



Application of surface polaritons in thin SiC films to sub-wavelength lithography and compact particle acceleration

Gennady Shvets, University of Texas at Austin

Collaborators:

Alan Feinerman (UIC) Chris Zorman (Case Western) Herman Geraskin (UT) Yaroslav Urzhumov (UT)

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What is proposed?

Test several near-field concepts at the ATF facility using its short-pulse CO_2 lasers:

- -Sub-micron patterning of surfaces with a CO₂ laser
- -Development of a SiC "single-wafer" accelerator/electron injector, also driven by a CO_2 laser

Who funds (or may fund) this?

•NSF's Nanoscience and Engineering Research (NSER) Program → sub-wavelength lithography

- •Argonne's LDRD → cold tests and fabrication of SiC high-gradient accelerator structure
- •DOE HEP \rightarrow proposal in preparation

Collaborators and the team for this proposal

PI: Gennady Shvets, University of Texas at Austin

Senior researchers:

Alan Feinerman (UIC) \rightarrow fabrication facility Chris Zorman (Case Western) \rightarrow growth of SiC films Dmitry Korobkin (postdoc, joint UT/Argonne) \rightarrow experiment (fabrication, thermoresist, laser testing)

Graduate students:

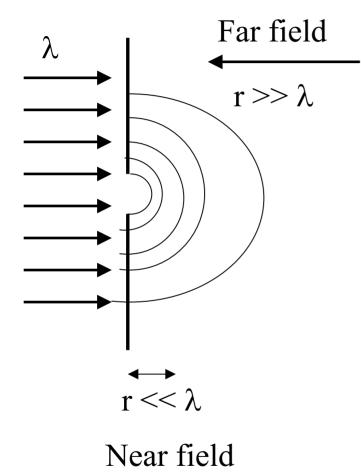
Herman Geraskin (UT) \rightarrow experiment, simulations

ATF: Igor Pogorelsky, Vitaly Yakimenko, Igor Pavlishin

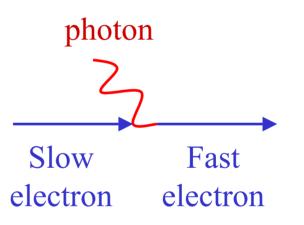
LDRD-funded effort at Argonne:

Wei Gai, John Power

What's wrong with far field, or What do nanolithography and accelerators have in common?

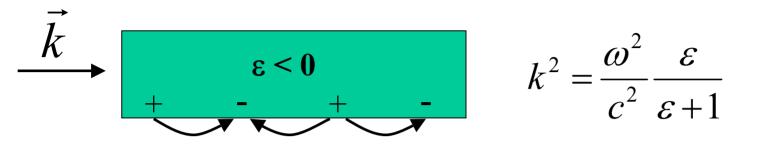


•Impossibility of linear in electric field acceleration (Woodward-Palmer's theorem)



•Impossibility of subwavelength resolution in far field (Abbe's resolution limit)

Near-field example: surface waves



vacuum, $\varepsilon = 1$

•Exponentially decays away from the interface
•Can accelerate non-relativistic charged particles
•Can be resonantly excited in the near-field

We propose to excite surface waves in mid-IR with a CO₂ laser and use it for lithography and accelerator applications

Outline of the talk

- Motivation: why near field in mid-IR?
- Beating diffraction with a "super-lens": Enhanced Near-Field Imaging
 - 1. PENFIL: Phonon-Enhanced Near-Field Infrared Lithography with silicon carbide
- SiC-based miniature surface wave accelerator :
 - 1. Principle of operation and advantages
 - 2. Wafer fabrication
 - 3. Deposition of SiC films and testing

Why sub- λ resolution in mid- to far-IR?

 Many materials (plastic, ceramics, polymers) strongly absorb in IR → "sculpting" various surfaces on a nanoscale by delivering IR to small (100 nm)³ volumes (laser ablation)

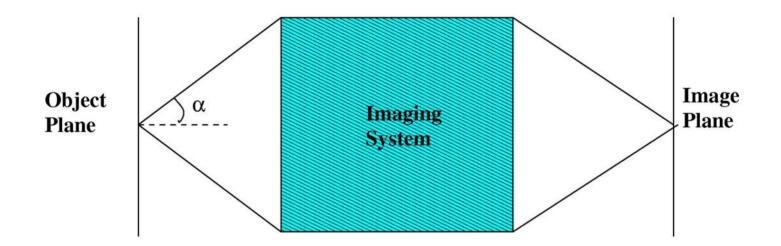
•Small biological objects (e.g., spores) can be imaged one at a time, by studying infrared Raman resonances.

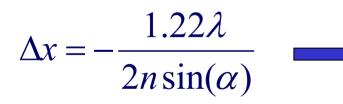
•IR imaging and manipulation of biological molecules can be preferable to shorter wavelengths because it does not alter chemical composition.

•Ultra-compact accelerators driven by CO₂ lasers

EM fields predominantly longitudinal for very small (sub- λ) gaps

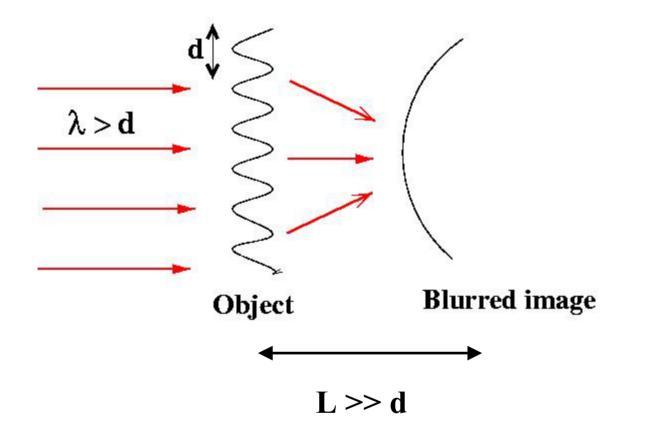
Abbe's Resolution Limit for Optical Microscope





Hard to resolve objects much smaller than $\frac{1}{2} \lambda$ using conventional far-field optics.

Imaging small-feature objects in far field



Small features (or large wavenumbers) of the object are lost because of the exponential decay of evanescent waves

Exponential decay of large-k waves

•Magnetic and electric fields in the object plane are represented by a sum of plane waves:

$$H_z(x=0,y) = \int dk_\perp A(k_\perp) e^{ik_\perp y} \qquad f$$

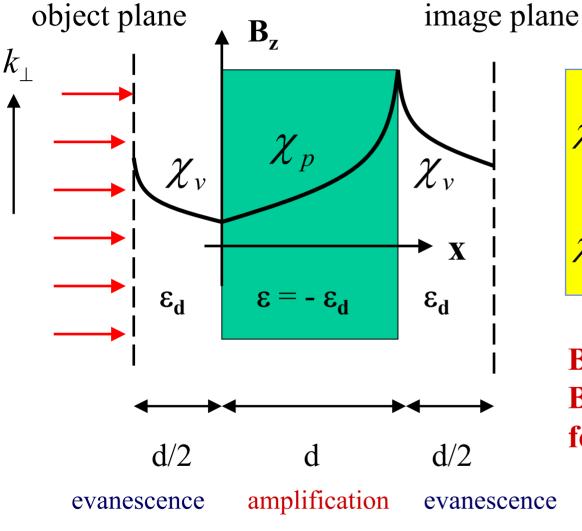
Spectral width $\Delta k \rightarrow$ feature size $\Delta y = \Delta k^{-1}$

•Image reconstruction → propagation of A(k) through vacuum, medium, and lenses:

 $\begin{aligned} & \underset{\mathbf{M}}{\text{Intropy}} \begin{cases} A(k_{\perp}) \propto e^{i\sqrt{k_0^2 - k_{\perp}^2}x} & \text{For smooth features with } \mathbf{k} < \mathbf{k}_0 = \omega_0/\mathbf{c} \\ & A(k_{\perp}) \propto e^{-\sqrt{k_{\perp}^2 - k_0^2}x} & \text{For sharp features with } \mathbf{k} > \mathbf{k}_0 = \omega_0/\mathbf{c} \end{aligned}$

•How to overcome decay? Go near-field! (a) get real close (L < d)(b) amplify evanescent waves

Surface waves enable super-lensing

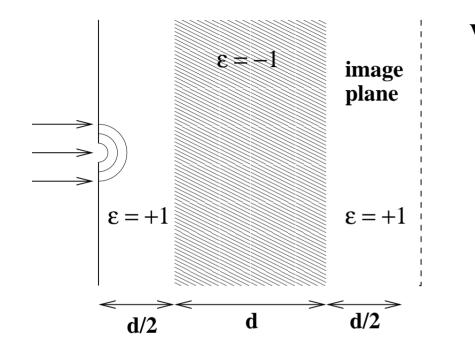


 $\chi_{\nu} = \sqrt{k_{\perp}^2 - \varepsilon_d \mu_d \frac{\omega_0^2}{c^2}}$ $\chi_p = \sqrt{k_{\perp}^2 - \varepsilon \mu \frac{\omega_0^2}{c^2}}$

 $B_z(object plane) = B_z(image plane)$ for $\mu = -1$, $\varepsilon = -\varepsilon_d$

 $\mu = +1$ works too!

Poor man's super-lens: $\varepsilon < 0$, $\mu = 1$



What determines the accuracy of image reconstruction in the image plane?

 Retardation (finite lens thickness) → Important for μ = 1 even if ε = -ε_d

2. Losses
$$\rightarrow \varepsilon = -\varepsilon_d + i \sigma$$

$$T = \frac{4(\chi_v / \varepsilon_v)(\chi_p / \varepsilon_p)\exp(-\chi_v d)}{(\chi_p / \varepsilon_p + \chi_v / \varepsilon_v)^2 \exp(\chi_p d) - (\chi_p / \varepsilon_p - \chi_v / \varepsilon_v)^2 \exp(-\chi_p d)}$$

Ideally, want T = 1, but issues (1) and (2) interfere...

Choose a low-loss negative ε materials and design a near-field "super-lens"!

•Plasma:
$$\varepsilon = 1 - \omega_p^2 / \omega^2$$

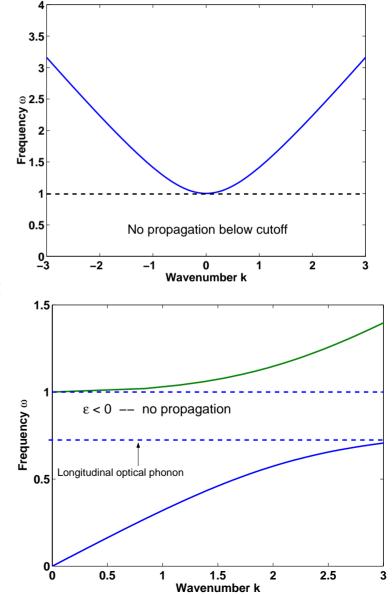
 $\omega_p^2 = 4\pi e^2 n/m$

Typically, ω_p is in 10¹⁵ Hz range in metals \rightarrow surface plasmons

•Polaritonic crystals (ZnSe, SiC):

$$\varepsilon = \varepsilon(\infty) \frac{\omega_L^2 - \omega^2 + i\gamma\omega}{\omega_T^2 - \omega^2 + i\gamma\omega}$$

Typically, ω_T and ω_L are in 1-30 THz range \rightarrow surface phonons



Proven properties of phononic materials

•Smaller damping of IR phonons (reduced by cooling)

$$\varepsilon = \varepsilon(\infty) \frac{\omega_L^2 - \omega^2 + i\gamma\omega}{\omega_T^2 - \omega^2 + i\gamma\omega}$$

For SiC
$$\gamma/\omega_L = 0.005 \ll 1$$

For ZnSe $\gamma/\omega_L = 0.001 \ll 1$

•Higher sensitivity to surface quality and laser frequency is explored in surface science! (Hillenbrand et. al., Nature 2002)

•Desirable $\varepsilon = -1$ regime is not accessible in mid-IR using metal films

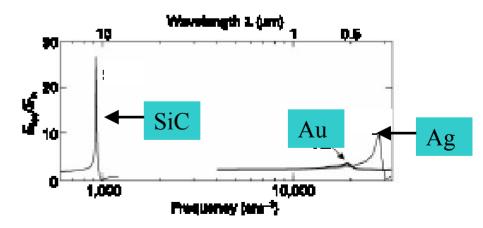
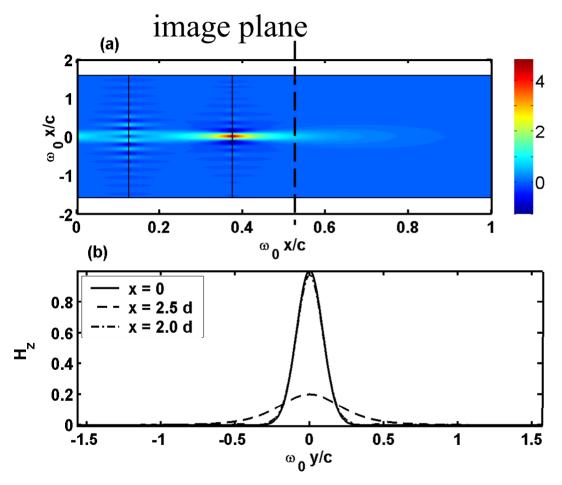


Figure 1 Calculated field enhancement E_{ind}/E_{in} due to Fröhlich resonance at the surface of a 10-nm-diameter sphere. The infrared lattice vibration (phonon) of a polar detectric (SIC) induces a considerably stronger near-field amplitude than the green/blue electronic excitation (plasmon) of a noble metal (Au, Ag).

Enhanced Near-Field Imaging (from Shvets, Phys Rev B 2003)



•Object: thin slit of width $w = 0.21 \ \mu m = \lambda/50 \ at x=0$

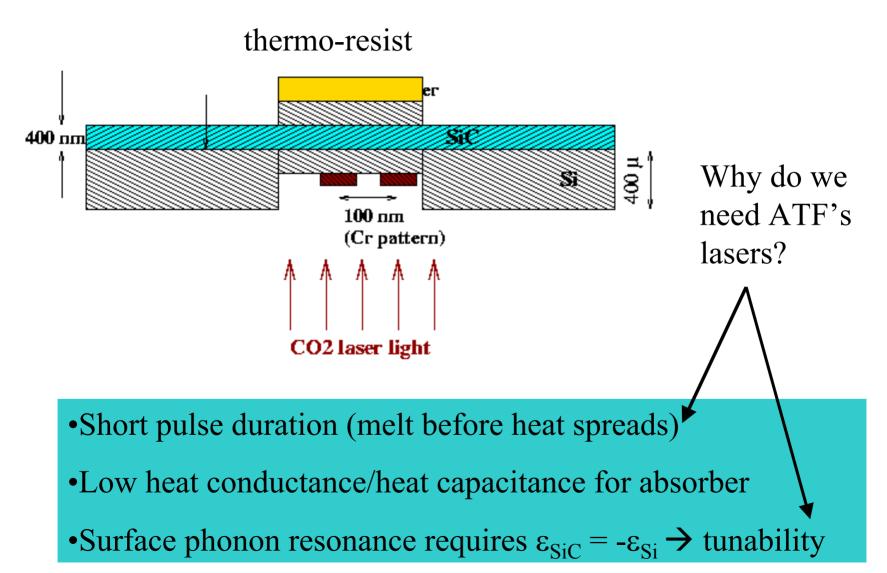
Lens: between x=d/2 and x=1.5d planes
(SiC with ε = -1)

•Image: in x=2d plane

•Blurred image (conventional near-field): in x=2.5d

Observe significant image improvement with SiC lens

Phonon Enhanced Near Field Infrared Lithography (PENFIL)



Applications of miniature accelerators



•Miniature X-ray sources for medical imaging

- •Novel high-gradient HEP accelerators
- •Injectors into plasma accelerators
- •Tiny cathodes exist for tiny accelerators: CNTs in 50nm bundles

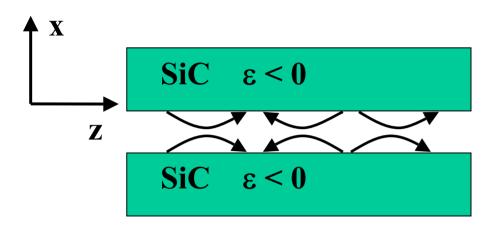
Otto Zhou (2002): 14keV X-ray image using CNT cathode

-Durable (field emission v.s thermal)

Challenge: increase electron energy (100's keV for medicine, 10's MeV/wafer for HEP)

Solution: single-wafer surface-wave accelerator

Novel miniature <u>surface-wave</u> accelerator

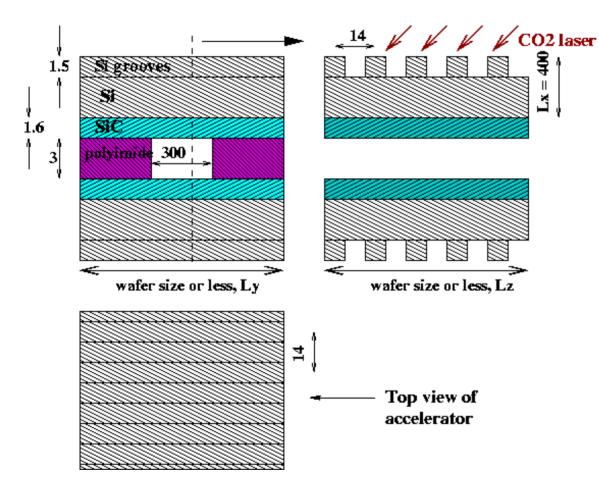


•Supports ω = kc mode

•Near field (small gap) \rightarrow attractive ratio E_z/E_x

- •Can support high fields → no breakdown (for W=300 Mev/m DC) because of large band gap: 2.3 eV for 3C-SiC
- •Supports high frequency surface waves ($\epsilon = -1.2$ for $\lambda = 10.6 \mu m$) excitable by CO₂ laser.

Schematic of SiC Surface Wave Accelerator



Accelerating surface mode: $\omega = \omega(k)$

•Grating necessary to couple into Si with appropriate wavenumber

•Accelerating mode "leaks" from cavity into Si → high-Q coupling

•Cold tests: study reflection and transmission coefficients as functions of incident angle and frequency



Initial fabrication was done at

Microfabrication Applications Laboratory

at University of Illinois at Chicago

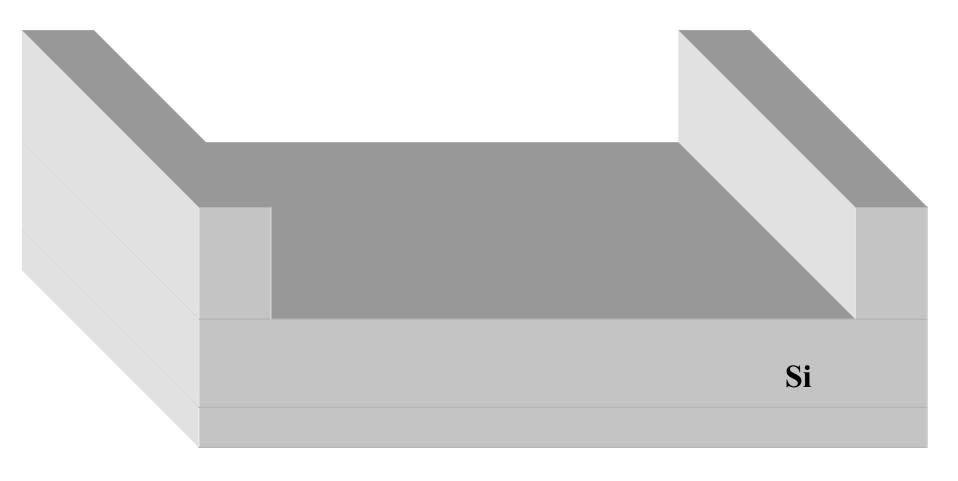
http://www.mal.uic.edu/

Summer 2003

As a base of acceleration structure we use conventional Silicon wafers:



3-4 microns deep cavity is made on one side of the wafer:



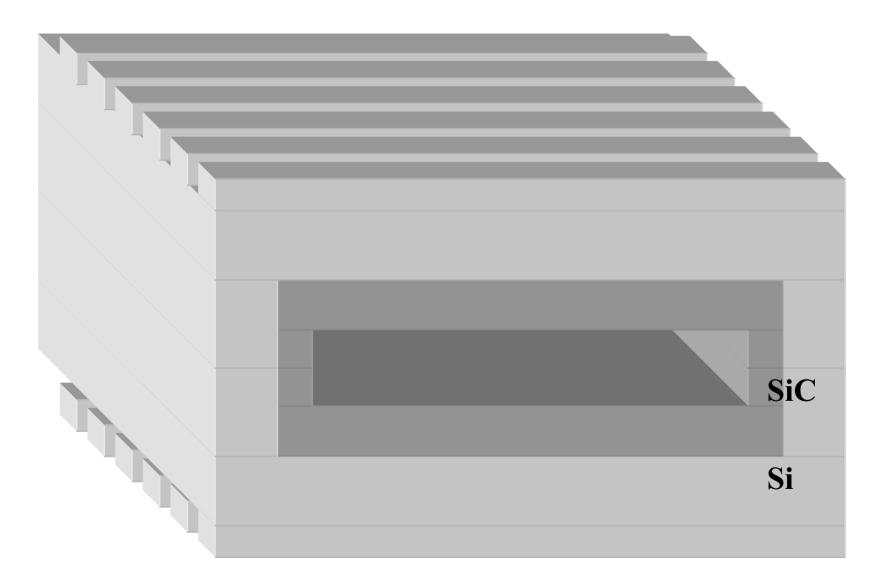
1.5 - 2 microns SiC film is deposited on the internal surface of cavity (done at CWRU)



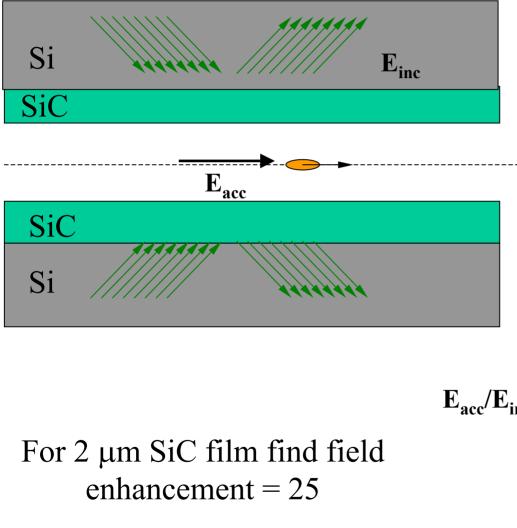
Diffraction grating is made on the other side of wafer:



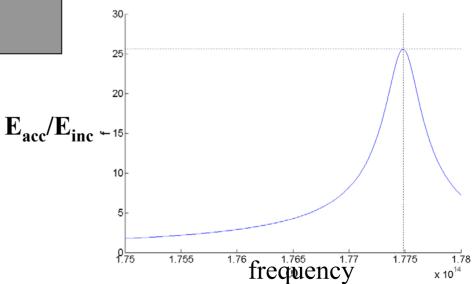
Wafers are bonded together – accelerator is ready!



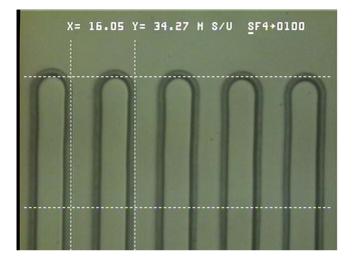
What's so great about this accelerator?



- •Can maximize the ratio $E_{acc}/E_{inc} \rightarrow$ resonance!
- Seek this maximum by
 varying laser frequency
 or incidence angle →
 drop in reflectivity

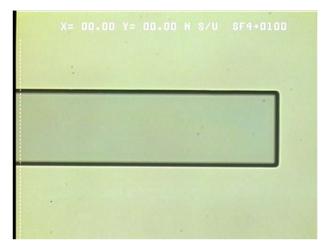


Fabrication Progress



Coupling grating, spacing = $16 \mu m$, depth = $1.6 \mu m$





Accelerating cavity, depth = $3.0 \ \mu m$

$1.5\ \mu m$ layer of SiC at the bottom of the Si cavity has been deposited

Conclusions

"Explosive" combination of a low-loss polaritonic material (SiC) and a short-pulse high-power driver (CO₂ laser) can be explored at the ATF for two promising applications

•Enhanced near-field nano-lithography in IR

•New advanced accelerator concept: not metallic, not dielectric, not plasma \rightarrow single-wafer surface wave accelerator. Uses standard wafer processing techniques \rightarrow potentially cheap!

We already know how to fabricate the structures. Preliminary proof-of-principle cold tests will be done at UT and Argonne.

We need short-pulse (30 - 100 ps) CO2 laser to demonstrate lithography and to test the accelerating structure. We are undemanding users: don't require a beam, don't ask for 2 ps pulses.