

# High-intensity laser interactions with near-critical density plasmas: NP-315758

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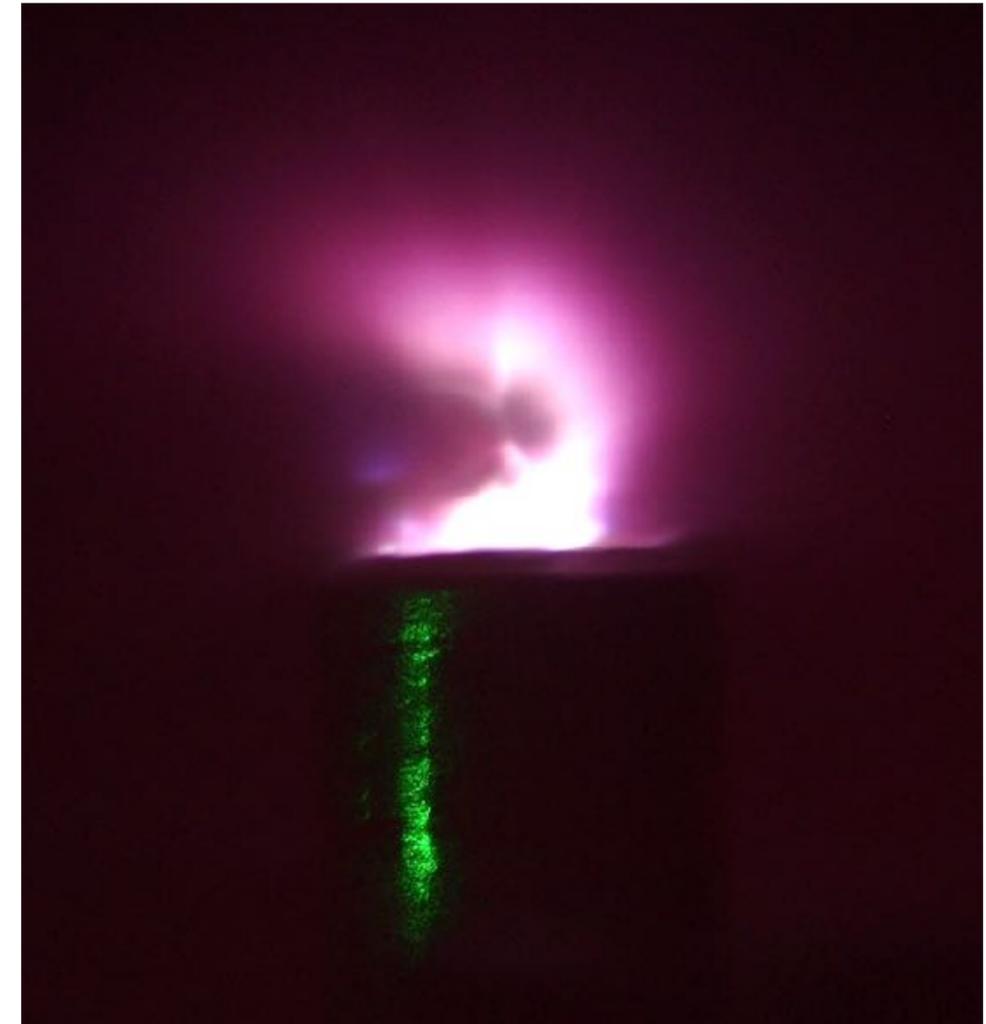
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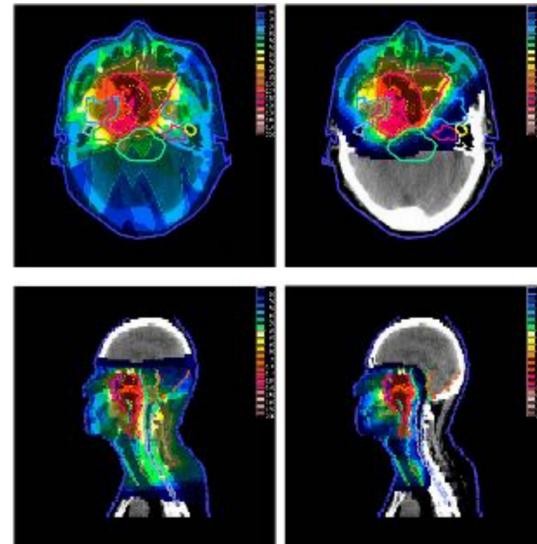
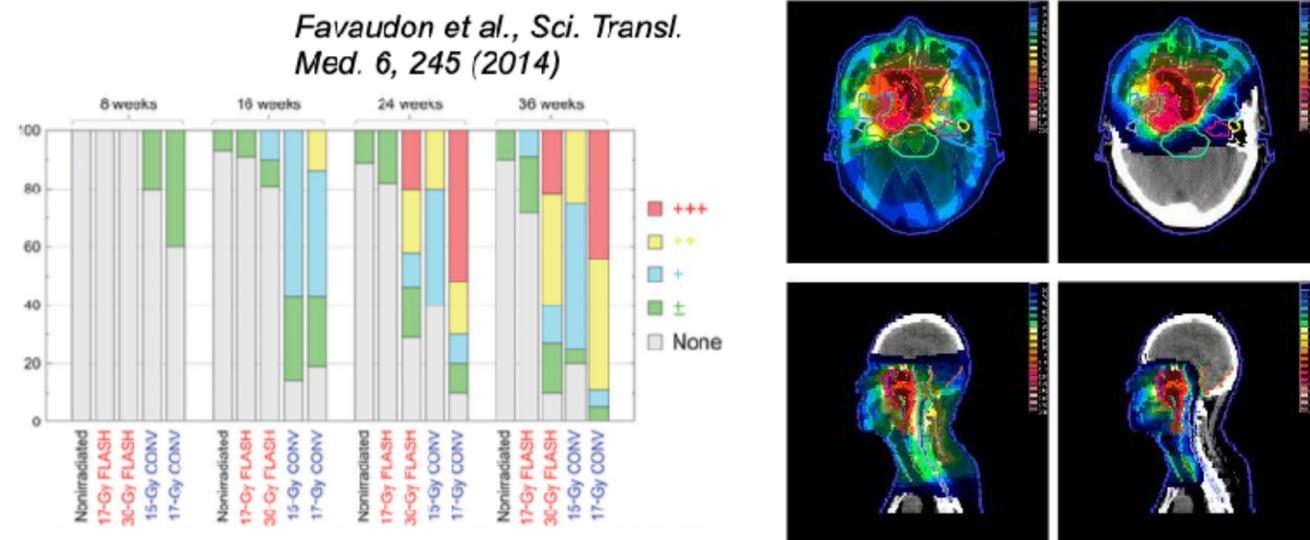
**ATF User Meeting,  
27th March 2024**

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# Motivation for laser-driven ion sources

## Why laser driven ion sources?



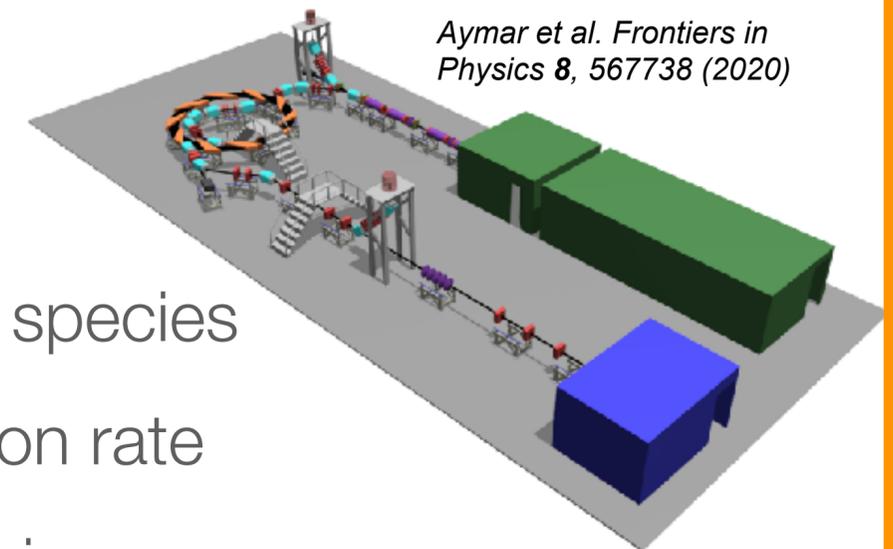
*Z. Taheri-Kadkhoda et al. Radiation Oncology 3 (2008)*

Laser driven ion sources increasingly attractive due to high source energy and short bunch length

For example, these sources are well suited for high dose rate radiobiology - e.g. FLASH

## Important characteristics of laser driven source for applications

- High energy
- High flux
- Different ion species
- High repetition rate
- Minimal debris

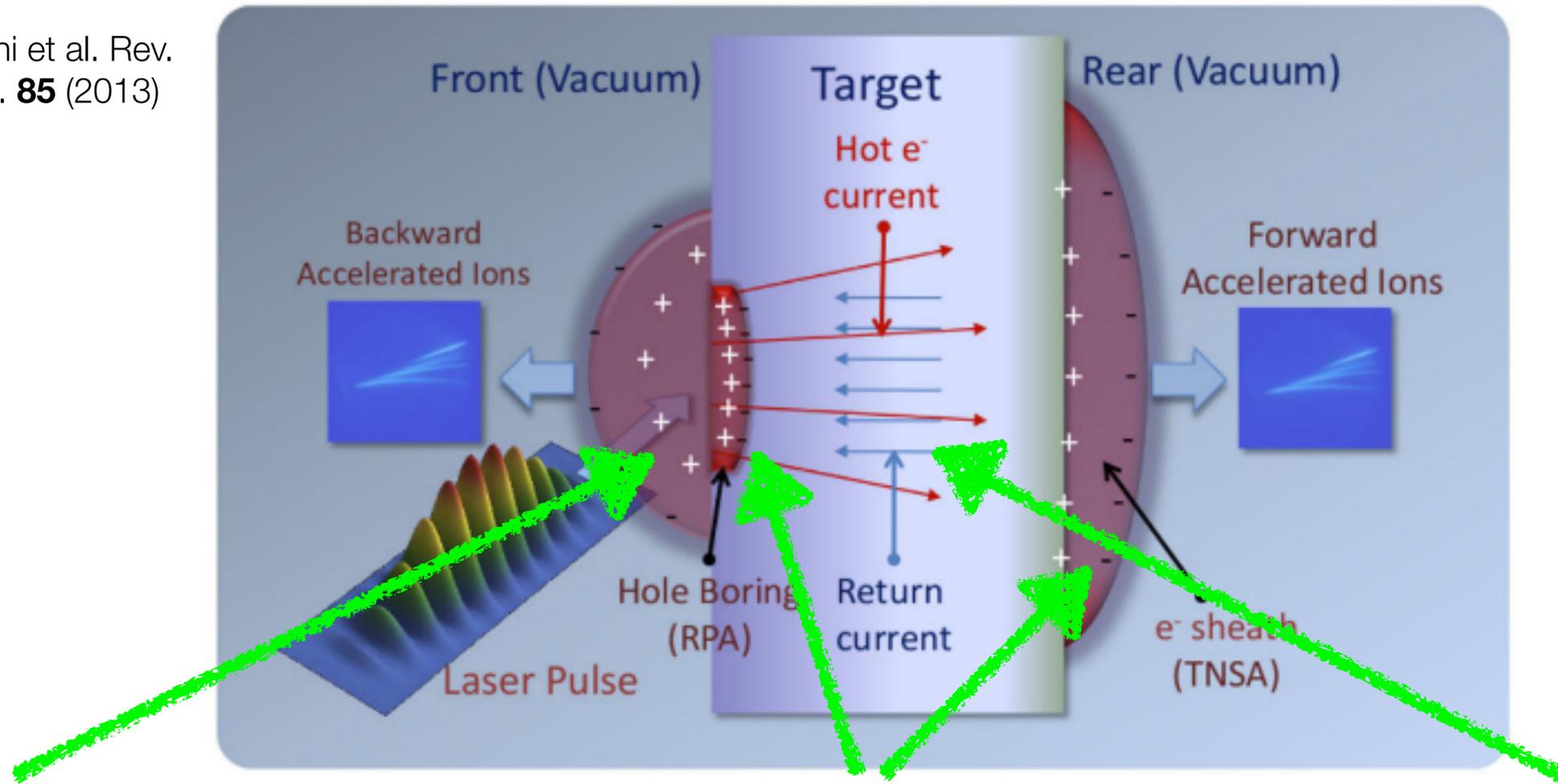


*Aymar et al. Frontiers in Physics 8, 567738 (2020)*

Gaseous targets are a great choice, if high energy, high flux ions can be produced...

# Physics of laser driven ion sources difficult to diagnose directly

From Macchi et al. Rev. Mod. Phys. **85** (2013)



Laser propagation in underdense plasma

Acceleration of ions at critical density surface and plasma boundary

Propagation of "fast" electrons in the target

Ion sources undergo multiple nonlinear and dynamic processes, near-impossible to see experimentally

# Diagnosing laser driven ion sources - a new approach?

Nearly all laser driven ion source experiments performed in the near-IR

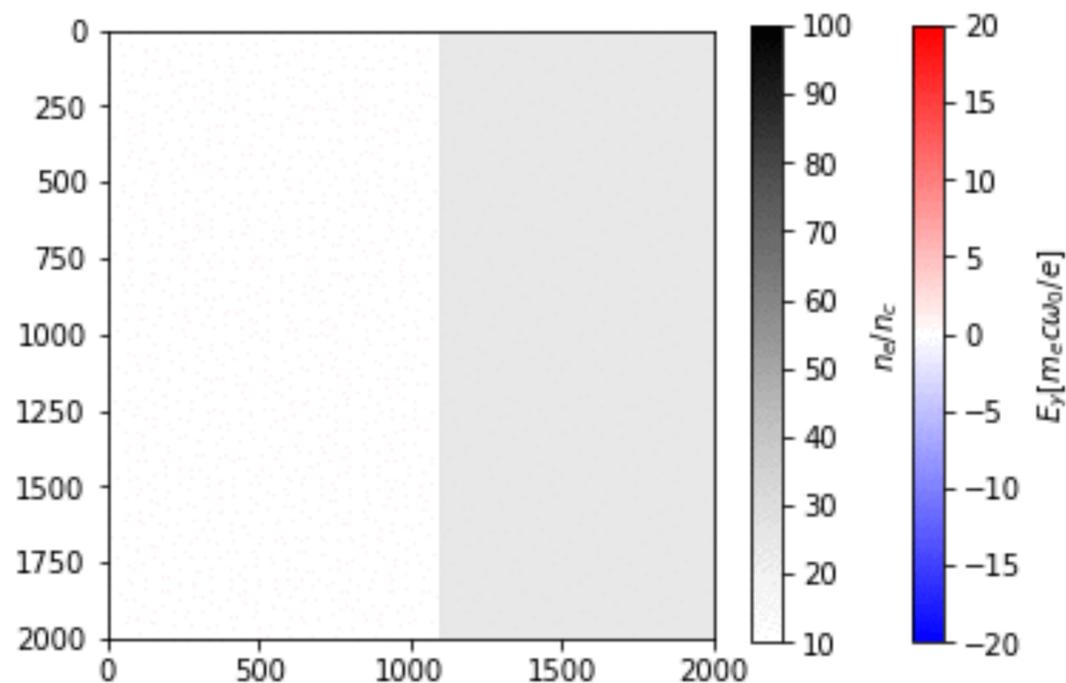
Typical dynamical scales	<i>Time</i>	<i>Length</i>	<i>Density</i>
	~10 fs	~1 $\mu\text{m}$	$>\sim 10^{21} \text{ cm}^{-3}$

Can we diagnose it?

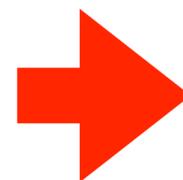
Too quick

Too short

Too dense



Rely on simulations, many assumptions



- Reduced dimensionality
- Uncertainty over experimental parameters
- Can only verify by looking at certain outputs e.g. ion beam

# Exploiting dimensional scaling of collisionless laser-plasmas

Collisionless laser plasmas can be defined using reference frequency\*:

*Time*

*Length*

*Density*

$$\tilde{t} = \omega_L t$$

$$\tilde{x} = \frac{\omega_L}{c} x$$

$$\tilde{n} = \frac{1}{n_c} n \propto \frac{1}{\omega_L^2} n$$

near-IR

~10 fs

~1 μm

>~10<sup>21</sup> cm<sup>-3</sup>



longwave-IR ~100 fs

~10 μm

>~10<sup>19</sup> cm<sup>-3</sup>



Resolvable

Resolvable

Ideal for optical probing

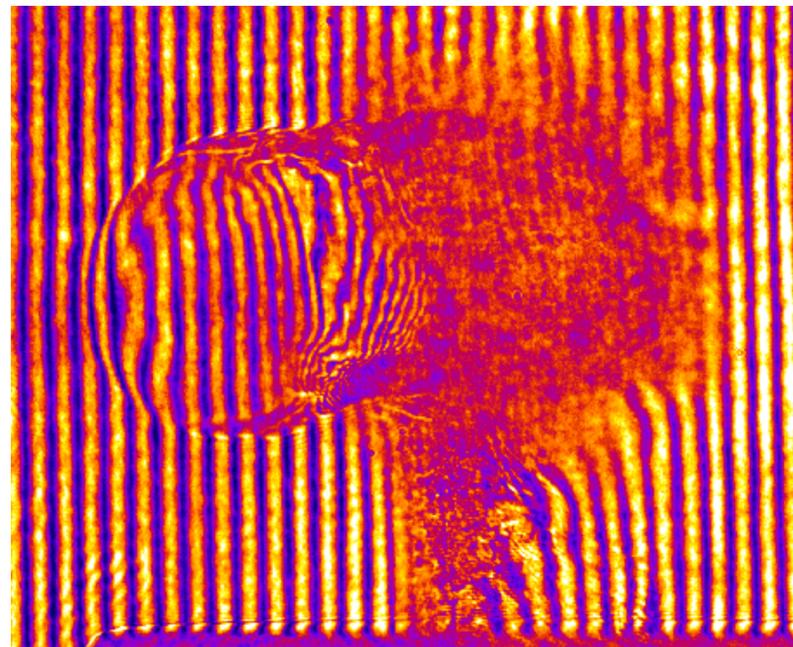
\*if e.g. ionisation/QED not important

# A unique facility at the ATF for investigating ion source physics

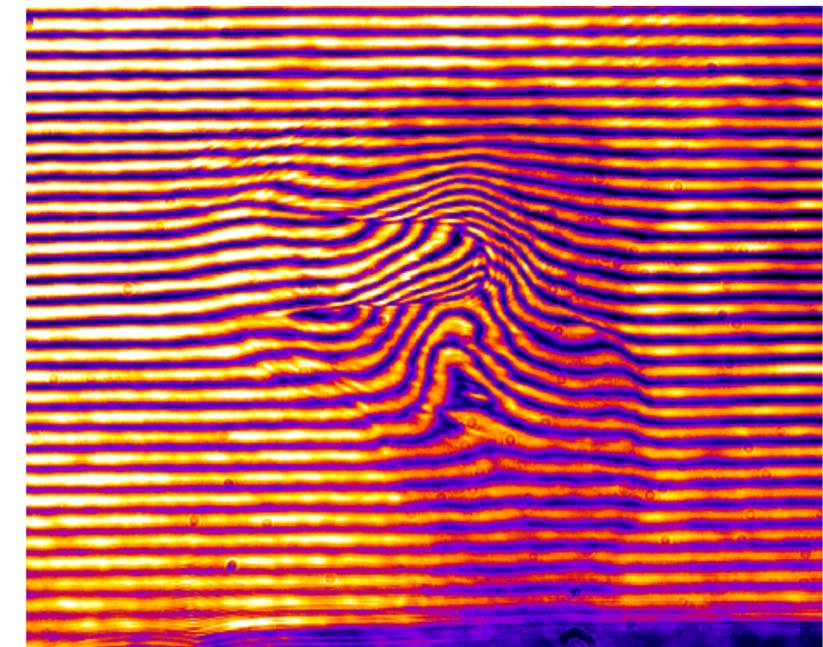
Utilising the ATF's NIR and MWIR laser facilities, we have a unique and exciting platform for investigating ion source physics dynamics

- CO<sub>2</sub> laser - 2 ps, 9.2  $\mu\text{m}$  wavelength drive laser for ion acceleration @  $10^{19} \text{ cm}^{-3}$
- TiS laser - <100 fs, 800 nm wavelength laser ideal for optical probing such densities

Enables *direct dynamic observation* of fundamental scale-independent processes driving all laser driven ion sources



*Previously: blur due to ionisation and plasma dynamics when temporal overlap between drive and probe*



*Now: clean images when overlapping drive and probe, allowing measurements of evolving overdense LPI*

# Proposed Objectives

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We plan to exploit this setup to investigate three regimes of importance to ion acceleration:

1. Laser propagation in underdense plasmas preceding the critical surface
2. The dynamics of ion acceleration by shock structures
3. Particle beam propagation through plasmas

# Proposed Investigation

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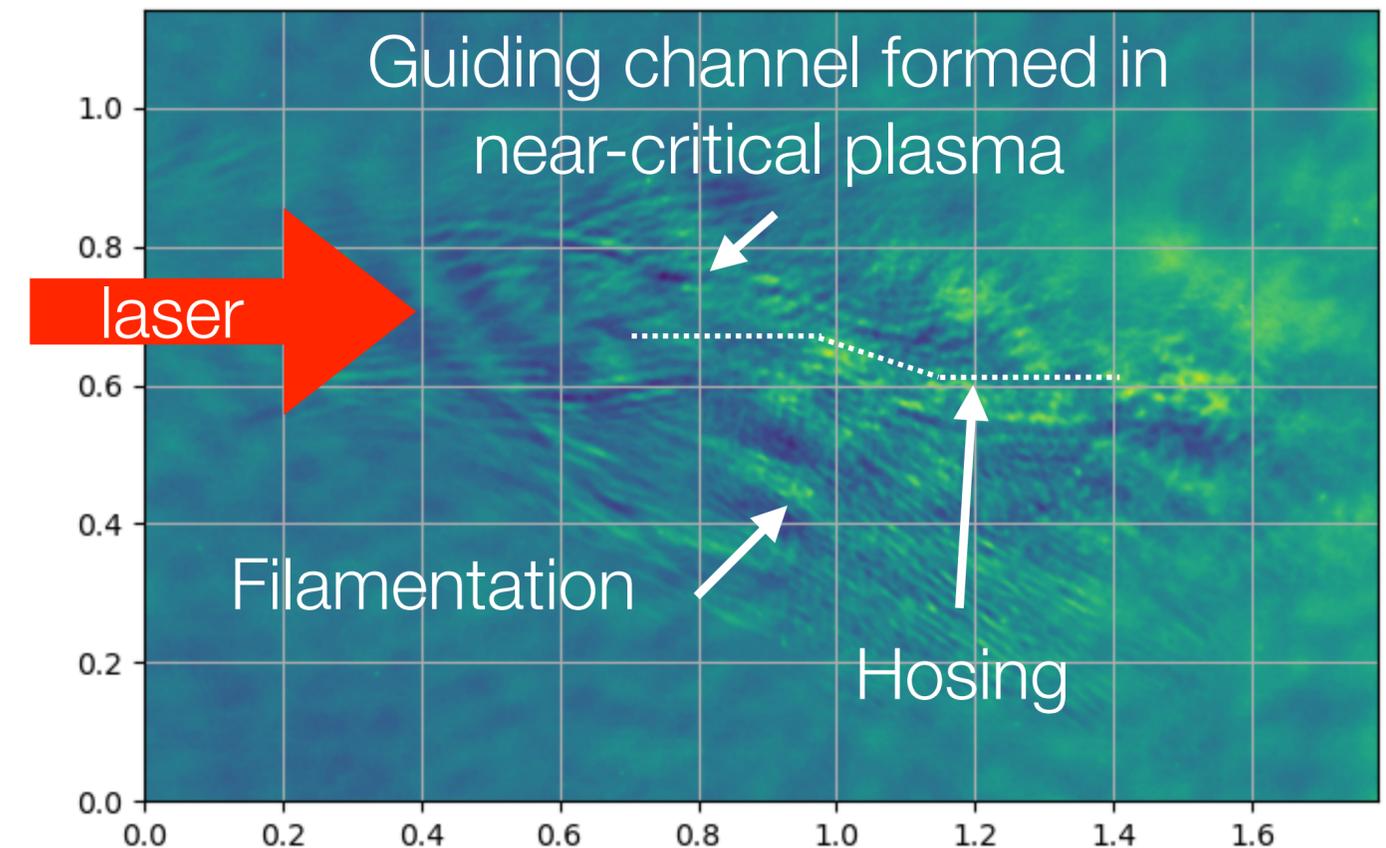
- 1. Laser propagation in underdense plasmas preceding the critical surface**
2. The dynamics of ion acceleration by shock structures
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# Laser propagation in underdense plasmas

To fully understand the acceleration process laser conditions at the critical surface must be fully understood

A number of effects can occur as a laser propagates an underdense plasma:

- relativistic self-focussing and dispersion
- laser hosing
- filamentation

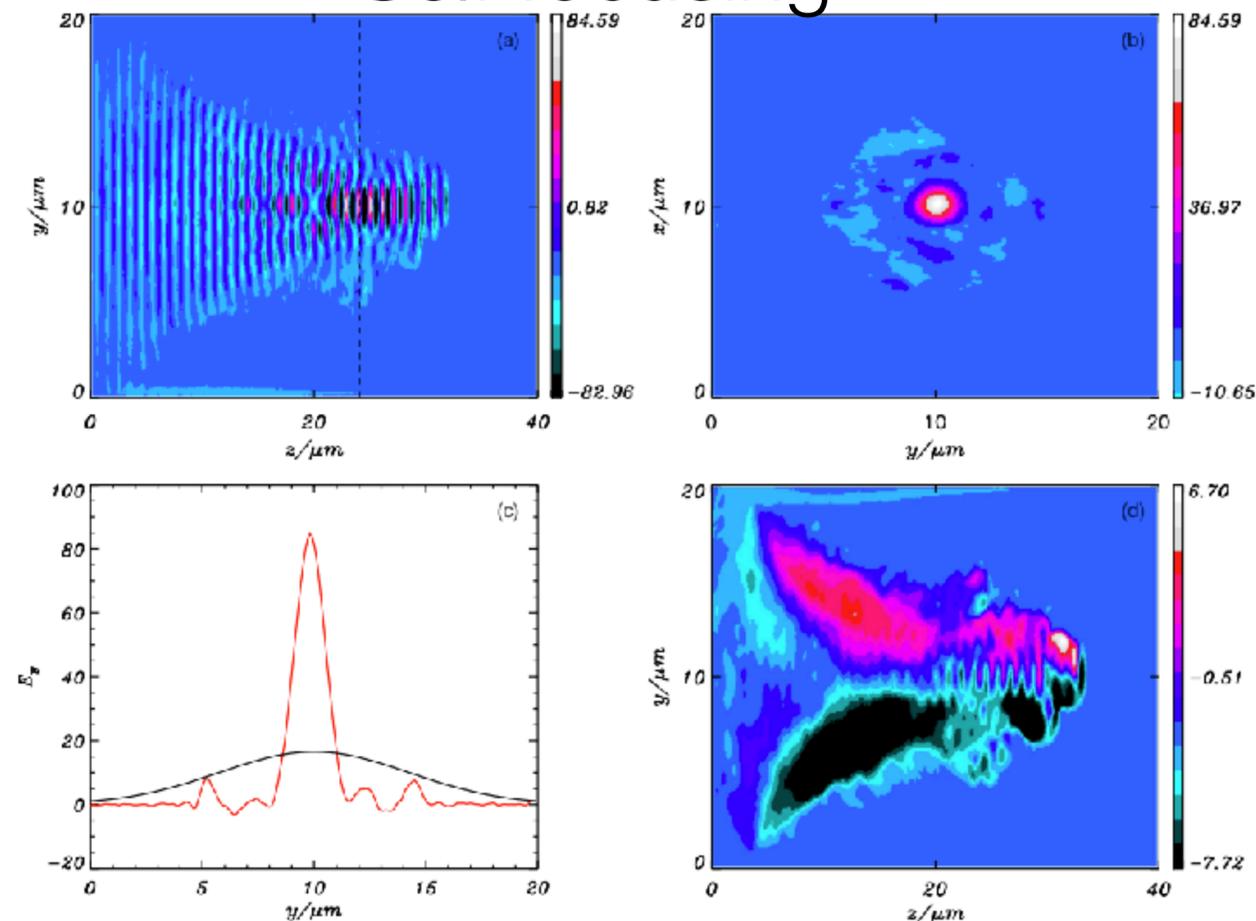


# Laser propagation in underdense plasmas

These processes affect laser conditions at critical density surface

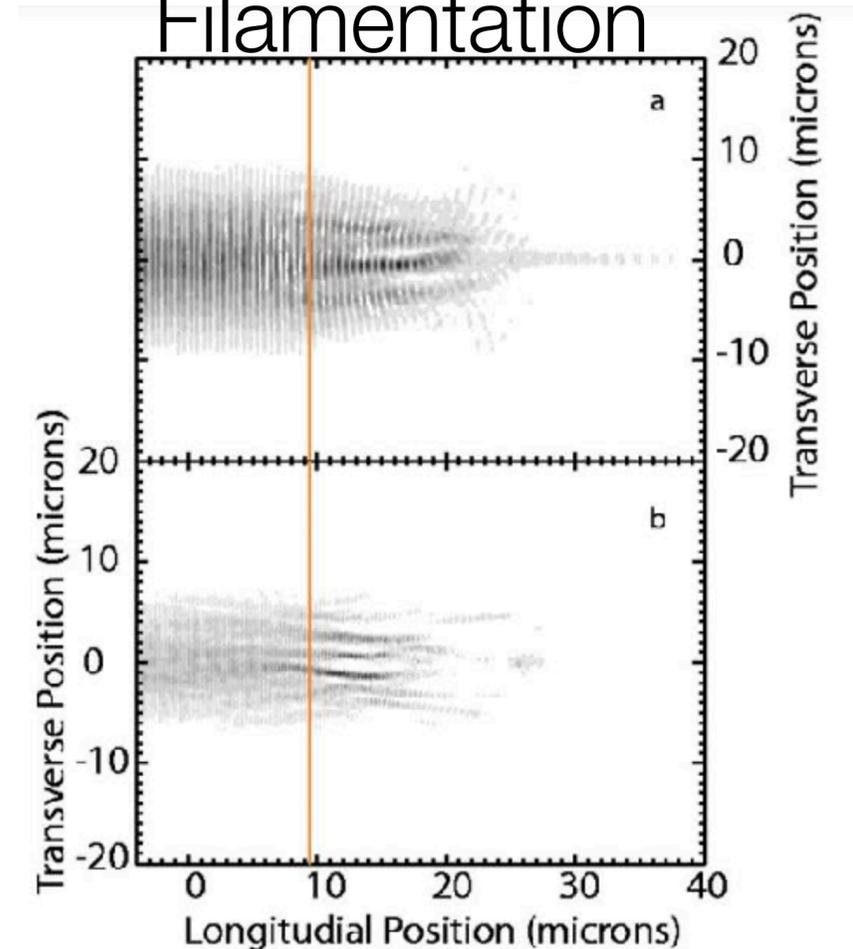
- ➔ uncertainty in laser energy, focal spot size and shape
- ➔ Can enhance or degrade acceleration performance

## Self focusing



From Jiang et al. PRL **107** (2011)

## Filamentation



From Brady et al. Phys. Plasmas **19** (2012)

# Experimental investigation of these concepts

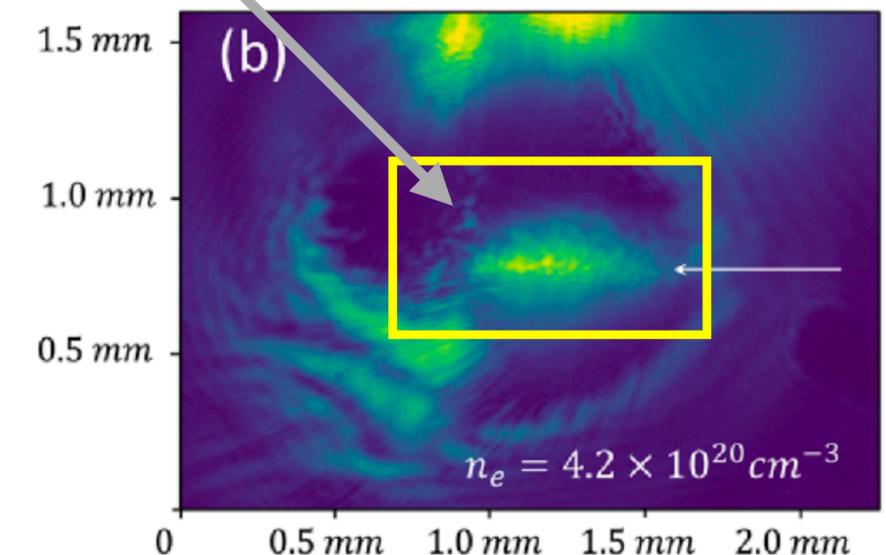
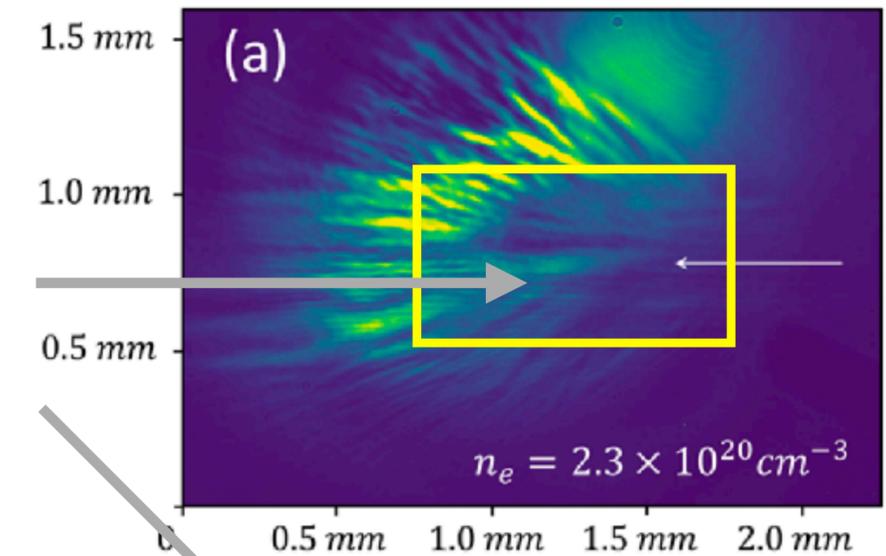
We will use the unique setup to investigate these processes:

- ➔ Measure plasma dynamics and transmitted laser properties
- ➔ Vary density and density gradients
- ➔ Vary laser parameters

Example with near-IR drive: LPI completely obscured due to high densities

We will directly image these phenomena at near-critical densities, not possible with near-IR drivers

From Singh et al. Sci. Rep. **10** (2020)



# Proposed Investigation

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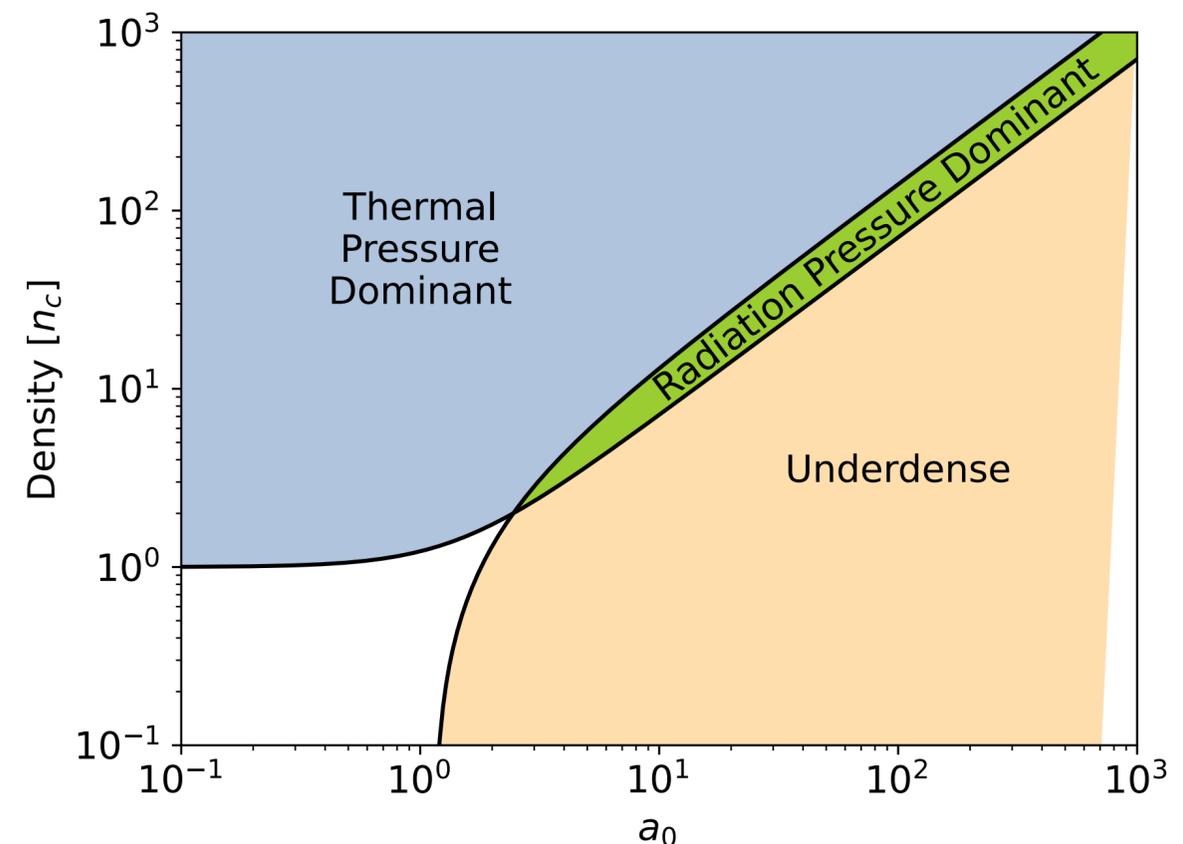
# ATF acceleration conditions

CO<sub>2</sub> laser allows near-critical investigations, enabling the study of shock acceleration mechanisms

Changing the laser and target conditions enables transition of hole-boring (HB) and collisionless shock acceleration (CSA)

$$P_R = n_c m_e c^2 a_0^2$$

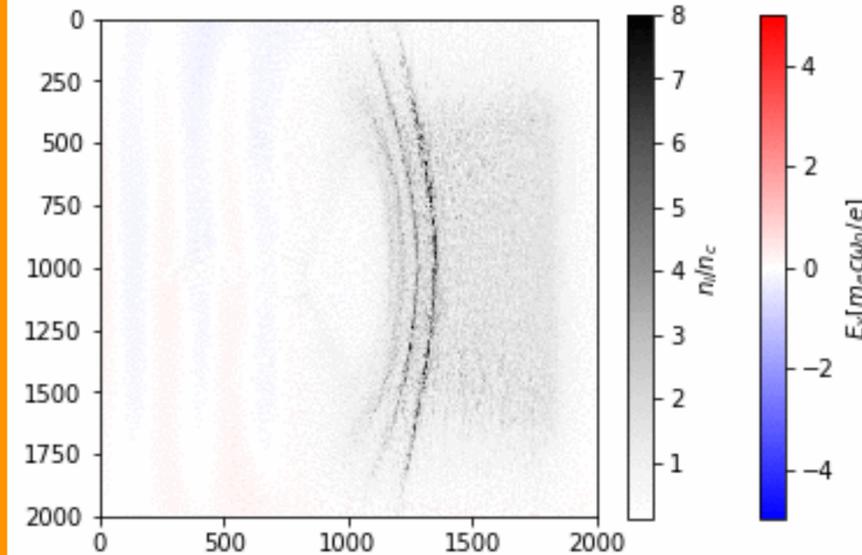
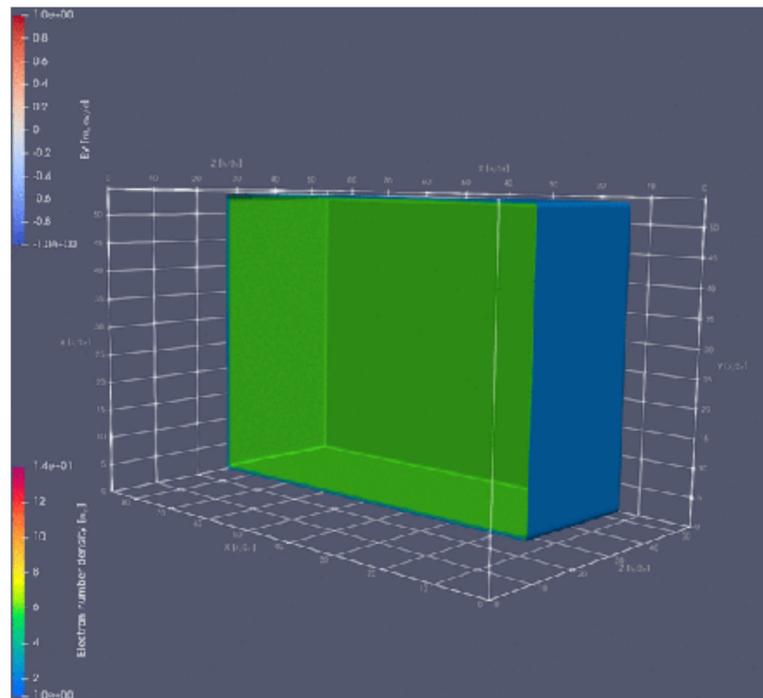
$$P_{Th} = n_e m_e c^2 \left[ \sqrt{1 + \frac{a_0^2}{2}} - 1 \right]$$



# Proposed areas of study

A number of interesting areas of study:

- Continue work of AE100 and make first detailed optical studies of the ion acceleration phase



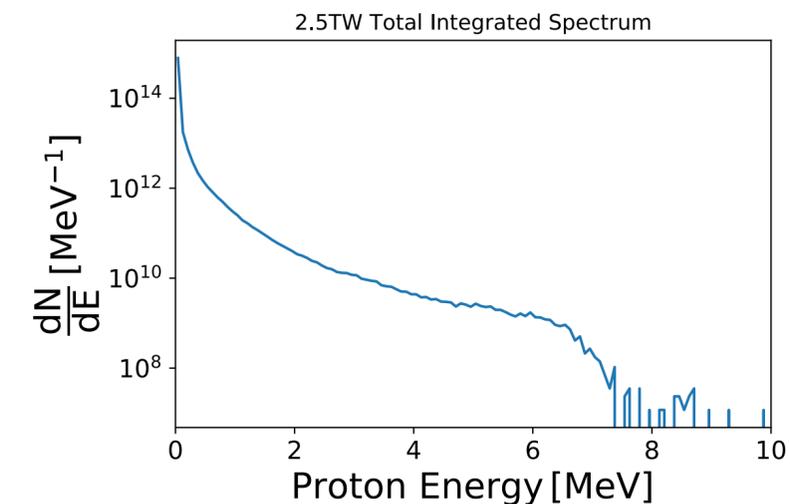
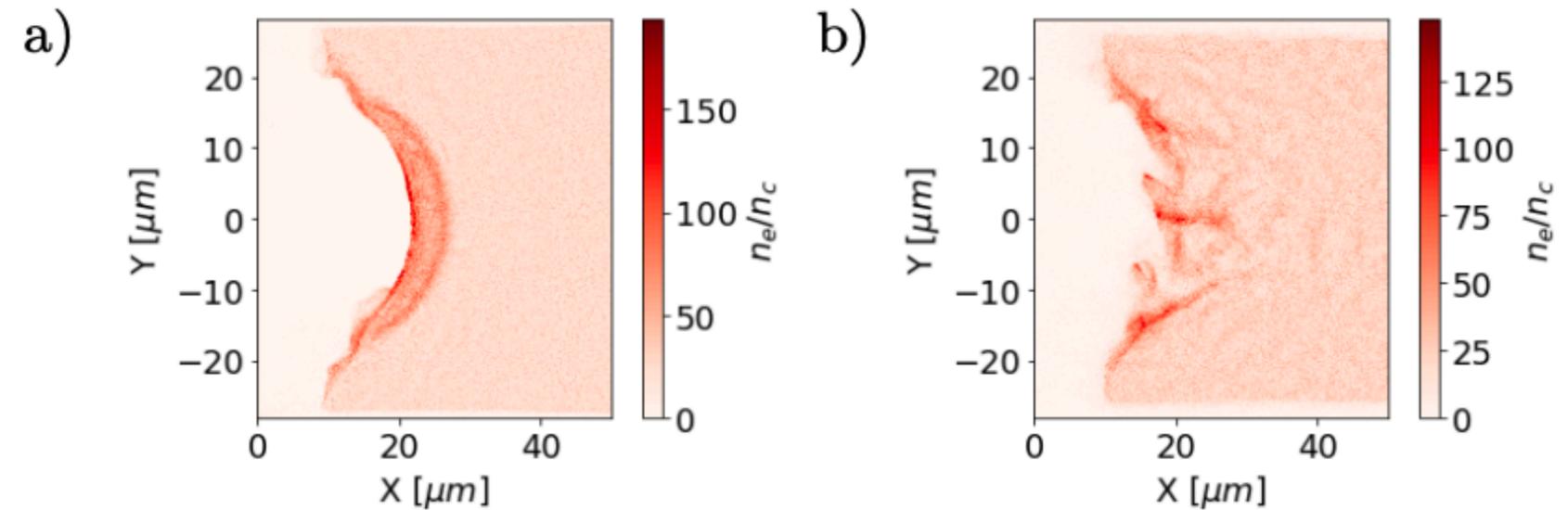
At  $T_e \sim 1$  MeV, and  $1e7$  shock velocity, feature size is  $\sim \mu\text{m}$  scale. In 2 ps, shock moves  $\sim 20 \mu\text{m}$

—> sub 100 fs probe needed

# Proposed areas of study

A number of interesting areas of study:

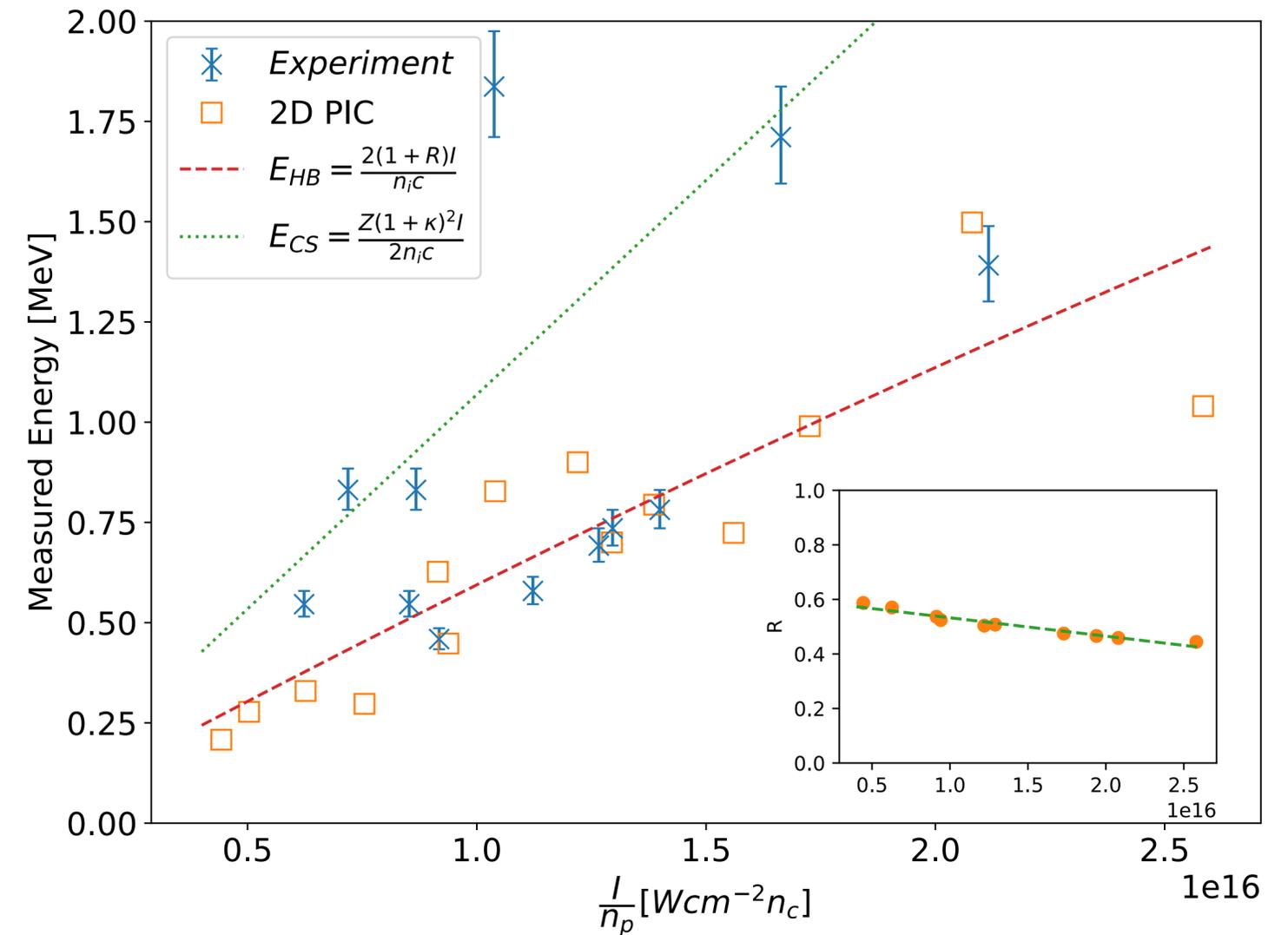
- First detailed optical studies of the ion acceleration phase
- Interplay between CSA and HB and how to control this



# Proposed areas of study

A number of interesting areas of study:

- First detailed optical studies of the ion acceleration phase
- Interplay between CSA and HB and how to control this
- Improved ion acceleration performance at multi-TW level



# Proposed Investigation

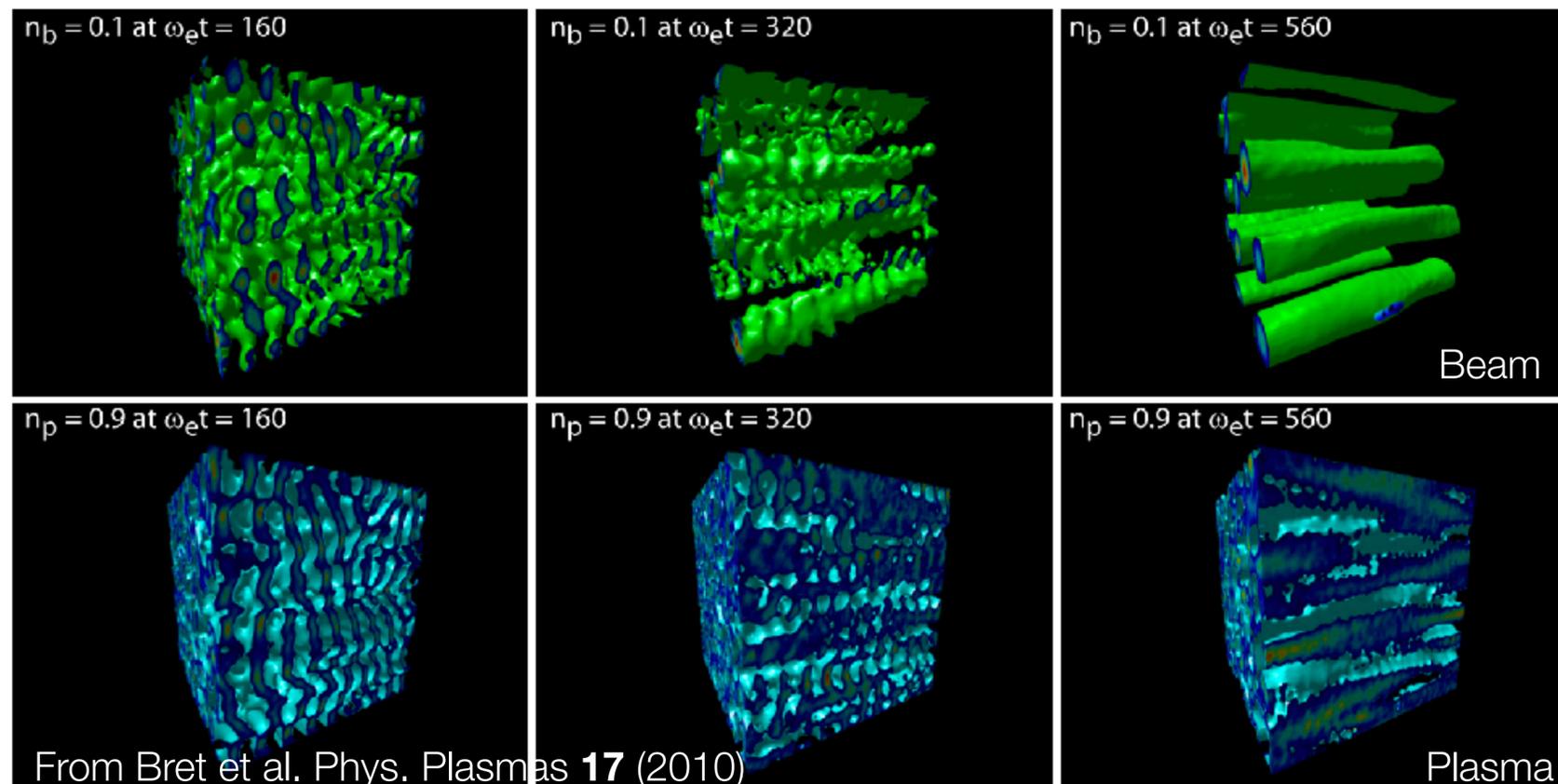
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# Beam propagation through plasmas

- Energetic electrons are generated in LPI and propagate into upstream plasma
- Their current greatly exceeds Alfvén limit  $\rightarrow$  balanced by return current
- These counter propagating populations are subject to collisionless instabilities, such as the current filamentation instabilities



- Affects electron beam transport
- Can degrade ion generation
- Proposed mechanism for magnetic field generation in some astrophysical scenarios

# Beam propagation through plasmas

Studies with near-IR cannot measure this in dense plasmas directly, only through subsequent impact on particles

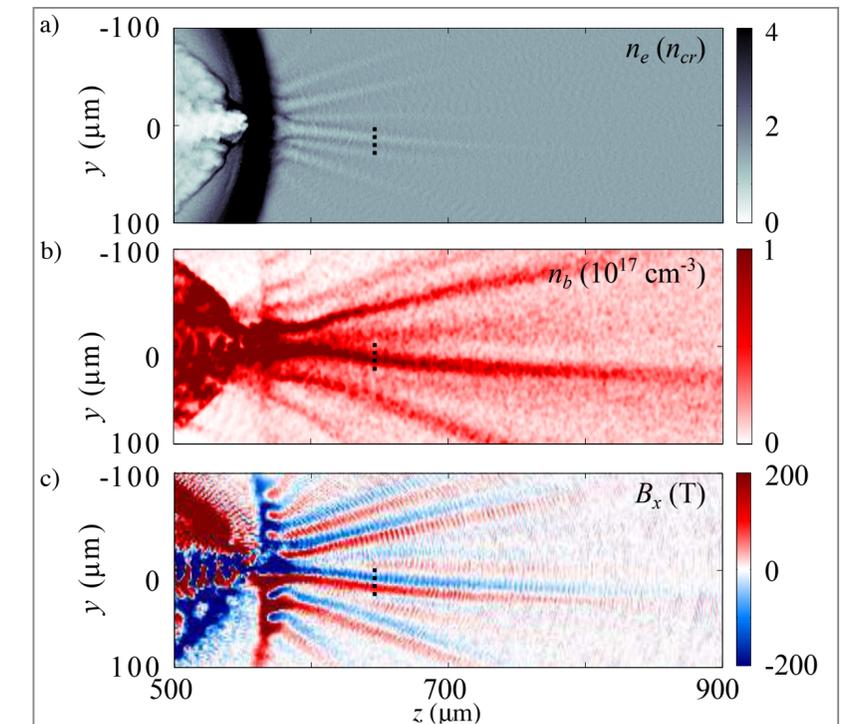
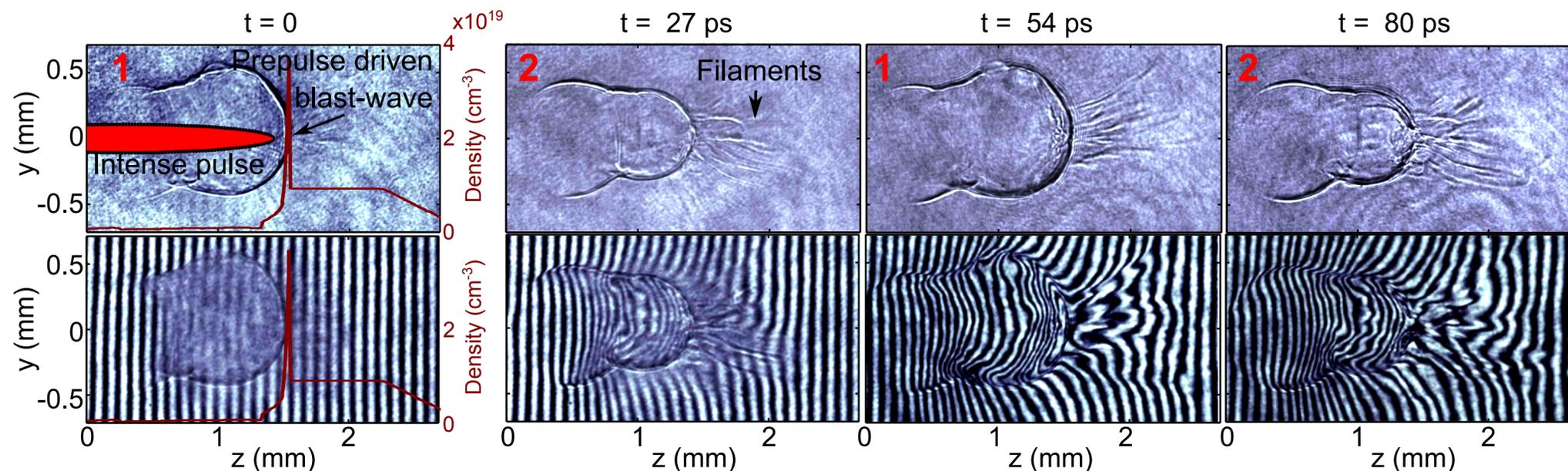
## The ATF can study the same physics, but on a different scale:

- 1) Density achievable by gas jets and ideal for optical probing
- 2) Filament 10x wider - resolvable using TiS
- 3) Instability driven over 10x longer time - time evolution can be captured using TiS

	CO <sub>2</sub> laser@ BNL	Equivalent NIR
$\lambda_L$ ( $\mu\text{m}$ )	9.2	0.8
$a_0$	3	3
$\tau_L$ (fs)	2000	200
$n_e$ ( $\text{cm}^{-3}$ )	$1.3 \times 10^{19}$	$1.7 \times 10^{21}$
$\gamma_e$	$\sim 3$	$\sim 3$
$\alpha = n_b/n_e$	$\sim 0.1$	$\sim 0.1$
$\lambda_{Fil} \sim 2\pi c/\omega_p$ ( $\mu\text{m}$ )	$\sim 10$	$\sim 1$
$\tau_f$ (fs)	$\sim 10$	$\sim 1$
e-folds	$\sim 200$	200

# Beam propagation through plasmas

Previously observed the *endpoint* of current filamentation instability, measuring filamentary density structures after the end of LPI



Previous measurements were limited because growth phase was not resolvable with old YAG probe. **Ti:S will enable time-resolved characterisation of filamentation.**

# Summary - NP-315758 Proposal

- Proposal builds on experience of AE66 and AE100 experiments, exploiting improved laser capabilities
- Objective is to investigate each stage of the ion acceleration process:
  - Laser propagation dynamics in underdense plasmas
  - Acceleration physics at the critical surface
  - Particle beam propagation in plasmas

Thank you for listening.  
Questions?

# Electron Beam Requirements

Parameter	Units	Typical Values	Comments	Requested Values
Beam Energy	MeV	50-65	<i>Full range is ~15-75 MeV with highest beam quality at nominal values</i>	
Bunch Charge	nC	0.1-2.0	<i>Bunch length &amp; emittance vary with charge</i>	
Compression	fs	Down to 100 fs (up to 1 kA peak current)	<i>A magnetic bunch compressor available to compress bunch down to ~100 fs. Beam quality is variable depending on charge and amount of compression required.  NOTE: Further compression options are being developed to provide bunch lengths down to the ~10 fs level</i>	
Transverse size at IP ( $\sigma$ )	$\mu\text{m}$	30 – 100 (dependent on IP position)	<i>It is possible to achieve transverse sizes below 10 <math>\mu\text{m}</math> with special permanent magnet optics.</i>	
Normalized Emittance	$\mu\text{m}$	1 (at 0.3 nC)	<i>Variable with bunch charge</i>	
Rep. Rate (Hz)	Hz	1.5	<i>3 Hz also available if needed</i>	
Trains mode	---	Single bunch	<i>Multi-bunch mode available. Trains of 24 or 48 ns spaced bunches.</i>	

**Electron beam not required**

# CO<sub>2</sub> Laser Requirements

Configuration	Parameter	Units	Typical Values	Comments	Requested Values	
<b>CO<sub>2</sub> Regenerative Amplifier Beam</b>	Wavelength	mm	9.2	<i>Wavelength determined by mixed isotope gain media</i>		
	Peak Power	GW	~3			
	Pulse Mode	---	Single			
	Pulse Length	ps	2			
	Pulse Energy	mJ	6			
	M <sup>2</sup>	---	~1.5			
	Repetition Rate	Hz	1.5	<i>3 Hz also available if needed</i>		
	Polarization	---	Linear	<i>Circular polarization available at slightly reduced power</i>		
	<b>CO<sub>2</sub> CPA Beam</b>	Wavelength	mm	9.2	<i>Wavelength determined by mixed isotope gain media</i>	
<i>Note that delivery of full power pulses to the Experimental Hall is presently limited to</i>		Peak Power	TW	5	<i>~5 TW operation will become available shortly into this year's experimental run period. A 3-year development effort to achieve</i>	5TW
		Pulse Mode	---	Single		
		Pulse Length	ps	2		2ps
		Pulse Energy	J	~5	<i>Maximum pulse energies of &gt;10 J will become available within the next year</i>	5
		M <sup>2</sup>	---	~2		
		Repetition Rate	Hz	0.05		Highest available
		Polarization		Linear	<i>Adjustable linear polarization along with circular polarization can be provided upon request</i>	LP/CP

# Other Experimental Laser Requirements

<b>Ti:Sapphire Laser System</b>	<b>Units</b>	<b>Stage I Values</b>	<b>Stage II Values</b>	<b>Comments</b>	<b>Requested Values</b>
Central Wavelength	nm	800	800	<i>Stage I parameters should be achieved by mid-2020, while Stage II parameters are planned for late-2020.</i>	✓
FWHM Bandwidth	nm	20	13		✓
Compressed FWHM Pulse Width	fs	<50	<75	<i>Transport of compressed pulses will initially include a very limited number of experimental interaction points.</i>	≤75
Chirped FWHM Pulse Width	ps	≥50	≥50		
Chirped Energy	mJ	10	200		
Compressed Energy	mJ	7	100		7
Energy to Experiments	mJ	>4.9	>80		4.9
Power to Experiments	GW	>98	>1067		99

<b>Nd:YAG Laser System</b>	<b>Units</b>	<b>Typical Values</b>	<b>Comments</b>	<b>Requested Values</b>
Wavelength	nm	1064	<i>Single pulse</i>	
Energy	mJ	5		5mJ+
Pulse Width	ps	14		
Wavelength	nm	532	<i>Frequency doubled</i>	X
Energy	mJ	0.5		X
Pulse Width	ps	10		X

# Special Equipment Requirements and Hazards

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- Electron Beam **N/A**
- CO<sub>2</sub> Laser
  - Please note any specialty laser configurations required here:
- Ti:Sapphire and Nd:YAG Lasers
  - Please note any specialty non-CO<sub>2</sub> laser configurations required here:
    - **YAG amplifier for highest possible energies**
- Hazards & Special Installation Requirements
  - **New magnet installation for particle spectrometer - 0.6T (already ordered)**
  - **HV for time-of-flight diamond detector**

# Experimental Time Request

## CY2024 Time Request

Capability	Setup Hours	Running Hours
Electron Beam Only		
Laser* Only (in FEL Room)		40 120
Laser* + Electron Beam		

## Time Estimate for Full 3-year Experiment (including CY2024-26)

Capability	Setup Hours	Running Hours
Electron Beam Only		
Laser* Only (in FEL Room)	120	360
Laser* + Electron Beam		

\* Laser = Near-IR or LWIR (CO<sub>2</sub>) Laser

# Experimental Layout

