

Status of Transverse Spin Physics

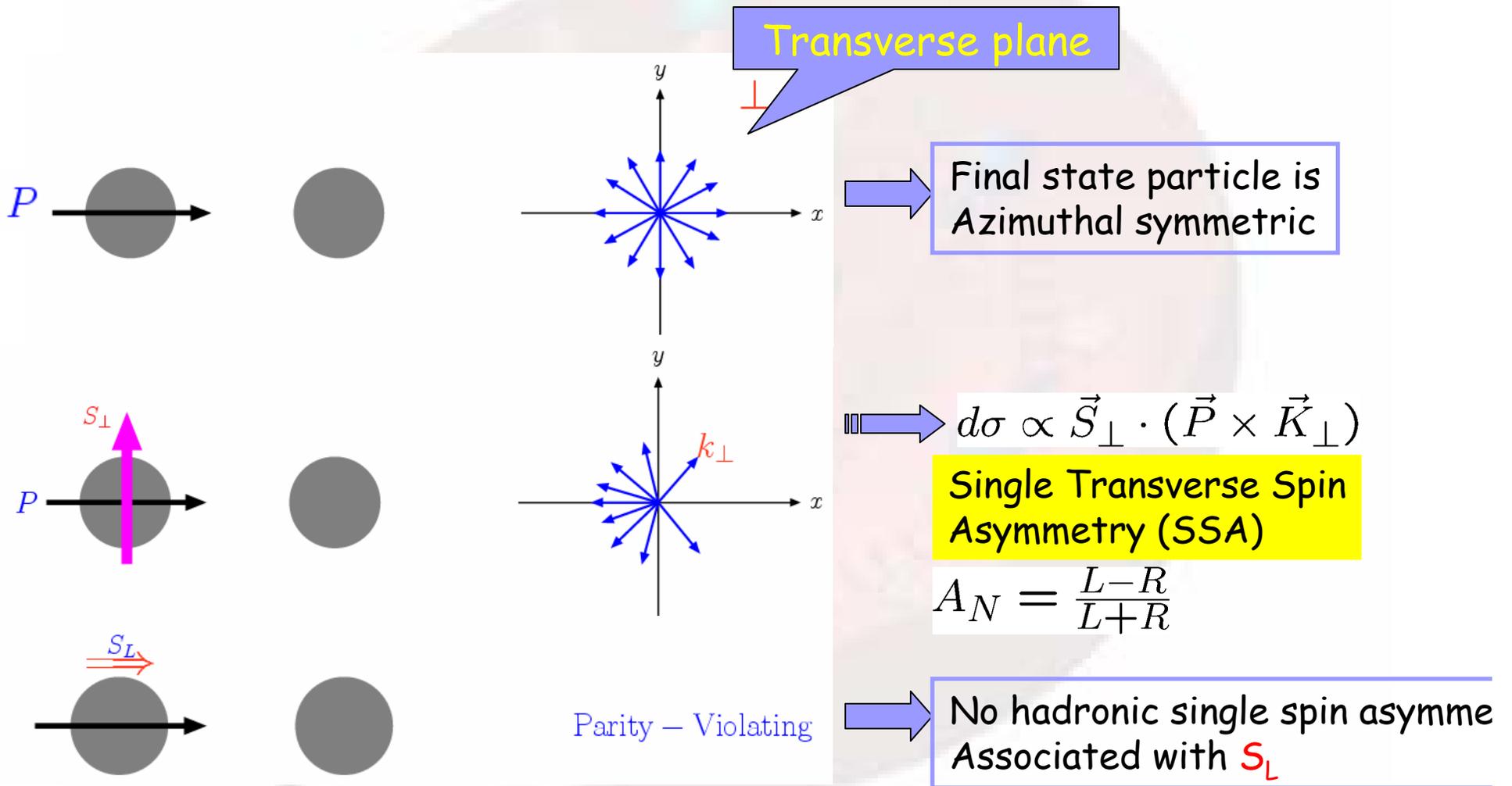
Feng Yuan

RBRC, Brookhaven National Laboratory

Transverse spin physics at RHICII and eRHIC

- Double transverse spin asymmetry
 - A_{TT} for Drell-Yan, jet, hadron, etc. to measure the quark transversity
- **Single transverse spin asymmetry**
 - Various TMD physics (additional information on nucleon structure) see more in Gamberg's talk
 - **Sivers function (PDF)**
 - **Collins function (FF)**
 - ...

What's Single spin asymmetry?



Single Spin Asymmetry

■ Motivations:

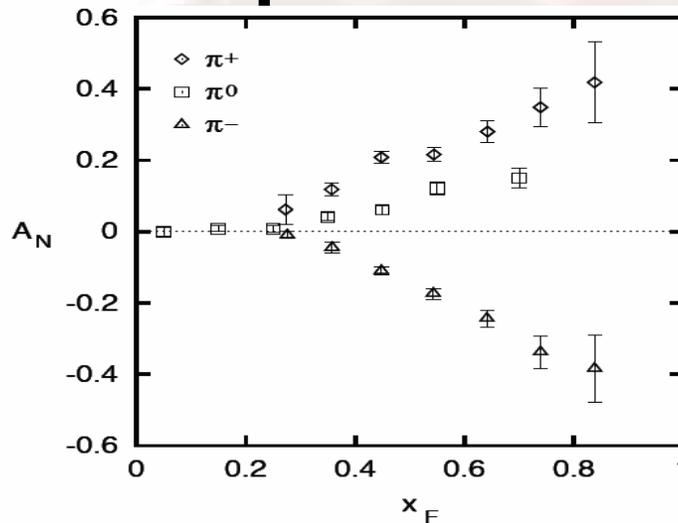
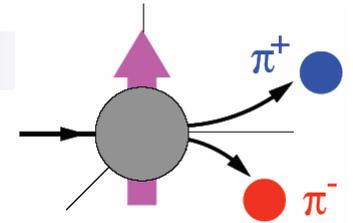
- Its strong tied with the quark orbital angular momentum
- We have beautiful data
- Nontrivial QCD dynamics, and fundamental test of the **factorization**, and the **universality** of PDFs, FFs,...

Transverse spin physics: a Renaissance



- Single transverse spin asymmetry has long been observed in high energy hadronic reactions (late 70's and 80's)
- Naive Parton model fails to explain large Single transverse spin asymmetry (*G. Kane et al, PRL41, 1978*), we have to go beyond to understand the large SSAs
- Many activities in recent years

SSAs in "Ancient" age: fixed target exp.

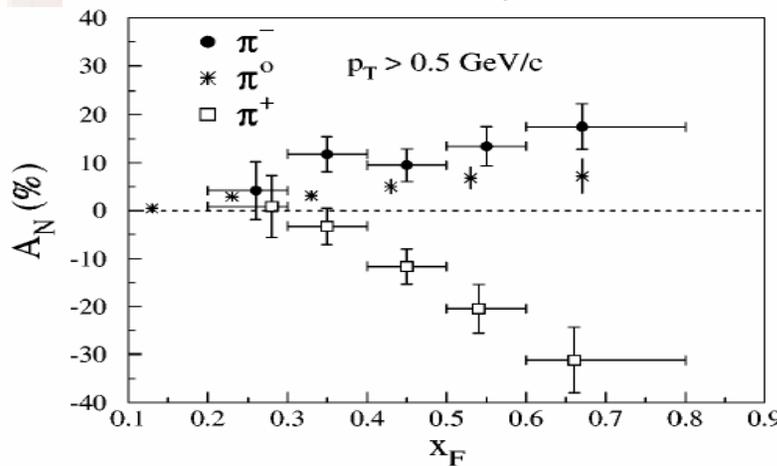


$$A_N = \frac{L-R}{L+R}$$

$$\sqrt{s} = 20 \text{ GeV}$$

Hadron production in p.p

Adams, et al., PLB264 (1991)

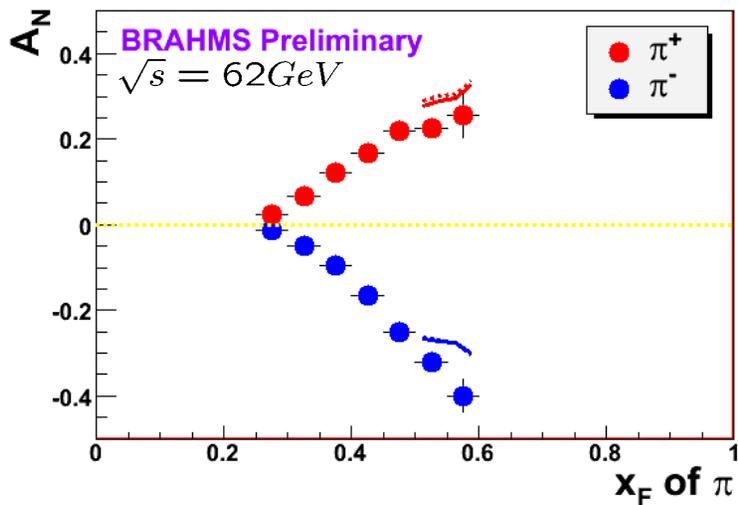
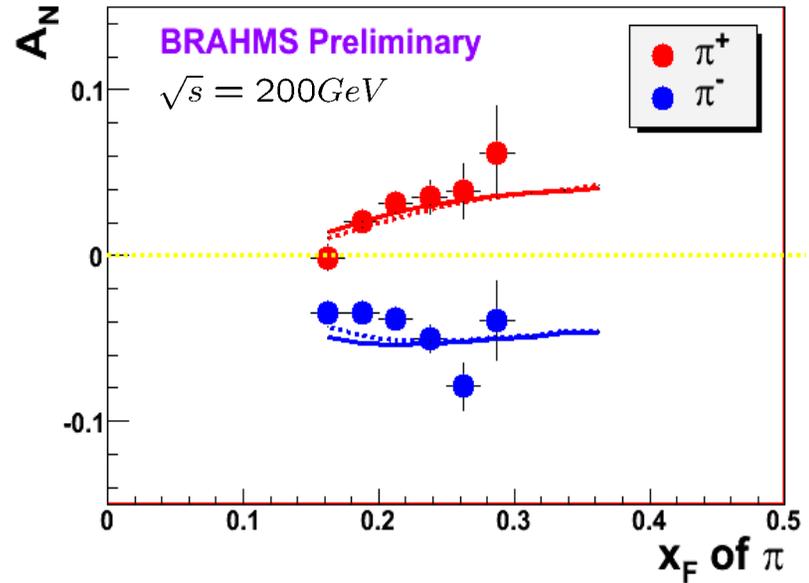
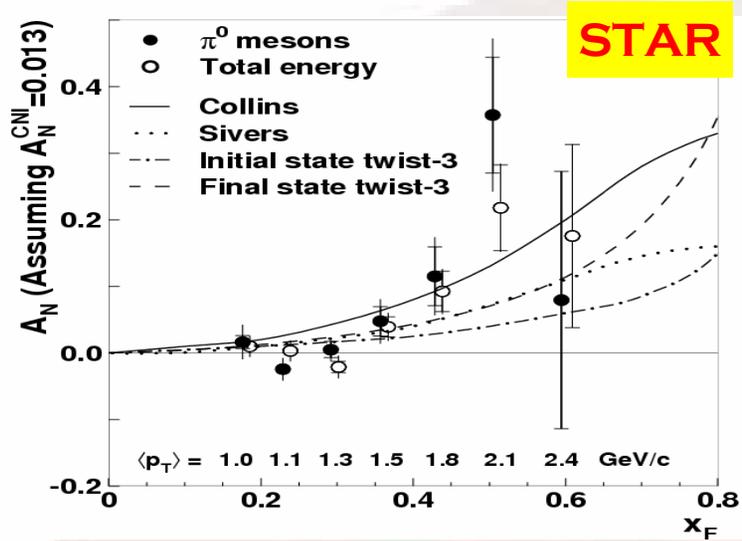
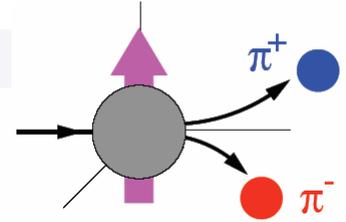


$$\sqrt{s} = 20 \text{ GeV}$$

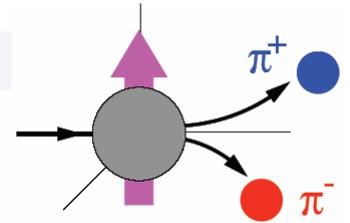
hadron production in $\bar{p}.p$

A. Bravar et al., E704, PRL77(1996)

SSAs in Modern era: RHIC

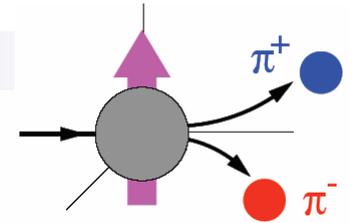


Large SSA continues at Collider experiments!!
(Also in ep DIS exp.)



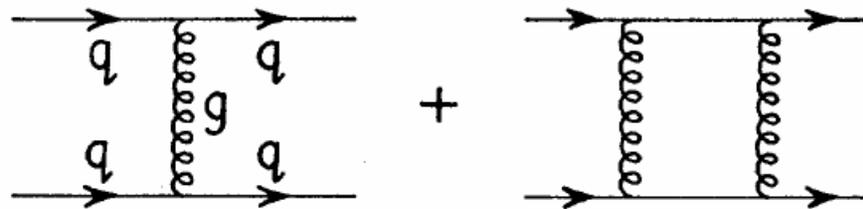
Why Does SSA Exist?

- **Single Spin Asymmetry requires**
 - **Helicity flip**: one must have a reaction mechanism for the hadron to change its helicity (in a cut diagram)
 - **A phase difference**: the phase difference is needed because the structure $S \cdot (p \times k)$ violate the naïve time-reversal invariance

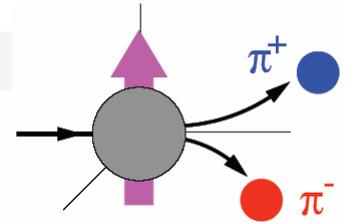


Naive parton model fails?

- If the underlying scattering mechanism is hard, the naive parton model generates a very small SSA: (*G. Kane et al, 1978*),



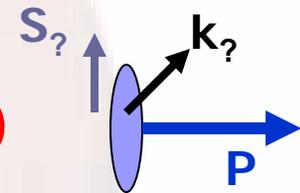
- It is in general suppressed by $\alpha_S m_q/Q$
- We have to go beyond this naive picture



Two mechanisms in QCD

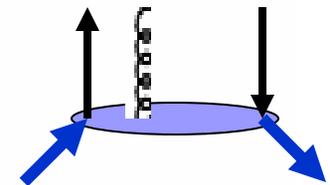
■ Spin-dependent transverse momentum dependent (TMD) function

- Sivers 90 **Sivers function** » $S_? \phi(P \times k_?)$
- Brodsky, Hwang, Schmidt, 02 (**FSI**)
- Gauge Property: Collins 02; Belitsky-Ji-Yuan, NPB03
- Factorization: Ji-Ma-Yuan, PRD04; Collins, Metz, 04



■ Twist-3 quark-gluon correlations (coll.)

- Efremov-Teryaev, 82, 84
- Qiu-Sterman, 91, 98



Regions

- Twist-three: the single inclusive hadron production in pp, require large P_T , SSA is suppressed by $1/P_+$
- TMD: low P_T , require additional hard scale like Q^2 in DIS and Drell-Yan, $P_T \ll Q$, SSA survives in Bjorken limit
- Overlap: $\Lambda_{\text{QCD}} \ll P_+ \ll Q$, unifying these two

Recent theoretical developments (twist-three)

- Complete formalism for single inclusive hadron production in pp collision has been derived, including the derivative and non-derivative terms

$$E_\ell \frac{d^3 \Delta\sigma(\vec{s}_T)}{d^3 \ell} = \frac{\alpha_s^2}{S} \sum_{a,b,c} \int_{z_{\min}}^1 \frac{dz}{z^2} D_{c \rightarrow h}(z) \int_{x'_{\min}}^1 \frac{dx'}{x'} \frac{1}{x'S + T/z} \phi_{b/B}(x')$$
$$\times \sqrt{4\pi\alpha_s} \left(\frac{\epsilon^{\ell s_T n \bar{n}}}{z \hat{u}} \right) \frac{1}{x} \left[T_{a,F}(x, x) - x \left(\frac{d}{dx} T_{a,F}(x, x) \right) \right] H_{ab \rightarrow c}(\hat{s}, \hat{t}, \hat{u})$$

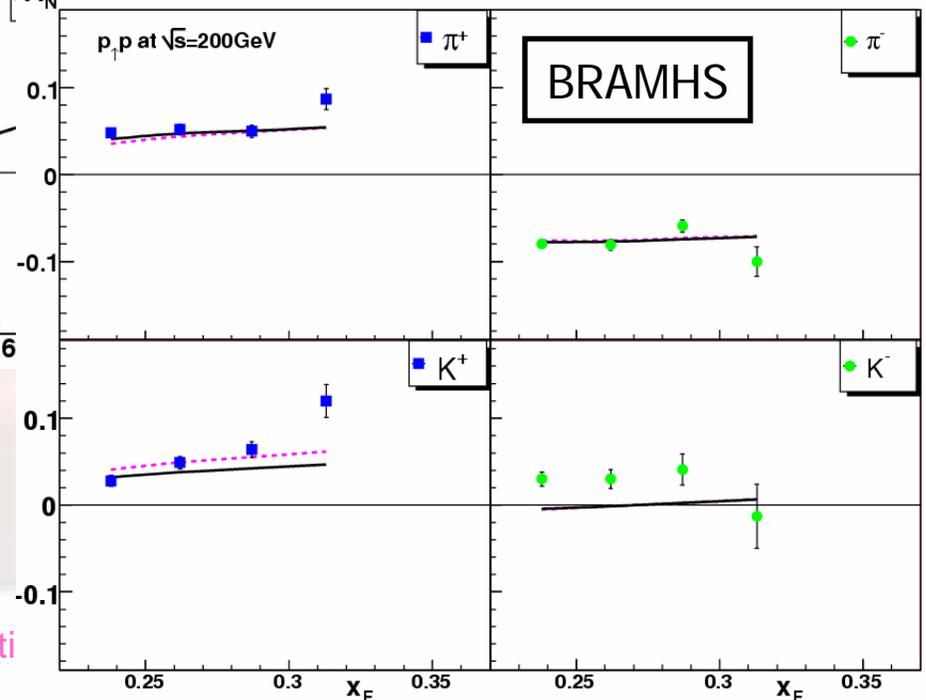
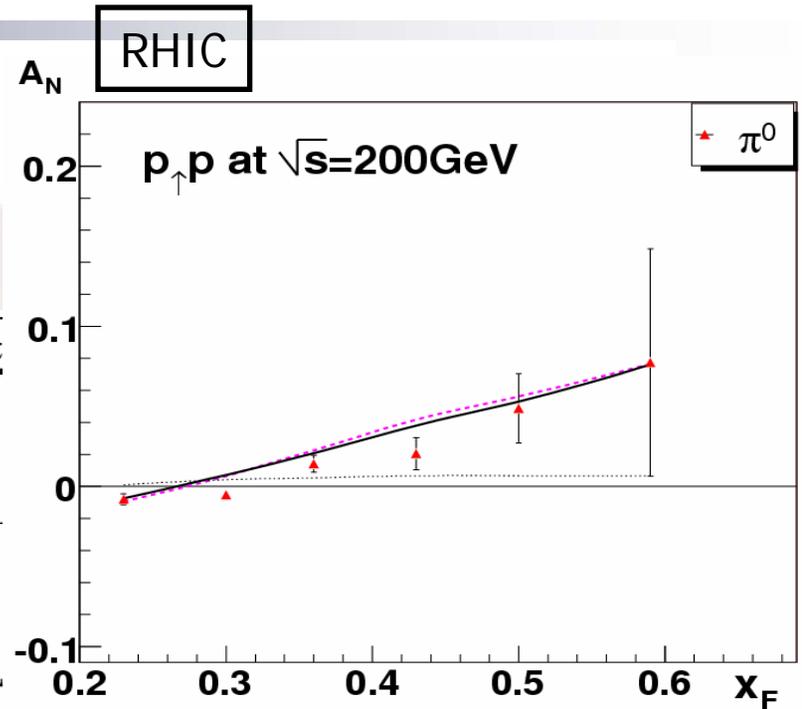
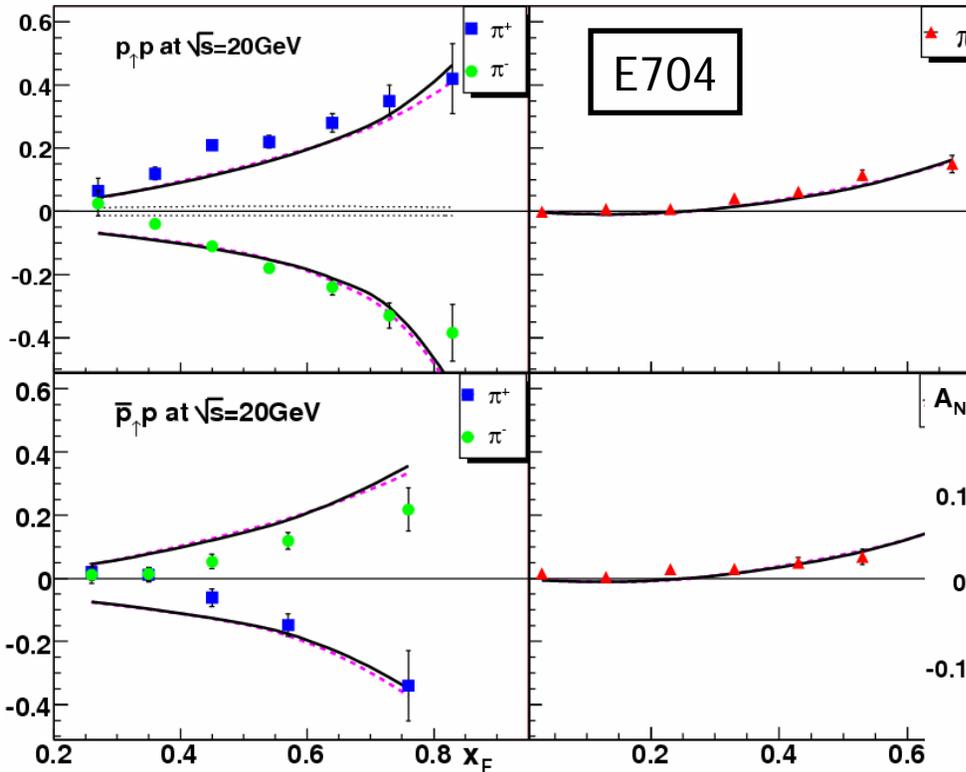
Qiu, Sterman, 91, 98

Kouvaris, Qiu, Vogelsang, Yuan, 06

See also Koike, et al., 06,07

Twist-3 Fit to data

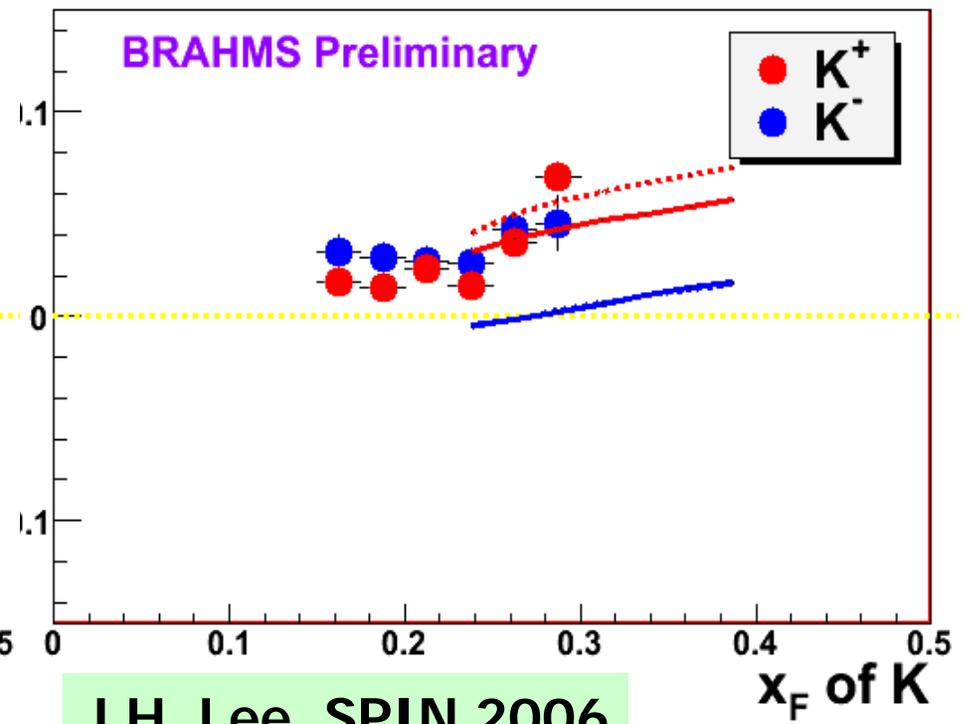
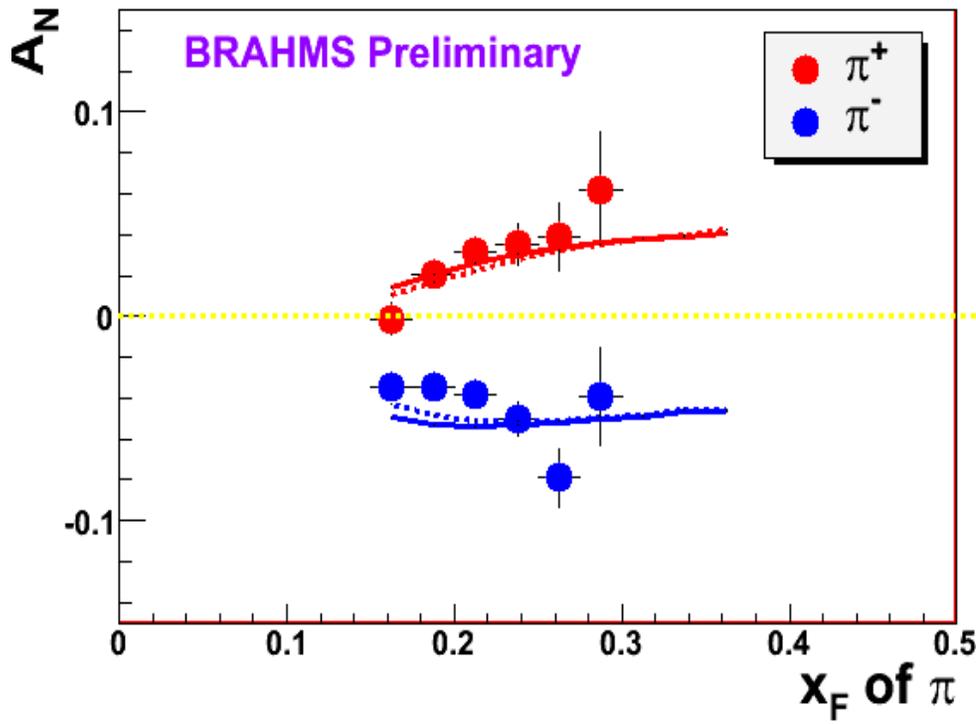
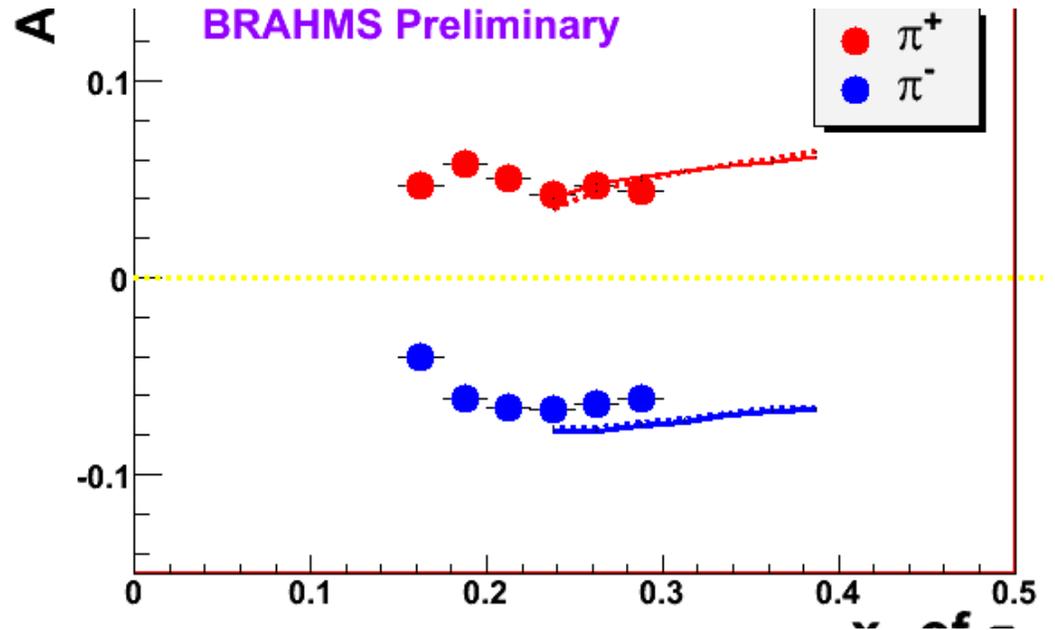
$$p \uparrow p \rightarrow \pi + X$$



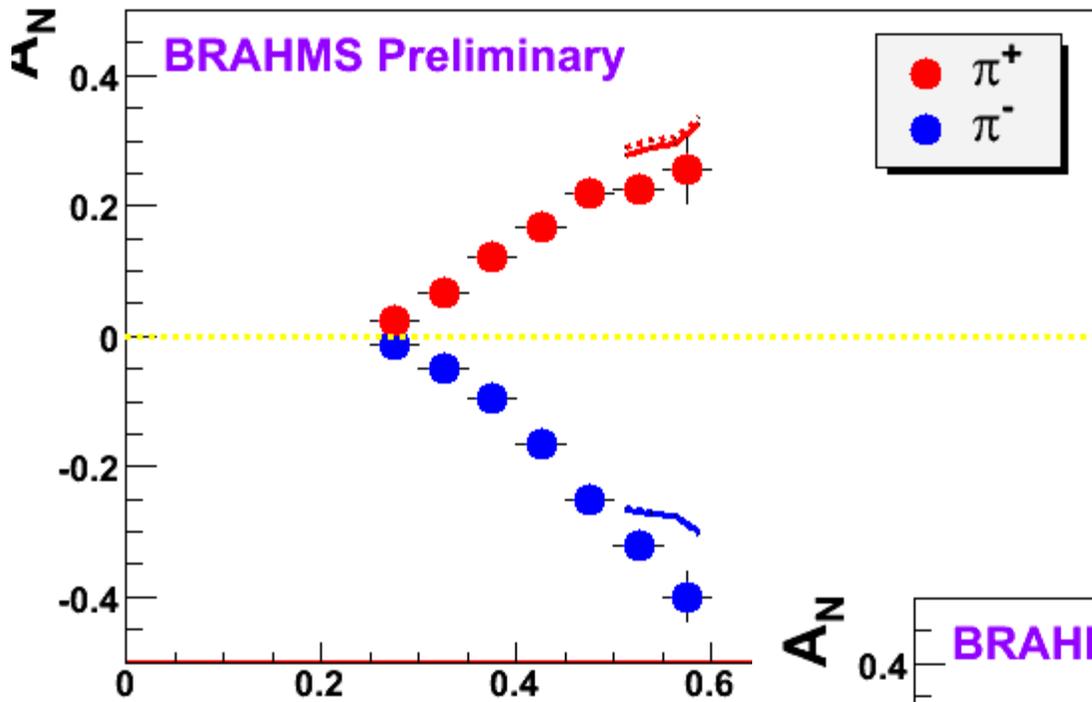
$$T_F^a(x) = N_a x^{\alpha_a} (1-x)^{\beta_a} f_a(x)$$

Kouvaris, Qiu, Vogelsang, Yuan, 06

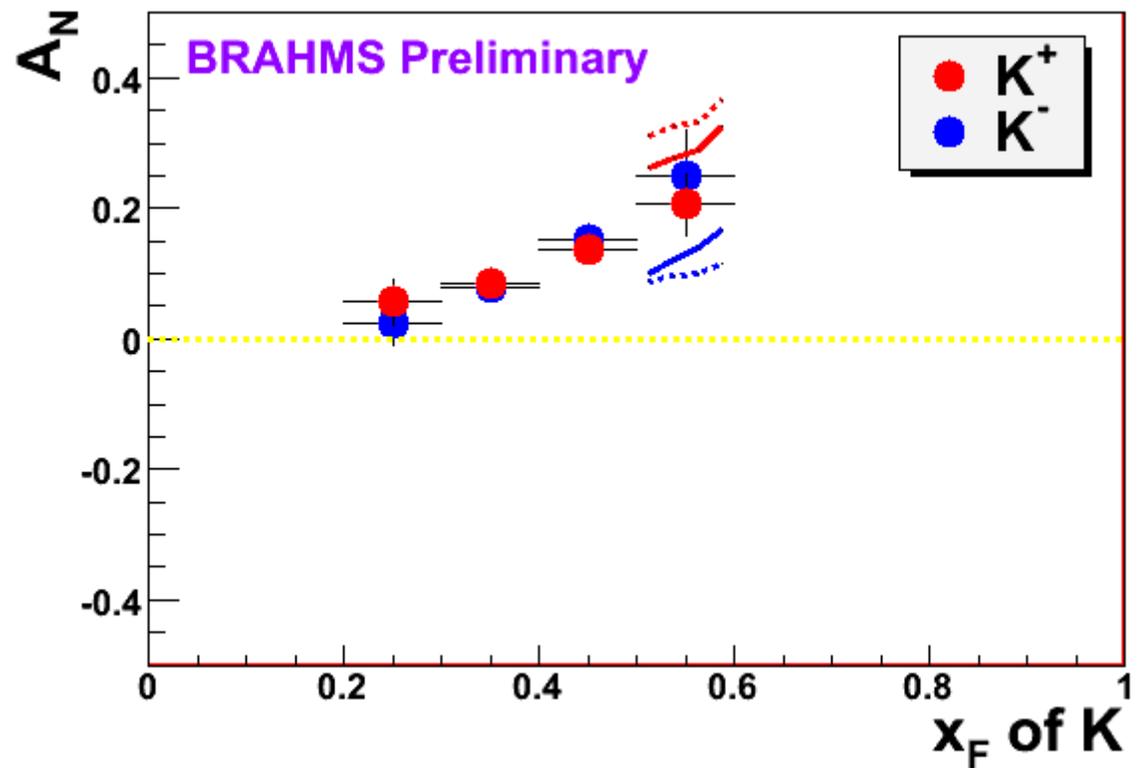
Compare to
2006 data
from RHIC



J.H. Lee, SPIN 2006



$$\sqrt{s} = 62\text{GeV}$$



CCO meeting, BNL

A unified picture for SSA

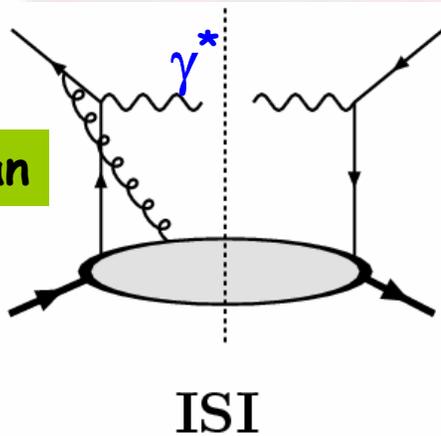
- In DIS and Drell-Yan processes, SSA depends on Q and transverse-momentum P_γ
 - At large P_γ , SSA is dominated by twist-3 correlation effects
 - At moderate P_γ , SSA is dominated by the transverse-momentum-dependent parton distribution/fragmentation functions
- The two mechanisms at intermediate P_γ generate the same physics!

Ji-Qiu-Vogelsang-Yuan, Phys. Rev. Lett. 97:082002, 2006

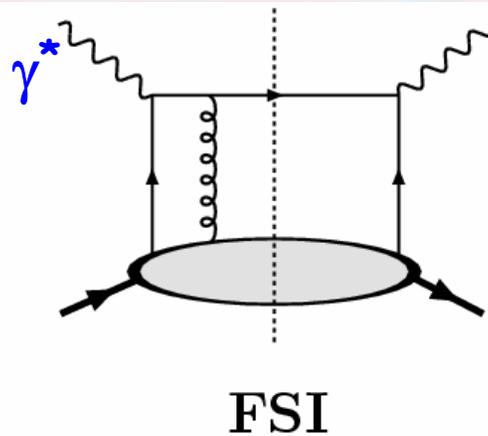
Testable prediction at RHIC

- Initial state vs. final state interactions

Drell-Yan



ISI



FSI

DIS

$$\text{Sivers}|_{\text{DY}} = -\text{Sivers}|_{\text{DIS}}$$

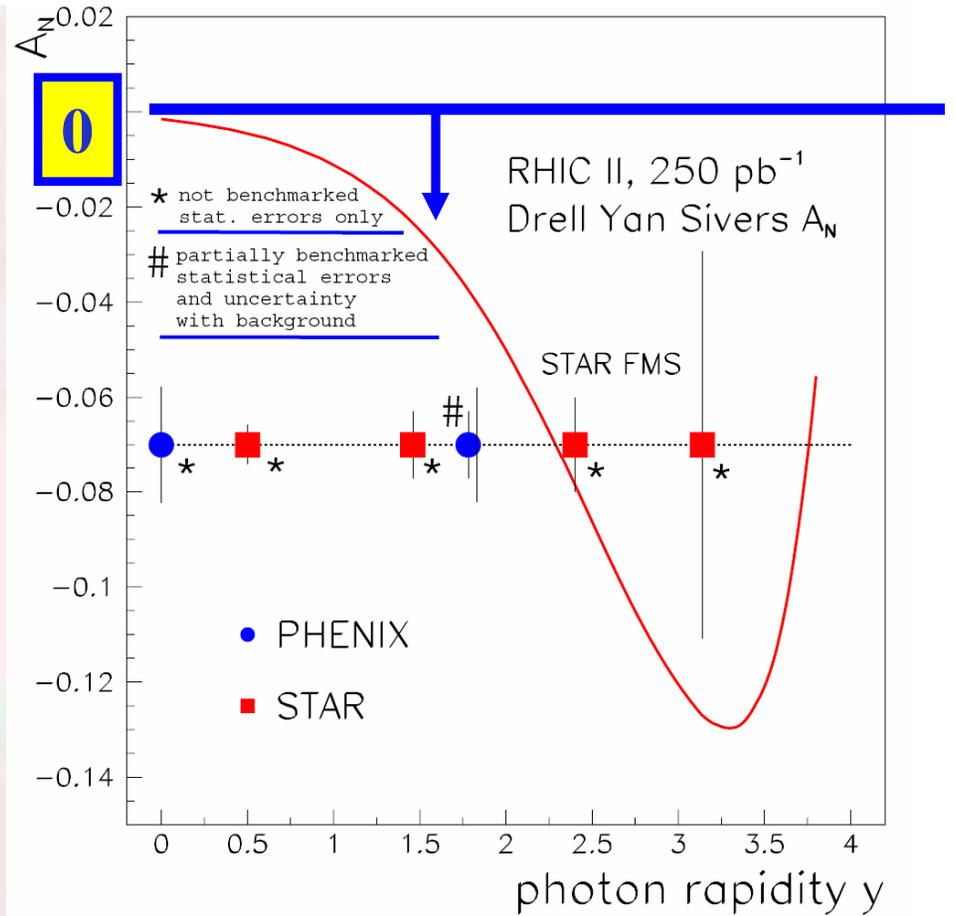
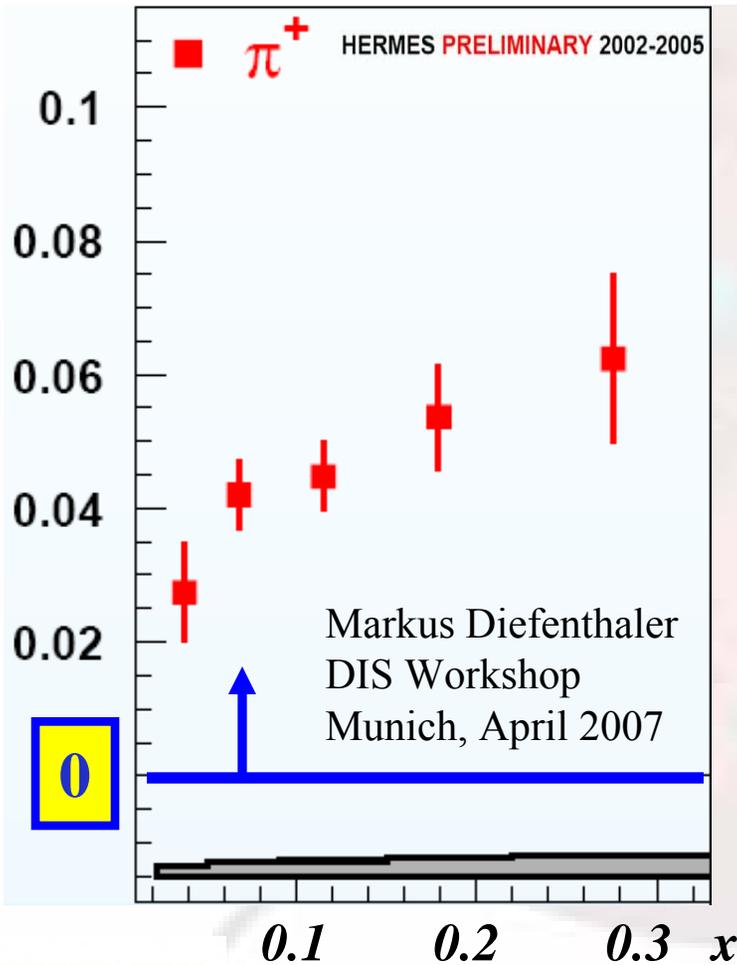
HERMES

- "Universality": fundamental QCD prediction

Experiment SIDIS vs Drell Yan

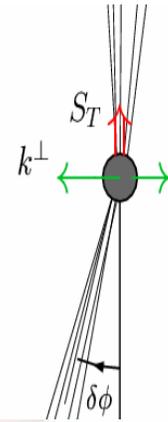
HERMES Sivvers Results

RHIC II Drell Yan Projections



Alternative: Dijet-correlation at RHIC?

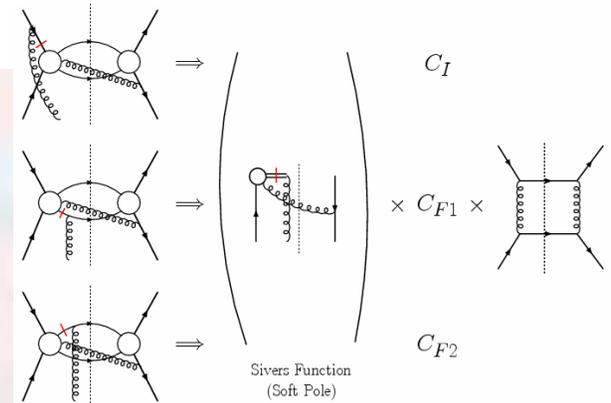
- Proposed by Boer-Vogelsang
 - Pheno. studies: Vogelsang-Yuan 05; Bomhof-Mulders-Vogelsang-Yuan 07
- Initial state and/or final state interactions?
 - Bacchetta-Bomhof-Mulders-Pijlman: hep-ph/0406099, hep-ph/0505268, hep-ph/0601171, hep-ph/0609206
 - Qiu-Vogelsang-Yuan, arXiv:0704.1153; 0706.1196
 - Collins-Qiu, arXiv:0705.2141
 - Mulders, et al., 07?
- Factorization? Universality?



The asymmetry can be related to that in DIS,
 in leading power of q_T/P_t ,
universality of the parton distributions?

$$\frac{d\sigma_{TU}}{d^2\vec{q}_\perp} = \sum_{ab} \int d^2k_{1\perp} d^2k_{2\perp} d^2\lambda_\perp \frac{\vec{k}_{1\perp} \cdot \hat{q}_\perp}{M} x_a q_{Ta}^{(\text{DIS})}(x_a, k_{1\perp}) x_b f_b(x_b, k_{2\perp}) H_{ab \rightarrow cd}^{\text{sivers}}(P_\perp^2) \delta^{(2)}(\vec{k}_{1\perp} + \vec{k}_{2\perp} - \vec{q}_\perp)$$

- q_T^{DIS} --- Sivers function from DIS
- q_T --- imbalance of the dijet
- H^{sivers} depends on subprocess



$$H_{ab \rightarrow cd}^{(\text{sivers})} = \frac{\alpha_s^2 \pi}{\hat{s}^2} \sum_i (C_I^i + C_{F1}^i + C_{F2}^i) h_{ab \rightarrow cd}^i(\hat{s}, \hat{t}, \hat{u})$$

Qiu, Vogelsang, Yuan, 07

Questions...

- To what extent do factorization breaks hold?
- Does a modified (weak) factorization hold if a general factorization breaks?
 - Universality will be an issue
- ...
- More rigorous works need to be done to draw any conclusion
- Experiments at RHIC and future eRHIC can help us to answer these questions

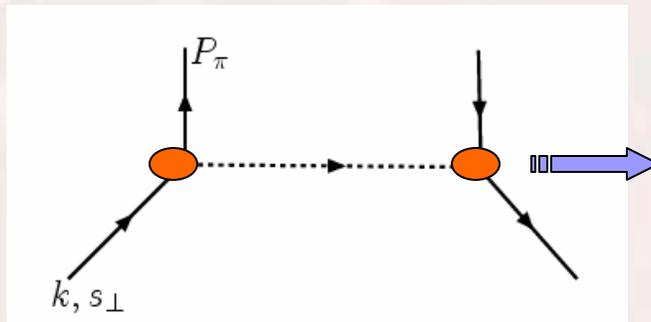
Collins effects

- An important ingredient to extract the quark **transversity** distribution

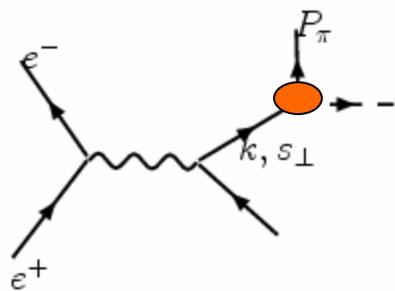
$$(k, s_?) \xrightarrow{\quad} (zk+p_?) \propto p_{\perp} \times s_{\perp}$$

- Its **universality** property has attracted many theoretical and experimental investigations
 - Metz 02, Collins-Metz 02, 04
 - Boer-Mulders-Pijlman, 03
 - Gamberg-Goldstein, ...
 - HERMES and Belle experiments, PRL 05-06

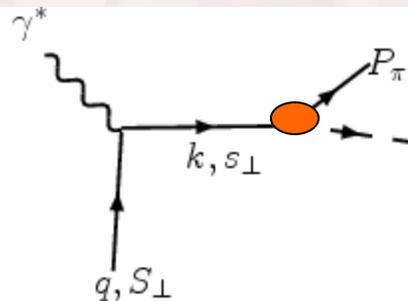
Simple model a la Collins 93



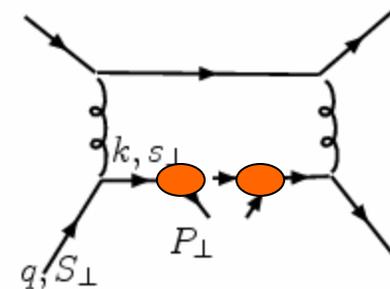
Phase information in the vertex
Collins-93



e^+e^- annihilation



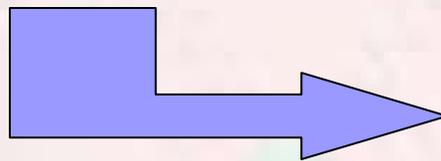
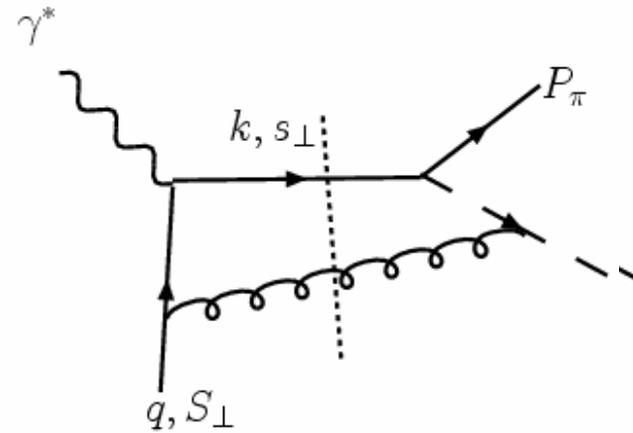
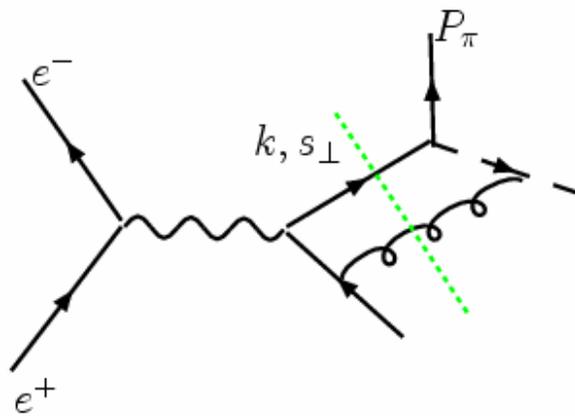
Semi-inclusive DIS



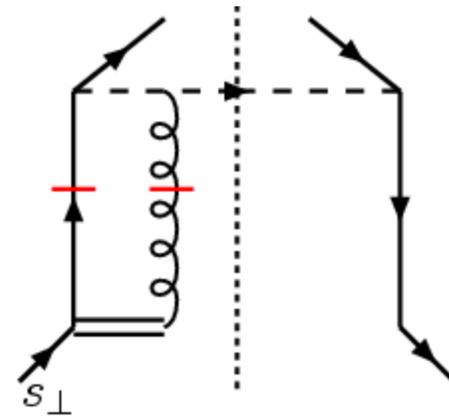
Hadron in a jet in pp

Universality of the Collins Function!!

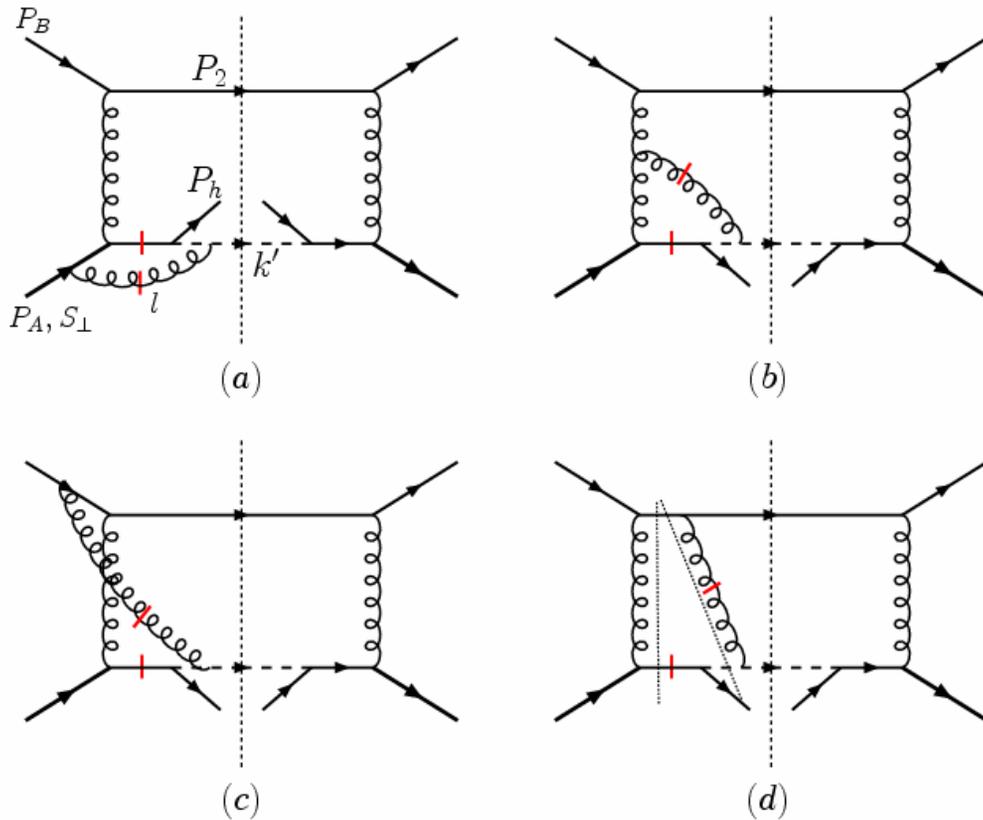
One-gluon exchange (gauge link)?



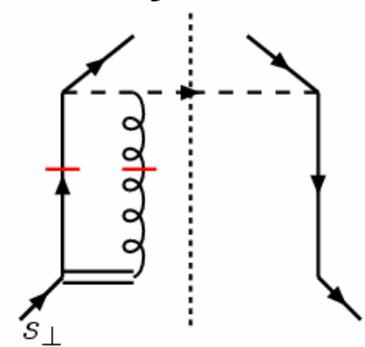
Metz 02, Collins-Metz 04:
Universality of the Collins function!!



Similar arguments for pp collisions



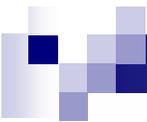
By using the Ward Identity:



same Collins fun.

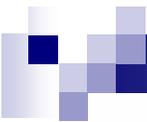
Vogelsang, Yuan, to appear

Conjecture: the Collins function will be the same as e^+e^- and SIDIS



Open questions

- Universality for the SSA in the **time-like** region (fragmentation)
 - Experiments at Belle, RHIC, eRHIC are crucial
- **Link with the quark-gluon correlation in the fragmentation process**
 - Koike et al., 00-06
- **Scale(Q^2)-dependence of all these SSAs**



Summary

- We are in the early stages of a very exciting era of transverse spin physics studies, where the future RHICII and eRHIC will certainly play very important roles
- We will learn more about nucleon structure from these studies, especially for the quark orbital motion

What can we learn from SSA

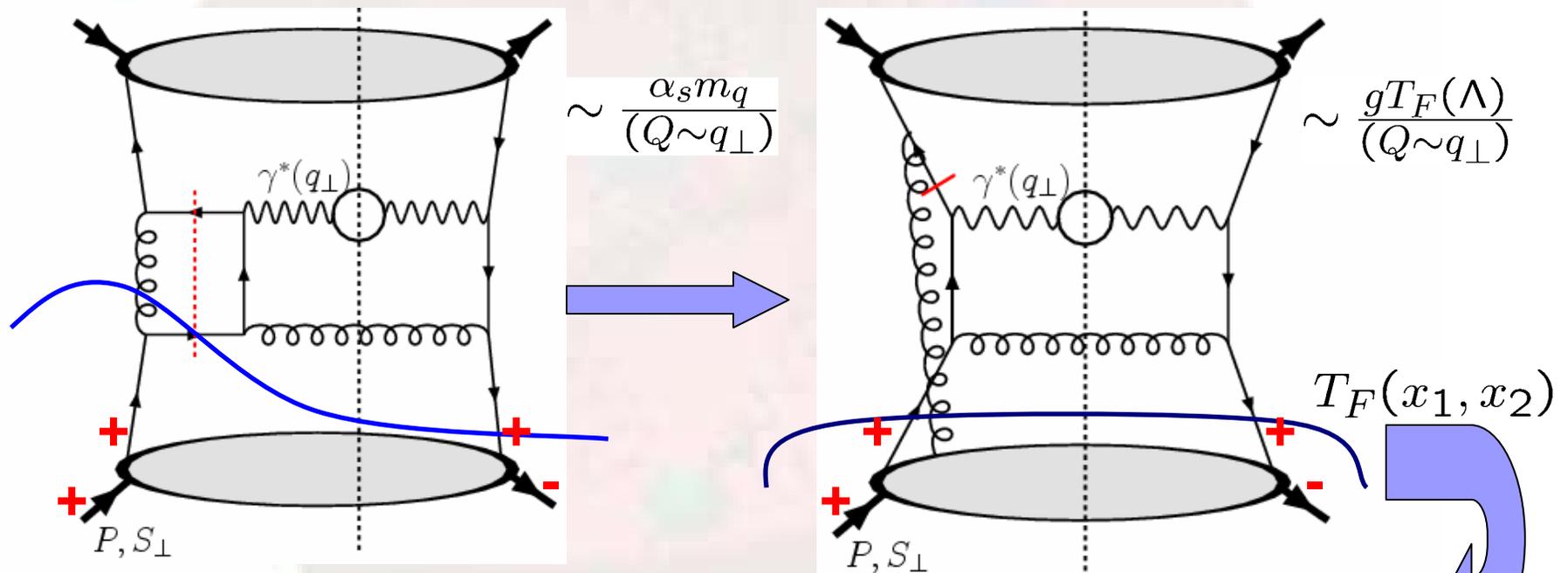
- Quark Orbital Angular Momentum
e.g, Sivers function \sim the wave function amplitude with nonzero orbital angular momentum!

Vanishes if quarks only in s-state!

Ji-Ma-Yuan, NPB03
Brodsky-Yuan, PRD06

Take Drell-Yan as an example (with non-zero transverse momentum q_{\perp})

- We need a loop to generate a phase

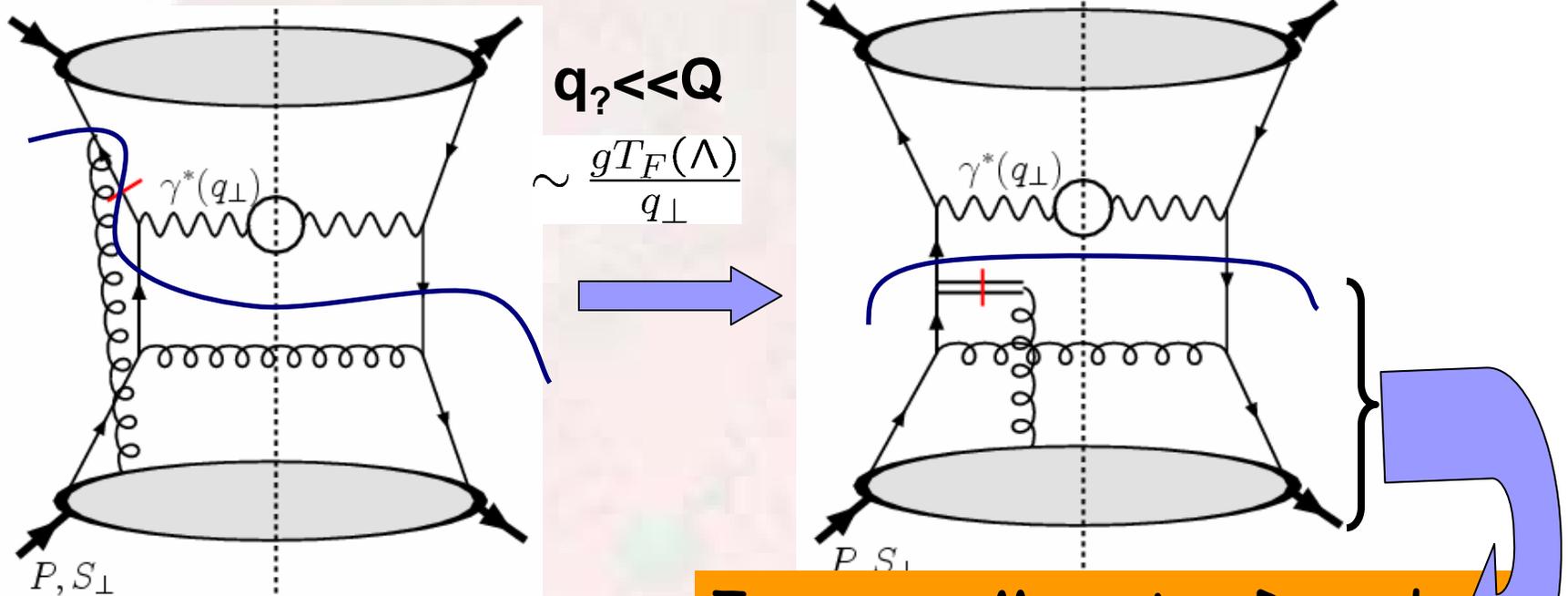


Kane et al., hard parton model

Twist-three Correlations
Efremov-Teryaev, 82, 84
Qiu-Sterman, 91,98

Further factorization ($q_{\parallel} \ll Q$)

- The collinear gluons dominate



Twist-three Correlations

Efremov-Teryaev, 82, 84

Qiu-Sterman, 91,98

Transverse Momentum Dependence distributions

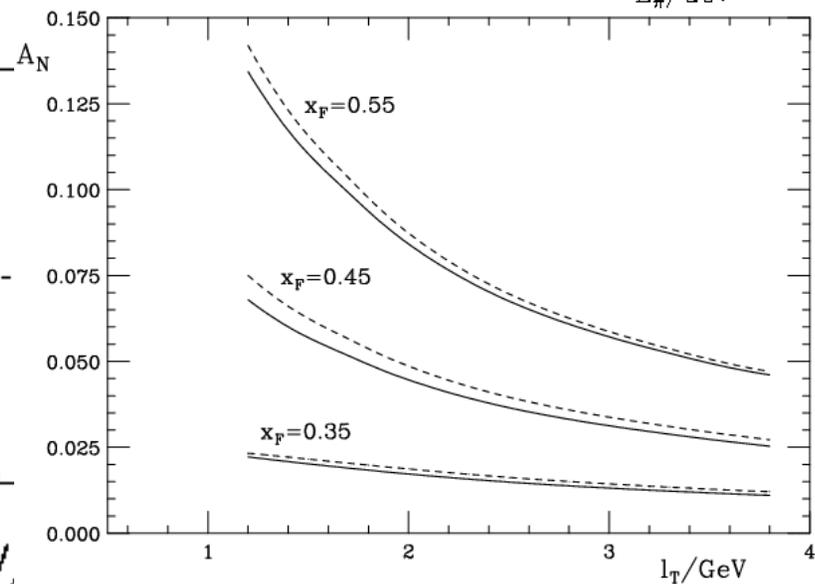
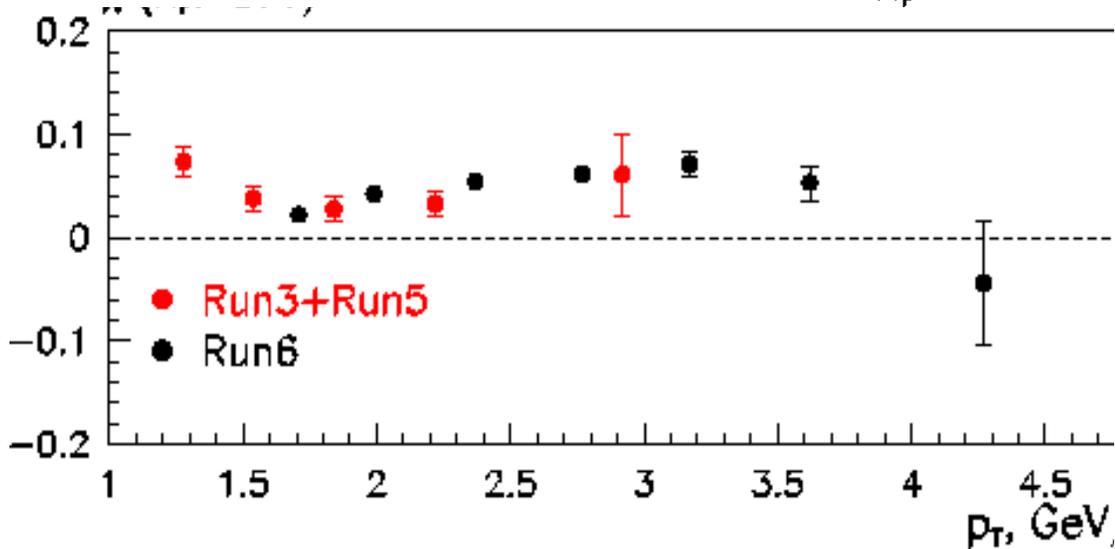
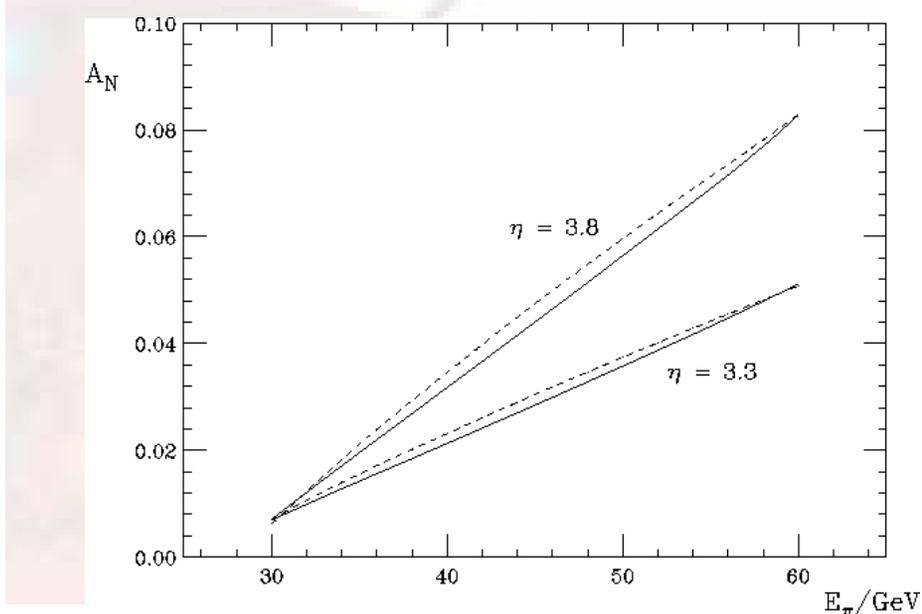
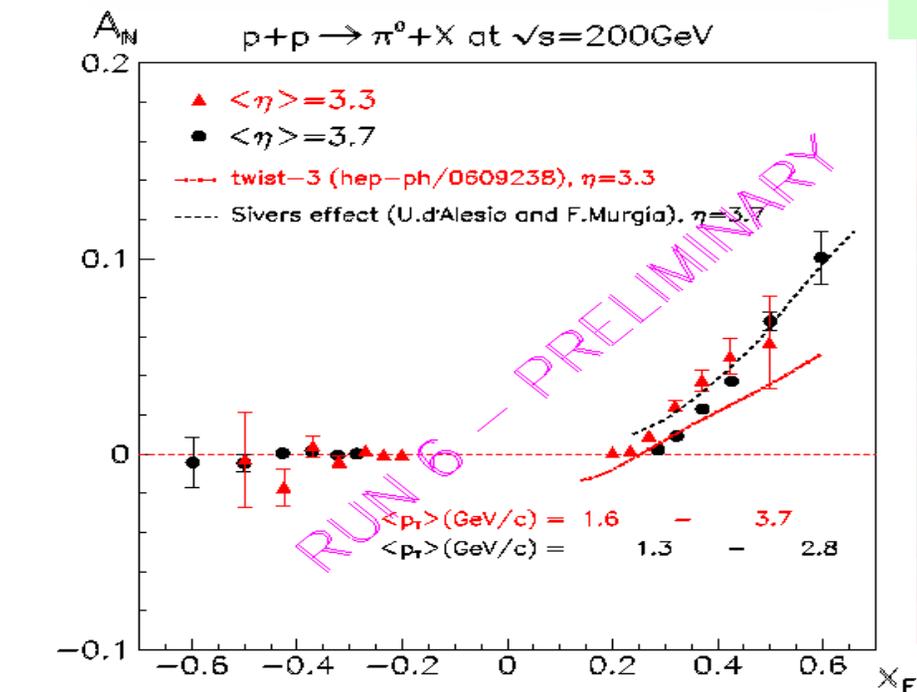
Sivers, 90, Collins, 93,02

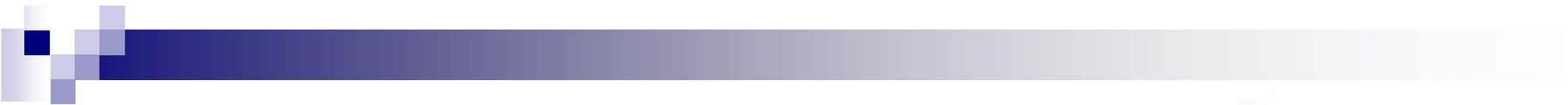
Brodsky-Hwang-Schmidt, 02

Ji-Qiu-Vogelsang-Yuan, 06

New challenge from STAR data (2006)

Talks by Ogawa and Nogach in SPIN2006





Some comments

- It's difficult to explain this pattern in the current theoretical approaches
- One possible reason could be the partial cancellation between the favored and disfavored contribution to Pi^0 production
 - Check with the charged pions?

Final Results

■ P_T dependence

$$\frac{d\Delta\sigma}{d^2q_\perp dy} = \int q_T(z_1, k_\perp) \bar{q}(z_2, k_\perp) + \left(\frac{d\Delta\sigma^{QS}}{d^2q_\perp dy} - \frac{d\Delta\sigma^{QS}}{d^2q_\perp dy} \Big|_{aspt.} \right)$$

Sivers function at low P_T

Qiu-Sterman Twist-three

- Which is valid for all P_T range
- SSA is suppressed by $1/P_t$ at large P_t

Extend to all other TMDs: large P_t power counting

- k_t -even distributions have the same dependence on k_t
- k_t -odd distributions are suppressed at large k_t
- Power Counting Rule

$$k_t\text{-even: } 1/k_t^2$$

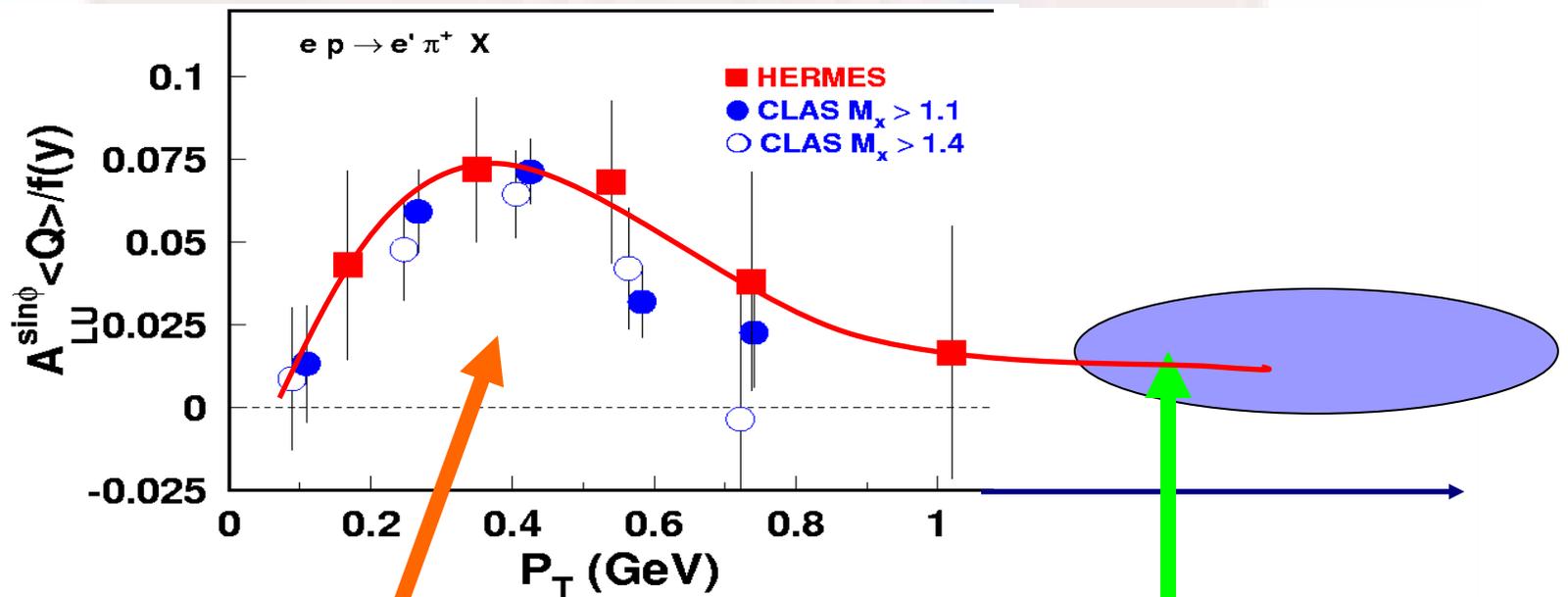
$$k_t\text{-odd: } 1/k_t^4$$

SIDIS cross sections at large P_t

$$\begin{aligned}
 d\sigma \propto & (1 - y + y^2/2)x_B F_{UU}^{(1)} && \xrightarrow{\text{blue}} && 1/P_t^2 \\
 & - (1 - y)x_B \cos(2\phi_h) F_{UU}^{(2)} && \xrightarrow{\text{green}} && \\
 & + \lambda_e \lambda y (1 - y/2)x_B F_{LL} && \xrightarrow{\text{blue}} && \\
 & + \lambda_e |S_\perp| y (1 - y/2)x_B \cos(\phi_h - \phi_S) F_{LT} && \xrightarrow{\text{green}} && 1/P_t^4 \\
 & + \lambda (1 - y)x_B \sin(2\phi_h) F_{UL} && \xrightarrow{\text{green}} && \\
 & + |S_\perp| (1 - y + y^2/2)x_B \sin(\phi_h - \phi_S) F_{UT}^{(1)} && \xrightarrow{\text{magenta}} && 1/P_t^3 \\
 & + |\vec{S}_\perp| (1 - y)x_B \sin(\phi_h + \phi_S) F_{UT}^{(2)} && \xrightarrow{\text{magenta}} && \\
 & + |\vec{S}_\perp| (1 - y)x_B \sin(3\phi_h - \phi_S) F_{UT}^{(3)}/2 && \xrightarrow{\text{cyan}} && 1/P_t^5
 \end{aligned}$$

Transition from Perturbative region to Nonperturbative region?

- Compare different region of P_T



Nonperturbative TMD

Perturbative region

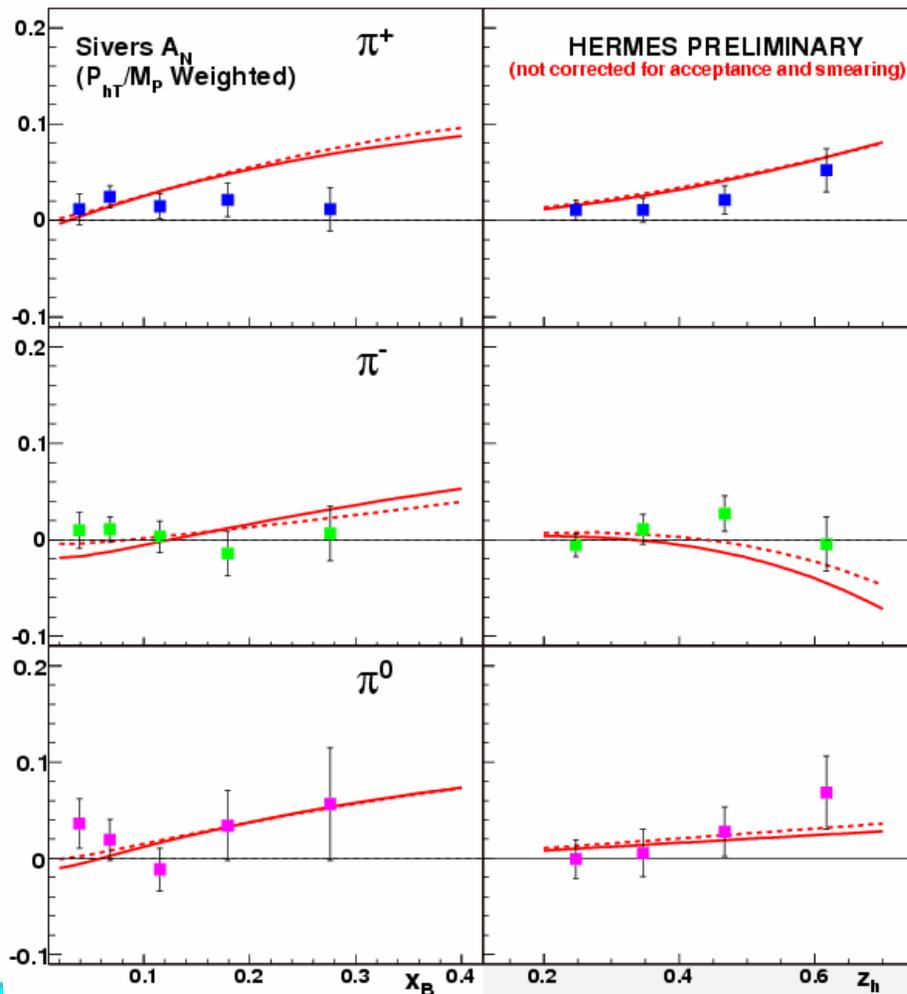
More over

- At low $P_{\perp} \ll Q$, the second term vanishes, use the Siverson function only
- P_{\perp} -moment of the asymmetry can be calculated from twist-three matrix element
 - In SIDIS, for the Siverson asymmetry

$$2 \left\langle \frac{P_{h\perp}}{M_P} \sin(\phi_h - \phi_S) \right\rangle = \frac{\int \frac{1}{Q^4} (1 - y + \frac{y^2}{2}) \frac{z_h}{M_P} x_B g_s T_F(x_B) D(z_h)}{\int \frac{1}{Q^4} (1 - y + \frac{y^2}{2}) x_B f_q(x_B) D(z_h)}$$

Boer-Mulders-Tangelmann, 96,98

Compare to the HERMES data



- T_F from the fit to single inclusive hadron data
 - [Kouvaris-Qiu-Vogelsang-Yuan](#), hep-ph/0609238
- This comparison is very nontrivial, because the SSA in DIS depends on final state interactions, whereas in hadronic collision both initial and final state interactions contribute
- Indicate the consistency of SSAs in DIS and hadron collisions

See also, Efremov, et al., PLB612,233 (2005)

A new way to study higher-twist quark-gluon correlation functions?

- Large p_t SIDIS certainly will provide information on higher-twist distributions
- P_t -weighted azimuthal asymmetries will also give constraints on these distributions
 - Systematic analysis has to be done
 - Evolution, NLO corrections, ...
 - It indeed opens a new window for studying higher-twist quark-gluon correlations in SIDIS