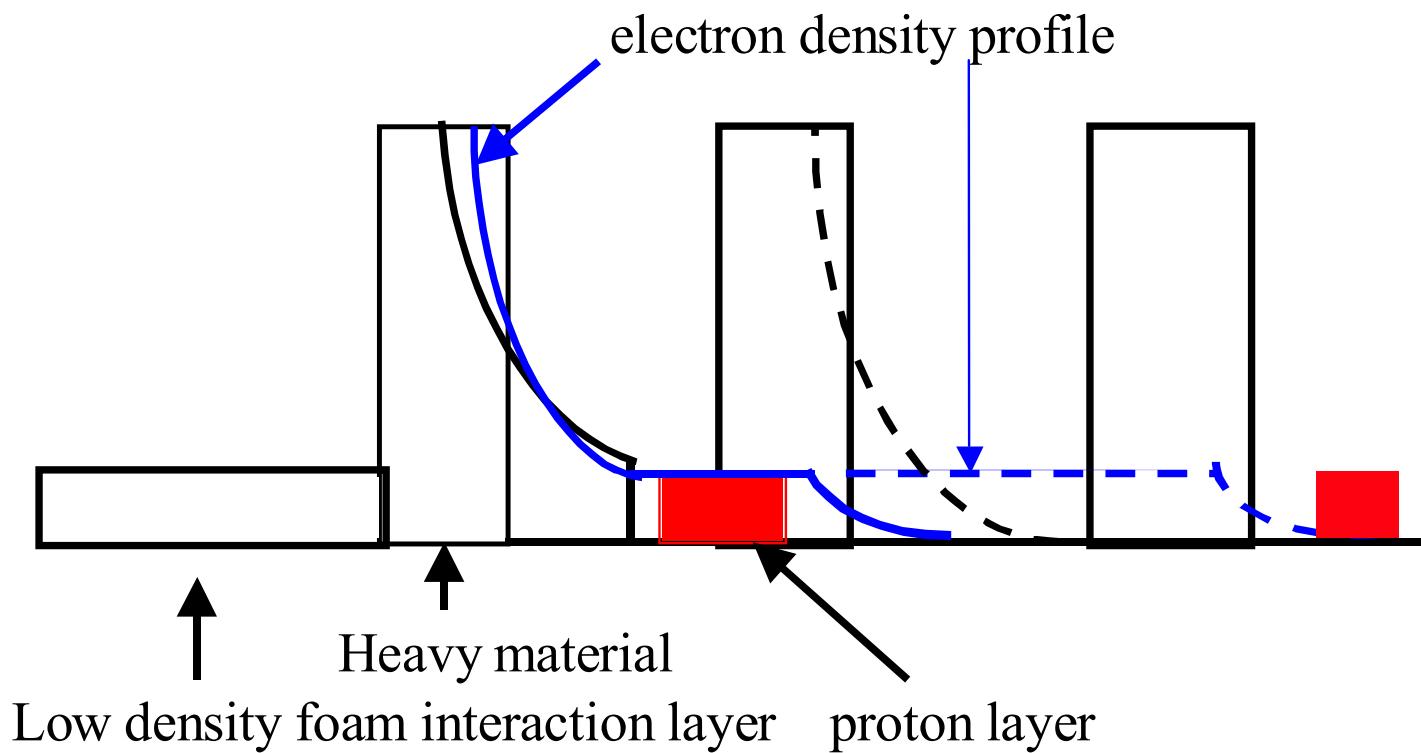
Target for high energy and monochromatic proton acceleration



$$Zn_{i, \text{foam}} < \gamma n_c , T_h \sim a^2 mc^2 / 2 \longrightarrow I_L > 10^{21} \text{ W/cm}^2$$

For 300 MeV proton

Multi-foil Expansion Processes



We heard some nice talks... (II)

- Vacuum Acceleration (Umstadter, Nakajima)
- Active Medium Acceleration (Shachter)
- Ferroelectrics (Kanareykin)
- Beam conditioning for FELs (Wurtele, Esarey, Schroeder)

Beam Conditioning for FELs

1. Jonathan Wurtele -- conv.
Approach to an exotic idea

LBNL/UCB

Co-workers: G. Penn, A. Wolski, A.
Sessler

2. E. Esarey et al. -- Plasma LWFA
approach
3. C. Schroeder et al.-- Thomson
Scattering laser approach

The Problem

FEL resonance: $\lambda = \left(1 - \frac{\bar{v}_z}{c}\right) \lambda_w \approx \frac{1+K^2}{2\gamma^2} \lambda_w$

$$\frac{\bar{v}_z^2}{c^2} = 1 - \frac{1}{\gamma^2} - \frac{\bar{v}_\perp^2}{c^2} = 1 - \frac{1+K^2}{\gamma^2}$$

Zero canonical momentum

$$\lambda = \frac{1}{2} \left(\frac{1+K^2}{\gamma^2} + \frac{2J_x}{\beta_x} + \frac{2J_y}{\beta_y} \right) \lambda_w$$

Typical particle

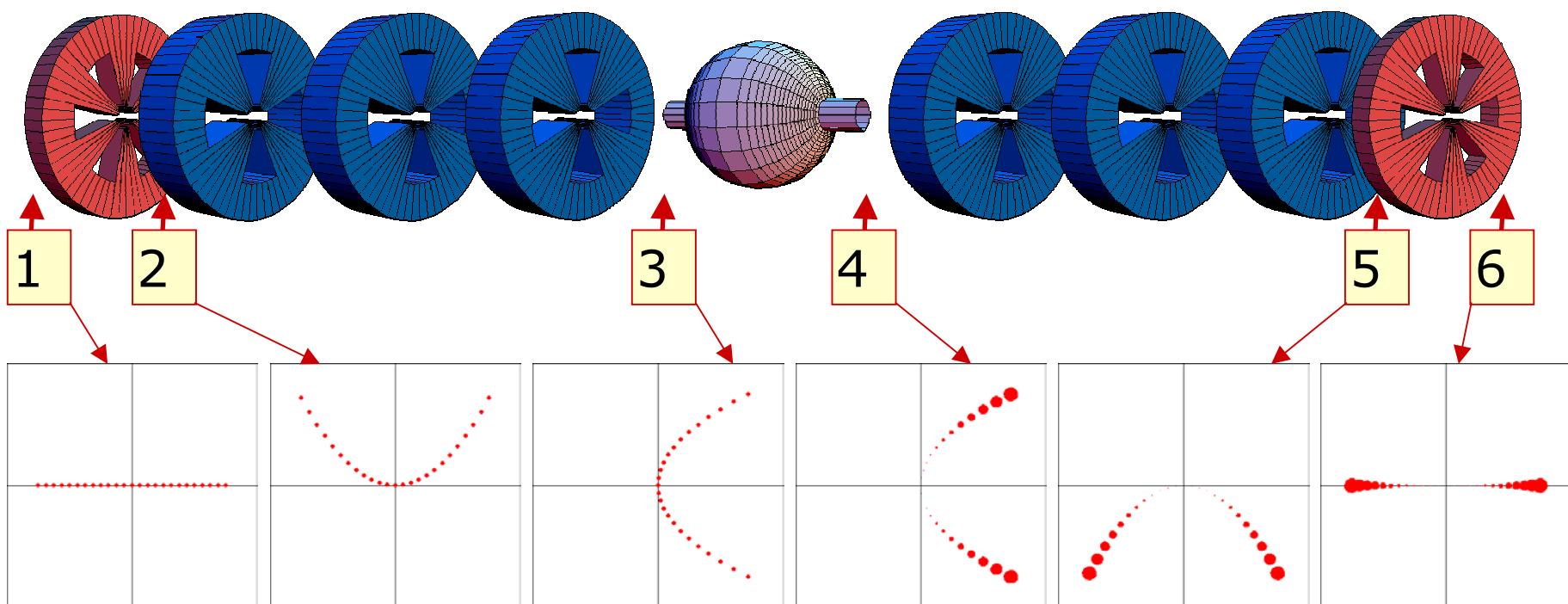
$$J_x = \frac{1}{2} \left[\frac{x^2}{\beta_x} + \beta_x \left(\frac{p_x}{P_0} + \frac{\alpha_x x}{\beta_x} \right)^2 \right]$$

Emittance is beam average of J

FEL requires:

- 1) Particle relative motion be less than one wavelength in a gain length a
- 2) Emittance be less than radiation wavelength (overlap).

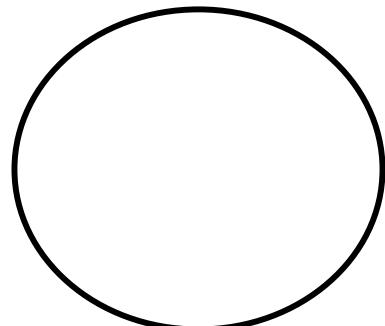
Example: Sextupole + TM_{110} cavity Conditioner



- 1 A “matched” beam enters the conditioner.
- 2 The first sextupole distorts the horizontal phase space.
- 3 The phase space is rotated through $\pi/2$.
- 4 The cavity gives a correlation between x and p_z .
- 5 The phase space is rotated through a further $\pi/2$.
- 6 The final sextupole removes the phase space distortion.

Works but weak E requires many stages

3 Alternate Approaches:



beam



profile of decelerator field

1. Esarey -- LWFA, good but wide laser => 700 TW
2. Schroeder -- Thomson laser => 300 TW
3. Working Group -- passive plasma (PWA deceleration)
0 TW!
but longitudinal dependence
 - use shaped bunch before compressing?
 - throw away all but small slice of bunch?

Vacuum Acceleration -- Umstadter, UM

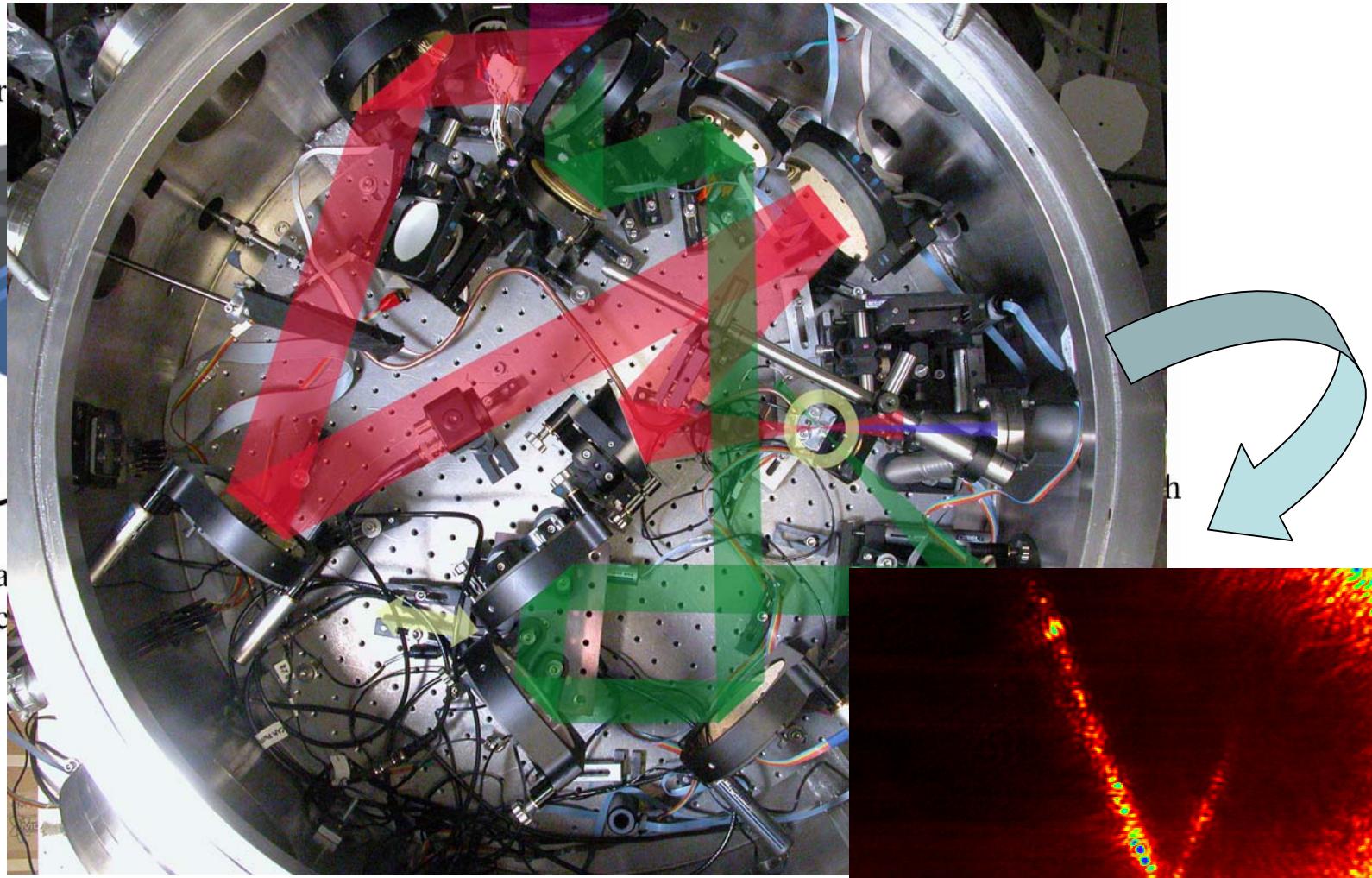
CCD Camera
Objective

Phosphor
Screen

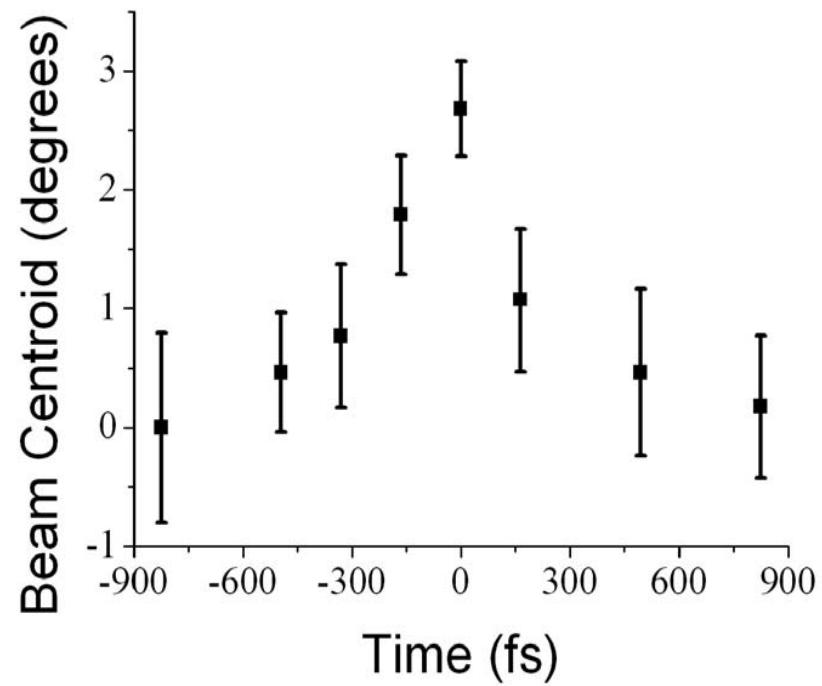
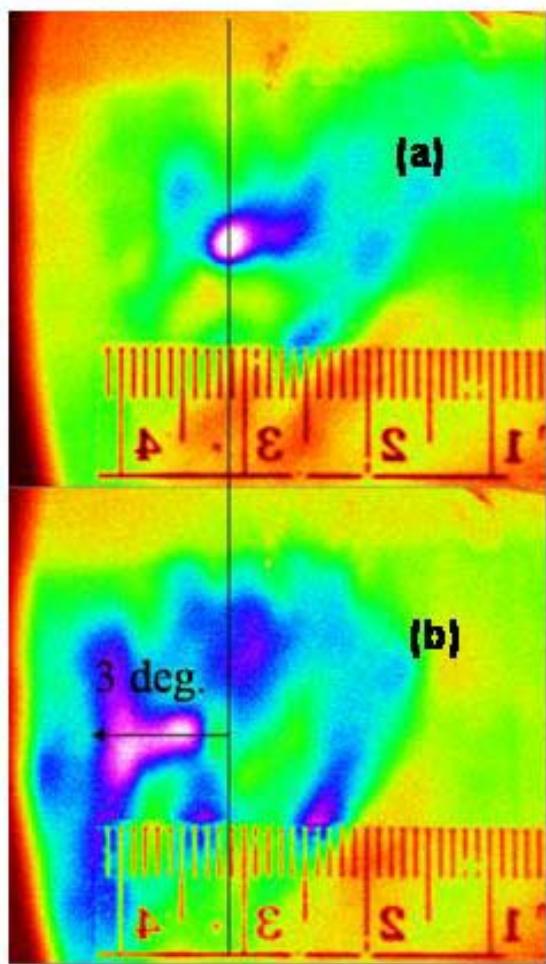
Micro
Channel Plate

Grating

Seydel
Spec

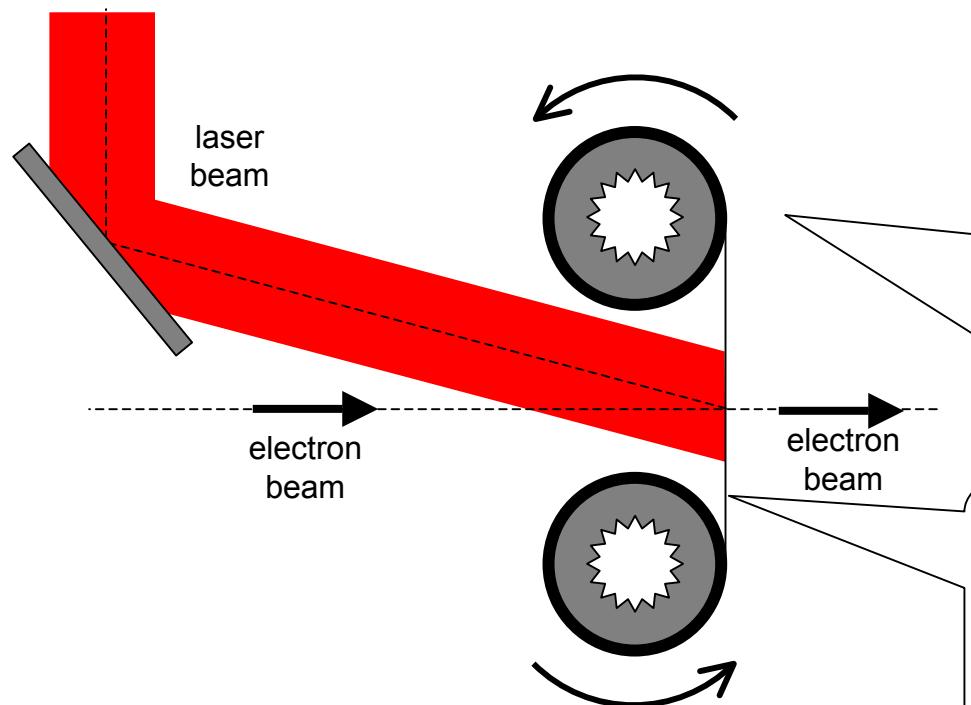


Ponderomotive deflection of electron beam



Cross-correlation of laser-wakefield
accelerated electron beams

The "Walkman" particle accelerator -- LEAP, Plettner



Sony SRF-86 Sports
AM/FM Radio Walkman



Properties of the tape

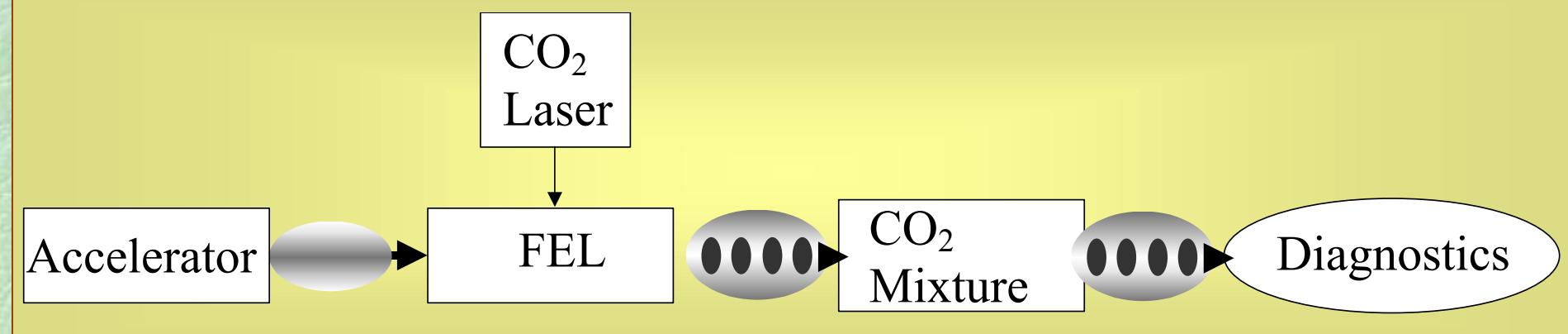
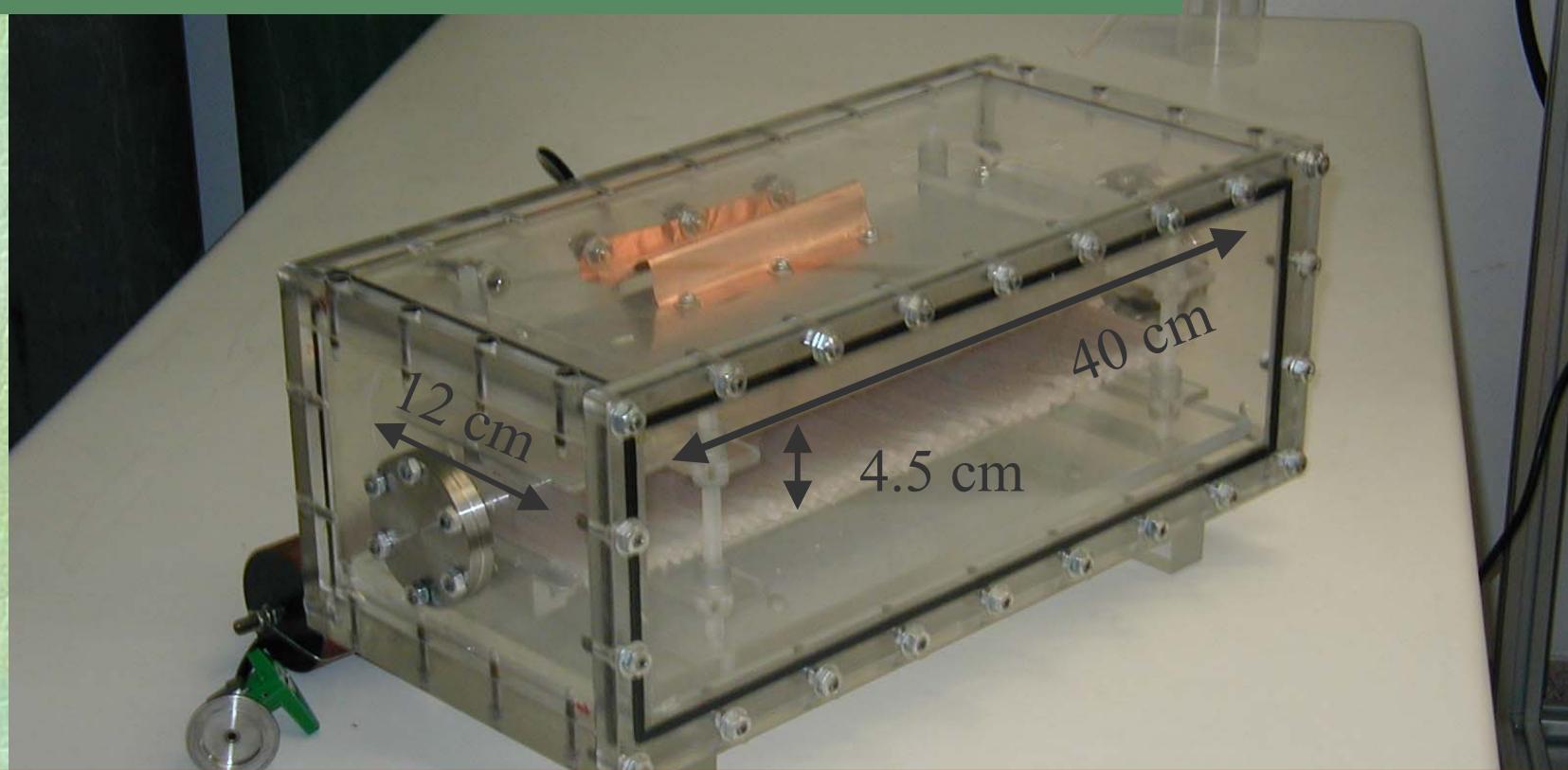
1 μm gold
8 μm kapton

Key modifications

1. Slit is lost (e^- traverses matter)
2. One laser beam focused on the surface of the tape
3. Damage threshold constraint is removed

New problems

1. e^- traverses matter
2. potential contamination from laser-ablated material
3. Potential interference of the acceleration effect due to the presence of plasma



Exotic Concepts Conclusion

“As our field matures, there is still room for innovation in all AAC areas: drivers, materials, combinations thereof.”

G. Shvets

“We came, we saw, we calculated.”

Anonymous

“I proclaim these the best games ever.”

Juan Antonio Samaranch

Thank you Ilan and Brookhaven National Lab!