

Neutron Physics

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5 lectures:

1. Physics/Technology of Cold and Ultracold Neutrons
2. Electroweak Standard Model Tests [neutron beta decay]
3. Nuclear physics/QCD [weak interaction between nucleons]
4. Physics Beyond the Standard Model [EDM/T violation]
5. Other interesting stuff that neutrons can do [NNN interaction, searches for extra dimensions,...]

1: Physics/Technology of Cold and Ultracold Neutrons

Neutron sources [reactors/spallation] and neutron moderation

Low energy neutron interactions with matter: neutron optics

Polarized neutrons

“Ultracold” neutrons

Examples

Thanks for slides to: Geoff Greene (Tennessee/ORNL),
K. Bodek (PSI), Chen-Yu Liu (LANL), Jen-chieh Peng (Illinois)

What is “Neutron Physics”?

neutron physics (Fr. physique de neutronique):
a subset of scientific studies begun at the
Institute Laue/Langevin (ILL) in Grenoble, France
which the major users of the facility [condensed matter
physics/chemistry/biology/materials science] found
difficult to classify.

Known to practitioners as **fundamental neutron physics**

Research which uses “low energy” neutrons from nuclear
reactors and accelerator-driven spallation sources to
address questions in nuclear, particle, and astrophysics

“Fundamental” Neutron Physics...

is “Nuclear” physics, but with an “isotope of nothing”

is Particle Physics: but at an energy of 10^{-20} TeV

employs a particle which, according to Big Bang Cosmology, is lucky to be alive

relies on experimental techniques and ideas from atomic and condensed matter physics

is pursued at facilities built mainly for chemistry, materials science, and biology

Neutron Properties

Electric charge: $q_n=0$, electrically neutral [$q_n < 10^{-22}e$]

Size: $r_n \sim 10^{-5}$ Angstrom = 1 Fermi [area $\sim 10^{-25}$ cm² = 0.1 “barn”]

Internal Structure: quarks [ddu, $m_d \sim m_u \sim$ few MeV] + gluons

Spin: $s_n = 1/2$ [Fermi statistics]

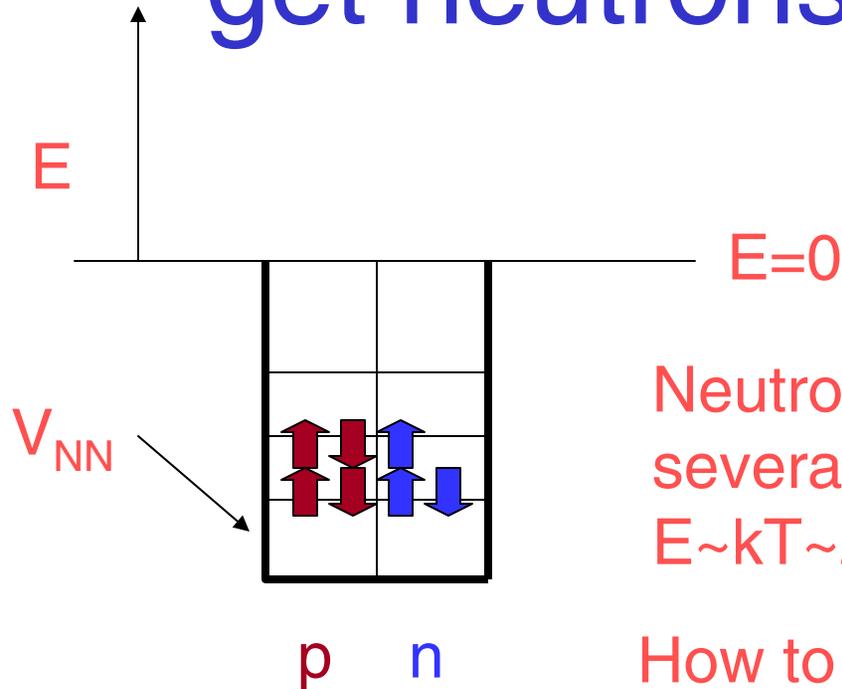
Magnetic Dipole Moment: $\mu_n / \mu_p = -0.68497935(17)$

Electric Dipole Moment: zero [$d_n < 10^{-26}$ e-cm]

Mass: $m_n = 939.566$ MeV [$m_n > m_p + m_e$, neutrons can decay]

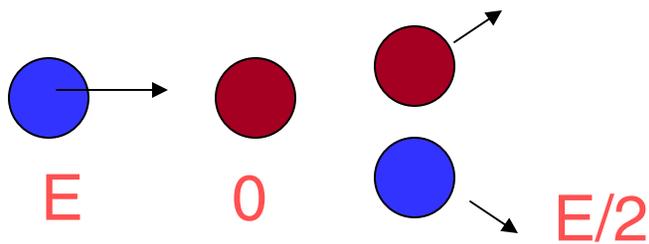
Lifetime: $\tau_n = 885.5 \pm 1.0$ seconds

Why is it such hard work to get neutrons to play with?



Neutrons are bound in nuclei, need several MeV for liberation. We want $E \sim kT \sim 25 \text{ meV}$ (room temperature)

How to slow down a heavy neutral particle with $M_n = M_p$? Lots of collisions...

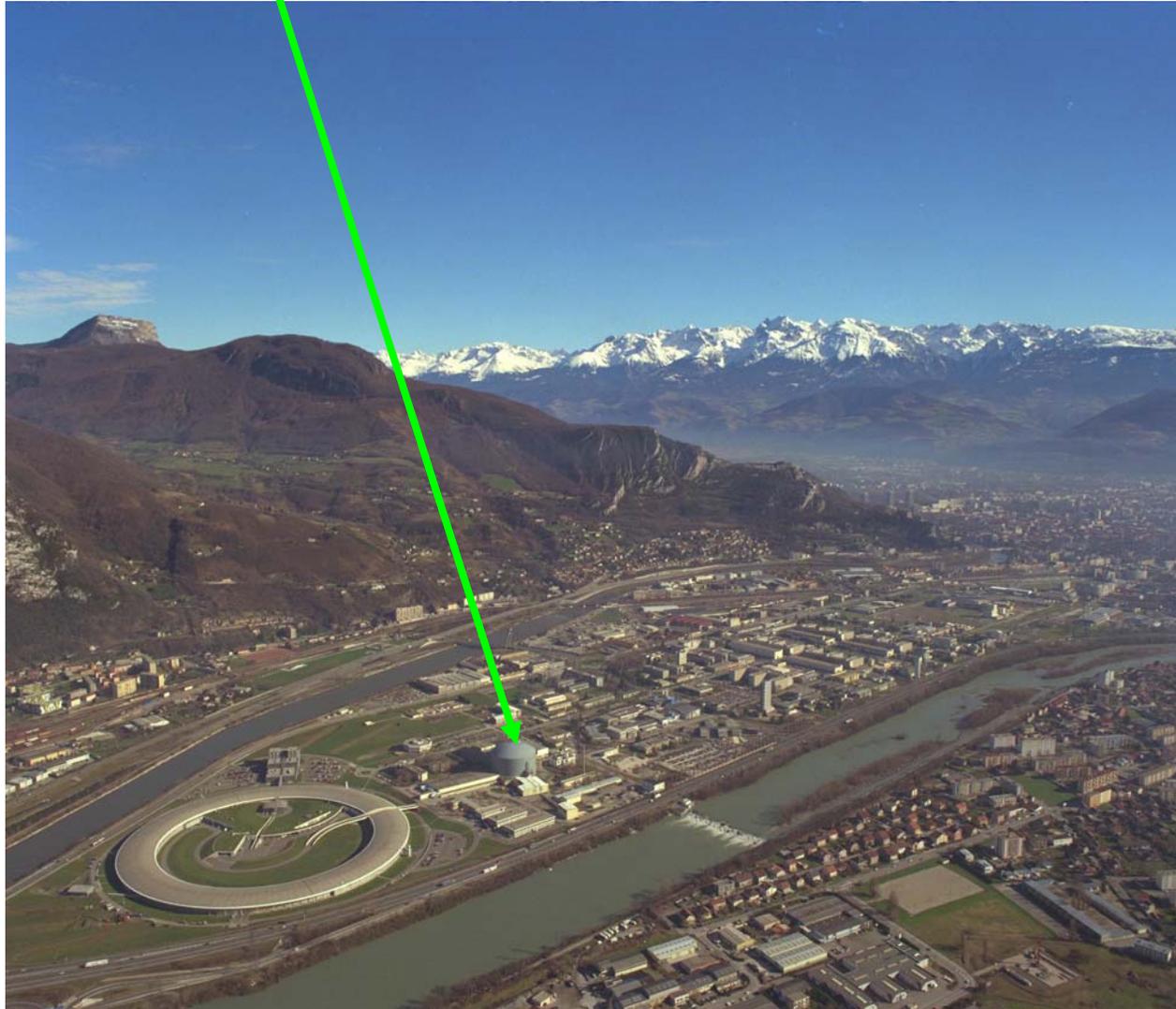


$$\left[\frac{1}{2}\right]^N = (1 \text{ MeV}) / (25 \text{ meV}) \text{ for } N \text{ collisions}$$

Neutrons are unstable when free \rightarrow they can't be accumulated easily

ILL, Grenoble

60 MW reactor

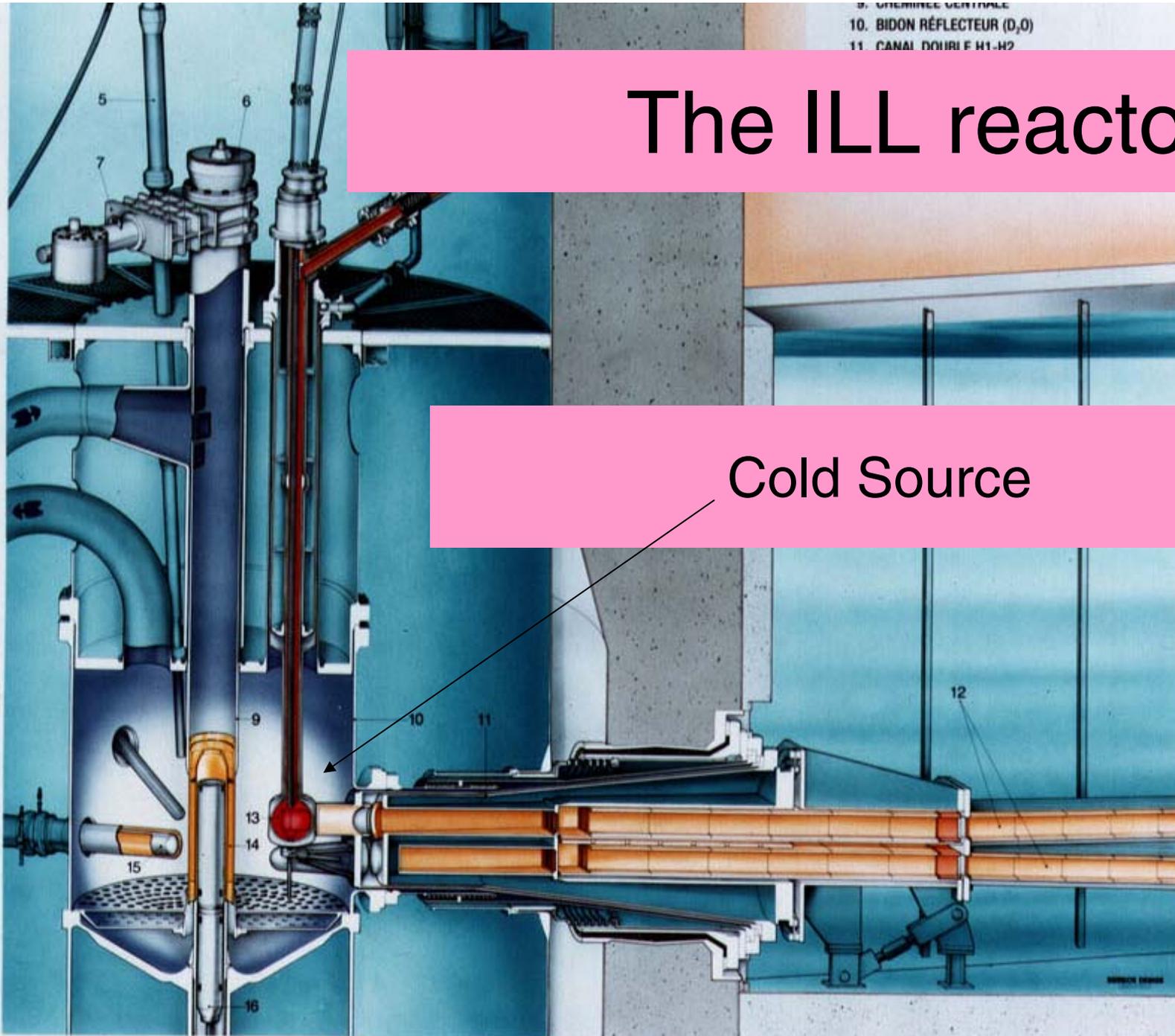


World's most intense source of neutrons for scientific research since 1973

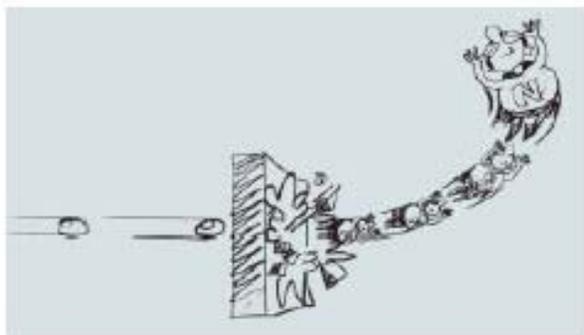
France, UK, Germany founding members

The ILL reactor

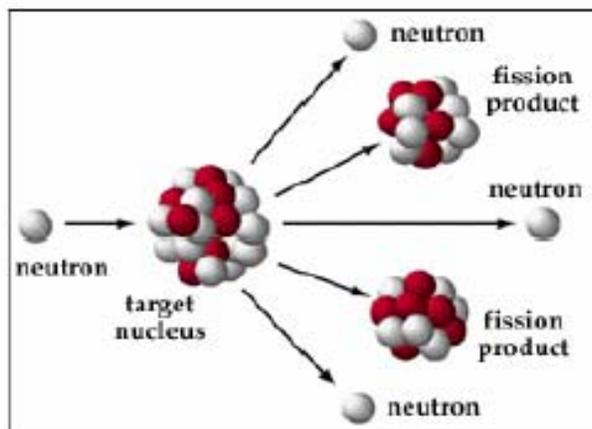
Cold Source



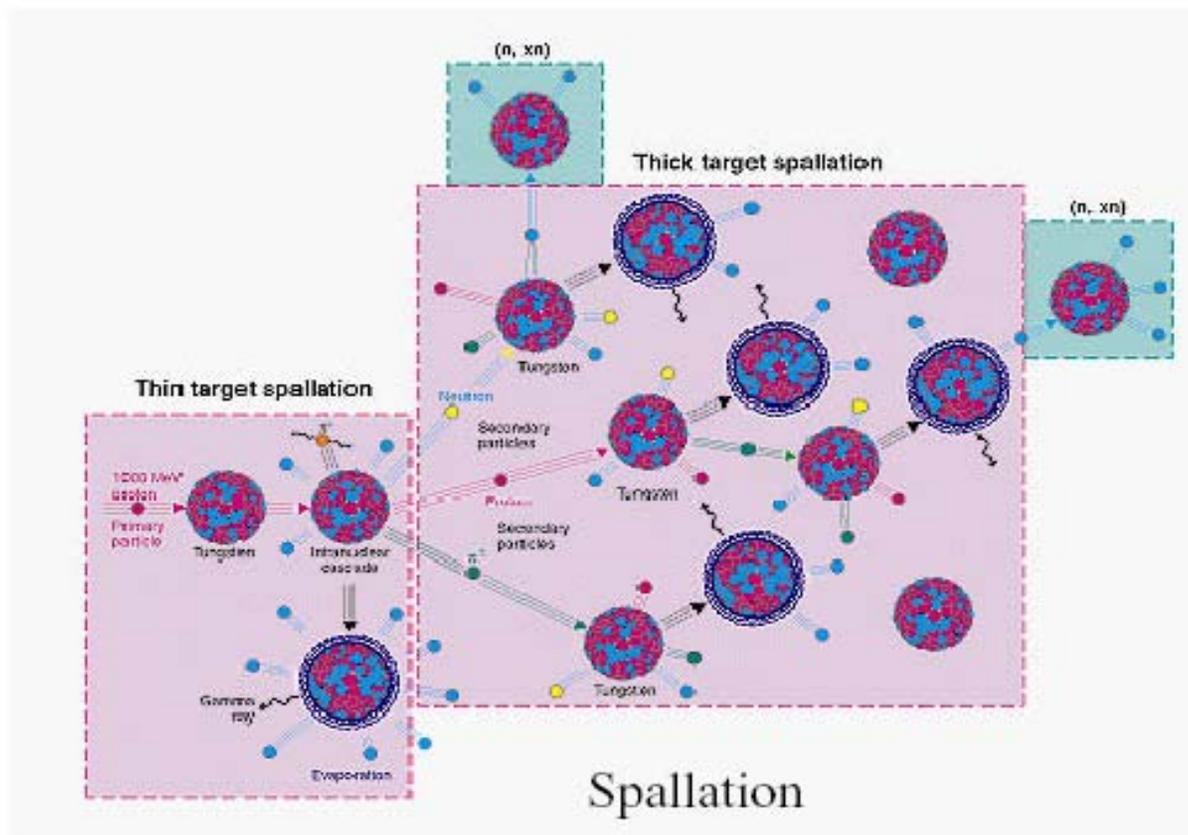
About 1.5 Useful Neutrons Are Produced by Each Fission Event in a Nuclear Reactor Whereas About 25 Neutrons Are Produced by spallation for Each 1-GeV Proton Incident on a Tungsten Target



Artist's view of spallation



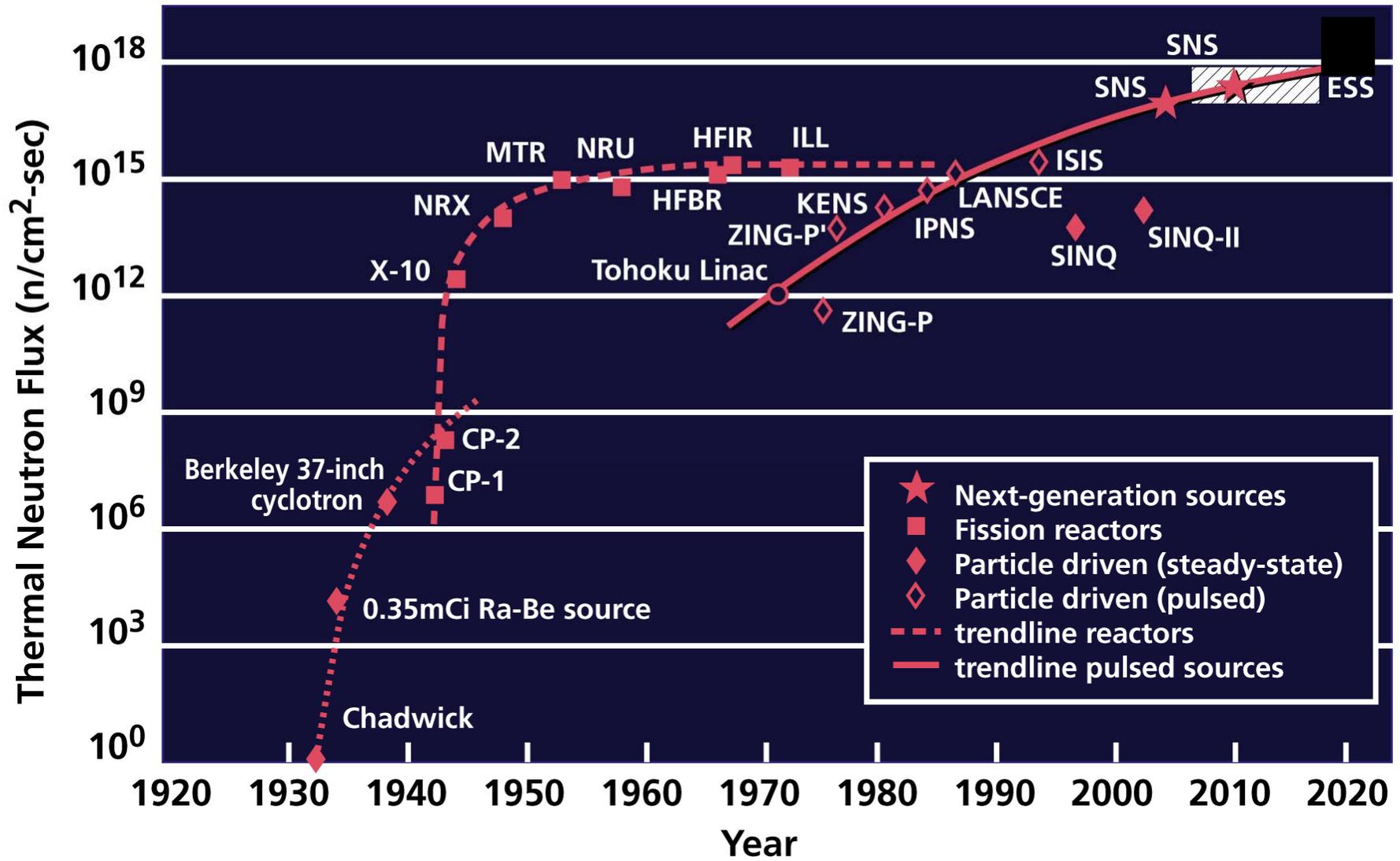
Nuclear Fission



From R. Pynn, UCSB

Reactors are limited by heat removal from the core

Pulsed sources have not yet reached that limit



(Updated from *Neutron Scattering*, K. Skold and D. L. Price: eds., Academic Press, 1986)

The Spallation Neutron Source



SNS Site Overlay

- ◆ \$1.4B--1GeV protons at 2MW, ready in 2007.
- ◆ Short (~1 usec) pulse— mainly for high TOF resolution
- ◆ >10 times the power of ISIS

The Spallation Neutron Source

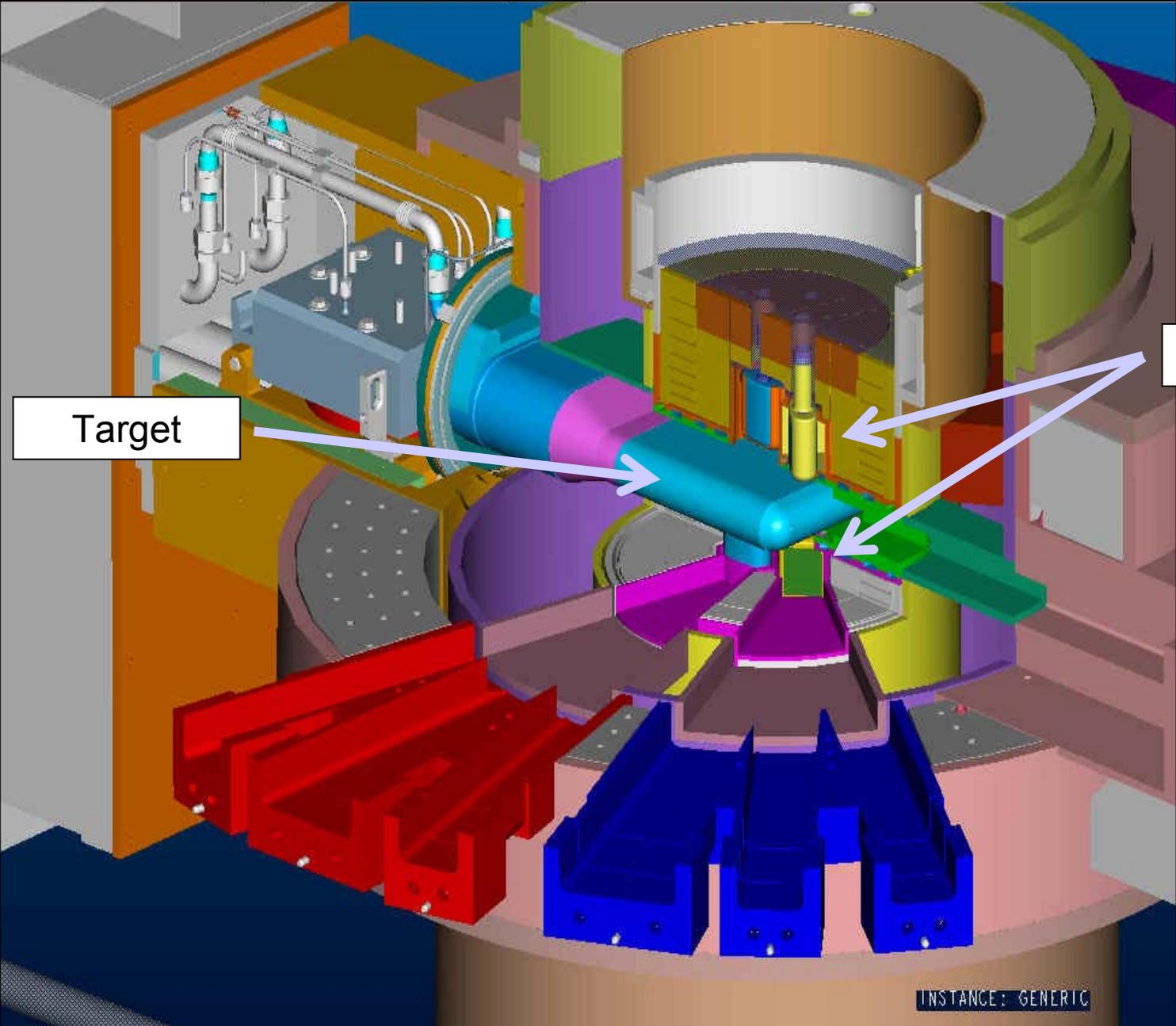
Photo, Oct 2003



SNS Construction: on schedule for 2007 Operation



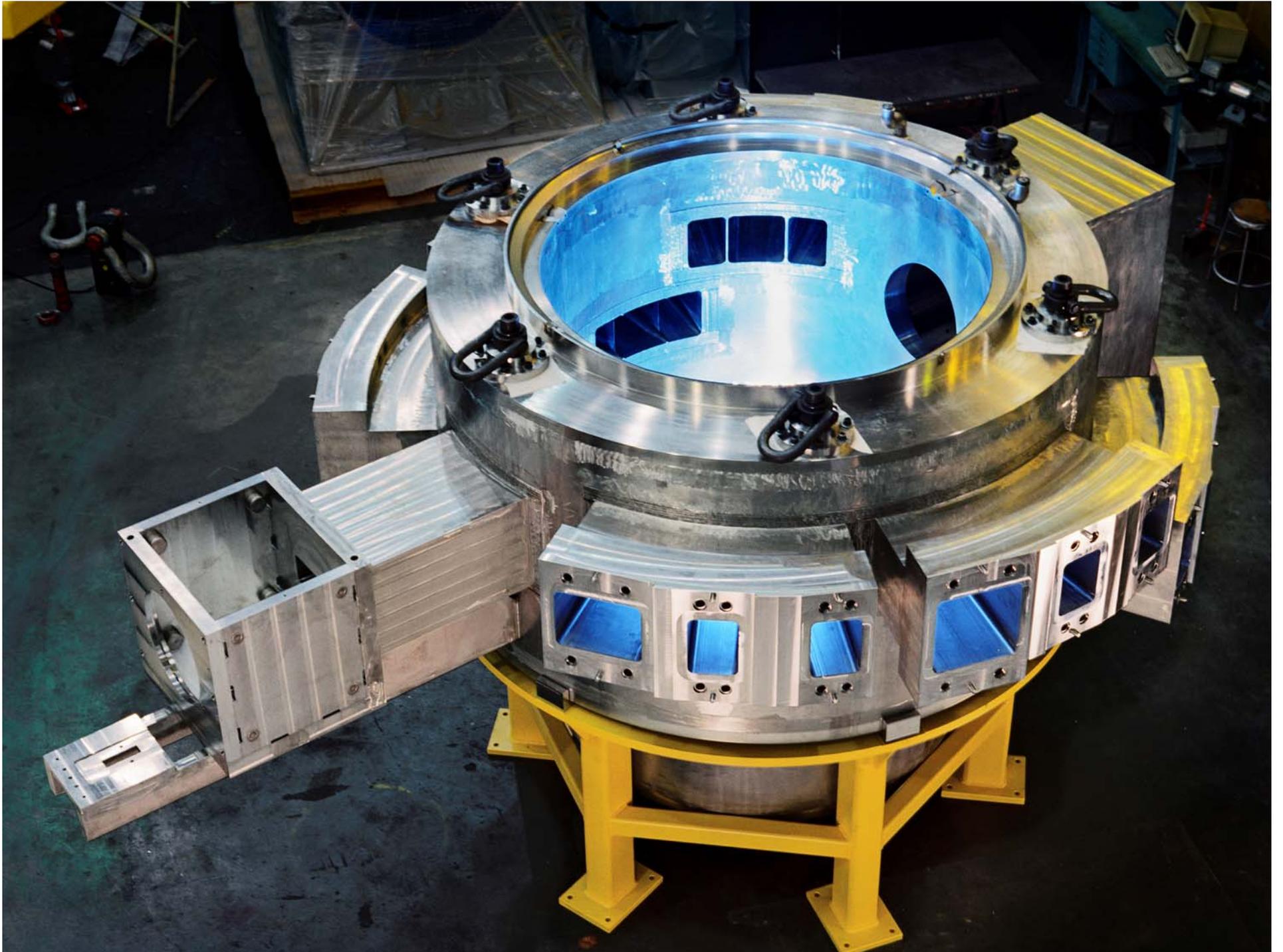
photos, April 2003



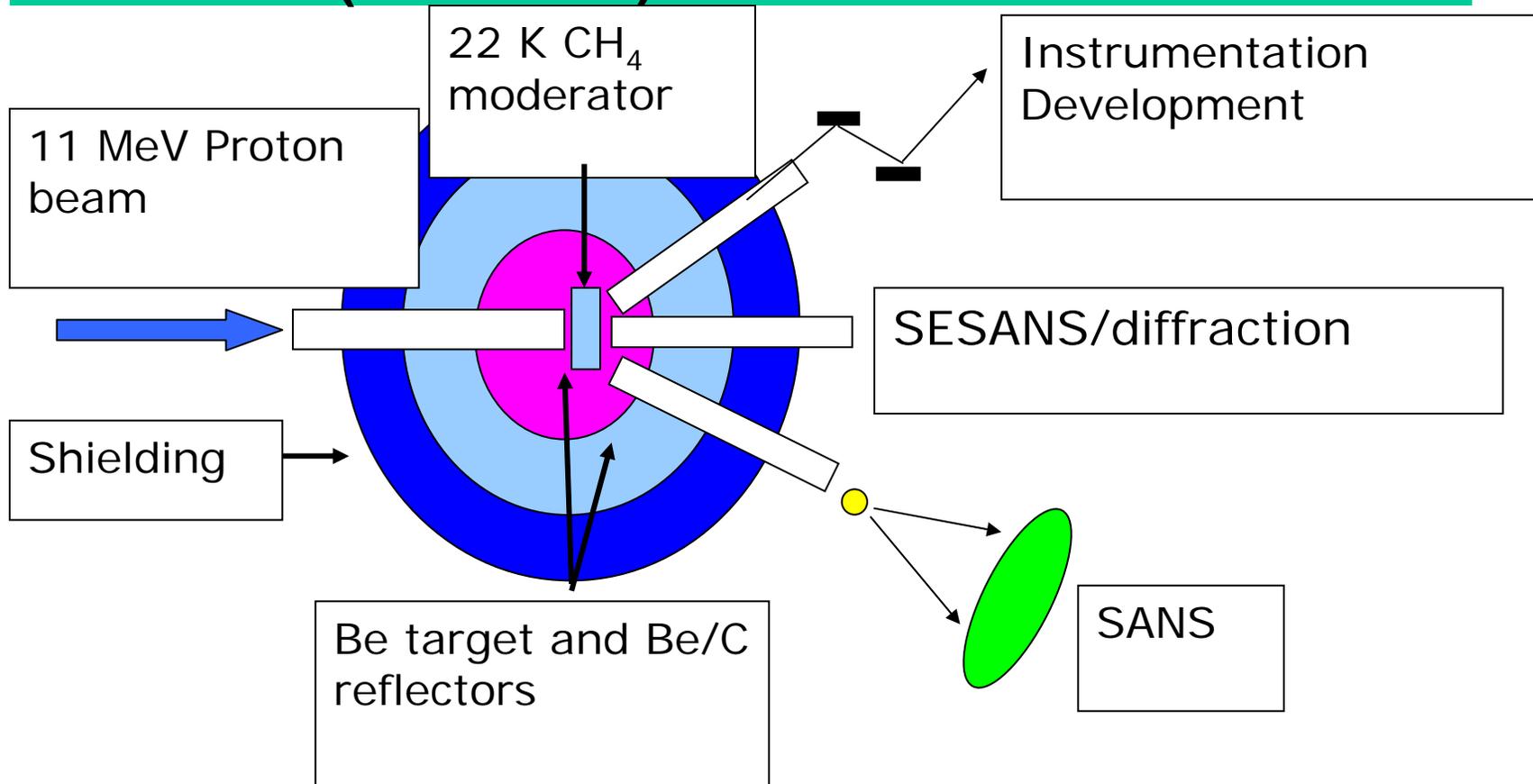
Target

Moderator

INSTANCE: GENERIC

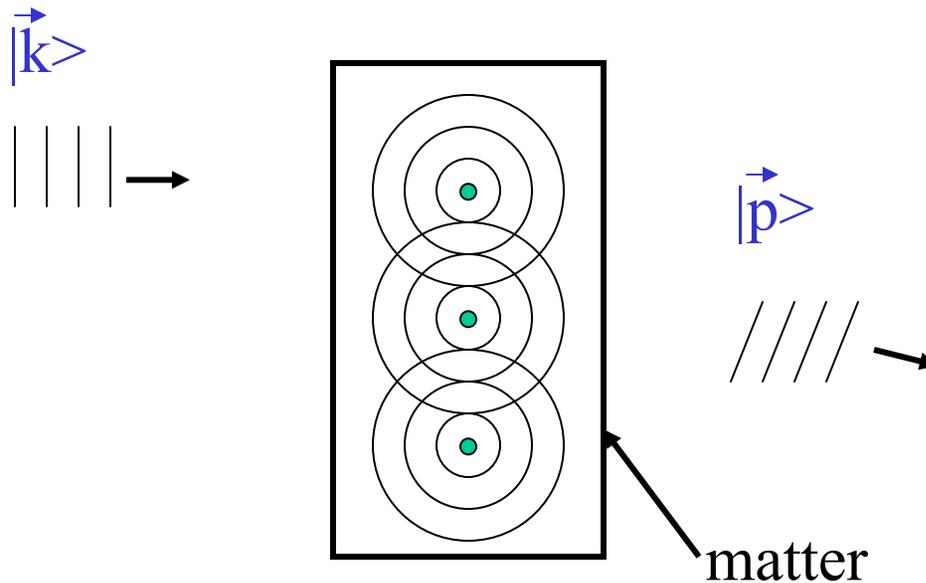


Low Energy Neutron Source (LENS) at Indiana



First university-based pulsed cold neutron source!

Neutron-Matter Interaction



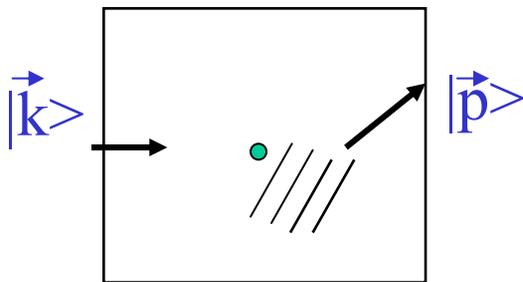
Thermal neutrons: $\lambda \sim 2\text{\AA}$
 $R_{\text{nuc}} \sim 10^{-5} \text{\AA}$
 $kR \gg 1 \rightarrow$ elastic scattering
 $f_{\text{scatt}} = b =$ s-wave scattering length
 b can be positive or negative

Inelastic scattering: $|\vec{k}| \neq |\vec{p}|$, energy transferred to localized excitations in matter (phonons, magnons,...)

Elastic scattering: $|\vec{k}| = |\vec{p}|$, energy transferred to object as a whole \rightarrow quantum mechanical state of matter unchanged...

Neutrons in Condensed Matter

for a “thermal” neutron ($E_K = mv^2/2 = 3/2 k_B T$, $T = 300\text{K} \rightarrow E_K = 25\text{ meV}$)
the de Broglie wavelength of the neutron is $\lambda \approx 2\text{ Angstroms}$



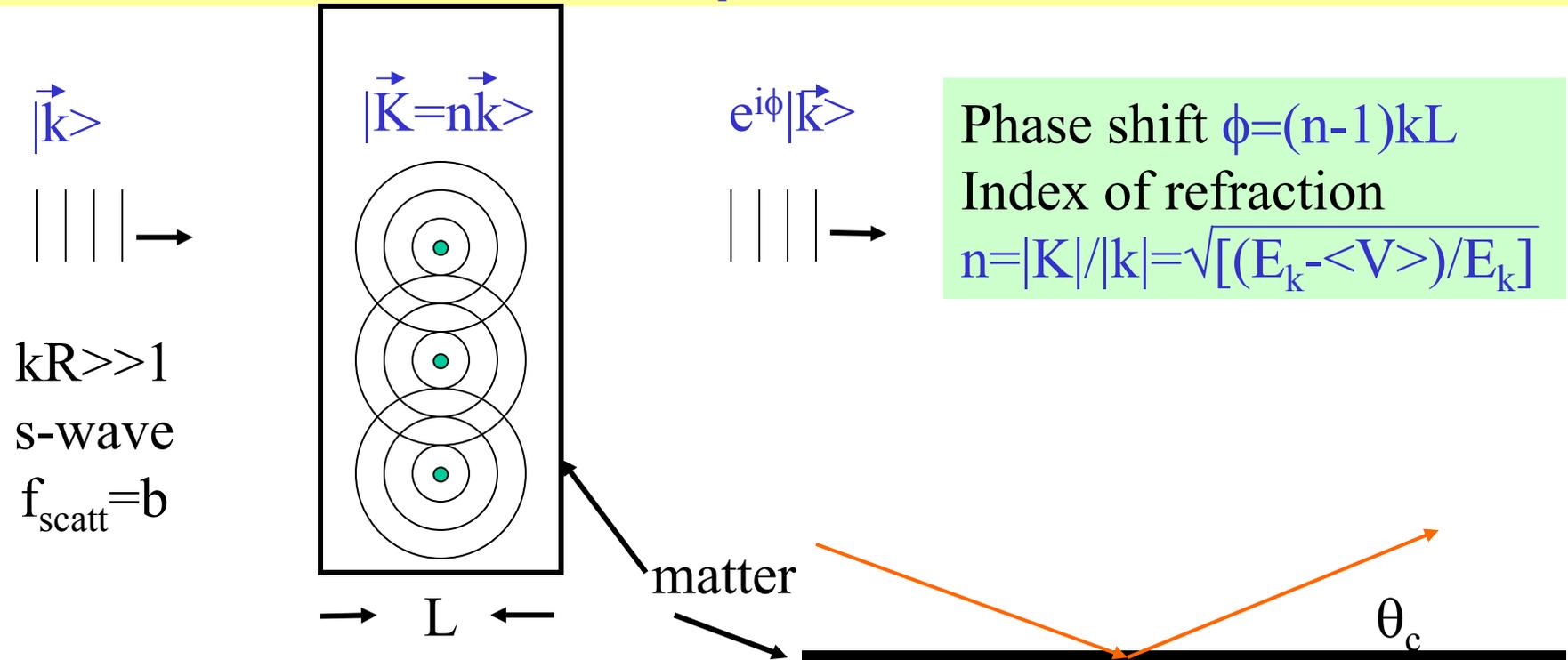
phonon



diffraction

Thermal neutrons have the right energies and momenta to match excitations (phonons, spin waves, molecular rotations...) and **static structures** (crystals, molecular shapes,...) in condensed media

Neutron Optical Potential



$$\theta_c = \lambda \sqrt{[\rho b / \pi]} \text{ critical angle}$$

$$\langle V_{\text{strong}} \rangle = 2\pi h^2 \rho b_s / m, \sim \pm 100 \text{ neV}$$

$$\langle V_{\text{mag}} \rangle = \mu B, \sim \pm 60 \text{ neV/Tesla}$$

$$\langle V_{\text{grav}} \rangle = mgz \sim 100 \text{ neV/m}$$

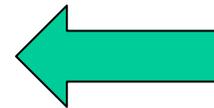
$$\langle V_{\text{weak}} \rangle = [2\pi h^2 \rho b_w / m] \vec{s} \cdot \vec{k} / |k| \sim 10^{-7} \langle V_{\text{strong}} \rangle$$

For $E_k - \langle V \rangle$ negative,
neutron reflects
from the optical potential



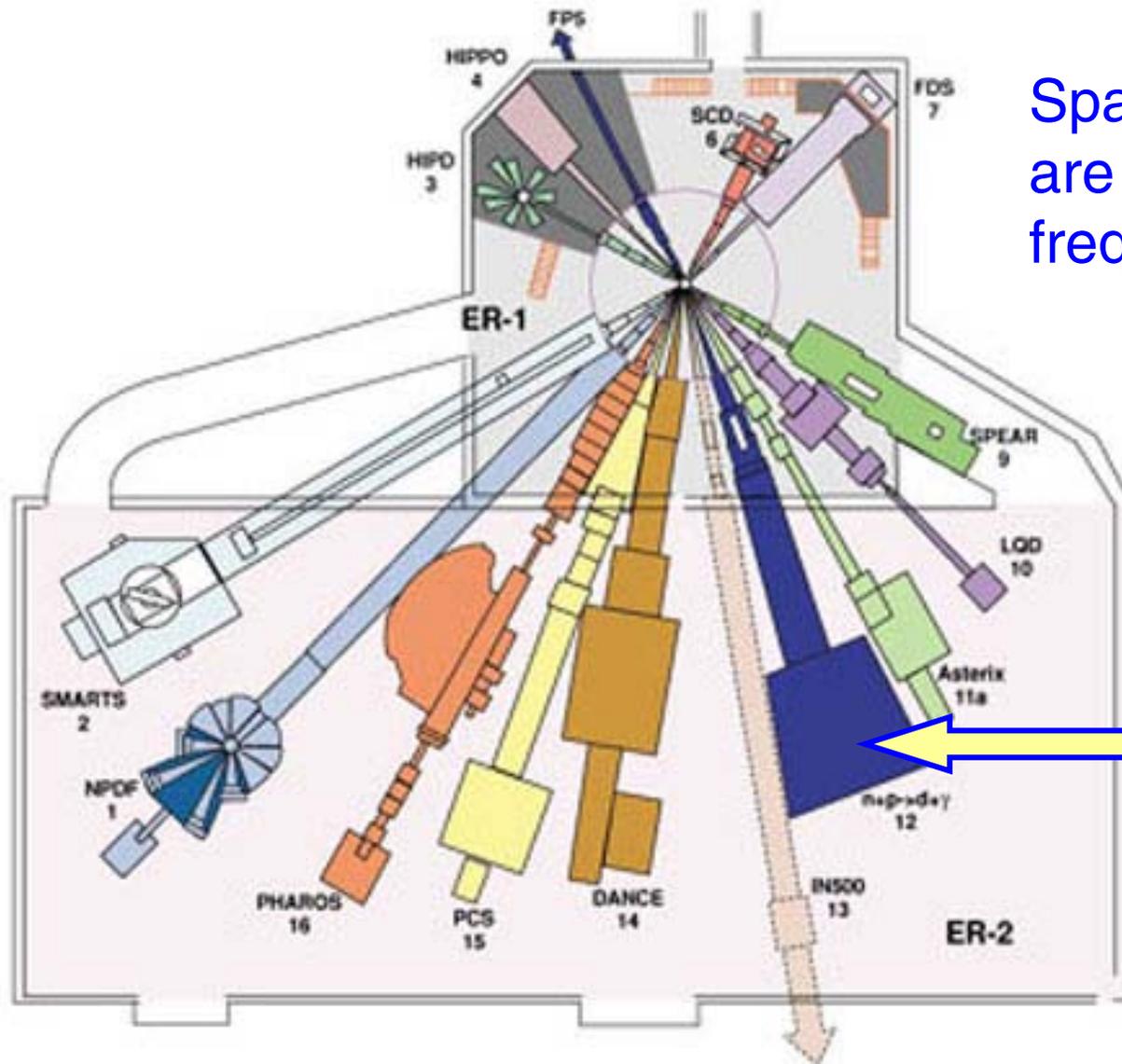
NIST Cold Neutron Guide Hall

Reactor makes neutrons, cooled to $\sim 20\text{K}$ by liquid hydrogen



Neutron mirrors (“guides”) conduct the neutrons ~ 100 meters with small losses.

LANSCCE Guide Hall

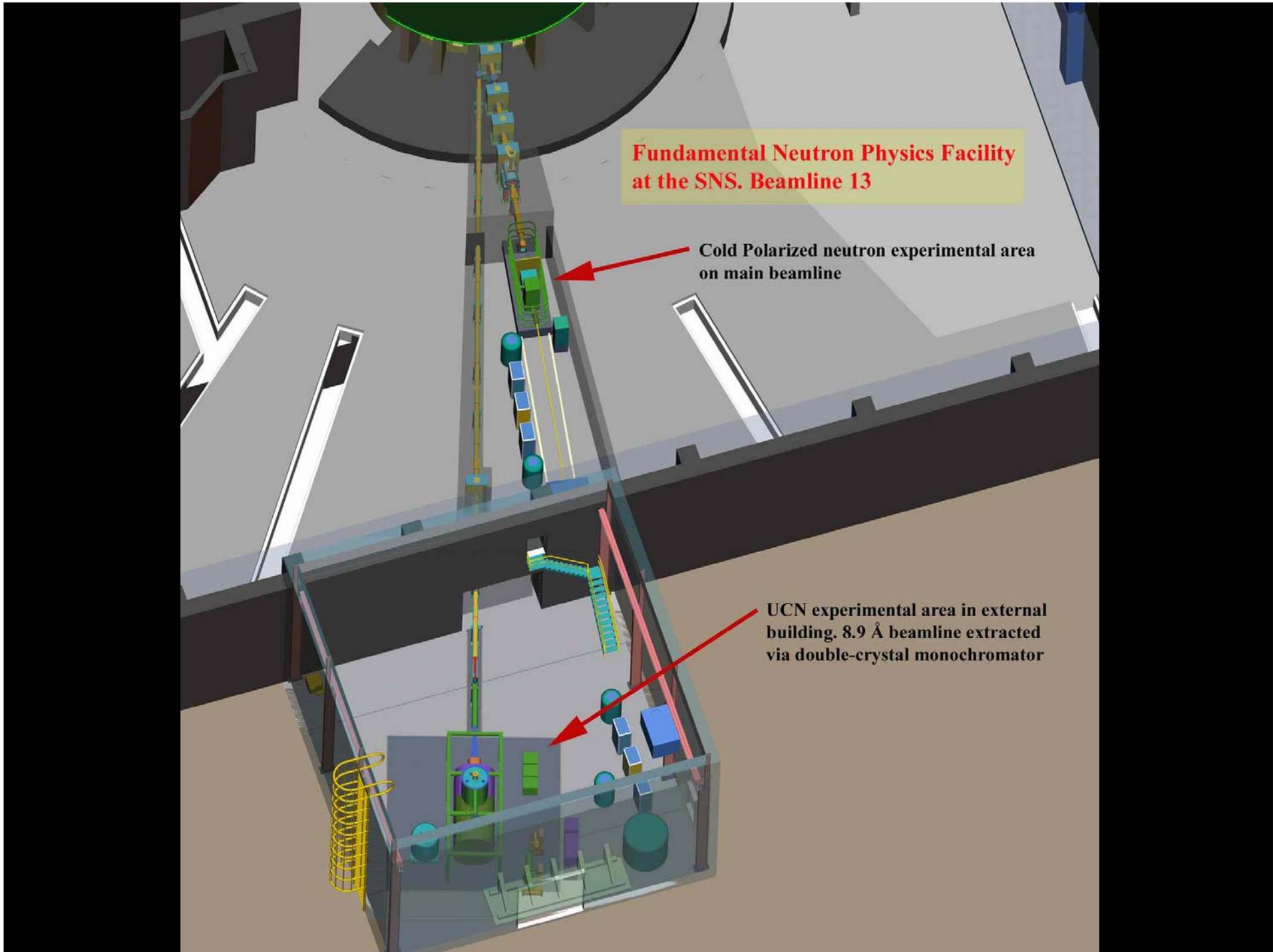


Spallation neutron sources are typically pulsed at frequencies of 10-60 Hz

**Fundamental Neutron Physics Facility
at the SNS. Beamline 13**

Cold Polarized neutron experimental area
on main beamline

UCN experimental area in external
building. 8.9 Å beamline extracted
via double-crystal monochromator



Comparison of Source Brightness

	time-averaged	peak
LANSCCE @80kW	2.5	~1250
SNS @1.4MW	~70	~12000
NIST @20MW	~150	~150
ILL @57MW	~450	~450
HFIR @80MW	~450	~450

From G. Greene, UT/ORNL

Cold Neutron Beams @ SNS

Capture Flux:

The estimated density of neutrons in an SNS cold beam will be: $\rho \approx 10^4 n \cdot cm^{-3}$

Absolute Flux

The estimated flux of neutrons in an SNS cold beam will be: $d^2 / dA dt \approx 10^9 n \cdot cm^{-2} \cdot s^{-1}$

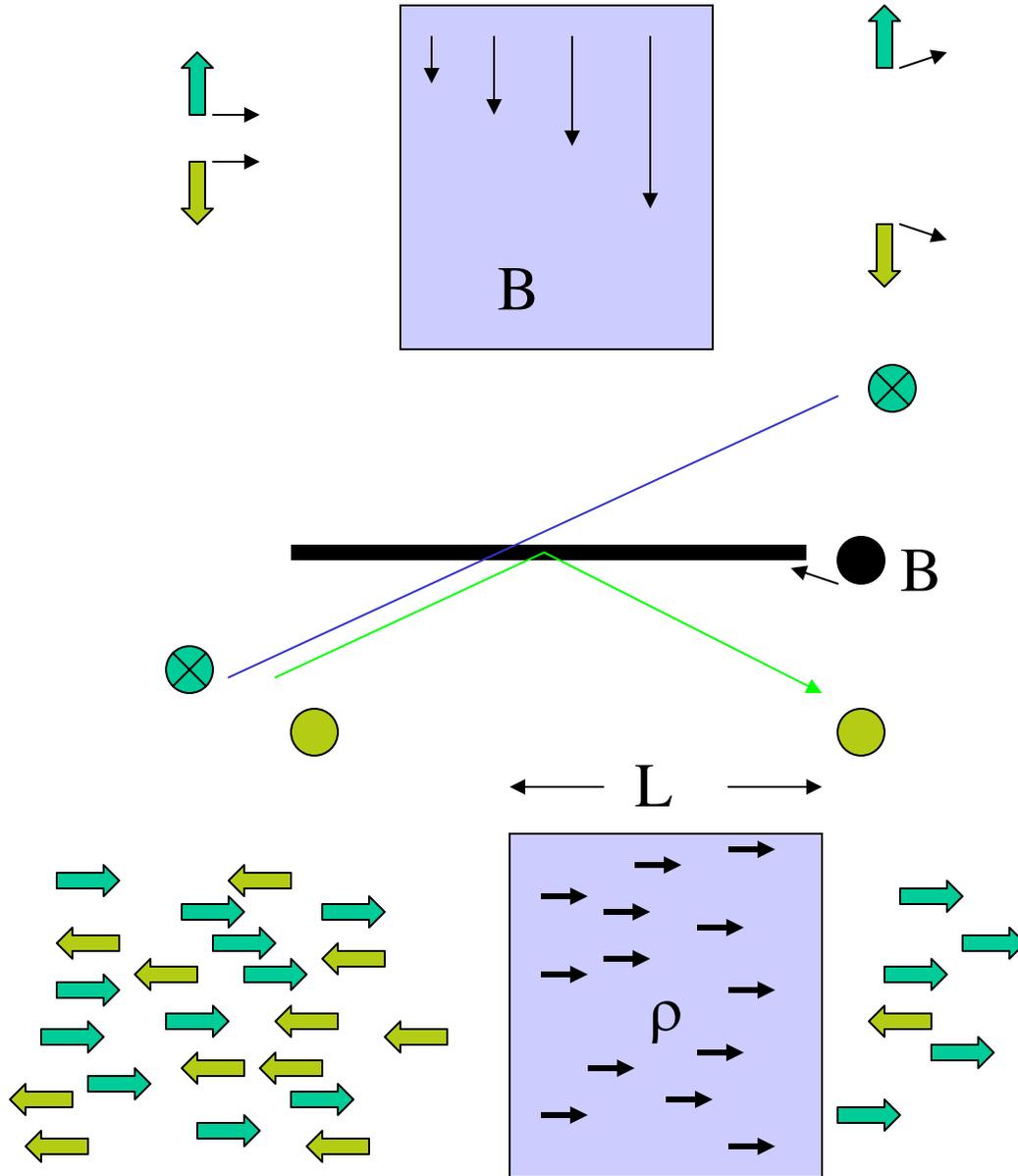
Fluence

The estimated fluence of neutrons in an SNS cold beam will be: $d^2 / dt \approx 10^{11} n \cdot s^{-1}$

How to Polarize Neutrons?

Force of nature	Does it depend on neutron spin?	How big is it?
Gravity	Yes if the other object has spin	Very, very small
Weak interaction (causes beta decay)	Yes, even if the other object has no spin	Only very small (=measurable)
Electromagnetism	Yes, due to magnetic moment in B	$\mu \cdot B \sim 300 \text{ neV/T}$
Strong interaction	Yes	$\sim 100 \text{ neV}$ (varies with nucleus)

What methods have been used?

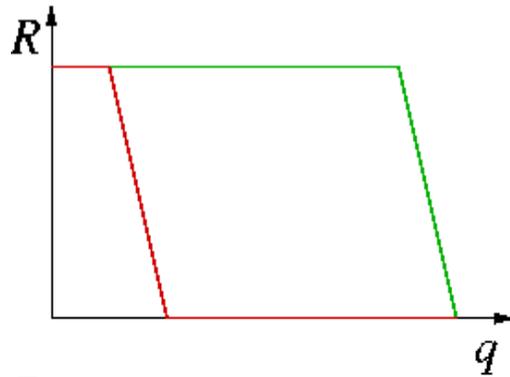


B gradients (Stern-Gerlach,
sextupole magnets)
electromagnetic
 $F=(\mu \bullet \nabla)B$

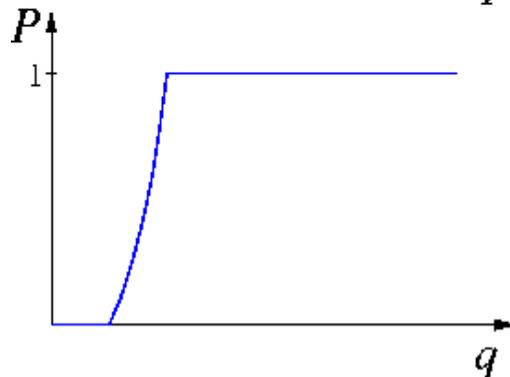
Reflection from magnetic
mirror: electromagnetic+
strong
 $f_{\pm}=a(\text{strong}) \pm a(\text{EM})$
with $|a(\text{strong})|=|a(\text{EM})|$
 $\Rightarrow f_{+}=2a, f_{-}=0$

Transmission through
polarized nuclei: strong
 $\sigma_{+} \neq \sigma_{-} \Rightarrow T_{+} \neq T_{-}$
Spin Filter: $T_{\pm}=\exp[-\rho \sigma_{\pm} L]$

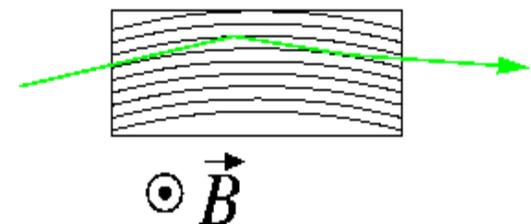
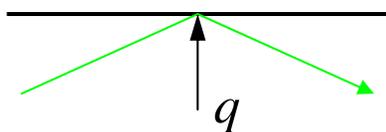
Polarisers for Neutrons: Super Mirrors



Magnetized mirror has large optical potential for one spin state, small for other spin state
-> spin states possess different critical angles for total reflection



To cover the whole beam the mirrors are stacked so that every neutron in the beam reflects once



FUNSPIN – Polarized Cold Neutron Facility at PSI

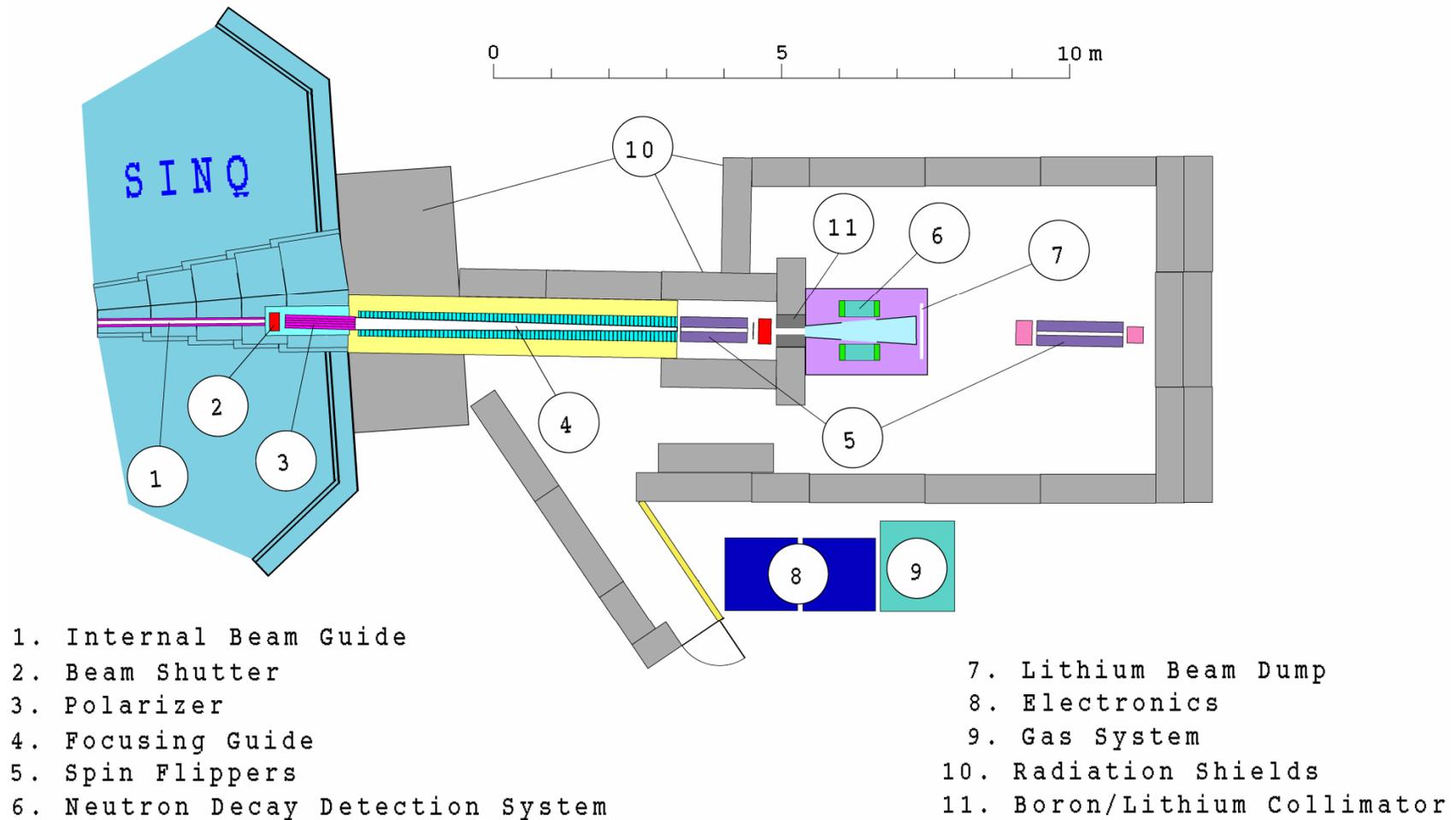


Figure 4: Layout of the Polarized Cold Neutron Facility at PSI.

FUNSPIN Polarized Cold Neutron Facility at PSI

Large momentum acceptance neutron beam guide

Super-mirrors:	450 Ni/Ti layers, $m \approx 3$
Flux density (unpolarized):	$\approx 10^9 \text{ cm}^{-2}\text{s}^{-1}\text{mA}^{-1}$
Thermal equivalent:	$\approx 3 \times 10^9 \text{ cm}^{-2}\text{s}^{-1}\text{mA}^{-1}$
Total unpolarized:	$\approx 10^{11} \text{ s}^{-1}\text{mA}^{-1}$
Polarizer:	multi-slit, magn. super-mirrors
Focusing section:	$8 \times 15 \text{ cm}^2 \rightarrow 4 \times 15 \text{ cm}^2$
Flux density (polarized):	$\approx 2 \times 10^8 \text{ cm}^{-2}\text{s}^{-1}\text{mA}^{-1}$
Total polarized:	$\approx 1.2 \times 10^{10} \text{ s}^{-1}\text{mA}^{-1}$
Average polarization:	$> 97\%$
Spin guiding field:	permanent magnets
Equipment:	RF-spin flippers, beam chopper

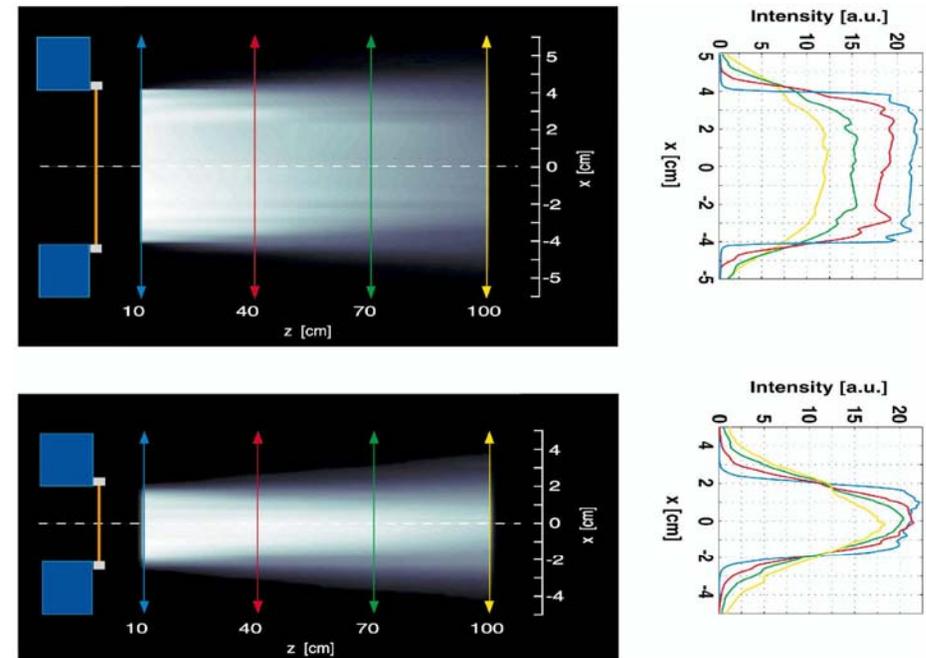


Figure 6: Beam profiles (right) and a representation of the beam intensity obtained by interpolation (left). Upper panel: unpolarized beam (SINQ insert). Lower panel: polarized beam (experimental position).

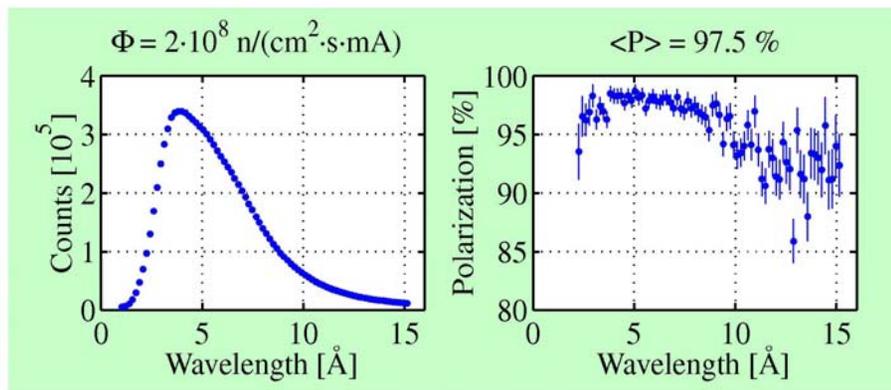
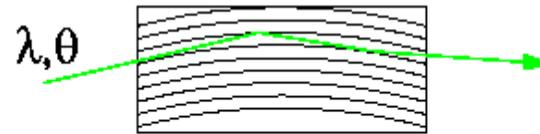
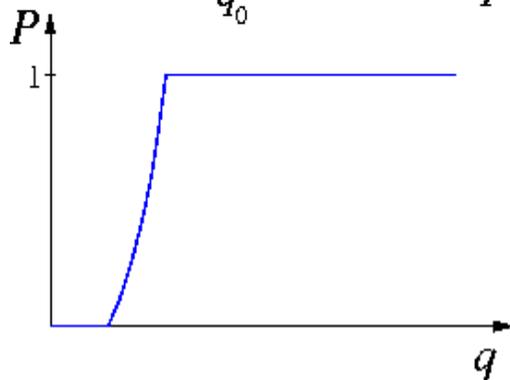
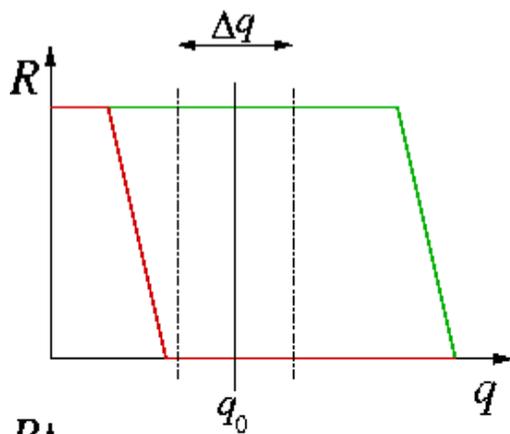


Figure 5: (a) Neutron wavelength spectrum measured using the TOF technique and a thin ³He transmission detector. (b) Wavelength dependence of the beam polarization.

$$P_{av} = (89.9 \pm 0.8) \%$$

(over fiducial volume of the beam)

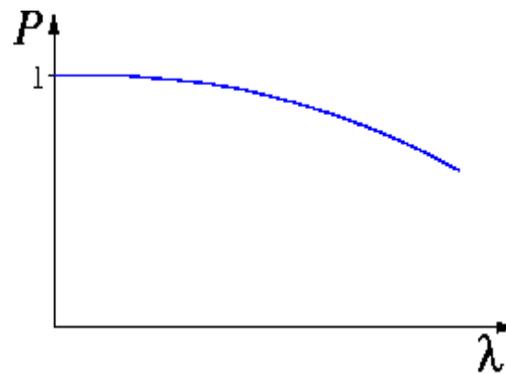
Problems with Supermirror Polarizers



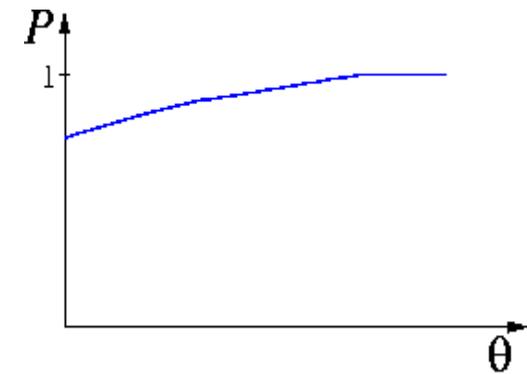
$$\odot \vec{B}$$

$$\Delta q = \Delta q(\lambda, \theta)$$

$$q_0 = q_0(\lambda, \theta)$$



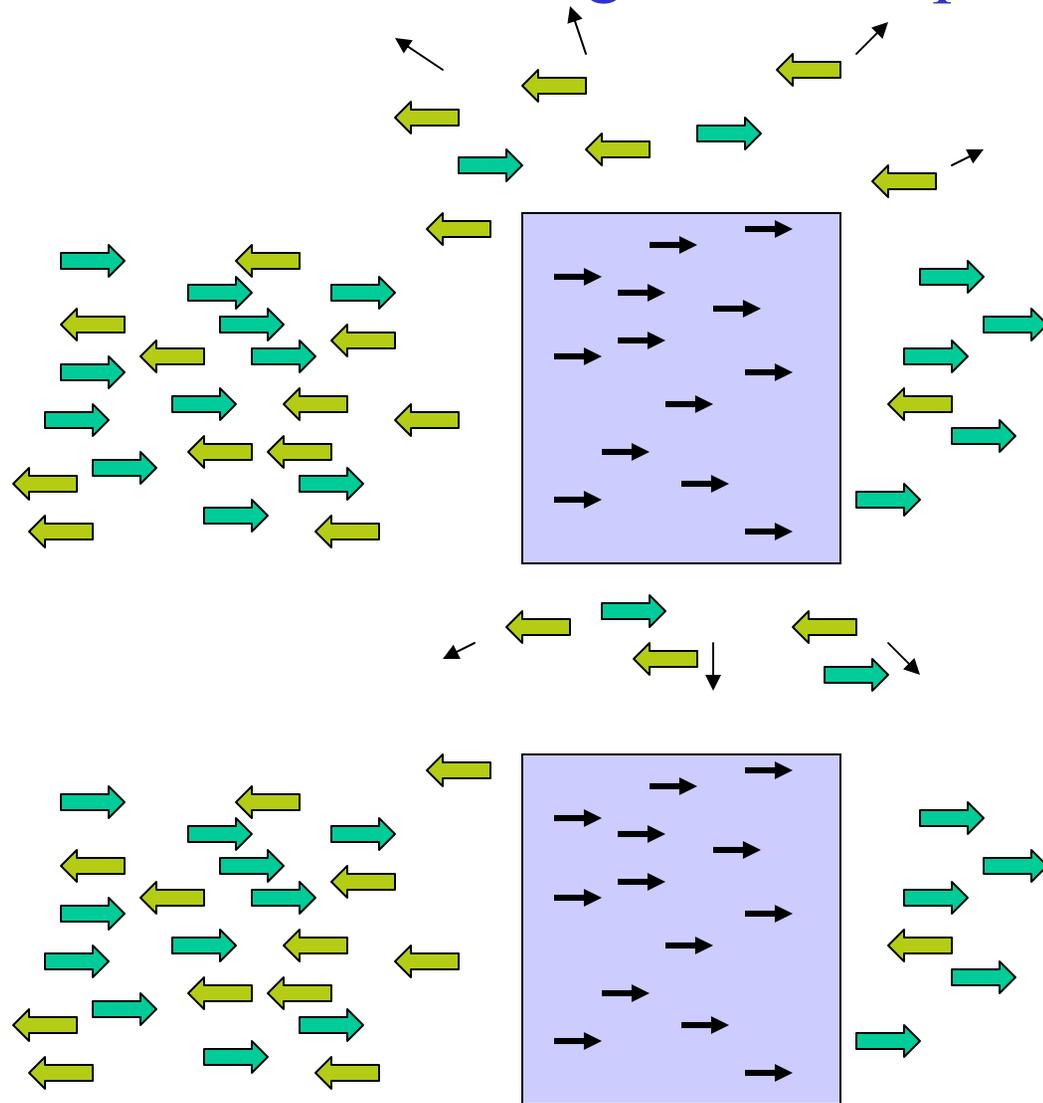
⇒ Wavelength
dependence



⇒ Angular
dependence

What Can Cause $\sigma_+ \neq \sigma_-$?

$$\sigma_{\text{total}} = \sigma_{\text{scattering}} + \sigma_{\text{absorption}}$$



Scattering: hard to get
 $\sigma_+ \text{scatter} \gg \sigma_- \text{scatter}$
(Hydrogen is good)

Absorption: can be very
large at resonance for one
spin state

$\sigma_+ \text{absorb} \gg \sigma_- \text{absorb}$ easier

Reaction products:

- (1) Gammas (almost all)
- (2) Charged particles
- (3) ^3He , ^6Li , ...

Properties of a perfect polarizer

Formula for σ_{\pm} for neutron (spin=1/2) + nucleus (spin=I, P_N):

$$\sigma_{\pm} = \sigma_0(1 \pm \rho P_N), \quad \rho = \{-1 \text{ for } J=I-1/2 \dots, I/(I+1) \text{ for } J=I+1/2\}$$

\Rightarrow we want a $J=I-1/2$ resonance so that $\sigma_- = 2\sigma_0$ and $\sigma_+ = 0$

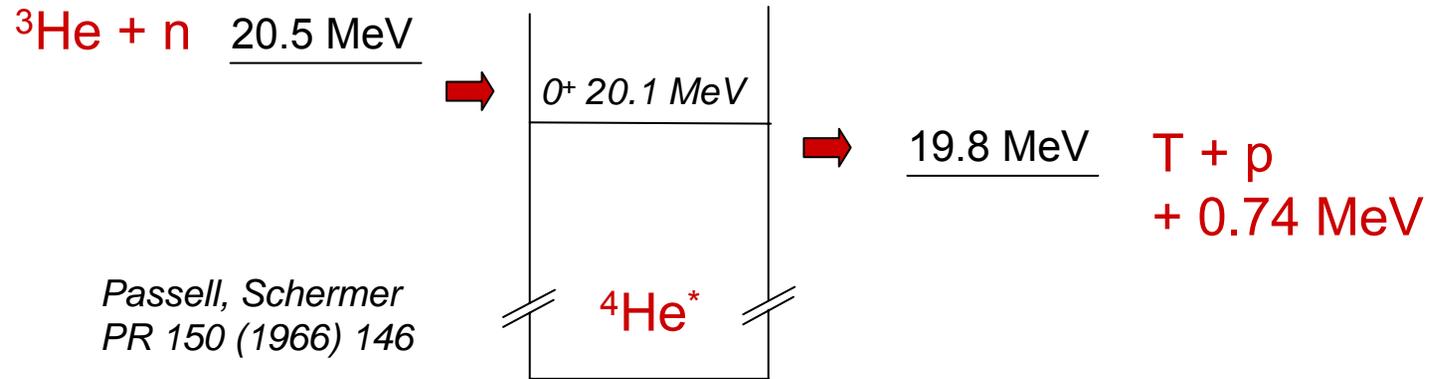
Once we polarize the nucleus, we want it to stay polarized!

An object with $I=1/2$ cannot have a quadrupole moment (quantum mechanics) \rightarrow it does not feel tensor fields (electric field gradients, ...). Closed electron shells are also good to isolate nucleus from the external world

\Rightarrow we want $I=1/2$, noble gas element

Final wish list: we want a noble gas element with an $I=1/2$ nucleus with large absorption resonance in $J=0$ channel and charged particles, not gammas, for reaction products

3He is Almost Perfect!



Noble gas

$I=1/2$

$n+{}^3\text{He}\rightarrow{}^4\text{He}^*(J=0 \text{ resonance!})$

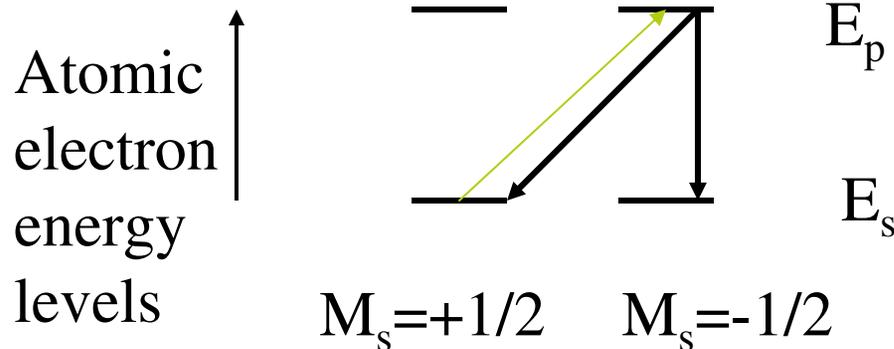
$\sigma_{-} = 2 \sigma_0 = \sim 10,000$ barns at 25 meV

$\sigma_{+, \text{absorb}} = \sim 0$

$\sigma_{+, \text{scatter}} = \text{few barns}$

$\sigma_{\gamma} = \text{few microbarns}$

Polarize by “Optical Pumping”



One photon is absorbed

Only $\Delta M = \pm 1$ allowed

Atom in excited state can decay back to either ground state

Circularly polarized photons, $S=1$, on atomic resonance

$E_p - E_s$

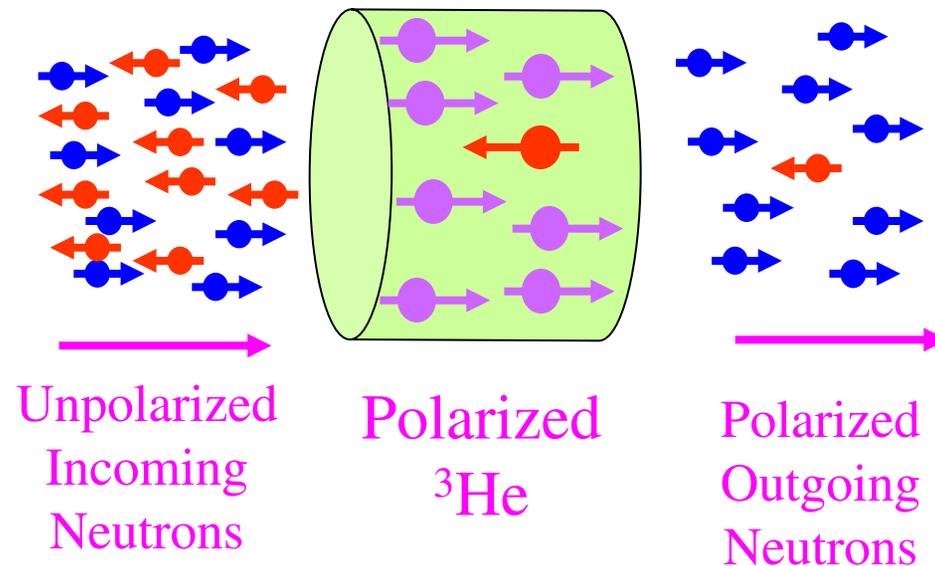
Keep absorbing photon again and again, eventually all ground state atoms in $M_s = -1/2$ substate

Finally: hyperfine interaction between electron and nucleus polarizes nucleus

-> electron is polarized!

Polarized ^3He as Neutron Spin Filter

- **strongly spin-dependent** neutron absorption cross section.



$$P_n = \frac{T^+ - T^-}{T^+ + T^-} = \tanh[\sigma(\lambda)N_{He}LP_{He}]$$

$$T_n = T_E \exp(-\sigma N_{He}L) \cosh(\sigma N_{He}LP_{He})$$

$$= \sqrt{1 - \left(\frac{T_0}{T_n}\right)^2}$$

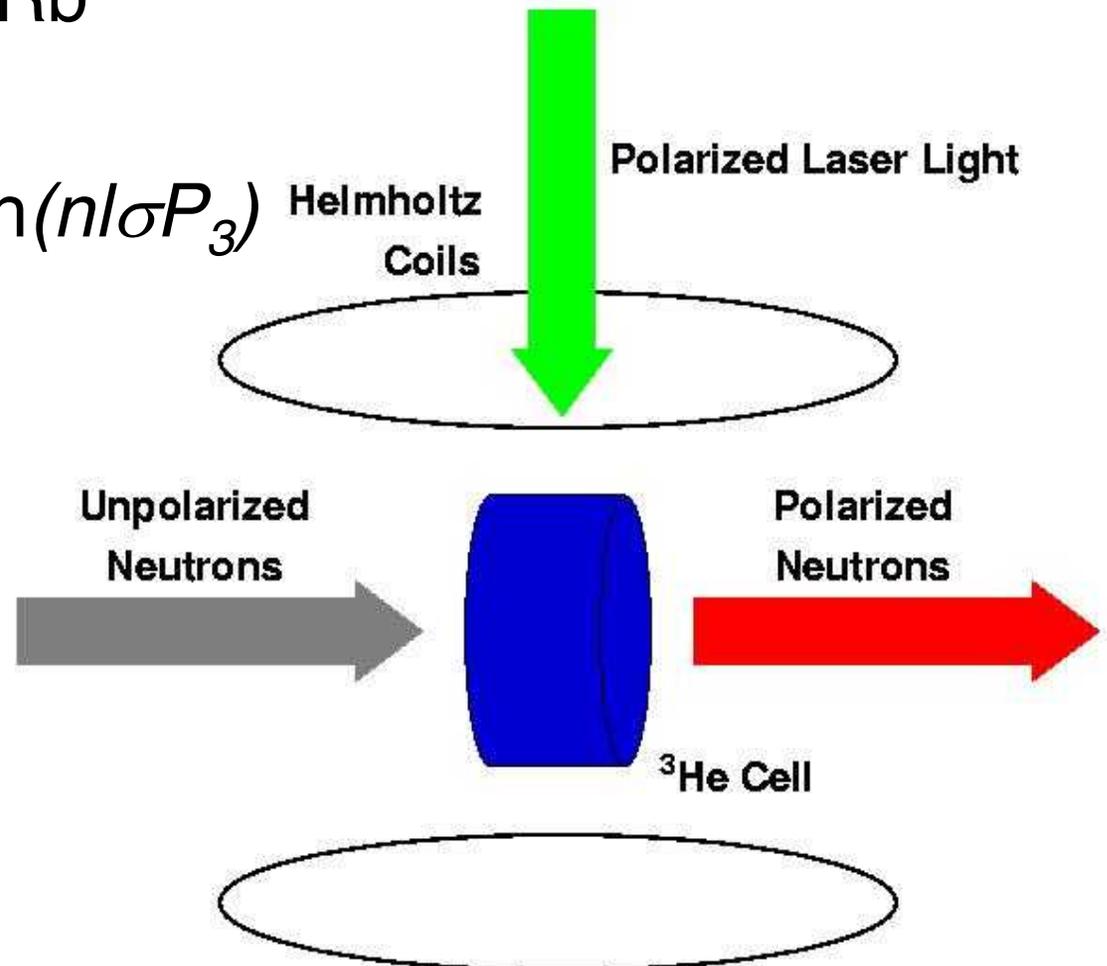
Neutron Polarizer: Atomic Physics

- ◆ Polarized ^3He Neutron Spin Filter

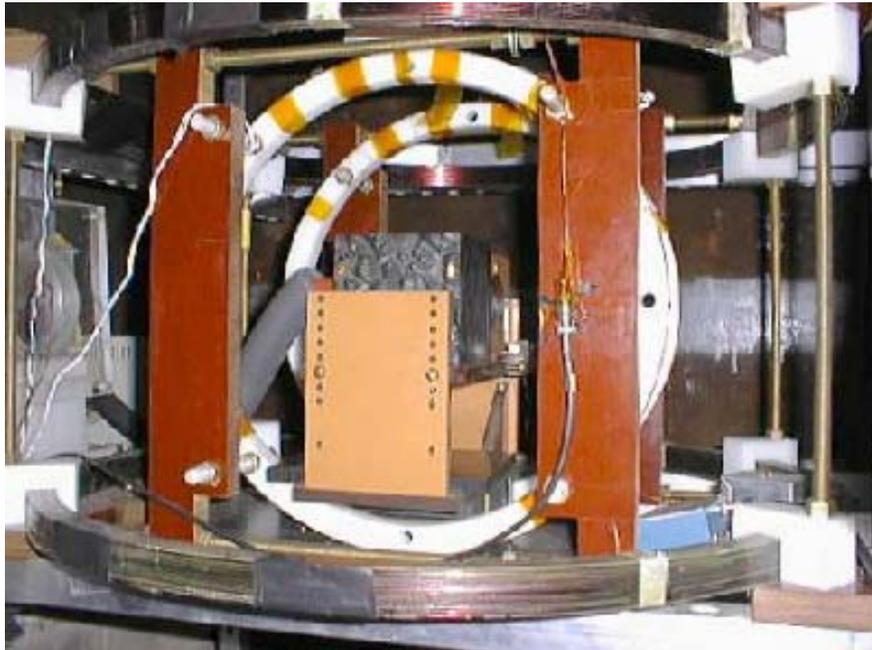
- Single cell Optical Pumping
- Spin Exchange w/ Rb

- $P_n = \tanh(nI\sigma P_3)$

- $t_n = t_0 \exp(-nI\sigma) \cosh(nI\sigma P_3)$



^3He Spin-Filter Setup



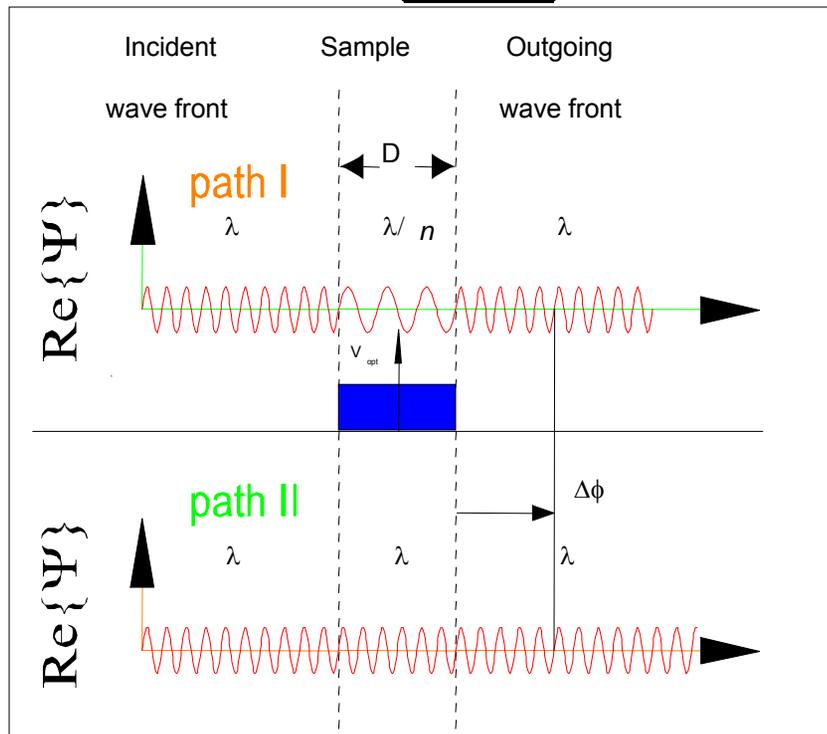
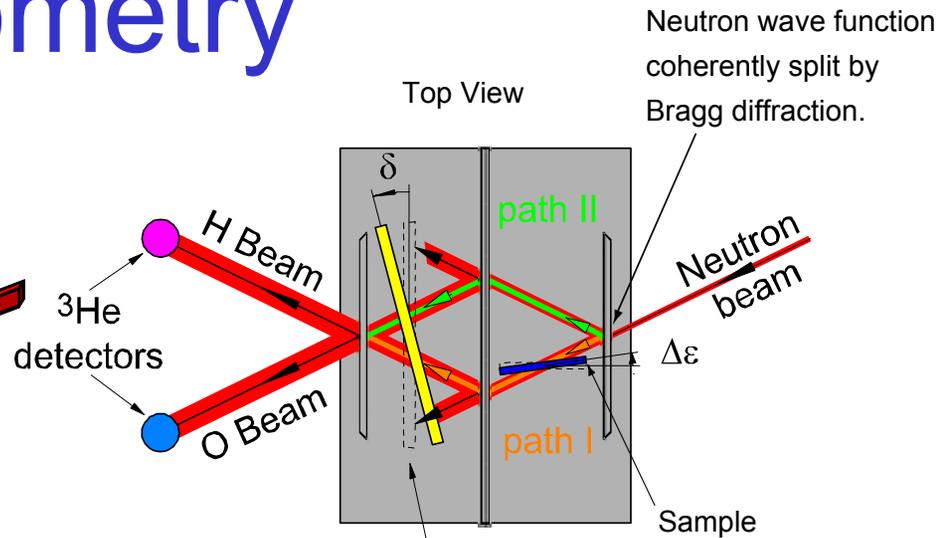
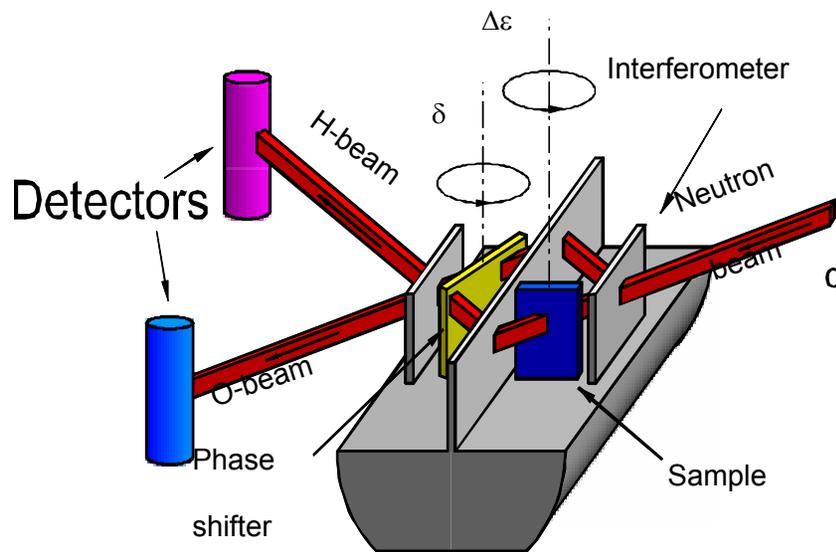
- ◆ Large area, 11 cm in dia, ^3He cells are required to cover the beam.
- ◆ ^3He spin filter allows a compact experimental setup.
- ◆ ^3He spin filter offers an extra spin flip without a change of the B field.

A 11-cm in diameter cell has $T_1 > 500$ hr. ^3He polarization 65% has been measured.

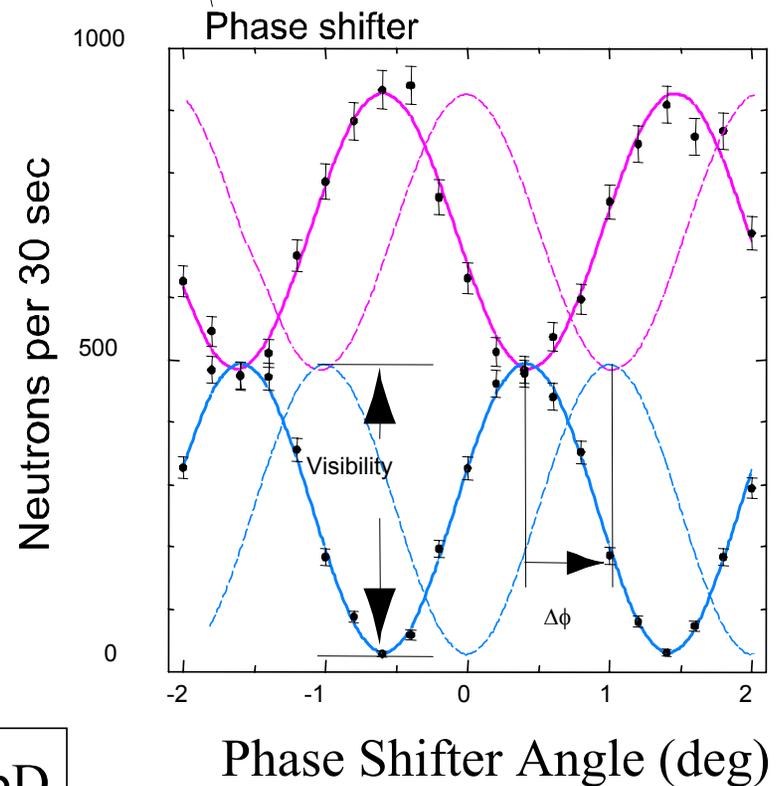


100 W light at 795 nm

Neutron Interferometry

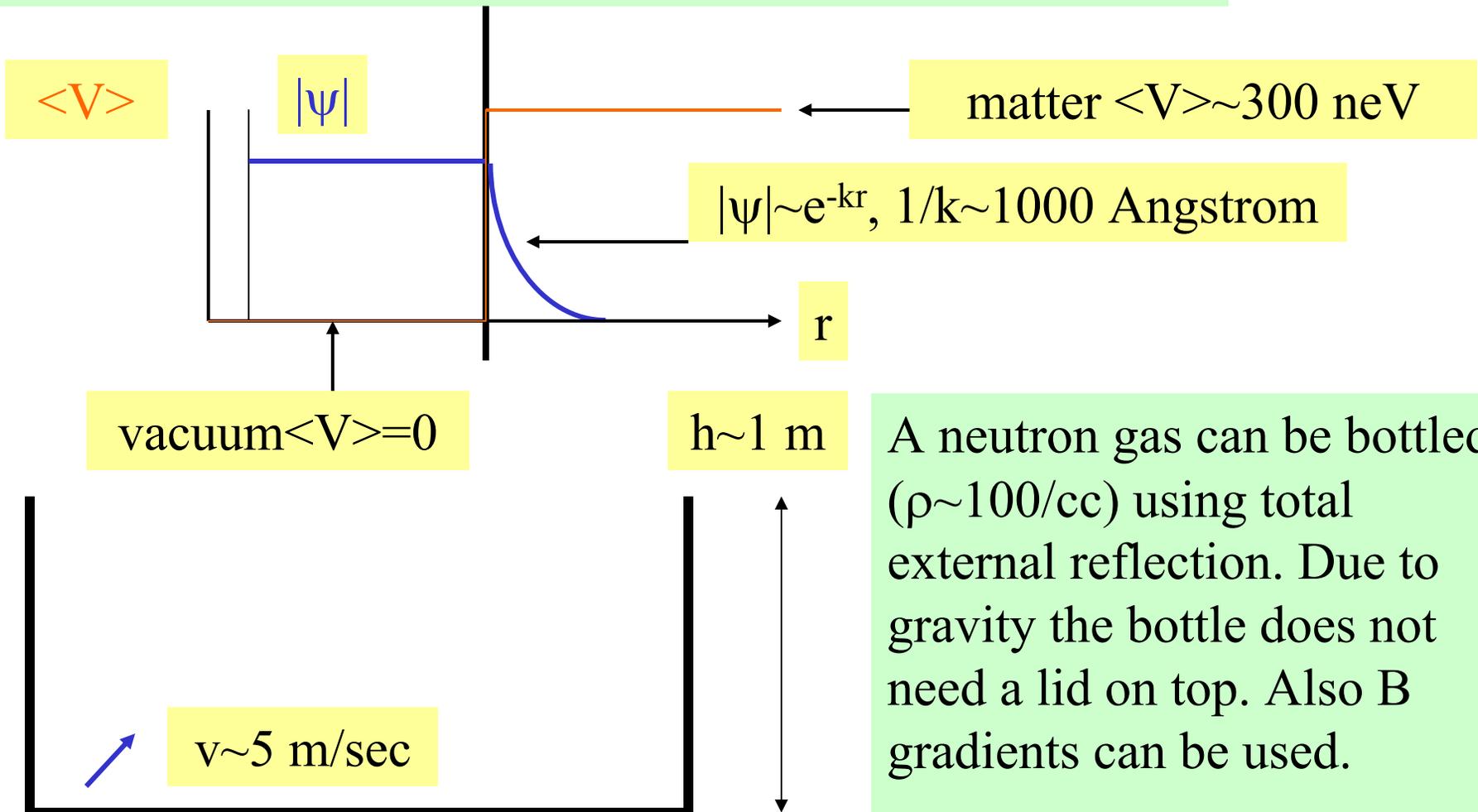


$$\Delta\phi = \lambda N b D$$



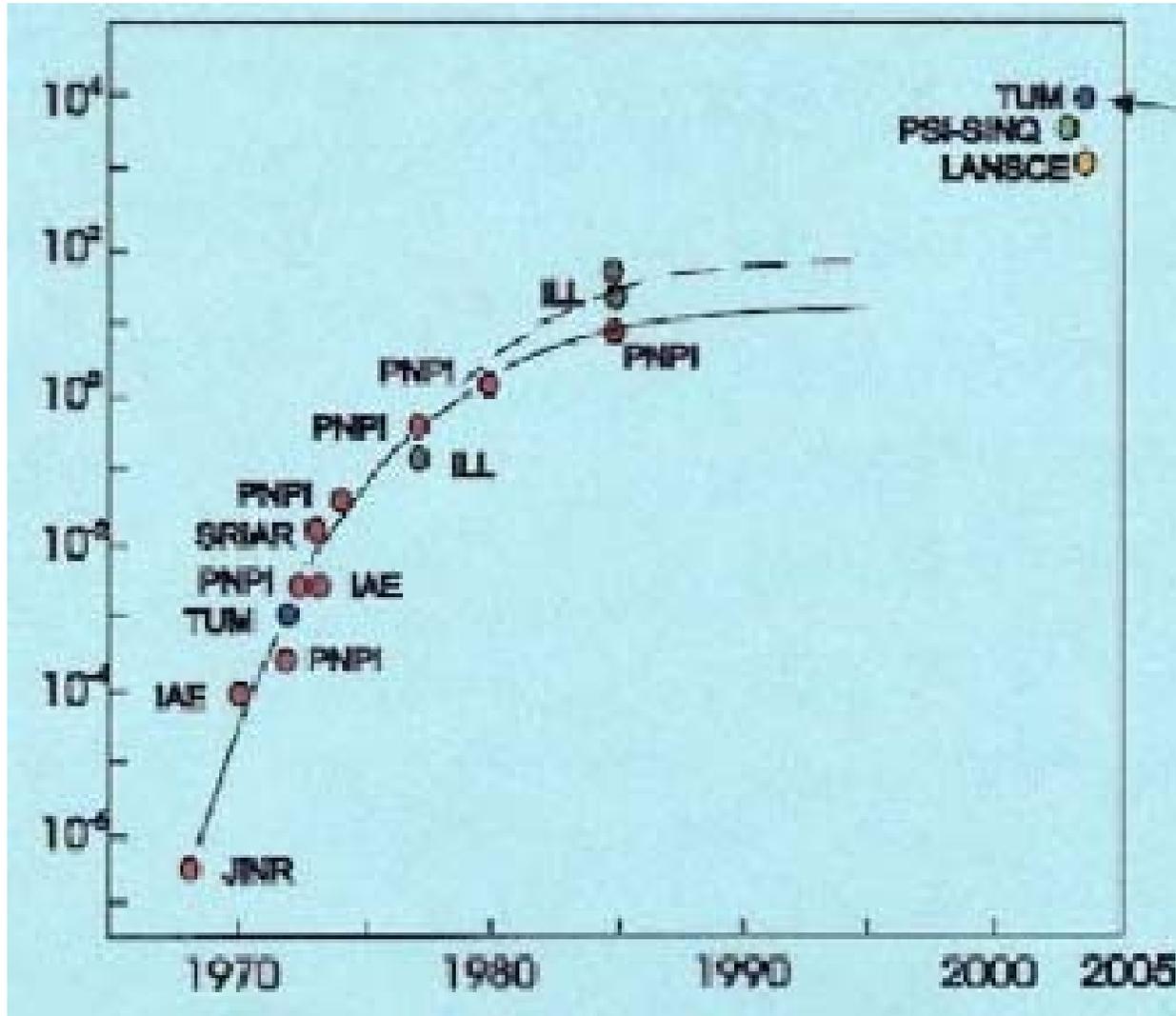
“Ultracold” Neutrons (UCN)

For very low energies ($E_k - \langle V \rangle$ negative, $\langle V \rangle \sim 300$ neV), matter forms a potential barrier for neutrons.



A neutron gas can be bottled ($\rho \sim 100/\text{cc}$) using total external reflection. Due to gravity the bottle does not need a lid on top. Also B gradients can be used.

Overview of existing UCN Sources



- ◆ Neutron moderation
 - Tail of Maxwell-Boltzman distribution
 - $N_{ucn} = 10^{-13} \Phi_0$
- ◆ Conservative force
 - Gravity deceleration
 - Turbine deceleration
 - could not increase the phase space density.
- ◆ **Superthermal source.**

The ILL reactor UCN Turbine

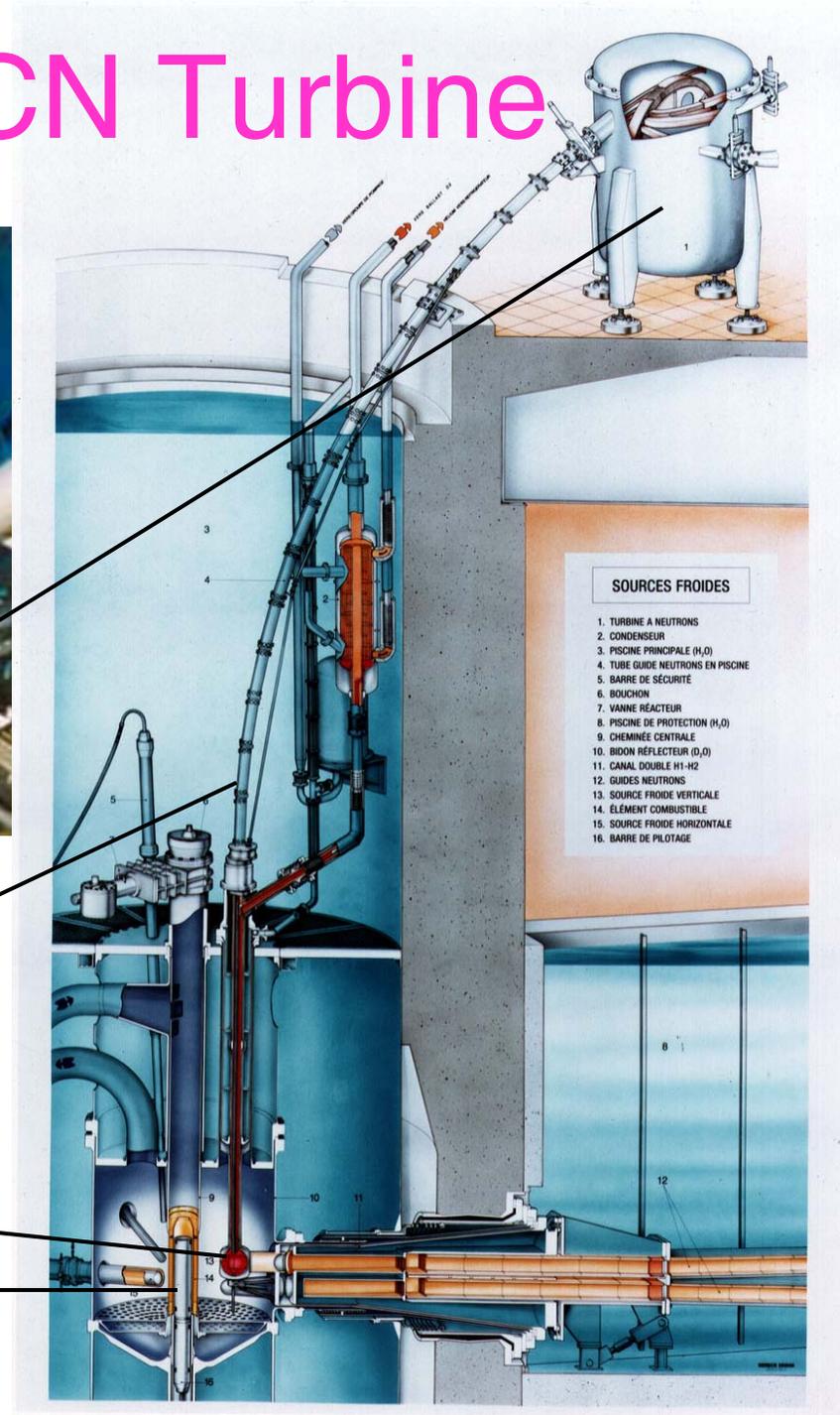


UCN Neutron turbine
A. Steyerl (TUM - 1986)

Vertical guide tube

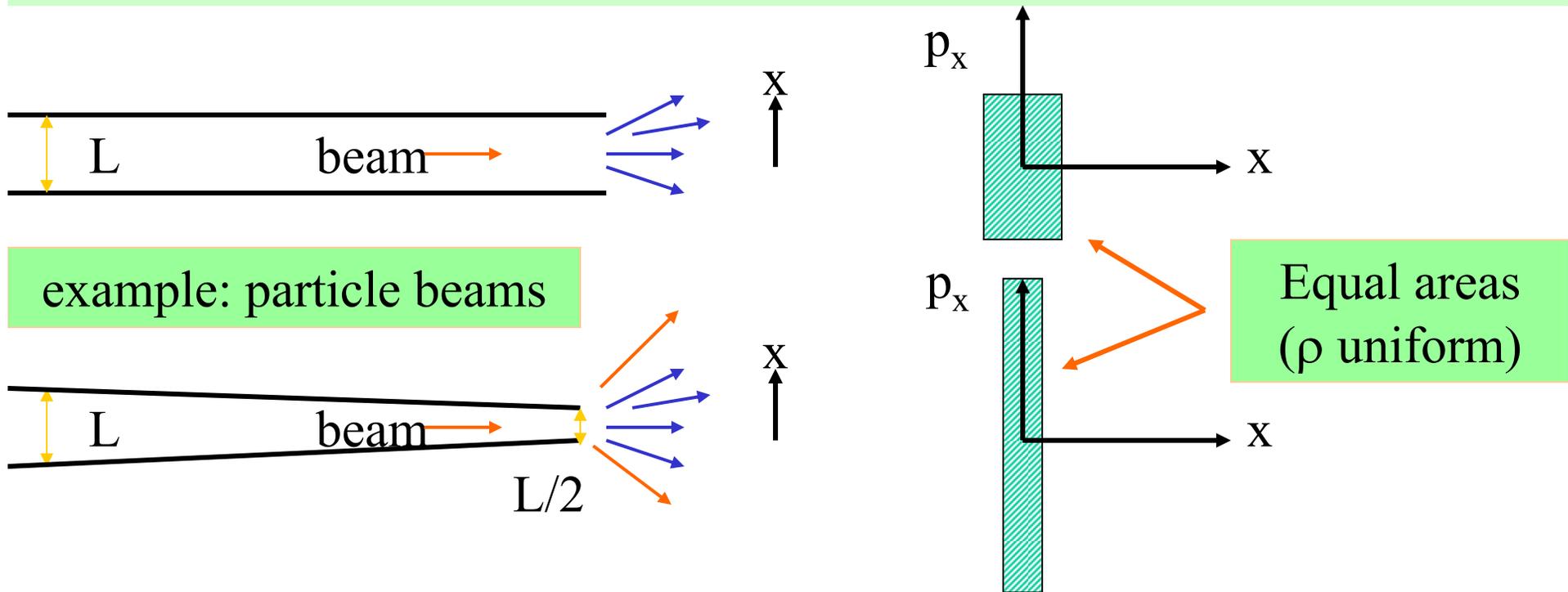
Cold source

Reactor core



Making More UCN: Liouville's Theorem

Describe an ensemble of neutrons by a density in phase space:
 $\rho(\mathbf{r},\mathbf{p})d^3r d^3p$. Liouville's theorem: phase space volume occupied by particles is constant if only "conservative" forces [=derivable from a potential] act



UCN gas in a bottle with d^3p filled $[0, p_{\max}]$: $N(\mathbf{r})d^3r$ constant.
Need dissipative (non-conservative) interaction to increase $N(\mathbf{r})$

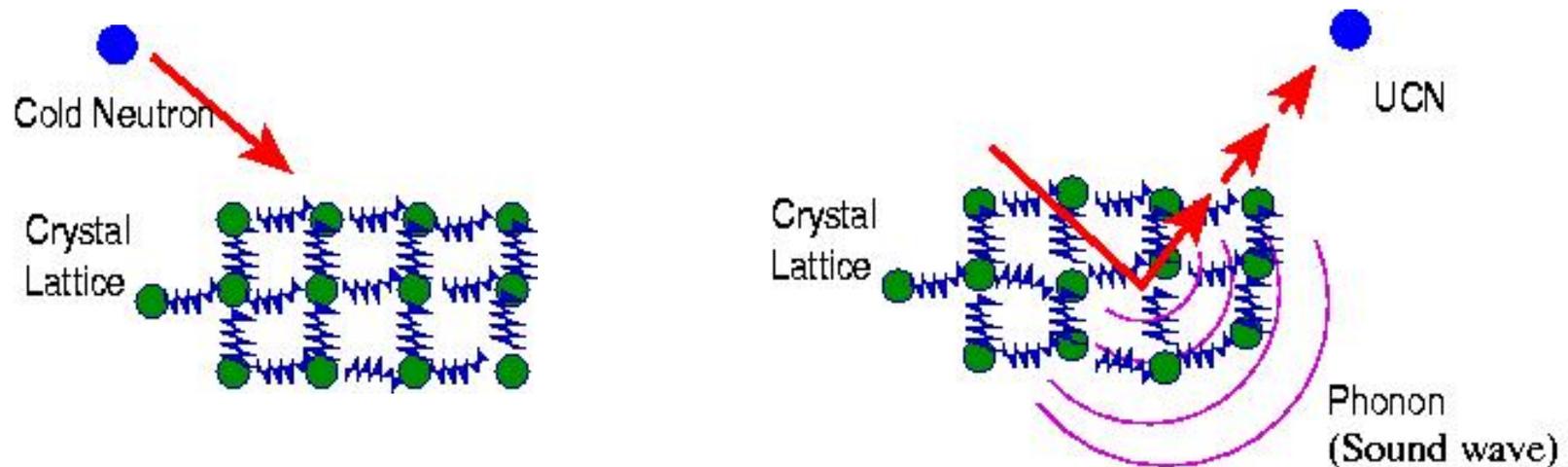
“Superthermal” UCN Sources

- ◆ **superfluid 4He:** highest densities, fed by 1 meV monochromatic neutrons using **phonon cooling**, need to extract UCN from 4He at <0.5K
- ◆ **solid D2:** high currents, fed by ~4 meV neutron spectrum using **phonon cooling**, needs to be maintained in ortho state at ~5K
- ◆ **solid O2:** highest potential currents, fed by ~1 meV neutron spectrum by **magnon cooling**, needs to operate ~2K

Superthermal Process: Beats Liouville a Dissipative Process

R. Golub and J. M. Pendlebury, Phys. Lett, A53, 133 (1975)

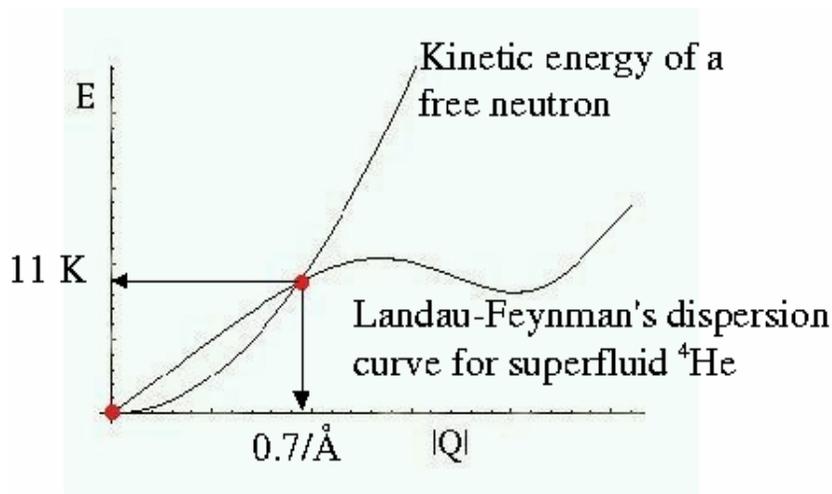
- ◆ Cold neutrons downscatter in the solid, giving up almost all their energy, becoming UCN.



- ◆ UCN upscattering (the reverse process) is suppressed by controlling the moderator at low temperatures.

Liquid ^4He as a UCN Source

- ◆ Isotropic liquid \Rightarrow Single degenerate dispersion curve for HeII.



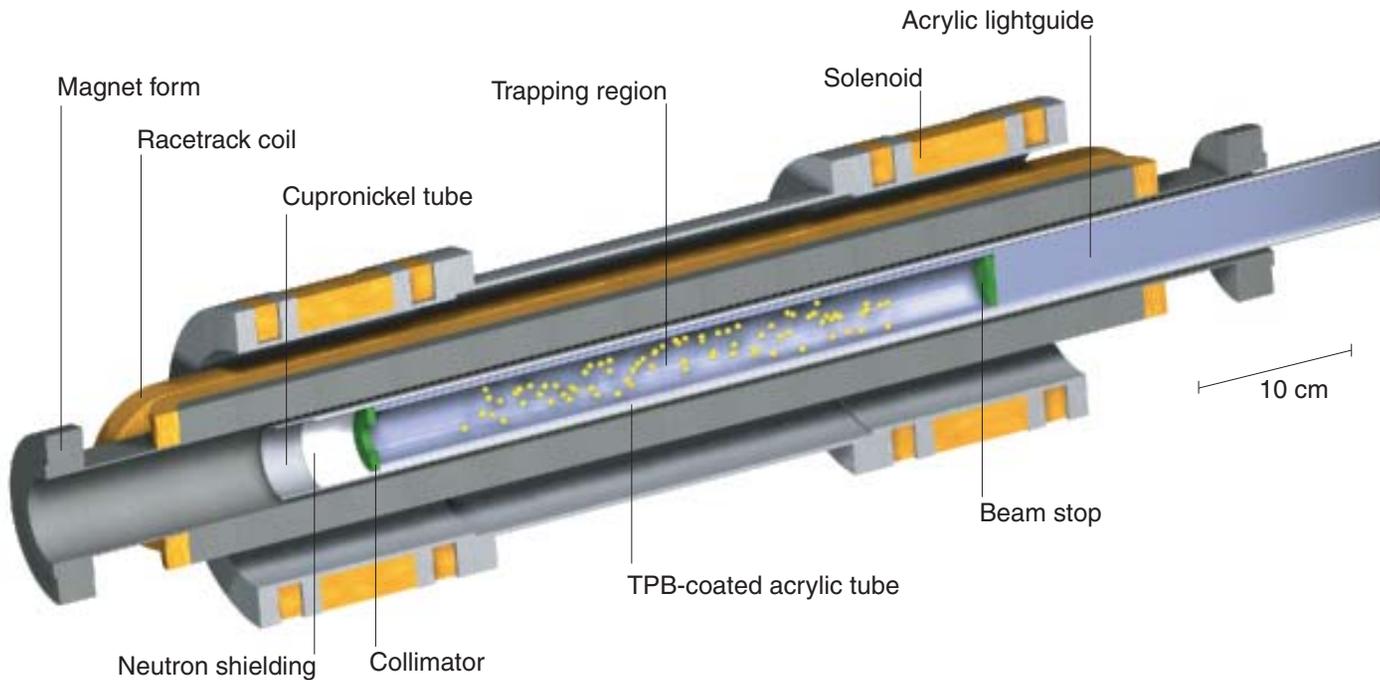
$$\sigma_{\text{coh}} = 1.34 \text{ barn}$$

$$\sigma_{\text{inc}} = 0 \text{ barn}$$

- ◆ Incident cold neutron with momentum of 0.7 \AA^{-1} (1 meV) can excite a phonon in ^4He and become an UCN
- ◆ No nuclear absorption loss for neutrons ($\sigma_{\text{abs}} = 0$).
- ◆ UCN can accumulate up to the β -decay time.
 - Produce a large steady-state density.

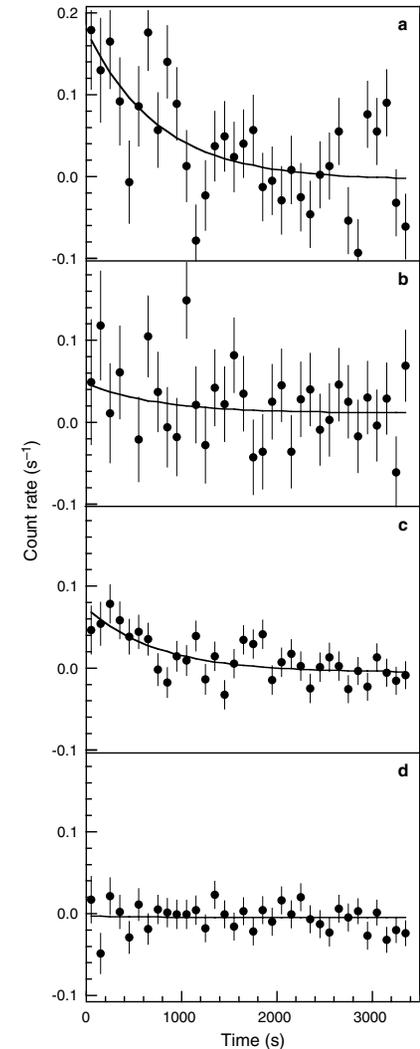
UCN Production in Superfluid ^4He

Magnetic Trapping of UCN
(Nature 403 (2000) 62)



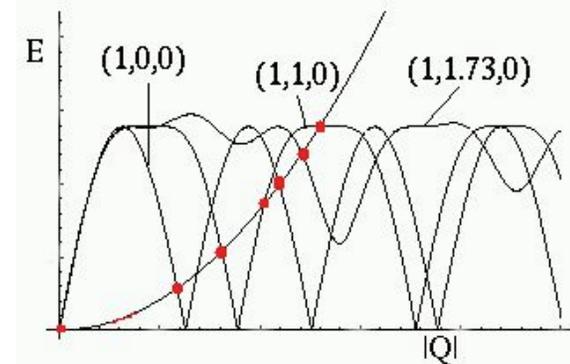
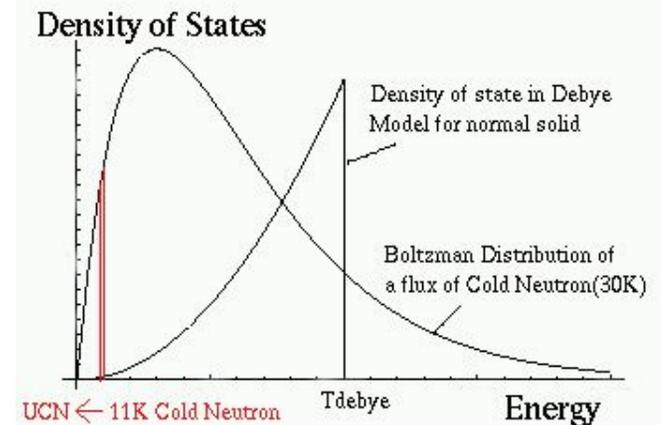
560 ± 160 UCNs trapped per cycle (observed)

480 ± 100 UCNs trapped per cycle (predicted)

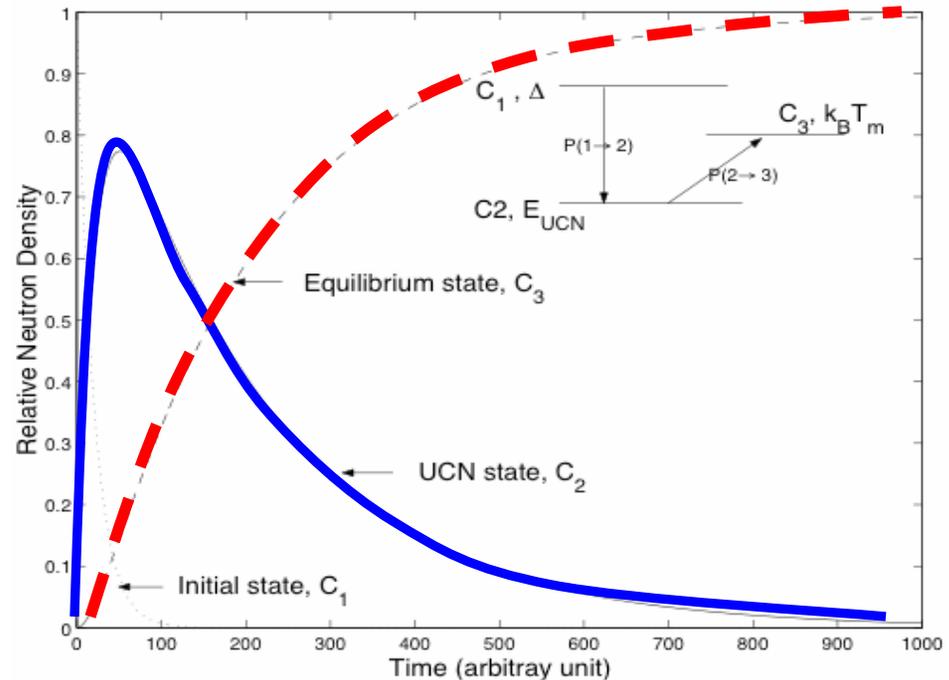
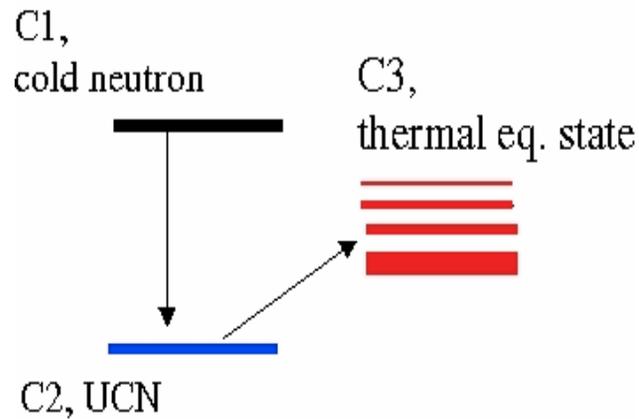


Solid Deuterium as a UCN Source

- ◆ $\sigma_{\text{inc}} = 2.04$ barn
 - Momentum is not conserved.
 - CN with energy smaller than the T_{debye} participate in the UCN production.
- ◆ $\sigma_{\text{coh}} = 5.59$ barn
 - The anisotropic dispersion curves broaden the inelastic scattering criteria for UCN production.
- ◆ More efficient use of cold neutrons.
 - Produce a large current(flux) of UCN.
- ◆ However, large UCN loss due to the nuclear absorption.

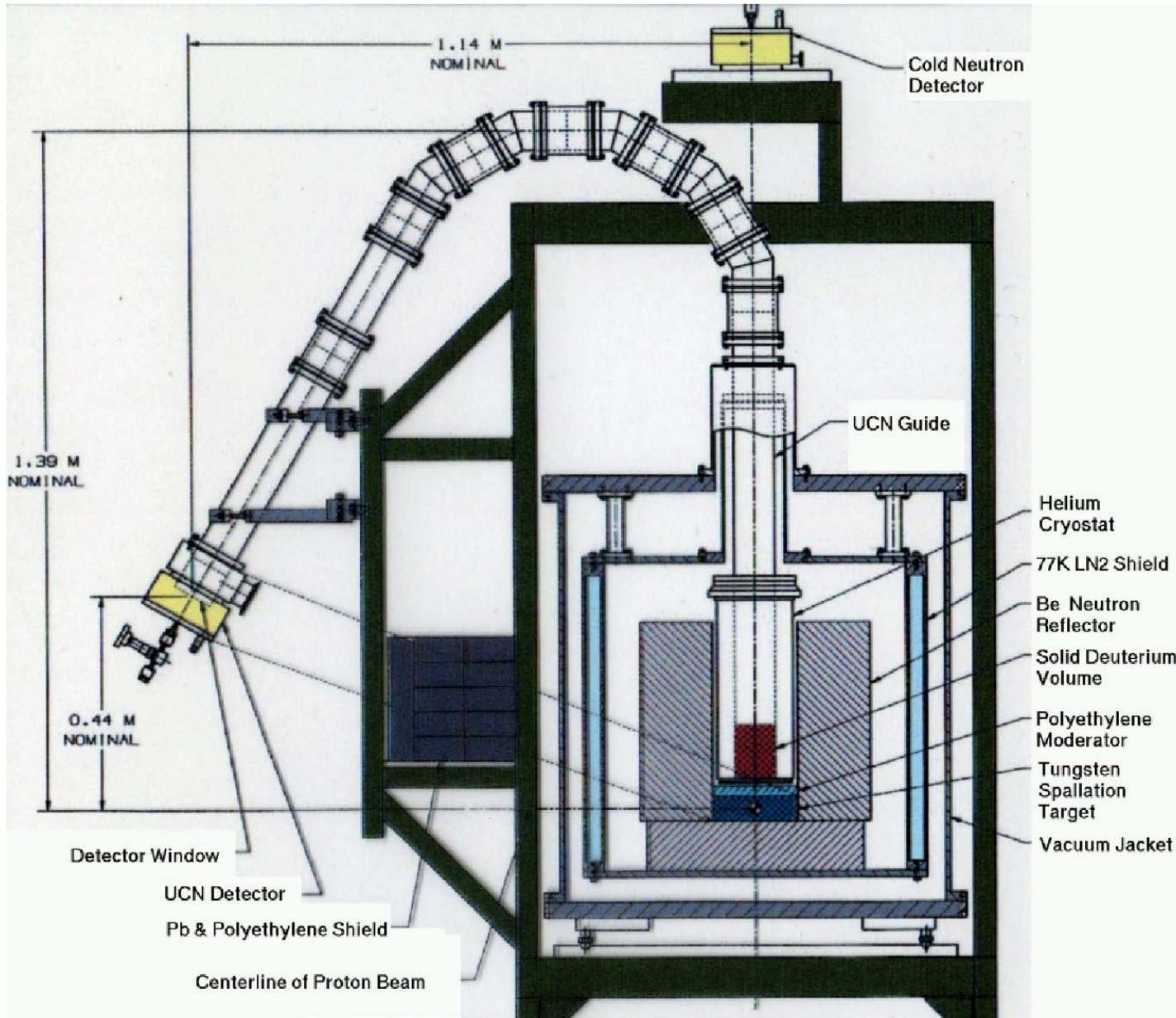


Dynamics of UCN Production



- ◆ Lifetime of UCN in the source material is a critical parameter in the establishment of large UCN densities.
- ◆ Extract UCN out of the source before it is thermalized \Rightarrow Spallation N source + Separation of the source and the storage + a UCN Valve

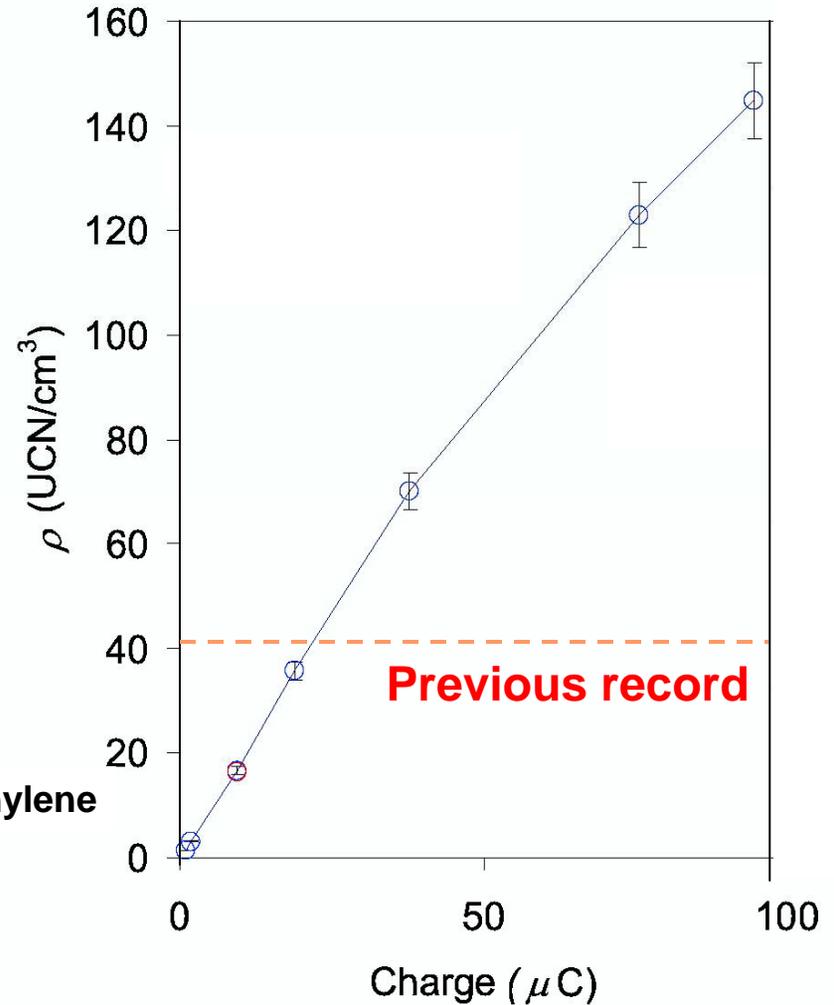
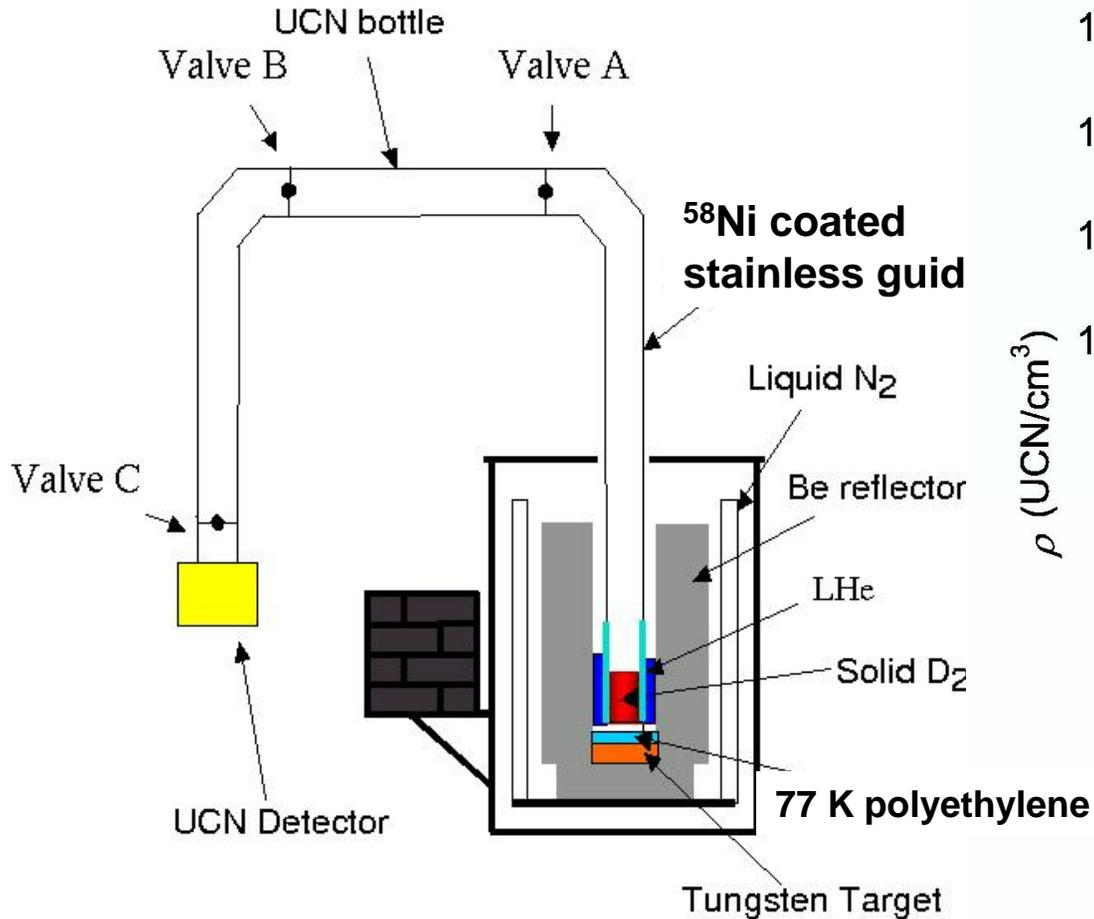
Solid Deuterium Source at LANL



SD₂ UCN Source

Based on R.Golub and K.Böning Z.Phys.B51,95,(83)

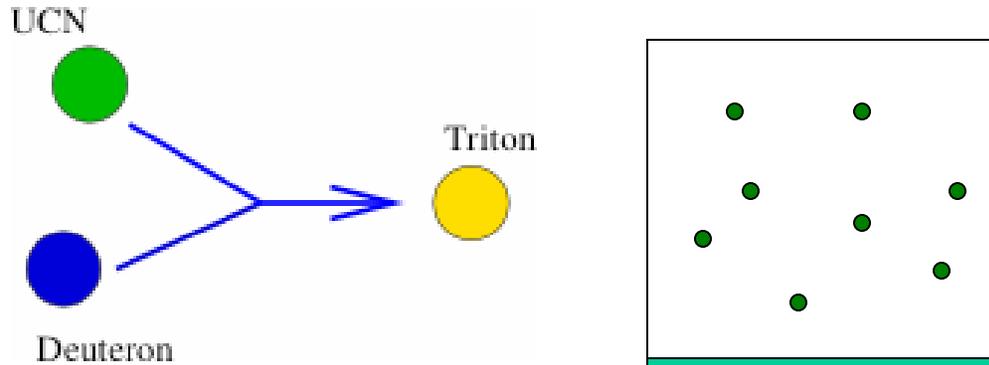
World record UCN density



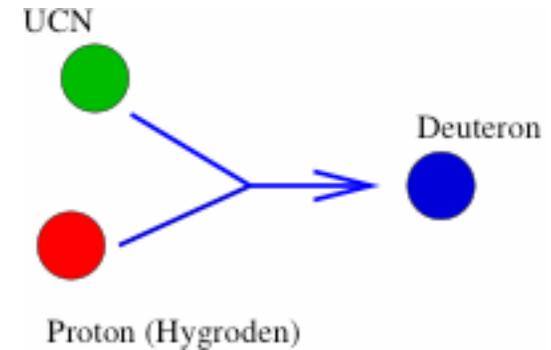
C.L.Morris *et al.* PRL **89**, 272501 (2002)

A. Saunders *et al.* nucl-ex/0312021

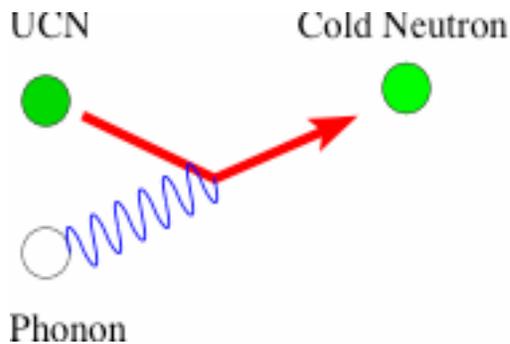
UCN Loss in Solid Deuterium



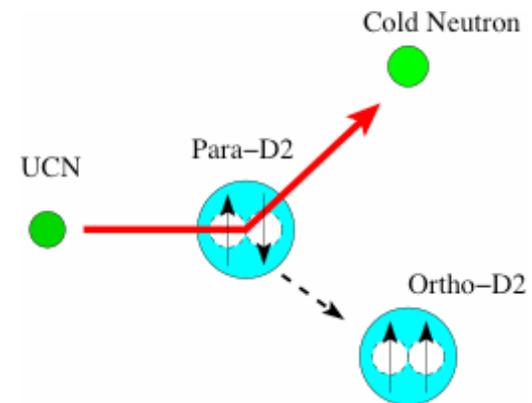
Nuclear absorption by S-D₂
 $\tau \sim 150$ msec



Nuclear absorption by Hydrogen
Impurities, $\tau \sim 150$ msec/0.2% of H

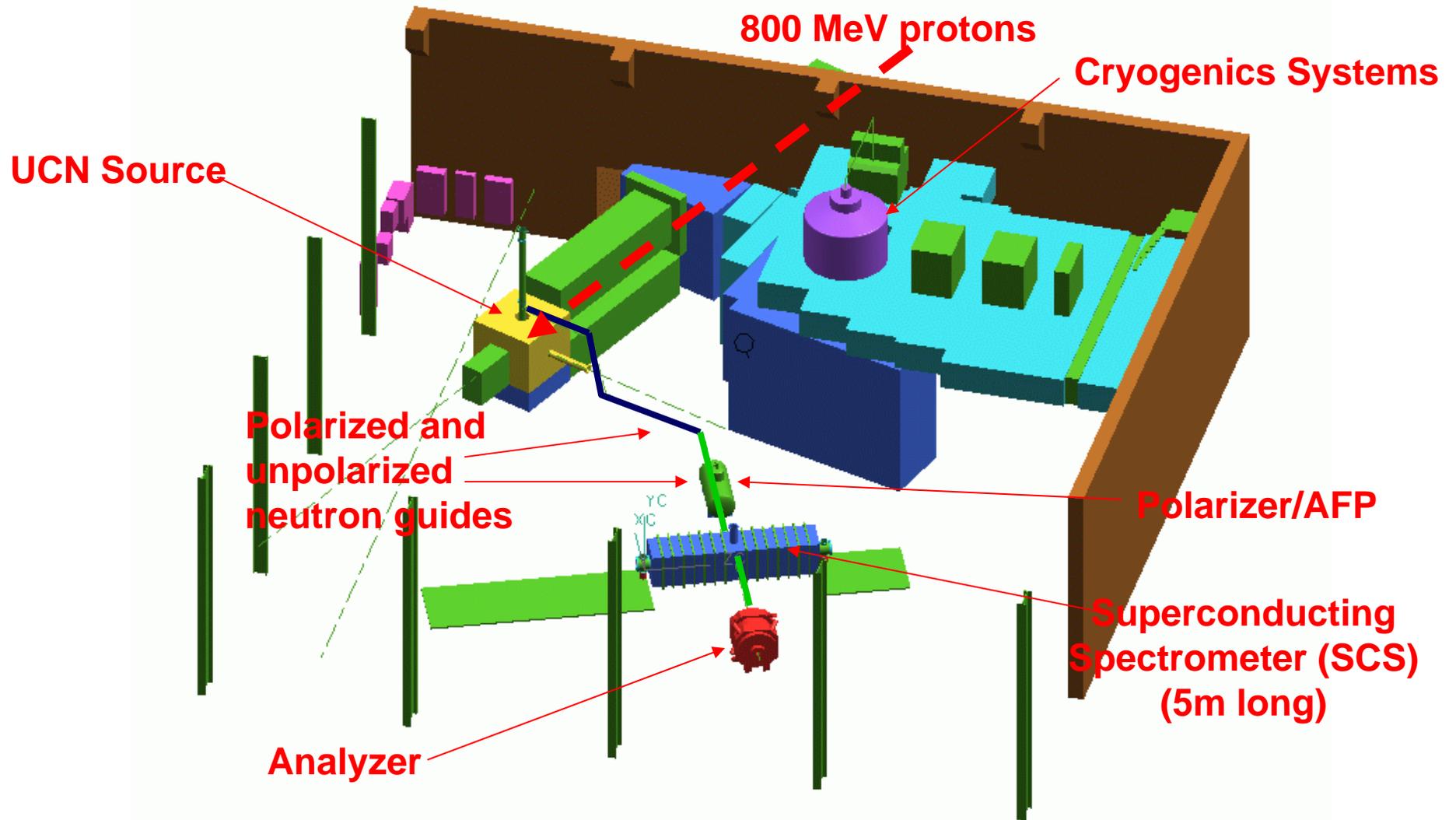


UCN upscattering by phonons
 $\tau \sim 150$ msec at T = 5K



UCN upscattering by para-D₂
 $\tau \sim 150$ msec/1% of para-D₂

Experiment Layout at LANSCE



Physics/Technology of Cold and Ultracold Neutrons

Uses techniques and concepts from atomic physics, condensed matter physics, low temperature physics, optics

Neutrons can be highly polarized, guided, and trapped

Energy range: ~ 100 neV to ~ 25 meV

CW and pulsed sources available