



Superconducting Magnet Division

Magnet Note

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Date: September 19, 2003
Topic No: 631-39 (AM-MD-331)
Topic: Superconductor R&D
Title: Critical Current Measurements of Bi-2212 Wires

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Critical Current Measurements of Bi-2212 wires

Arup K. Ghosh and W. B. Sampson

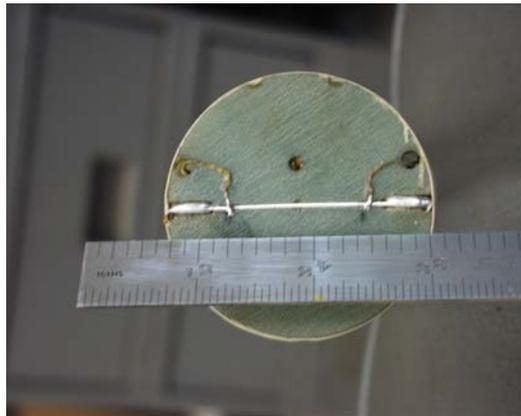
Abstract

Short lengths of Bi-2212 wires manufactured by SHOWA and OST were measured at 4.2K in fields up to 8T. The data have been analyzed in terms of an exponential field function at low fields and pinning-force functions at higher fields. In addition strands extracted from a 30-strand SHOWA cable were also measured. These are compared to the critical currents measured for the cable samples at BNL's Short Sample Test facility. A 0.8mm diameter strand from OST showed the highest J_c of 4350 A/mm² at zero-field and ~ 1900 A/mm² at 8T. The high J_c wires tend to be O₂-overdoped and have T_c 's ~ 80 K.

Introduction

The critical currents of HTS conductors are usually quoted for an electric field gradient, $E_c=1.0$ μ V/cm. This would seem appropriate at temperatures where the currents are small. However, at 4.2K, the currents for wires or tapes with cross-sectional area ~ 0.5 - 0.6 mm² can be several hundred amps so that the resistive transitions should be measured at a smaller electric field gradient or effective resistivity. In fact HTS conductors seem to have V-I characteristics similar to LTS wires of low n -value. Multifilamentary Bi-2212 is of particular interest for magnet builders as it can be processed as a wire, and Rutherford cables can be fabricated for making magnets using the "react-and-wind" method. At present, straight wire testing at 4.2K is limited lengths of a few cm since the solenoid bore is only 60mm, so that the critical current I_c is defined at $E_c=1.0$ μ V/cm. In the future testing of wires using an appropriate barrel needs to be developed so that longer sample lengths can be used to measure the transition at much lower electric fields.

Tests of three wire samples are described here. One from Showa is a 1.0mm strand (B1152DE28) with Ag/SC ratio of ~ 2.8 . The others are a 0.82mm and 0.72mm strand (PMM030224) from OST with a Ag/SC ratio of ~ 2.5 . Typical sample lengths were ~ 5 cm and voltage-tap separation was ~ 30 mm. The picture below shows the arrangement.



Results

Figure 1 shows the $E-I$ measurements for B1152 wire for different fields applied perpendicular to the wire. A log-log plot is shown to display the characteristic power-law behaviour seen in LTS wires. However, the $E-I$ curve often does not have a unique n -value. For analysis the n -value is averaged for electric fields between 0.3 and 3.0 $\mu\text{V}/\text{cm}$. Figure 2 summarizes the critical current and n as a function of field. Note that J_c initially falls rapidly with field and then gradually decreases with field.

Figures 3 and 4 show the critical current and n -value for the 0.82mm and 0.72mm wires from OST. Although the J_c of these three strands vary, the J_c behaviour as a function of field is very similar. This can be seen in Figure 5, where the normalized critical current $J_c(H)/J_c(1\text{T})$ is plotted versus the applied field. The similar J_c - H dependence would imply that the basic material properties at these fields are similar for these wires. Indeed if one analyzes the J_c in the two field regimes, $H < 1\text{T}$ and $H > 1\text{T}$, it appears that for $H < 1\text{T}$, the dependence is exponential in H ,

$$J_c = J_{c0} \exp^{-bH} \quad (1)$$

and above 1T, the data fit a power law

$$J_c = J_{c1} H^{-m} \quad (2)$$

An example of the fit is shown in Figure 6, and the fitting parameters for the three wires are summarized in Table 1.

Table 1

Wire ID	No. of Fils	Diameter	Jc (1T)	Jc (5T)	H < 1T	H < 1T	H > 1T	H > 1T
					Jc0	b	Jc1	m
		mm	A/mm^2	A/mm^2	A/mm^2	T	A/mm^2	
B1152	127x7	1.00	2617	1863	3864	2.5	2572	0.21
PMM030224	85x7	0.82	3076	2080	4410	2.7	3052	0.23
PMM030224	85x7	0.72	2738	1872	3980	2.6	2740	0.24

The low field exponential behaviour suggests the possibility of weak link coupling between the filaments or between the grains. In this region the n -value also decreases rapidly. It would be interesting to compare the magnetization with the transport current at low fields. At fields $> 1\text{T}$, it is common in the literature to show that the critical current density of Bi-2212 wires varies exponentially with field [1, 2]. The data shown here can also be fit to a form of Eq.(1) with $(l/b) \sim 20\text{T}$. However in this limited field region at 4.2K, one would have to measure to high fields $\sim 20\text{T}$ to clarify which is the better functional dependence. Behaviour similar to Eq. (2) has been reported [1], where at 4.2K J_c has been found to be proportional to $H^{-0.19}$. In the next section we shall examine the pinning force where other considerations suggest a different dependence on field.

Pinning Force

The pinning force calculated for these wires are shown in Figure 7. In the range of $0.1\text{T} < H < 8\text{T}$, F_p is proportional to $H^{0.8}$. This field dependence can only be a limiting form of the full field

behaviour, otherwise this would be unphysical. In analogy with metallic superconductors, the pinning force must have a maximum and can be discussed in the context of a scaling law of the form:

$$F_p(H, T) = C(T)h^p(1-h)^q \quad (3)$$

where C is a proportionality constant which is temperature dependent, and $h=H/H_{c2}$. For HTS H_{c2} is replaced by the irreversibility field H^* . For NbTi, $p=q=1$, and for Nb₃Sn the Kramer-law [3] prediction is $p=1/2$, $q=2$. When $h \ll 1$ (which is the case for HTS at 4.2K), the observed J_c behaviour would imply that $p \sim 0.8$. However, From a study of Bi-2212 wires at different temperatures, R. Wesche [1] has shown that Eq. (3) does fit the observed data using the following values for the exponents, $p=0.9$ and q of either 3 or 5. For low temperatures $T < 10K$, the much smaller accessible values of h make it difficult to predict the true value of H^* . To study Eq.(3) at 4.2K, I_c measurements have to be extended to the range of 50T as the peak in F_p is indicated to be at $h \sim 0.2$. In Figure 8 the J_c for the SHOWA wire is shown with the various fits to J_c , all of which describe the J_c well in the range of applied field used for the present measurements. Some recent measurements of I_c from OST taken at 4.2K and to fields of 25T seems to have the best fit to $p=0.8$ and $q=3.0$ with $H^* \sim 400T$!

Extracted Strands

In an attempt to compare the critical current measured for a cable to the strands in the cable, individual strands were carefully separated from a section of the cable **HTS-S-R007B147-1**. The straight section of the wire between the two edges of the cable was used to measure I_c as a function of field at 4.2K. 5 such samples were tested and the I_c at 1T was found to vary from 150 to 382A. The large variation of I_c in the strands might be due to some degradation in I_c sustained during the extraction process. The strands are fairly well sintered in the cable making the extraction procedure difficult. Cable measurements when corrected for self-field give an average strand I_c of 360A. Although the critical currents vary significantly, the normalized I_c behaviour is very similar as shown in Figure 9. Details of the cable measurements appear in the appendix.

Summary

Although the wires tested were from two different manufacturers with their own heat-treatments, the J_c behaviour with field is very similar at 4.2K. Also these wires have a low $T_c \sim 80K$ and negligible I_c at 77K. The low T_c is indicative of wires which are O₂-overdoped. In contrast the wires extracted from the cable sample have T_c 's $\sim 90K$ and have measurable critical currents at 77K, and lower J_c at 4.2K

Below 1T, the J_c dependence on field is exponential, indicative of weak-link coupling of the grains. Decoupling fields are $\sim 2.5T$. In the field range studied, the J_c above 1T seems to fit a power law of $H^{0.8}$. This however is only a limiting form of the general pinning-force function. Measurements to fields $\sim 50T$ are necessary to clearly estimate the upper critical field at 4.2K

Acknowledgements

We thank Dr. T. Hasegawa of Showa Electric Wire and Cable Co., and Dr. K. Marken of Oxford Superconducting Technology for providing the heat-treated samples for test.

References

1. R. Wesche, "Temperature dependence of critical currents in superconducting Bi-2212/Ag wires", *Physica C*, 246 (1995) 186-194.
2. C.M. Friend, J. Tenbrink, D.P. Hampshire, "Critical current density of Bi₂Sr₂Ca₁Cu₂O_d monocoire and multifilamentary wires from 4.2 K up to T_c in high magnetic fields", *Physica C*, 248 (1995) 213-221.
3. E.J. Kramer, *J. Appl. Phys.* 44 (1993) 1360.

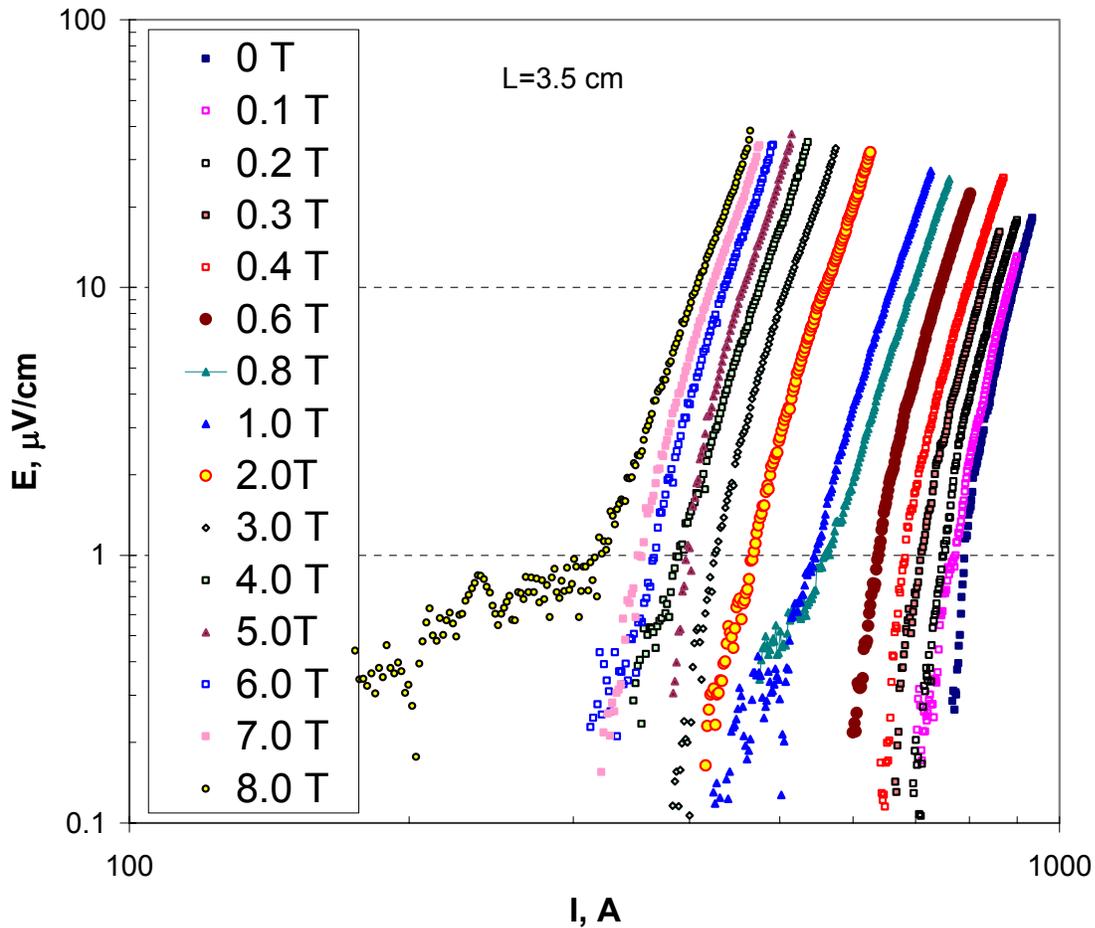


Figure 1 E-I curves for B1152DE28 strand at 4.2K and in applied fields to 8T.

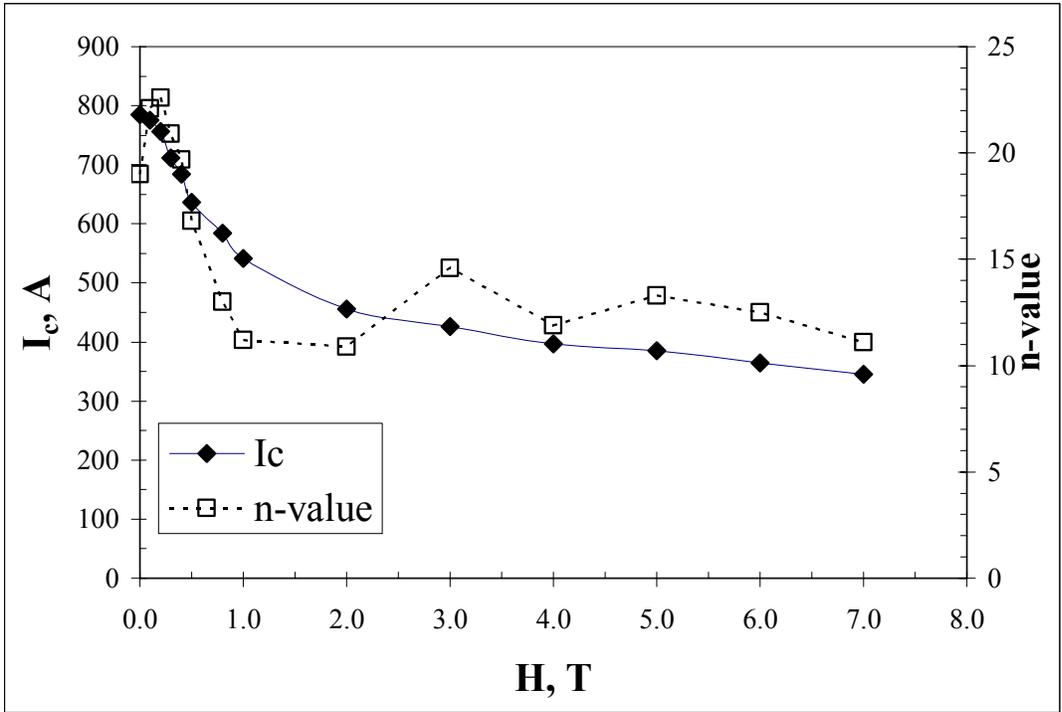


Figure 2 I_c and n-value for B1152DE28 strand from Showa

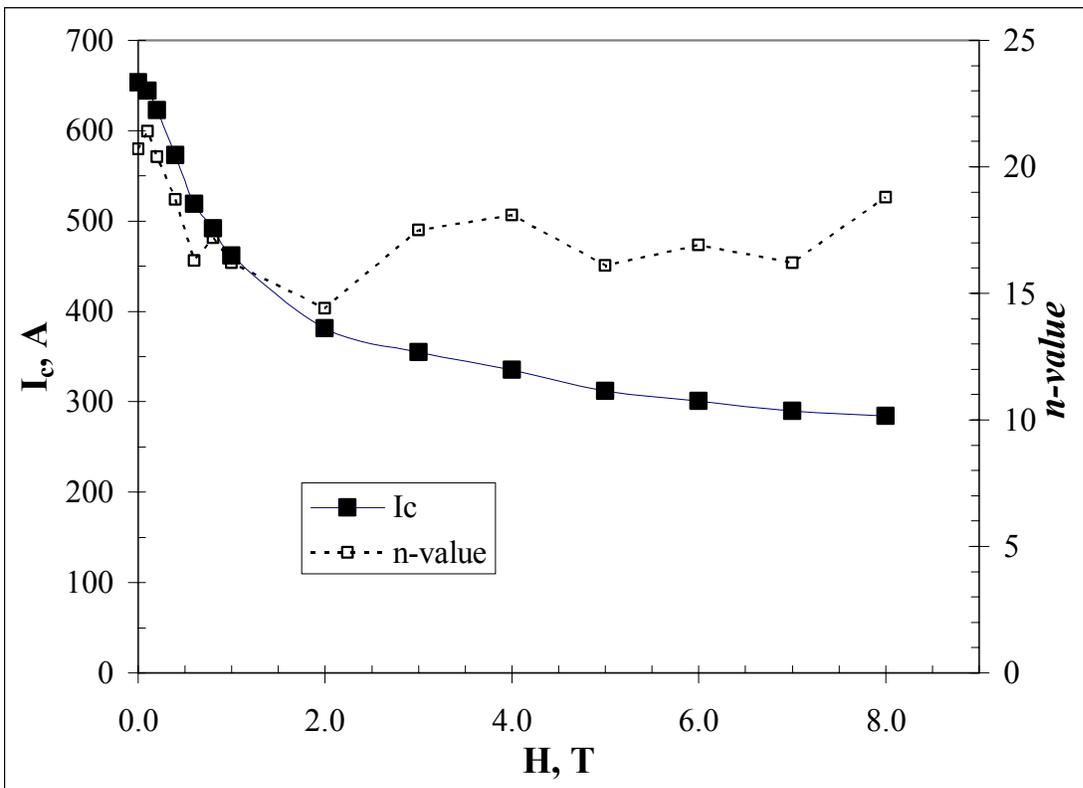


Figure 3 I_c and n-value for the 0.82 mm strand PMM030224 from OST

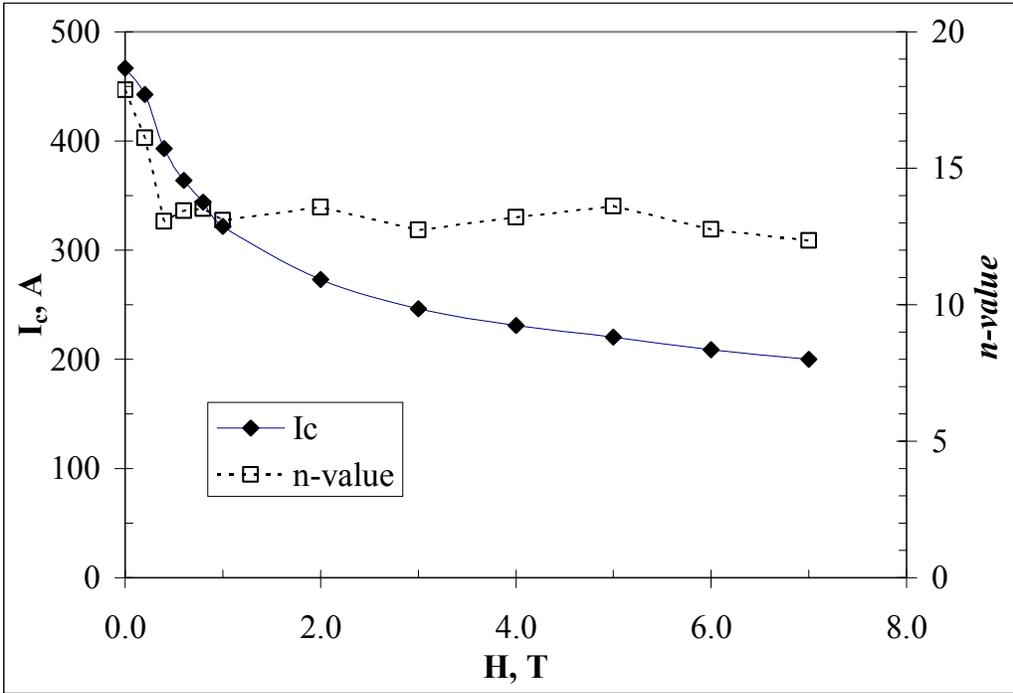


Figure 4 I_c and n-value for 0.72mm strand (PMM030224) from OST.

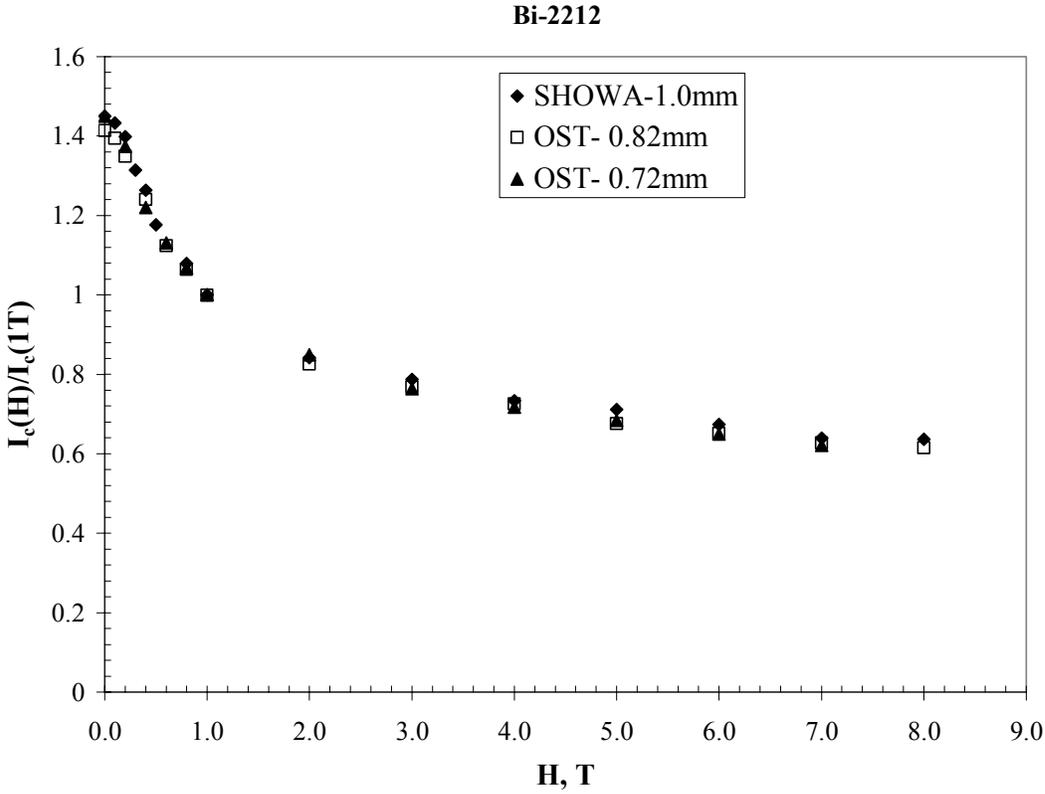


Figure 5 Normalized critical current as a function of field.

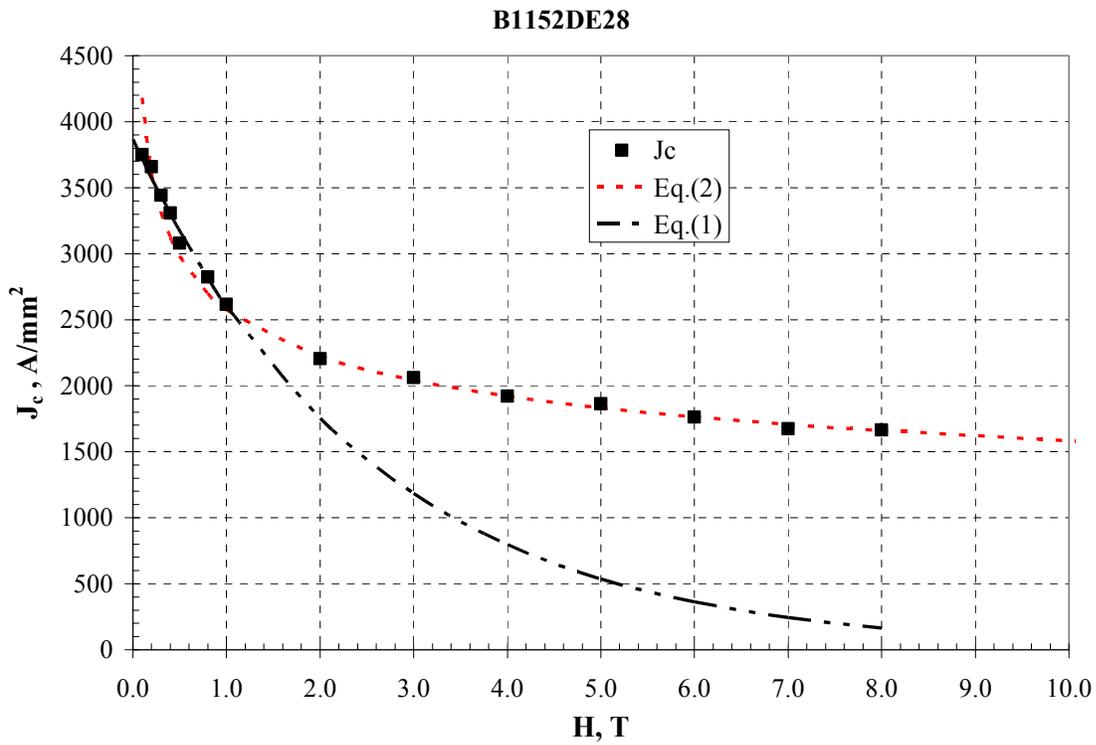


Figure 6. J_c vs. H for the SHOWA wire. The fit to Eq.(1) for $H < 1T$ and Eq.(2) for $H > 1T$ are also shown.

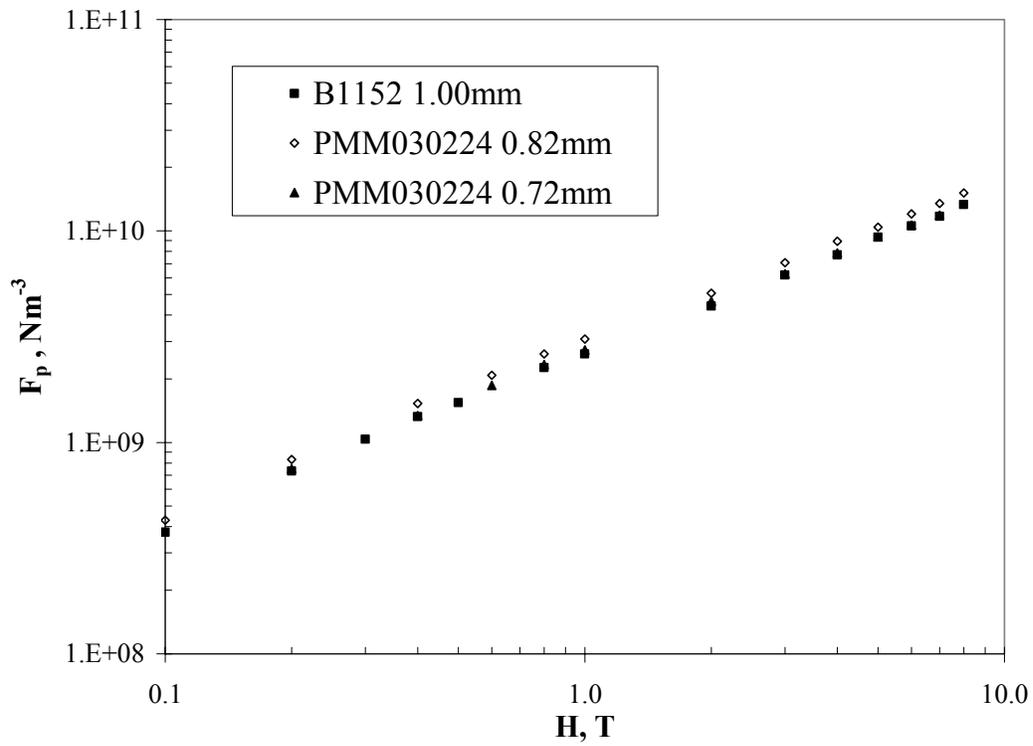


Figure 7. F_p as a function of H in the field range $0.1 < H < 8.0$ T

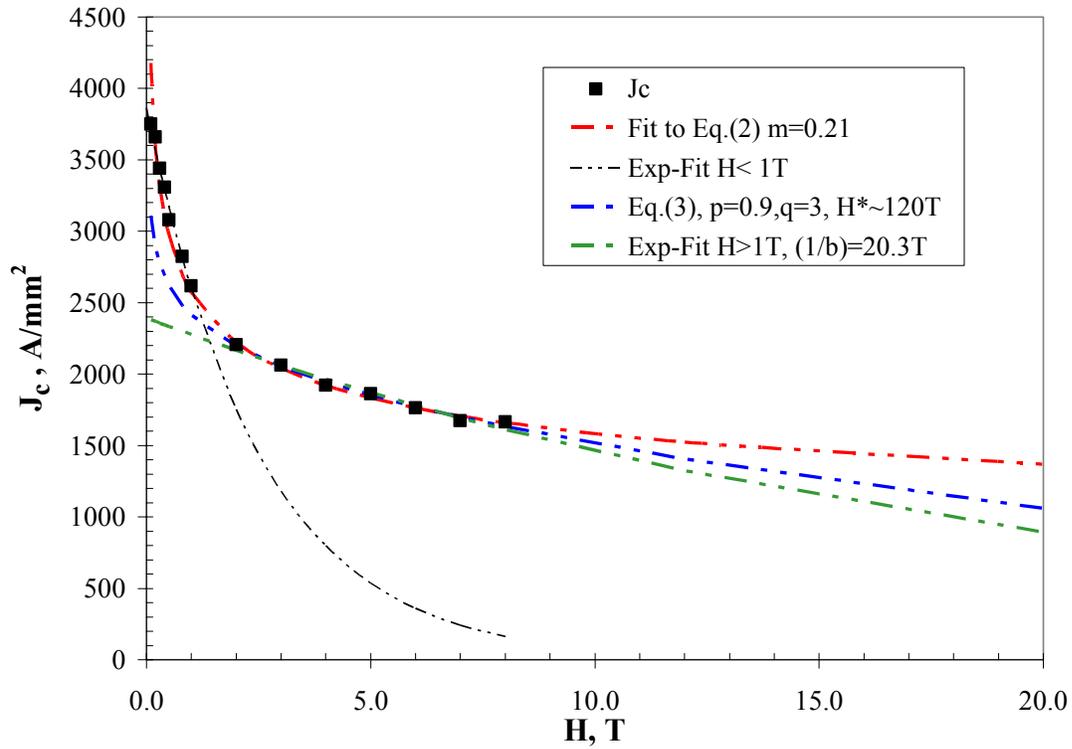


Figure 8. J_c - H for BII52 wire and the various fits are shown.

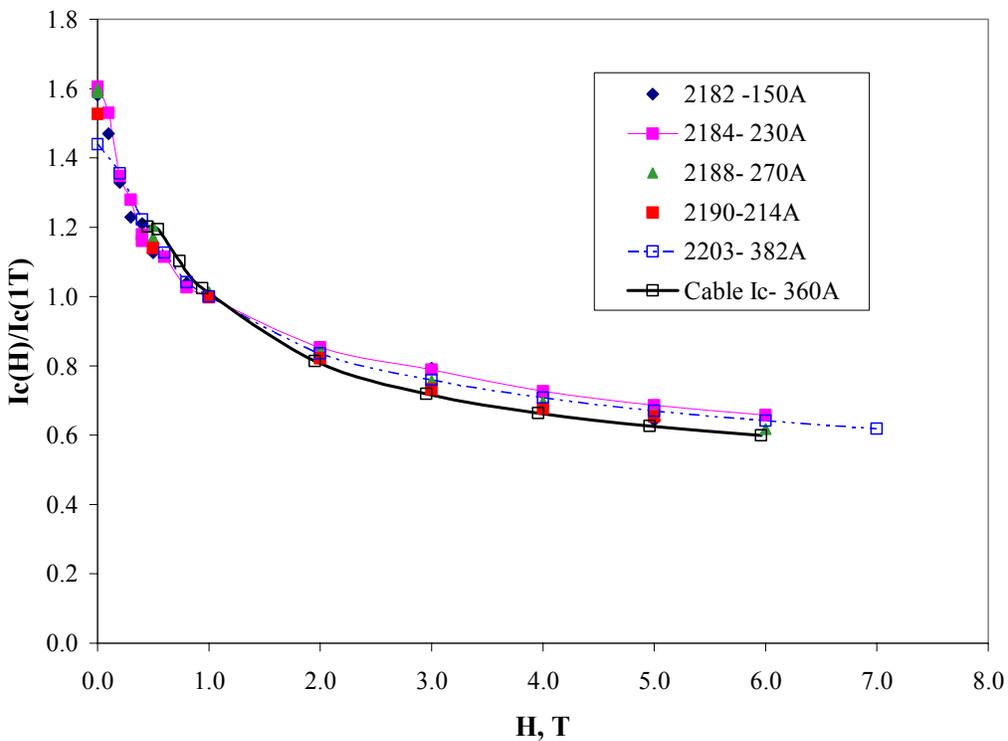


Figure 9 Normalized I_c as a function of field for strands extracted from cable HTS-S-R007B147-1

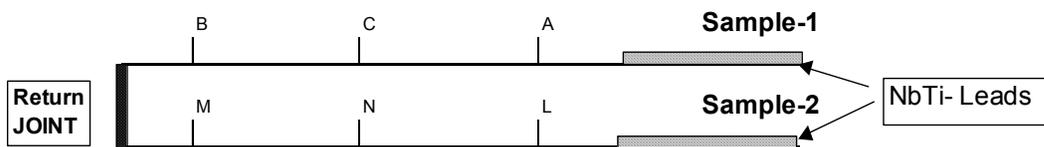
Appendix
Test Report for Cool-down # 4285

Cable Description

Strand Diameter 0.81 mm (nominal)
 No. of Strands 30
 Cable width 13.2 mm
 Cable thickness 1.52 mm
 No Keystone

The two lengths received Apr-3-03 are identified in the database as HTS-1-S-R007B147-1 and HTS-1-S-R013B147-1.

Sample configuration is shown below:

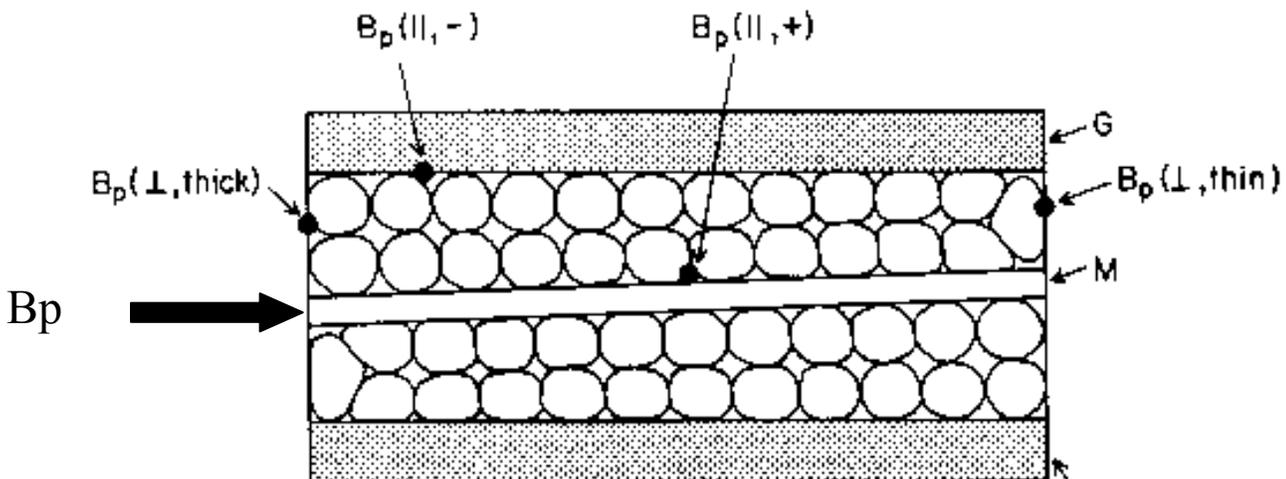


Voltage tap separation is as follows:

AC 300 mm
 BC 300 mm
 MN 300 mm
 LN 300 mm

Section AB is HTS-S-R007B147-1
 Section LM is HTS-S-R013B147-1

The samples were assembled into the test fixture with a transverse compression of 6 MPa. The test fixture was aligned in the dipole magnet with the field direction parallel to the flat face of the cable. A schematic of a cable pair is shown below:



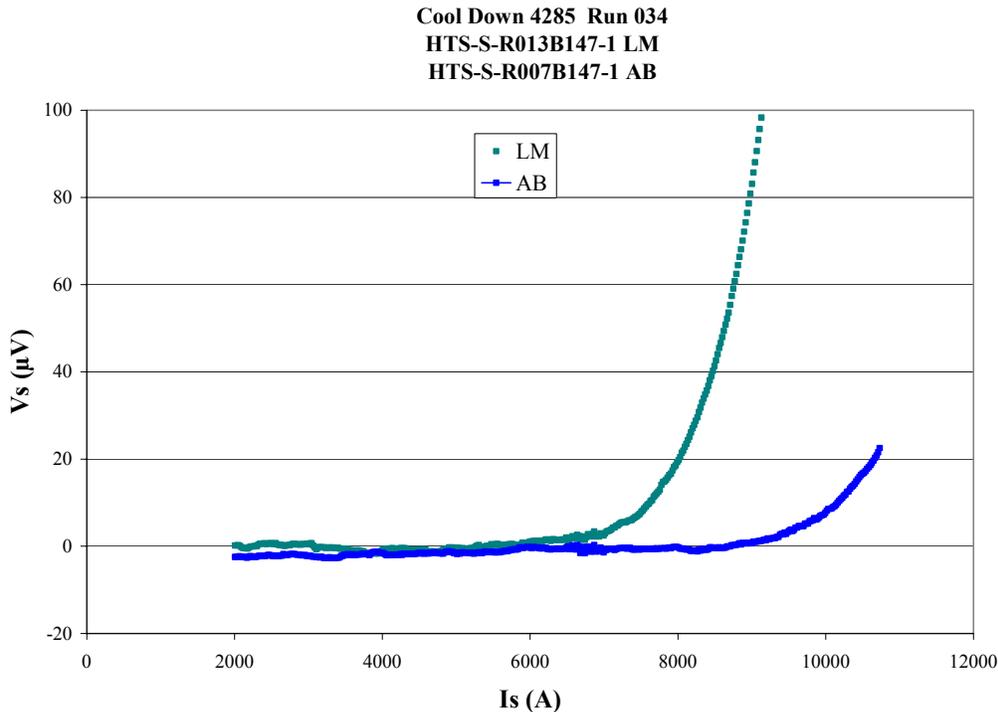
This picture depicts a keystone-cable rather than a rectangular cable. The purpose of this is to show, that with this orientation of the field, the peak field is at location indicated as Bp(//+) at zero applied field, but moves to Bp(//-) at fields greater than 0.4T. The other locations are for the case when the field is perpendicular to the wide face of the cable. The uniform field is 610mm covering the section AB and LM.

In analyzing the data, the V-I curve is fitted to the usual expression used to extract the Ic of LTS cables. After a base-line is subtracted, the data is fit to

$$\rho (\mu\Omega\text{-cm}) = 10^{-12} (I/I_c)^n$$

where $\rho = (V/I) L/A$, L=length of V-Tap and A is the cross-sectional area which is taken as 30 times wire area and Ic is defined for a resistivity of $10^{-12} \mu\Omega\text{-cm}$. We also show the Ic defined for 1 $\mu\text{V/cm}$ criterion and all plots are for this definition of Ic. The two Ic's are related with the n-value.

It appears that the critical currents of the two samples differ appreciably, with section LM showing much more voltage than AB. In fact the cable samples quench before AB reaches 1 $\mu\text{V/cm}$. *The data for the 1 $\mu\text{V/cm}$ is an extrapolation for the section AB and a measured value for section LM.* This is shown below for an applied field of 0.1T:



In the attached reports, the self-field has not been taken into account. With the field direction and current polarity used, the applied field is used to reduce the self-field thereby accessing the field region below the zero-applied-field region. It would be a lot easier to evaluate wires and cable if one moved away from zero-field and benchmarked conductors at a higher field, for instance 5T.

Below is the summary of the analysis for the two samples:

Superconductor Cable Test Report

10 Apr 2003

Cable: HTS-S-R007B147-1
Cool Down 4285

|| Low

$I_c(5\text{ T}) = 6748\text{ A}$
 $2\sigma I_c(5) = 0.021\text{ A}$

B (T)	T (K)	I _{min} Fit (A)	I _c (A)	σI _c (A)	I _q (A)	n	I _c (1.0μV/cm) (A)	File
0.00	4.443	2,004	8,906	9	10,399	10.6	11,406	4285_AB.002
0.05	4.441	2,002	9,274	10	10,524	12.3	11,473	4285_AB.008
0.05	4.436	2,006	9,075	22	10,673	11.2	11,464	4285_AB.009
0.10	4.439	4,066	9,104	26	10,249	10.2	11,740	4285_AB.014
0.10	4.436	2,006	9,550	12	10,792	14.4	11,454	4285_AB.016
0.20	4.445	2,008	9,702	20	11,258	11.7	12,080	4285_AB.017
0.30	4.444	2,008	9,921	18	11,447	11.0	12,489	4285_AB.018
0.40	4.442	2,008	10,225	30	11,215	12.4	12,537	4285_AB.019
0.50	4.434	4,101	9,747	36	11,256	8.8	12,944	4285_AB.020
0.60	4.444	2,007	9,716	36	11,675	8.9	12,870	4285_AB.021
0.80	4.446	2,007	8,974	19	11,072	9.2	11,881	4285_AB.022
1.00	4.437	2,007	8,130	11	10,285	8.7	11,033	4285_AB.023
2.00	4.434	2,070	6,363	8	8,238	8.9	8,797	4285_AB.024
3.00	4.442	2,070	5,646	17	7,370	9.5	7,750	4285_AB.025
4.00	4.444	2,070	5,106	14	6,842	9.2	7,145	4285_AB.026
5.00	4.434	2,219	4,680	21	6,454	8.6	6,748	4285_AB.033
6.00	4.440	2,118	4,494	17	6,215	8.8	6,459	4285_AB.034

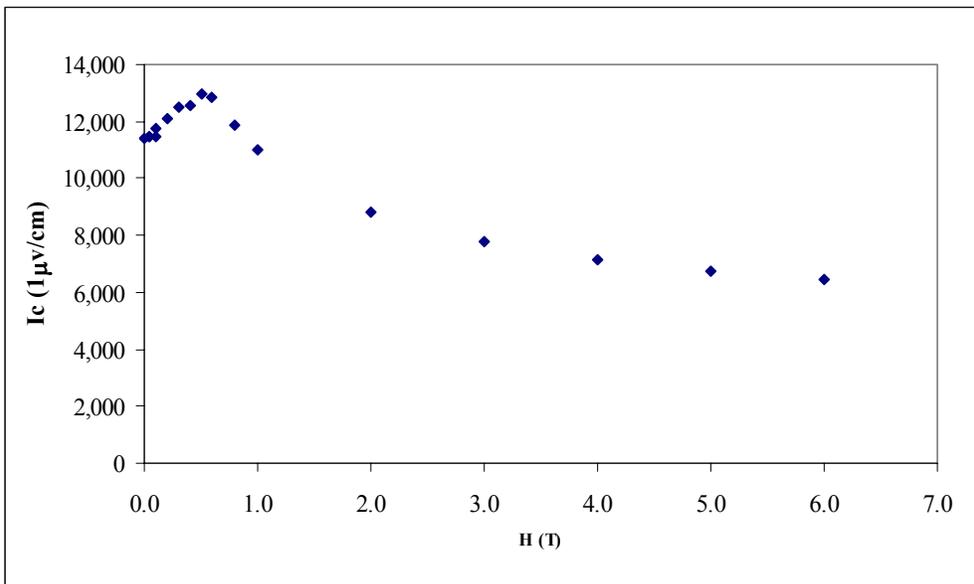
B (T)	I _c (A)	2σ (I _c) (A)	J _c (A/mm ²)
4.00	7145	14	1,861
5.00	6748	21	1,757
6.00	6459	17	1,682

R(293) = 18.958 μΩ/cm

Ag/Sc = 3.0

Comments: Short Length Sample
Paired with: HTS-S-R013B147-1

HT (rcvd): HT@Showa



Superconductor Cable Test Report

10 Apr 2003

Cable: HTS-S-R013B147-1
Cool Down 4285

|| Low

$I_c(5\text{ T}) = 4861\text{ A}$
 $2\sigma I_c(5) = 4\text{ A}$

B (T)	T (K)	I _{min} Fit (A)	I _c (A)	σI _c (A)	I _q (A)	n	I _c (1.0μV/cm) (A)	File
0.00	4.443	2,004	6,549	2	10,399	10.5	8,633	4285_LM.002
0.05	4.441	2,002	6,670	2	10,524	10.5	8,779	4285_LM.008
0.05	4.436	2,006	6,441	4	10,673	10.1	8,589	4285_LM.033
0.10	4.439	2,002	7,164	3	10,249	12.0	9,085	4285_LM.009
0.10	4.436	2,006	6,613	4	10,792	10.3	8,753	4285_LM.034
0.20	4.445	2,008	7,385	3	11,258	11.8	9,377	4285_LM.014
0.30	4.444	2,008	7,863	3	11,447	12.6	9,800	4285_LM.016
0.40	4.442	2,008	7,936	15	11,215	11.9	10,003	4285_LM.017
0.50	4.434	2,006	7,264	6	11,256	9.3	9,790	4285_LM.018
0.60	4.444	2,007	6,650	4	11,675	8.2	9,378	4285_LM.019
0.80	4.446	2,007	5,687	5	11,072	7.4	8,442	4285_LM.020
1.00	4.437	2,007	5,151	7	10,285	7.2	7,815	4285_LM.021
2.00	4.434	1,540	4,097	2	8,238	7.7	6,230	4285_LM.022
3.00	4.442	2,070	3,645	3	7,370	8.0	5,537	4285_LM.023
4.00	4.444	2,070	3,347	3	6,842	8.0	5,133	4285_LM.024
5.00	4.434	2,117	3,165	4	6,454	8.1	4,861	4285_LM.025
6.00	4.440	2,118	3,012	6	6,215	8.2	4,629	4285_LM.026

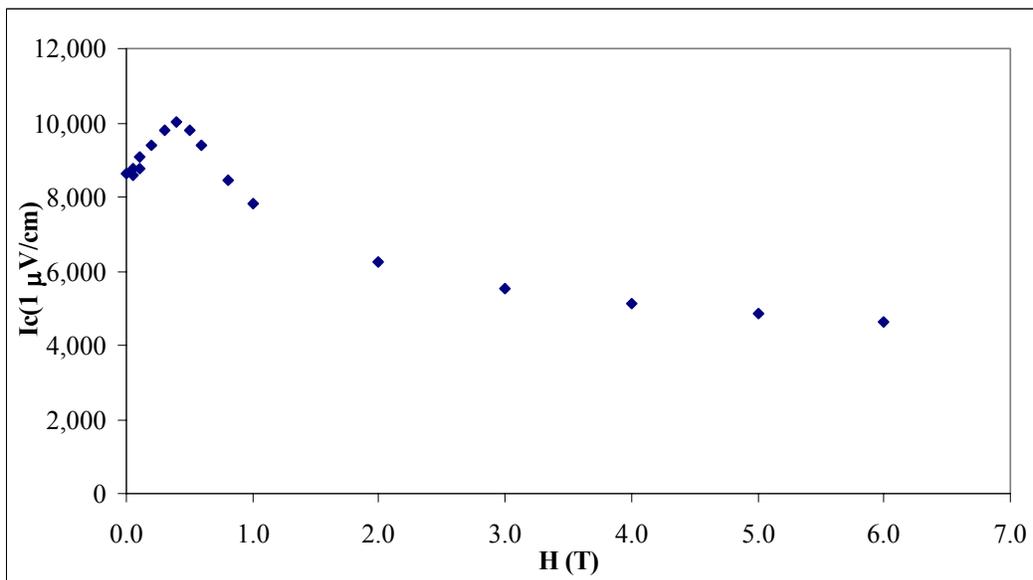
B (T)	I _c (A)	2σ (I _c) (A)	J _c (A/mm ²)
4.00	5133	3	1,337
5.00	4861	4	1,266
6.00	4629	6	1,206

R(293) = 19.755 μΩ/cm

Cu/Sc = 3.0

Comments: Long Coil Sample
Paired with: HTS-S-R007B147-1

HT (rcvd): HT@ Showa



Also summarized below are the tests in LN₂:

In Low-Temp Holder as a bifilar sample							
	Field (T)	Hsf T	T (K)	Ic (A)	Ic(1.0μv) (A)	n	Section
	0.00		77		148	6.4	AC
	0.00		77		149	6.7	BC
	0.00		77		130	6.2	LN
	0.00		77		132	6.6	NM