

# **High-brightness picosecond ion beam source based on BNL TW CO<sub>2</sub> laser: Proof-of-principle experiments**

*We report initial results and the status of the  
**first ever experiments** of proton and ion acceleration  
by ultra-short, high-intensity **far-infrared** (10  $\mu\text{m}$ ) laser pulses in metal foils.*

***Proton energy spectra** reveal a broad but pronounced **energy peak** at  $\sim 1$  MeV.*

*The peak, **never previously observed** with unstructured targets may be  
the **first experimental** indication of direct **Radiation Pressure acceleration**  
of protons by a circularly polarized laser.*

## *People*

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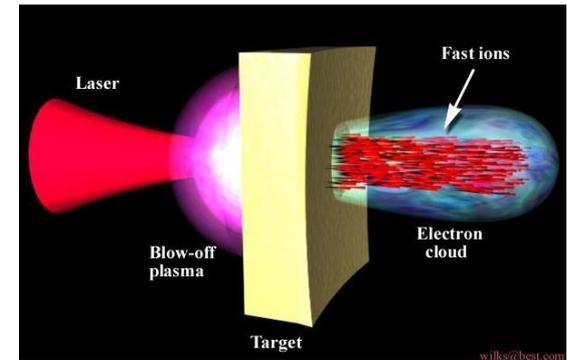
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*The work is supported by DOE HEP grant*

# Background

## The effect

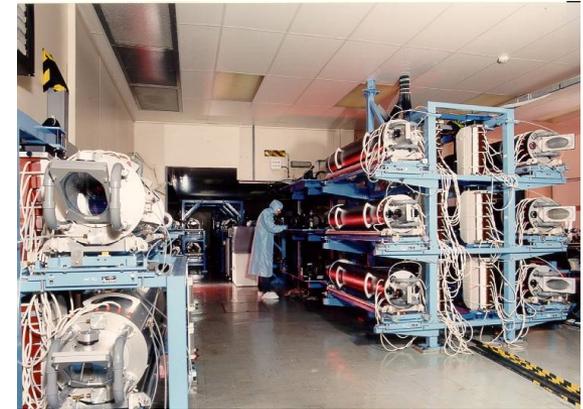
Thin foils irradiated by sub-picosecond laser pulses at relativistic intensities emit ions/protons at MeV energies in bright, well-collimated beams.



## The current status of the field

Ultra-fast solid state lasers, at  $\lambda \sim 1 \mu\text{m}$ ; usually **national facilities**.

Proton and heavier ions at **25-100 MeV** (potentially GeV) at **intensities 2 orders** of magnitude larger than those at conventional accelerators.



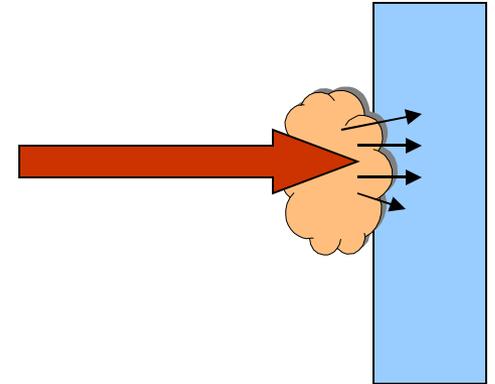
## We propose

To build and study a **high-brightness multi-MeV ion and proton beam source** driven by the unique **picosecond TW CO<sub>2</sub>** laser available at the ATF.

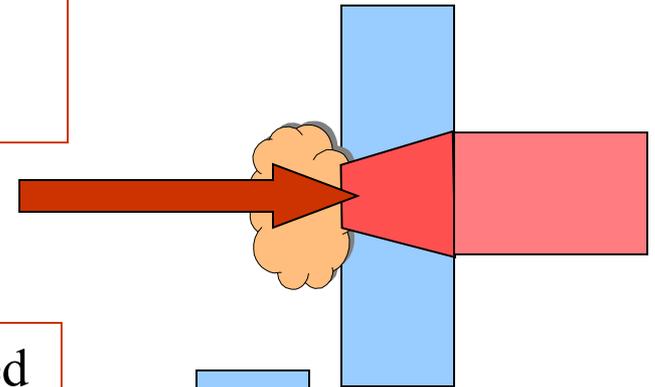


# *Physics of ion (proton) acceleration by a laser in a foil target: TNSA*

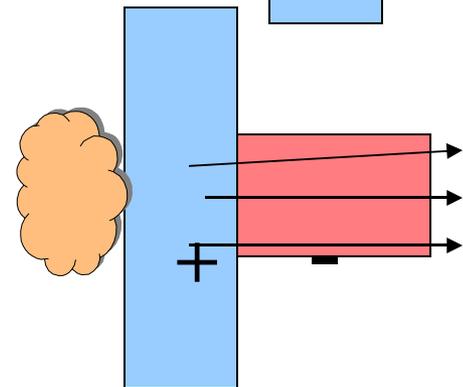
**Stage 1.** On the target surface, the laser beam creates plasma with **relativistic electrons moving into the target**



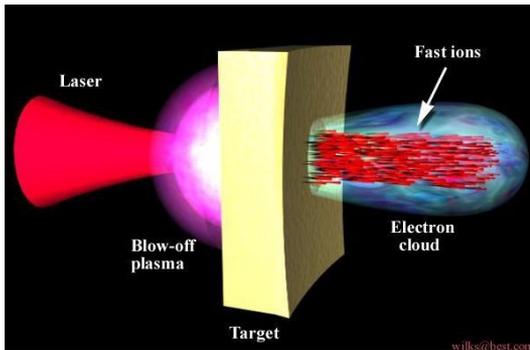
**Stage 2.** Relativistic electrons propagate through the target and form a dense cloud behind it



**Stage 3.** Electron cloud and the positively charged target create electric field ( $\sim 10$  GV/m), which ionizes atoms of the target back surface and pulls ions (protons) out, accelerating them



## *TNSA with CO<sub>2</sub> laser*



TNSA with CO<sub>2</sub> laser:

potential pros and cons as compared to solid-state lasers

**Pros:** larger wavelength and longer pulses →  
more electrons; relatively high ponderomotive potential

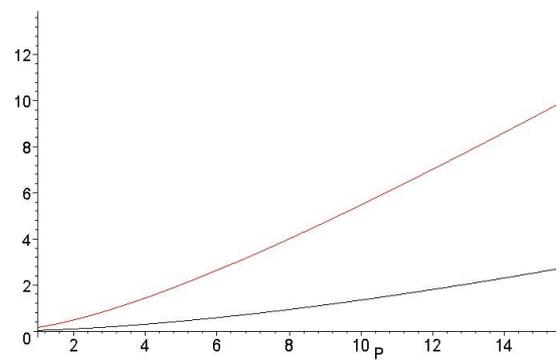
**Cons:** larger wavelength and longer pulses →  
larger electron volume; target heating

**Bottom line:**

- *significant promise of the CO<sub>2</sub> laser*
- *need to thoroughly investigate the physics in this new regime*

# Unique potentials of ATF CO<sub>2</sub> laser

**1. Uniform solid targets**  
Laser of moderate power  
New physics: very little is known



**Proton energy (MeV) vs laser power (TW). Red is for 10 μm, blue is for 1 μm**

## 2. Structured targets: H-rich microdots

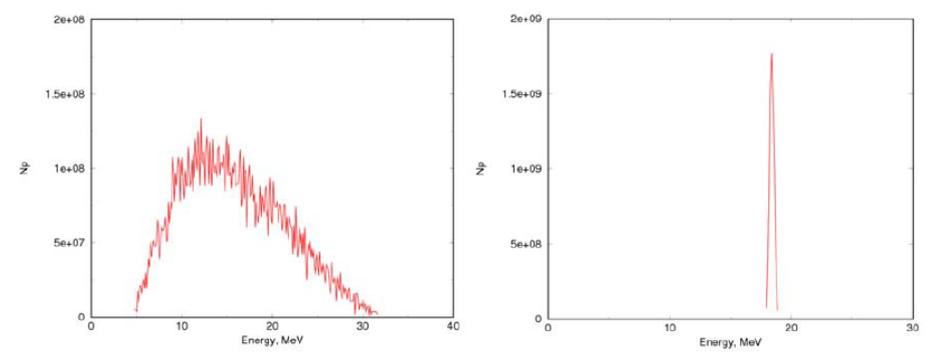
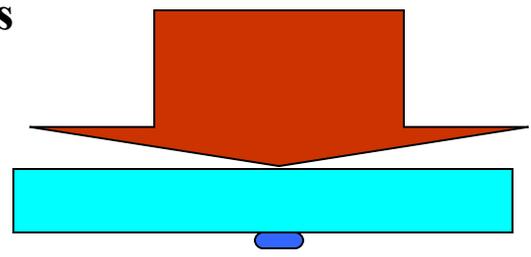
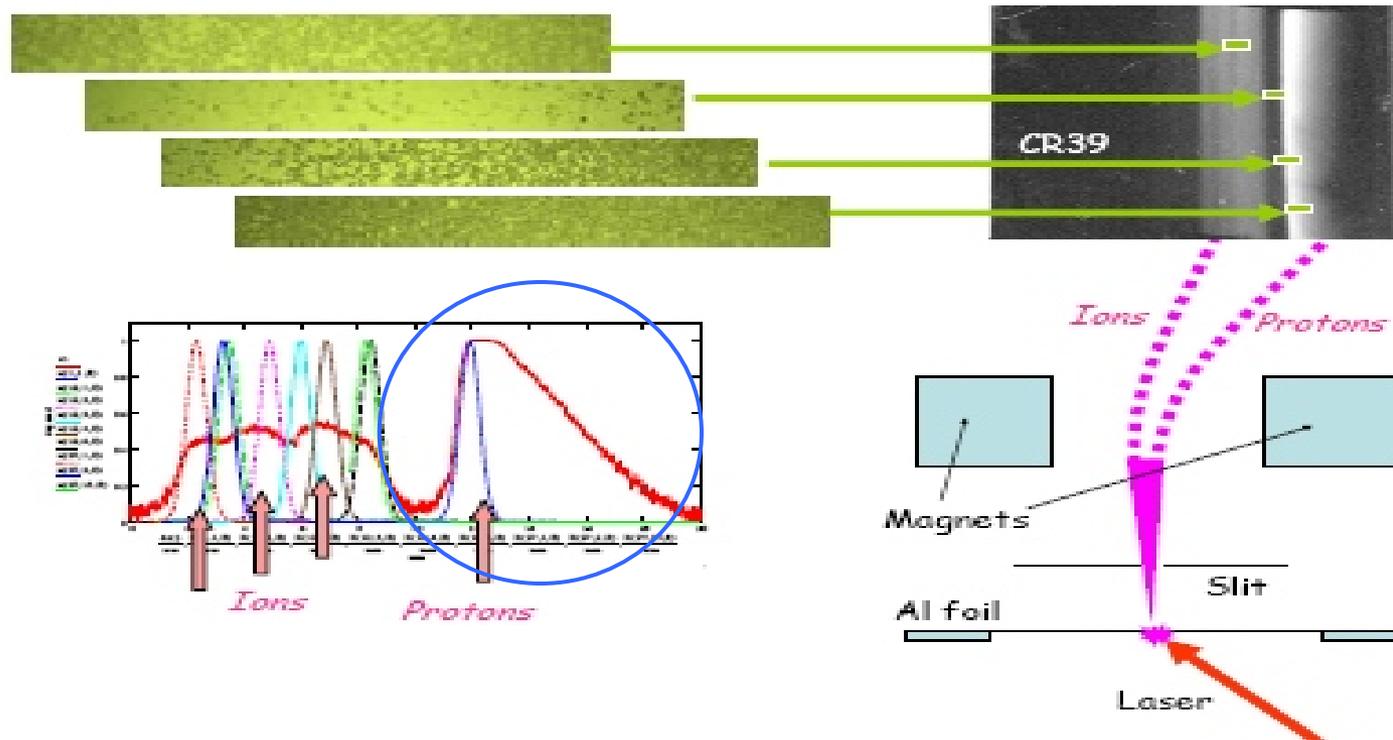


Fig. 8 Proton energy spectrum from a structured target. (a) Solid state laser with  $\lambda=1\mu\text{m}$ . (b) CO<sub>2</sub> laser with  $\lambda=10\mu\text{m}$ . The CO<sub>2</sub> laser produces a much narrower proton spectrum because of the narrower phase space fill.

**3. Low critical plasma density:  $10^{19} \text{ cm}^{-3}$**   
*Gas jets as targets*



# Experimental observations



Proton and ion energies were diagnosed with a compact magnetic spectrometer followed by a CR39 plate. A density plot obtained with a scanner is overlaid with artificial peaks that illustrate the expected position and spread due to a finite spectrometer slit for the 1-MeV protons and any possible  $\text{Al}^{+n}$  and  $\text{C}^{+n}$  ions accelerated in the same field.

## Circular vs linear laser polarization

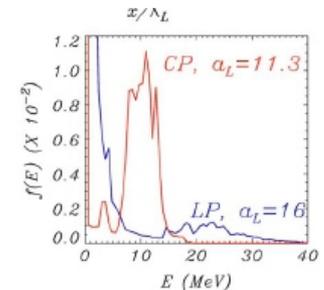
The *peak* nature of the proton spectra *highly unusual* for TNSA

Our laser radiation is **circularly** polarized → At normal incidence, electron heating suppressed → *TNSA suppressed*

Practically **all** experiments so far – with **linear** laser polarization

Recent theoretical predictions (*T. Liseikina et al*)  
Radiation pressure acceleration at the target's **front**

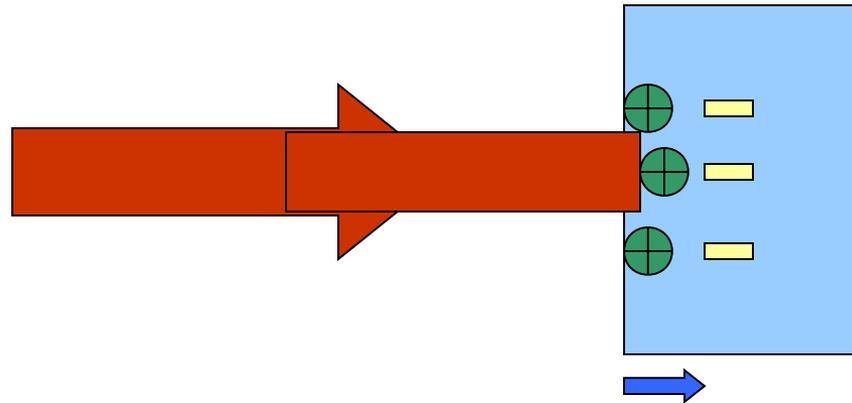
→ **Peak in energy spectra; Not yet observed experimentally**



Our case **more complicated: circular polarization, but at 45°**

## Radiation pressure acceleration by a circularly polarized laser

Normal incidence: *Very little electron heating by the laser directly*



*Maximum proton energy*  $E_{\max}^p (\text{MeV}) \approx a_0^2$  ,  $a_0^2 \approx 0.36 I_{18} \lambda^2 (\mu\text{m})$

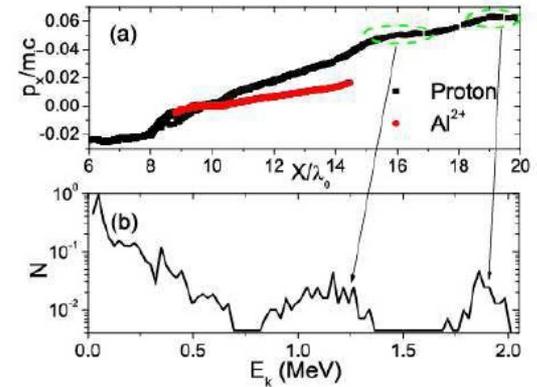
*Number of protons*  $N^p \approx v\pi\lambda / 4\pi$

# 1D and 3D PIC simulations

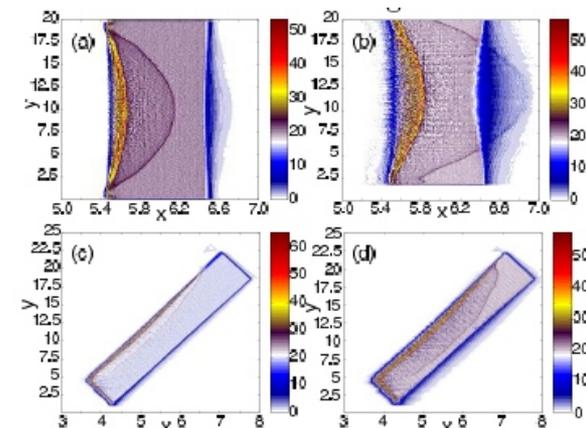
*Both mechanisms present*

*→ Need detailed computer simulations*

**1D** (a) Proton and Al<sup>2+</sup> distributions in the phase space at  $t/T_0 = 500$ . (b) Energy spectrum of the forward moving protons at  $t/T_0 = 500$



**3D** Distribution of ions at the time of  $t/T_0=40$  (a) and  $t/T_0=60$  (b) with the laser pulse normally incident onto the target; (c) and (d) correspond to the laser oblique incidence with the same time points



*Bottom line: interplay of two mechanisms retains quasi-monoenergetic peaks*

## ***Conclusion and near-term plans***

The first ever experiments of proton and ion acceleration ***by ultra-short, high-intensity far-infrared laser*** pulses interacting with metal foils.

Sub-TW, 6-ps, circularly polarized CO<sub>2</sub> laser pulses focused onto 12 μm-thick Al foils drive ion acceleration in the forward direction normal to the rear target surface.

The spectra of protons reveal a broad yet quite pronounced ***proton energy peak*** at ~1 MeV –never observed before with unstructured targets.

This peak may be the ***first experimental indication*** of direct RPA of protons by circularly polarized laser.

- Circular polarization at normal incidence
- Increase the laser output and further suppress pre-pulses
- Tighter focusing
- Thomson spectrometer for ion beam diagnostics
- Real-time beam diagnostics

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