

Modeling

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**Hydrodynamic
Modeling
for the search of
anomalous chiral effects**

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Outline

Modeling of chiral effects

Effects of dynamical EM fields

Outlook for future efforts

Anomalous chiral effects

$$\dot{\mathbf{j}}_{\text{anom}} = \kappa_B \mathbf{B} + \kappa_\omega \boldsymbol{\omega}$$

$$\dot{\mathbf{j}}_{5,\text{anom}} = \xi_B \mathbf{B} + \xi_\omega \boldsymbol{\omega}$$



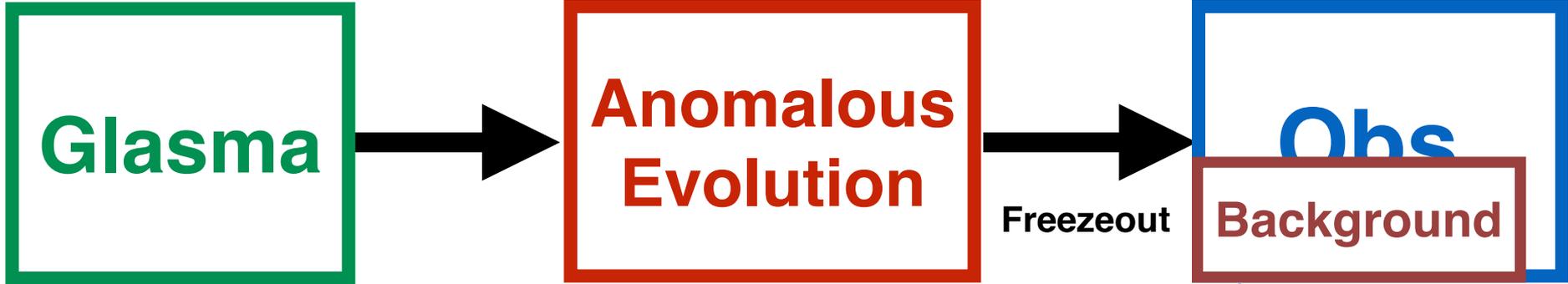
Theoretical frameworks

- Anomalous hydrodynamics

[Son-Surowka 2009; ...]

- Chiral kinetic theory

[Son-Yamamoto; Stephanov-Yin; ...]



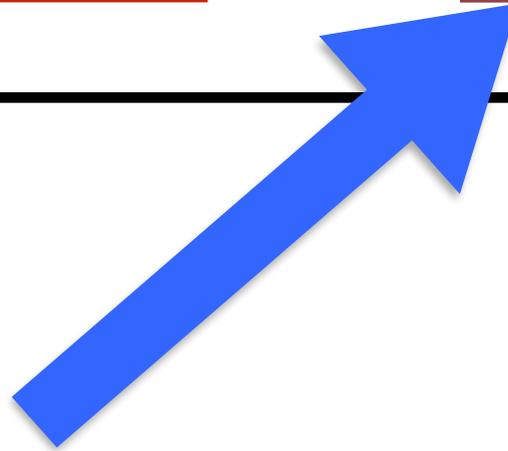
time

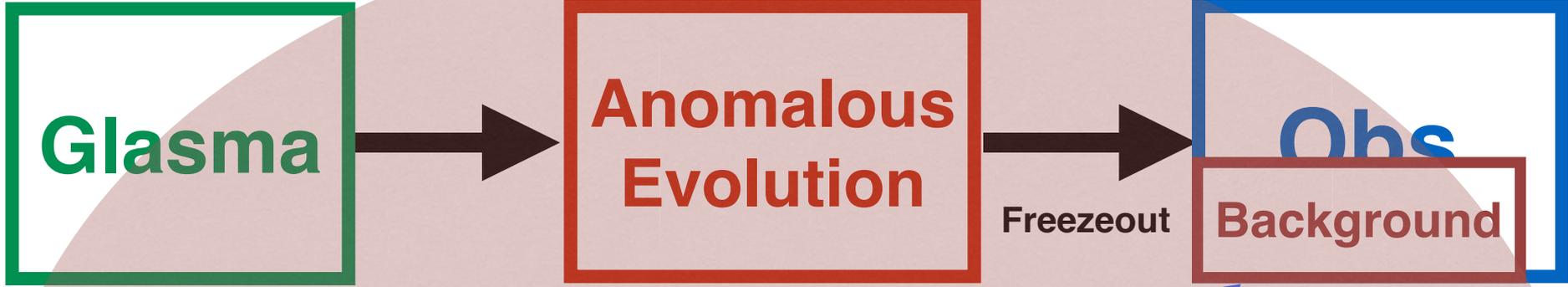
Magnetic field

CME/CMW

Vorticity

CVE/CVW



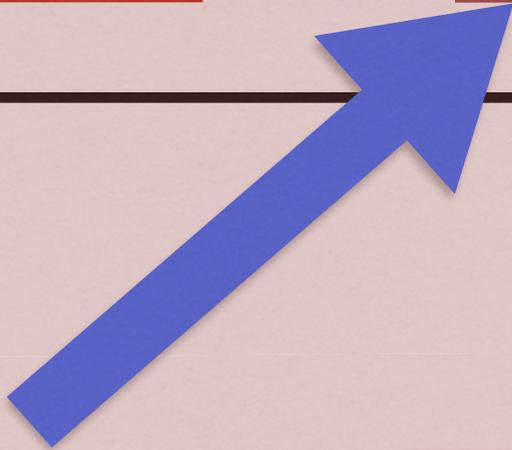


time →

Magnetic field →

CME/CMW →

Vorticity →
CVE/CVW →

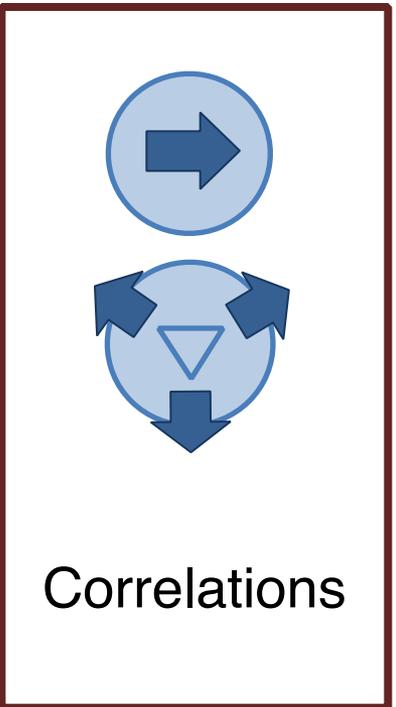
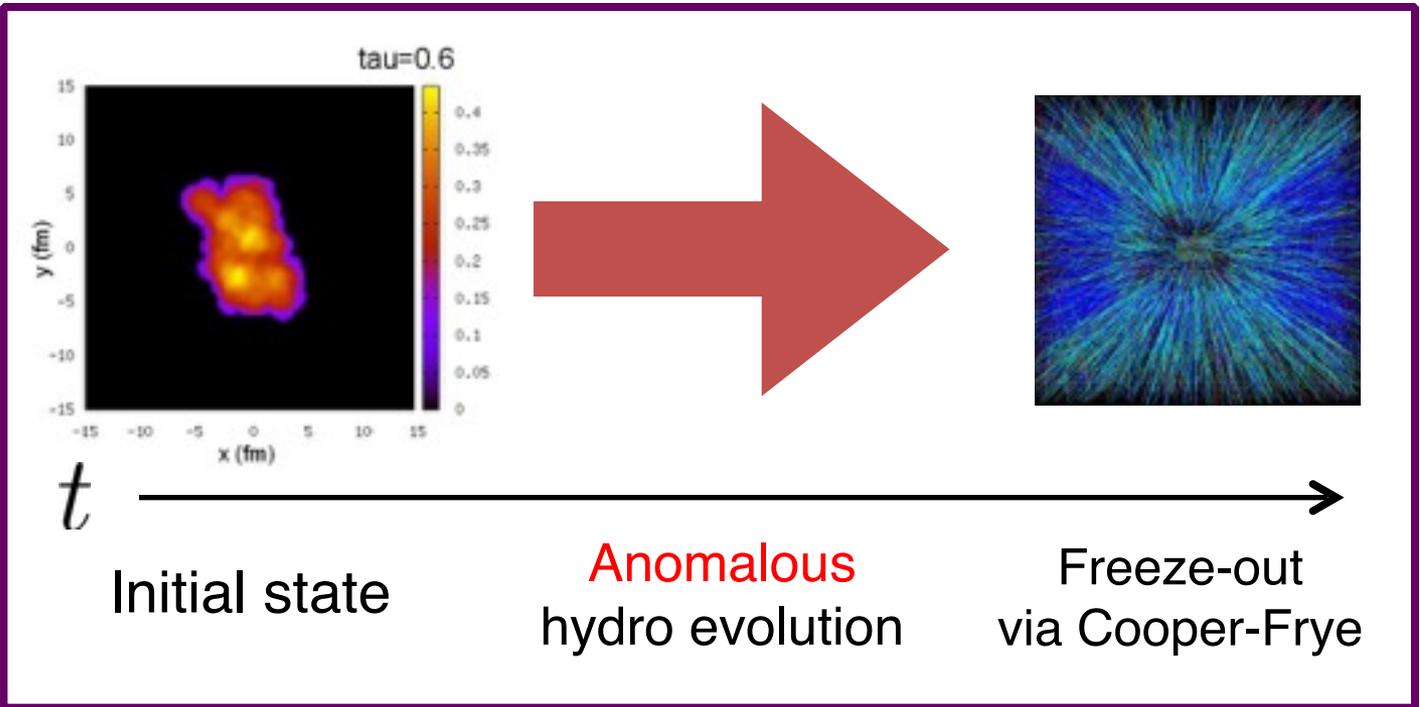


Dynamical Modeling

Event-by-event anomalous hydrodynamics for CME

[Hirono-Hirano-Khazeev 1412.0311]

Event-by-event anomalous hydrodynamic model



An initial cond.  Particle data for 1 event



Anomalous hydrodynamic equations

- Non-dissipative anomalous fluid in 3+1D
 - no viscosity/Ohmic conductivity
- Background electromagnetic fields

$$\partial_{\mu} T^{\mu\nu} = F^{\nu\rho} j_{\rho}$$

$$\partial_{\mu} j^{\mu} = 0 \quad \partial_{\mu} j_5^{\mu} = C E_{\mu} B^{\mu}$$

$$C = \frac{N_c N_f}{2\pi^2}$$

Anomalous hydrodynamic equations

- Constitutive equations

$$j^\mu = n u^\mu + \boxed{\kappa_B B^\mu} \quad j_5^\mu = n_5 u^\mu + \boxed{\xi_B B^\mu}$$

CME CSE

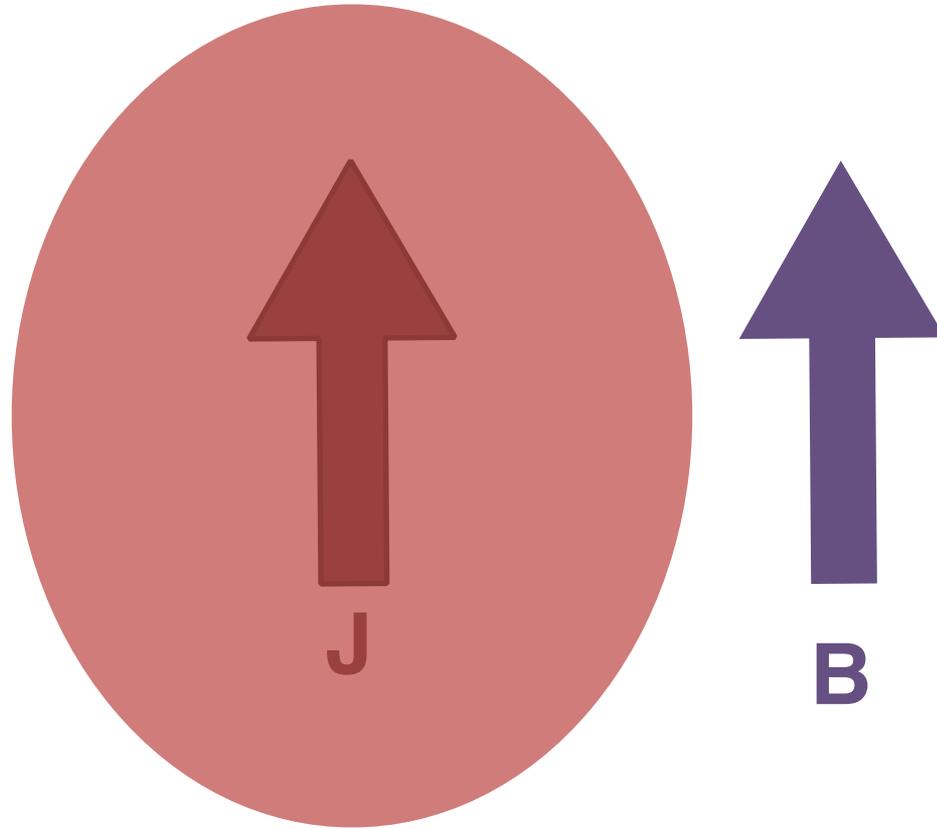
$$e\kappa_B = C\mu_5 \left(1 - \frac{\mu n}{e+p}\right) \quad e\xi_B = C\mu \left(1 - \frac{\mu_5 n_5}{e+p}\right)$$

[Son-Surowka 2009]
[Kalaydzhyan-Kirsch 2011]

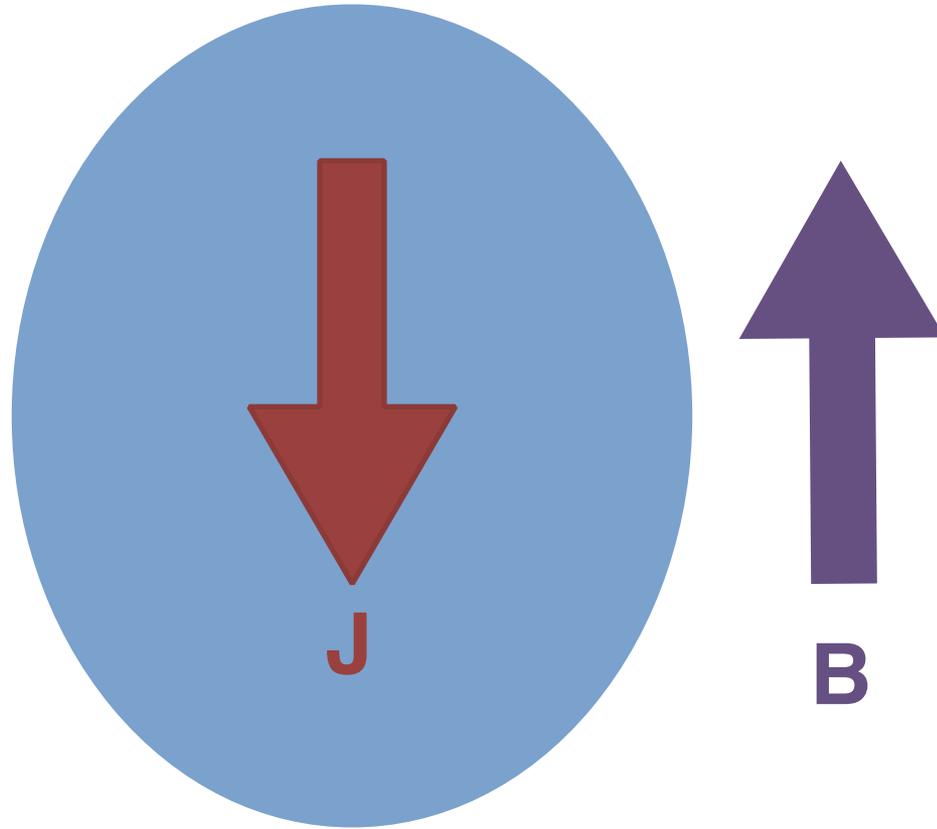
- Equation of state - conformal
 - massless quarks & gluons

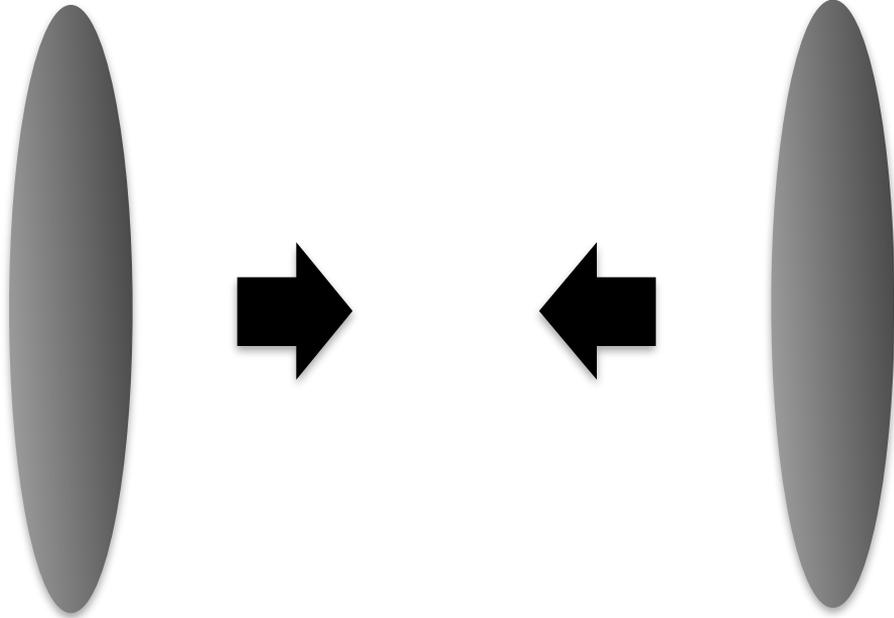
$$p(T, \mu, \mu_5) = \frac{g_{\text{qgp}} \pi^2}{90} T^4 + \frac{N_c N_f}{6} (\mu^2 + \mu_5^2) T^2 + \frac{N_c N_f}{12\pi^2} (\mu^4 + 6\mu^2 \mu_5^2 + \mu_5^4)$$

Initial axial charge?

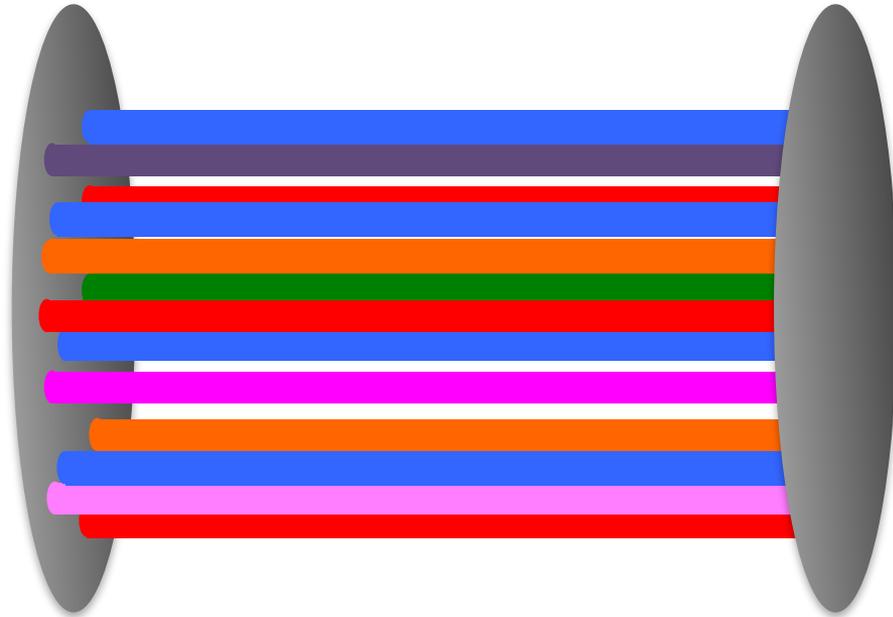


Initial axial charge?

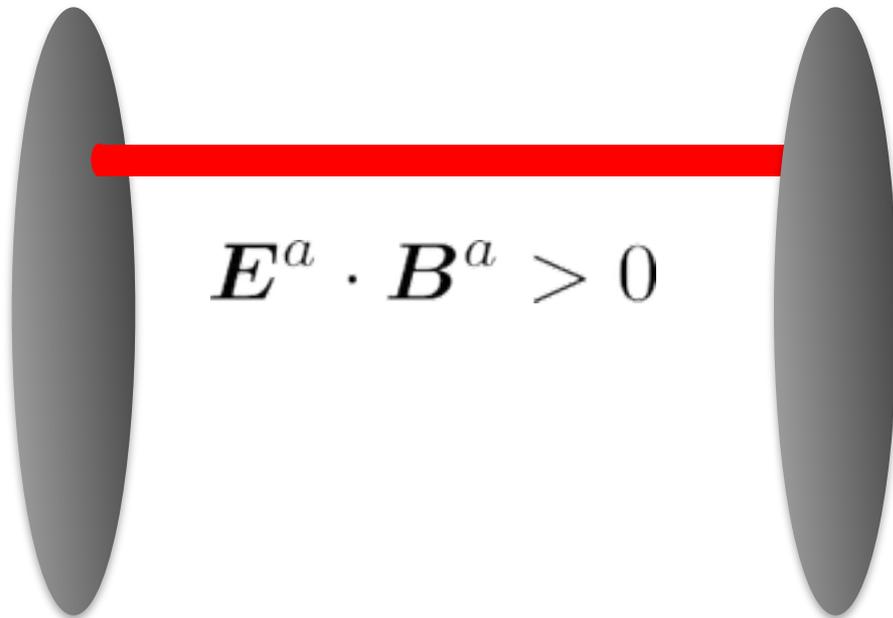




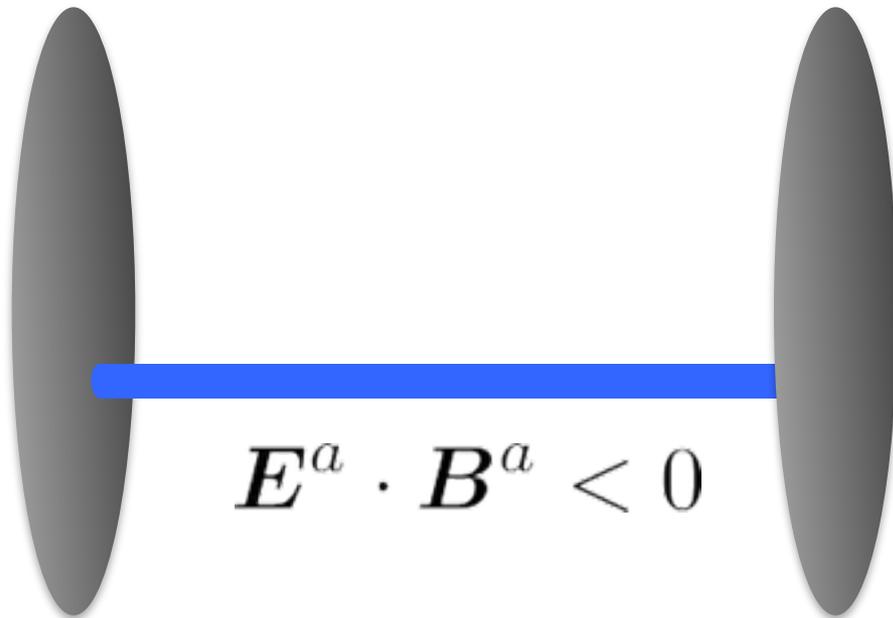
$$\partial_\mu j_5^\mu = \frac{g^2}{16\pi^2} \mathbf{E}^a \cdot \mathbf{B}^a$$



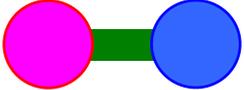
$$\partial_\mu j_5^\mu = \frac{g^2}{16\pi^2} \mathbf{E}^a \cdot \mathbf{B}^a$$

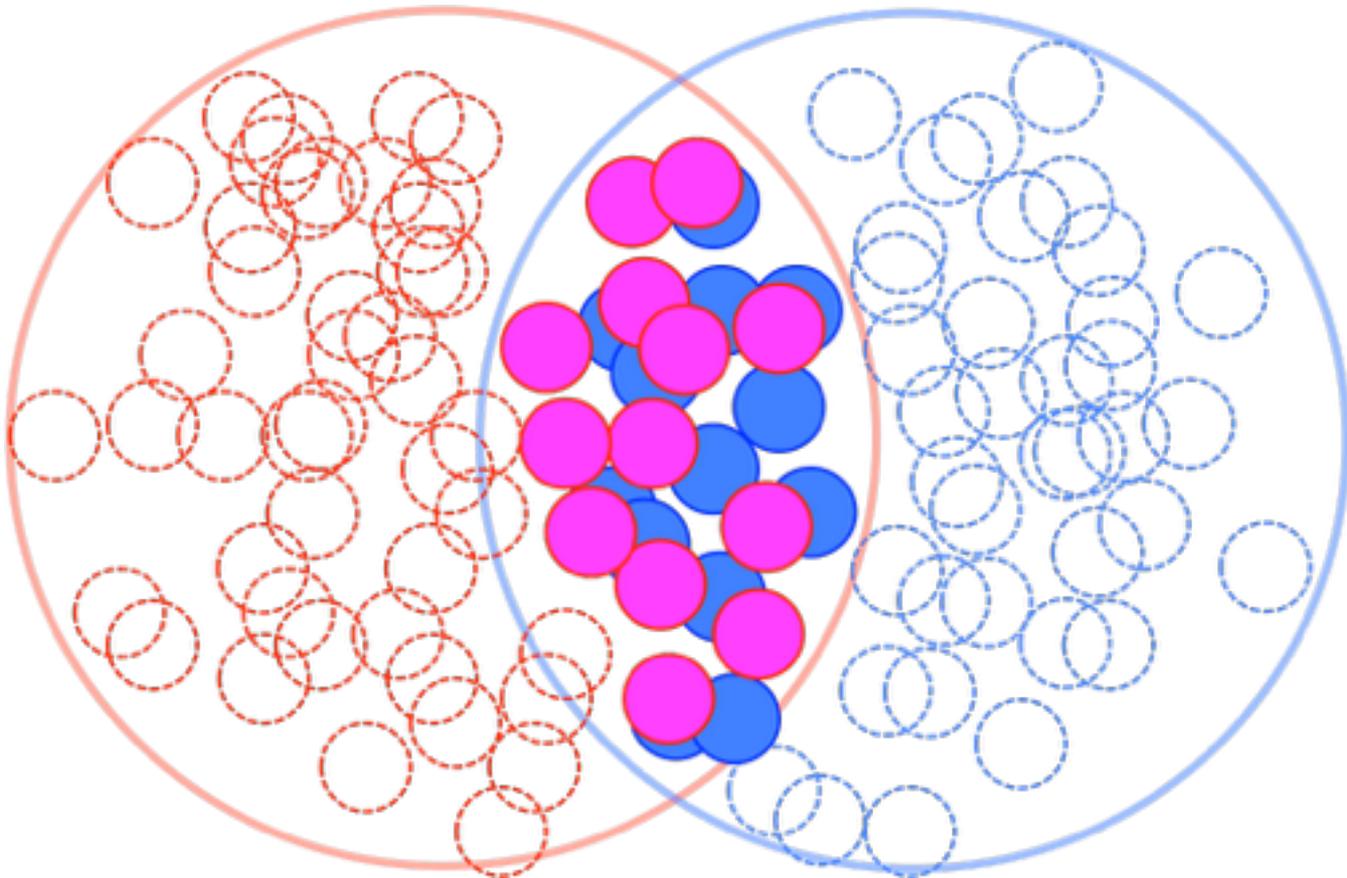


$$\partial_\mu j_5^\mu = \frac{g^2}{16\pi^2} \mathbf{E}^a \cdot \mathbf{B}^a$$



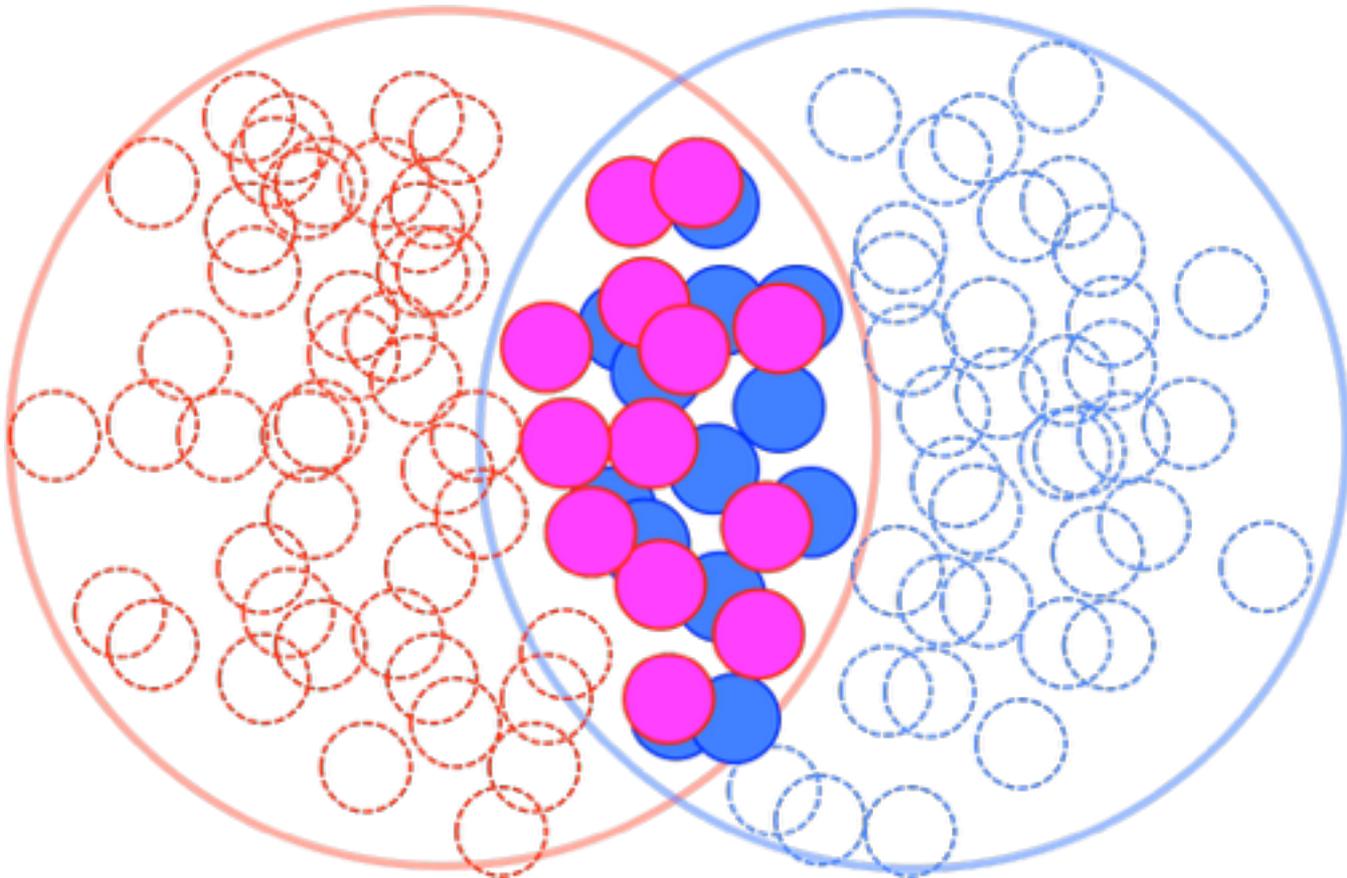
Axial charges from color flux tubes

$N_{\text{part}}^{A(B)}(\mathbf{x}_T)$: # of  () $N_{\text{coll}}(\mathbf{x}_T)$: # of pairs 



Axial charges from color flux tubes

$$s(\mathbf{x}_T, \eta_s) = Af(\eta_s) \left[\frac{1 - \alpha}{2} \frac{d^2 N_{\text{part}}}{dx_T^2} + \alpha \frac{d^2 N_{\text{coll}}}{dx_T^2} \right]$$

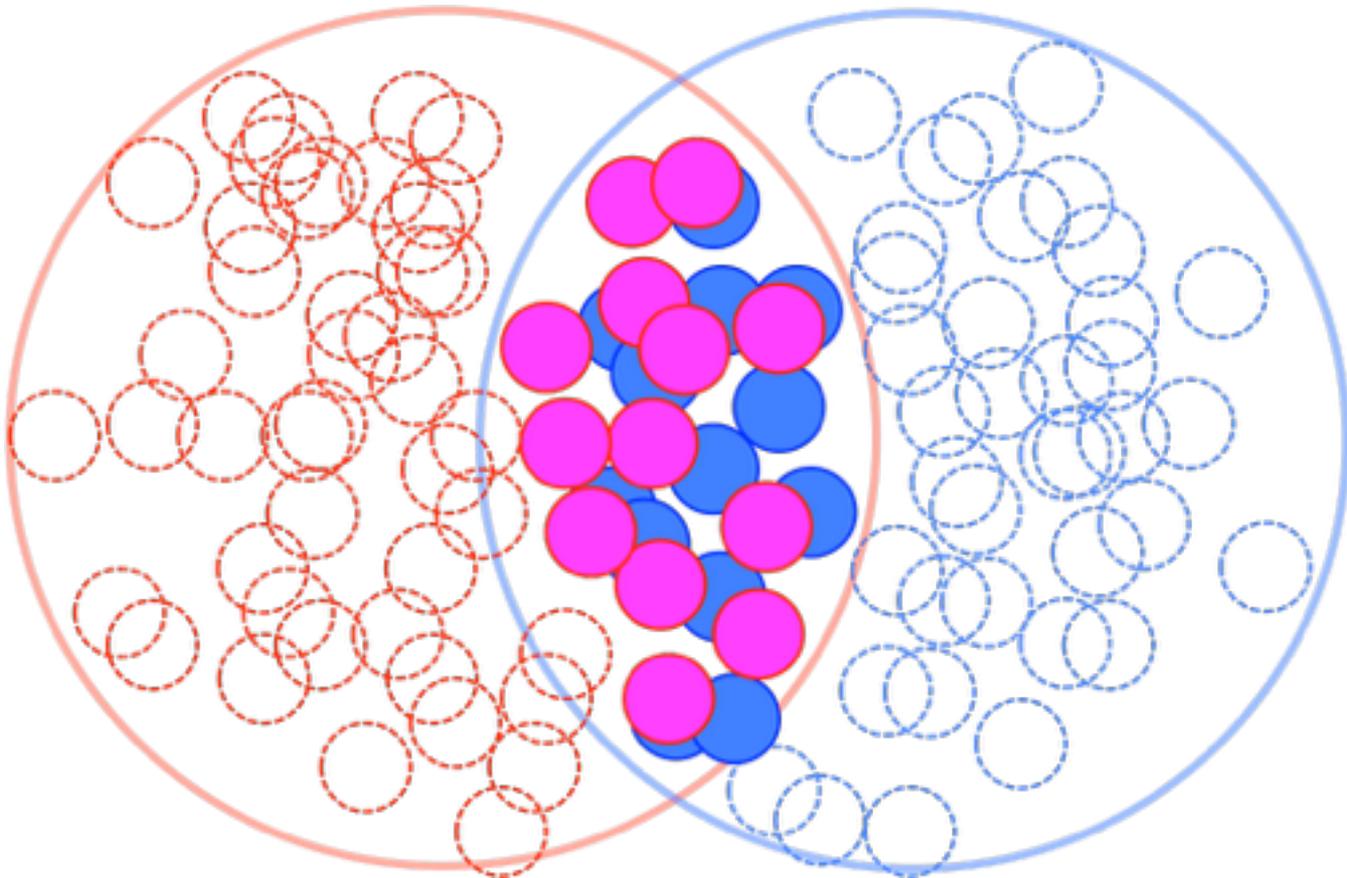


Axial charges from color flux tubes

$$X_j \in \{+1, -1\}$$

Sign of $\mathbf{E}^a \cdot \mathbf{B}^a$

$$\mu_5(\mathbf{x}_T) = C_{\mu_5} \sum_{j=1}^{N_{\text{coll}}(\mathbf{x}_T)} X_j$$

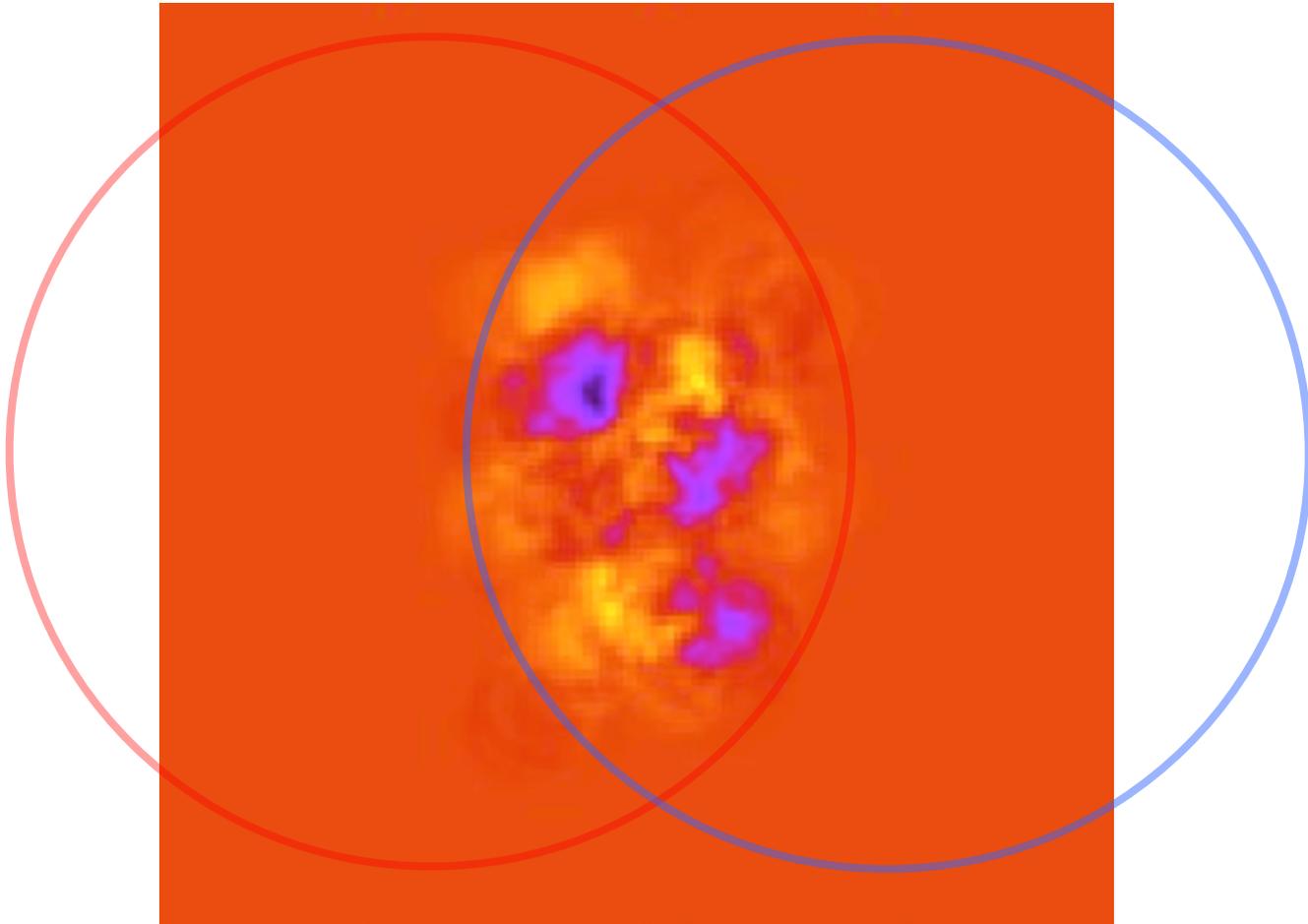


Axial charges from color flux tubes

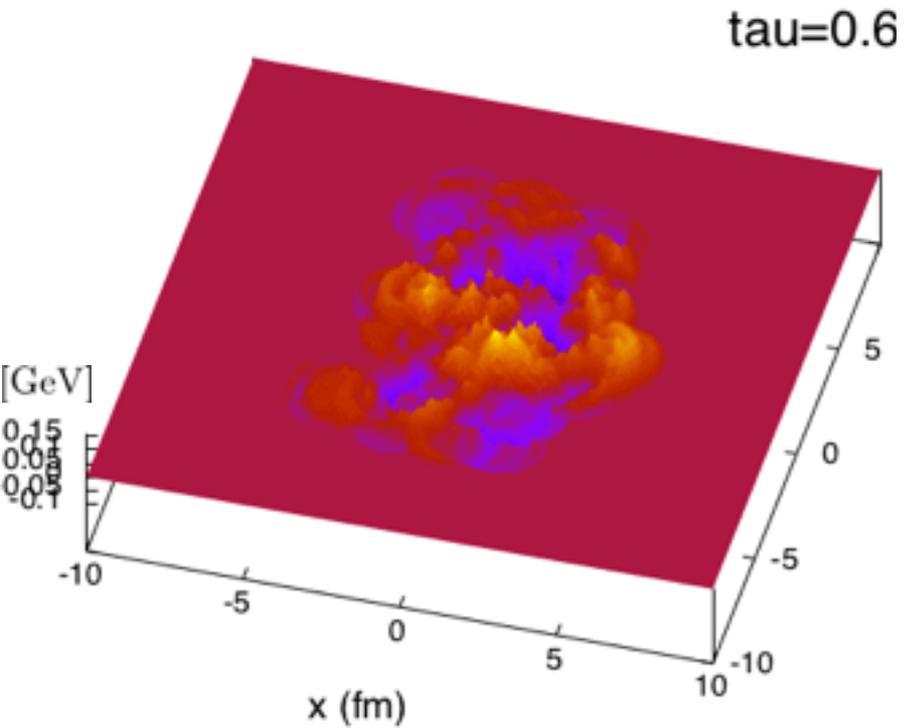
$$X_j \in \{+1, -1\}$$

Sign of $\mathbf{E}^a \cdot \mathbf{B}^a$

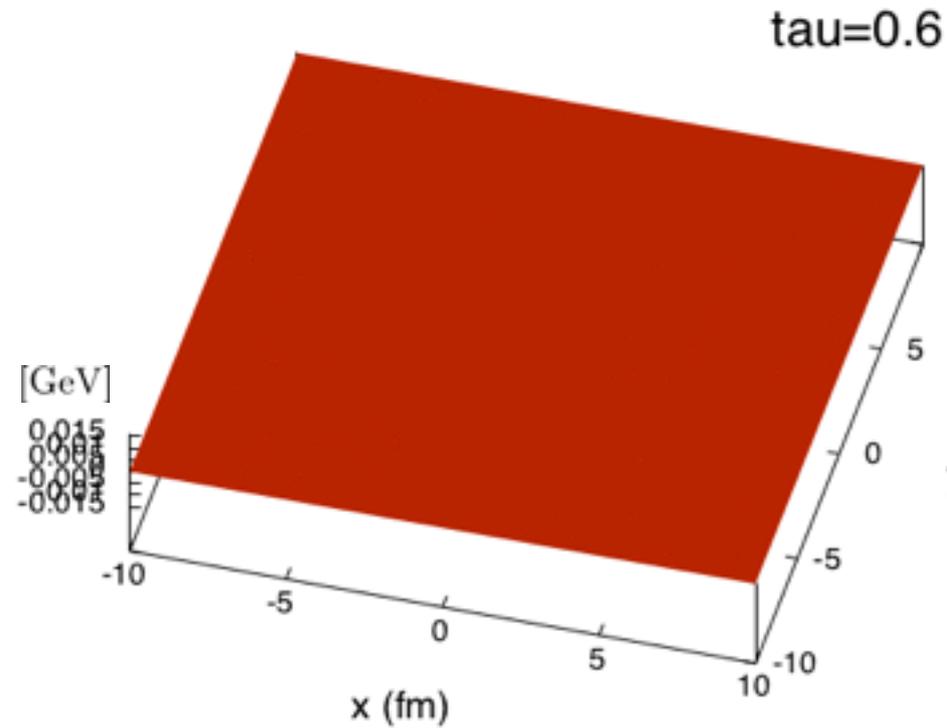
$$\mu_5(\mathbf{x}_T) = C_{\mu_5} \sum_{j=1}^{N_{\text{coll}}(\mathbf{x}_T)} X_j$$



Time evolution of chemical potentials

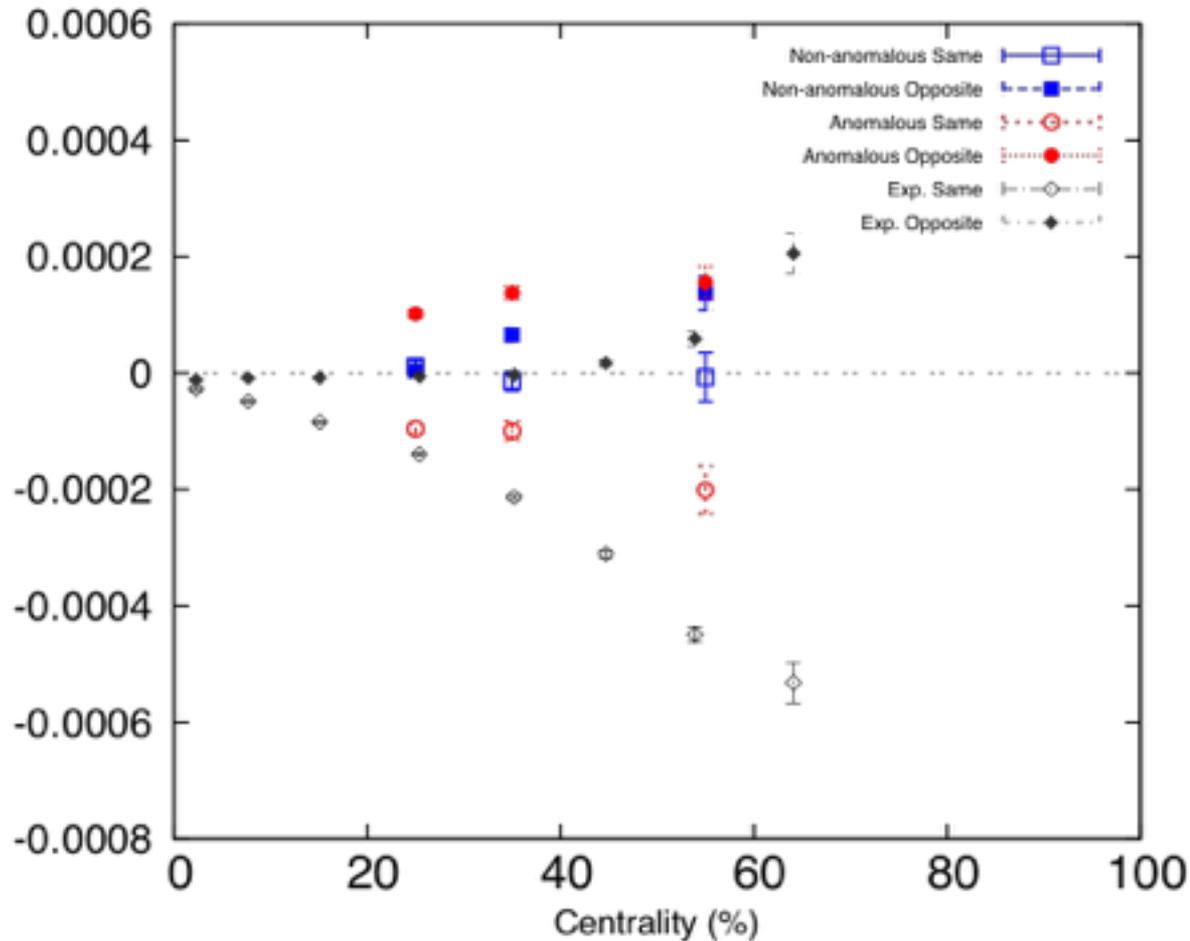


μ_5



μ

$$\gamma_{\alpha\beta} = \langle \cos(\phi_i + \phi_j - 2\psi_{\text{RP}}) \rangle_{\alpha\beta}$$



**Gamma correlations are indeed sensitive to CME/CSE
Values are similar to experimental ones**

Hydrodynamics with chiral anomaly and charge separation in relativistic heavy ion collisions

Yi Yin^{1,*} and Jinfeng Liao^{2,3,†}

[1504.06906]

- Charge transport from CME/CSE is solved on top of the solution of 2+1D viscous hydro (VISH)
- Quantify the effects of transverse momentum conservation (TMC)
- Same-charge correlation is calculated

Hydrodynamics with chiral anomaly and charge separation in relativistic heavy ion collisions

Yi Yin^{1,*} and Jinfeng Liao^{2,3,†}

[1504.06906]

$$\gamma_{\alpha\beta} = \langle \cos(\phi_i + \phi_j - 2\psi_{RP}) \rangle_{\alpha\beta} \quad \delta_{\alpha\beta} = \langle \cos(\phi_i - \phi_j) \rangle_{\alpha\beta}$$

$$\gamma_{\alpha,\beta}^{\text{data}} \simeq \boxed{\gamma_{\alpha,\beta}^{\text{CME}}} + \gamma_{\alpha,\beta}^{\text{TMC}}, \quad \delta_{\alpha,\beta}^{\text{data}} \simeq \boxed{\delta_{\alpha,\beta}^{\text{CME}}} + \delta_{\alpha,\beta}^{\text{TMC}}$$

$$(\gamma + \delta)^{\text{data}} \simeq (\gamma + \delta)^{\text{TMC}}$$

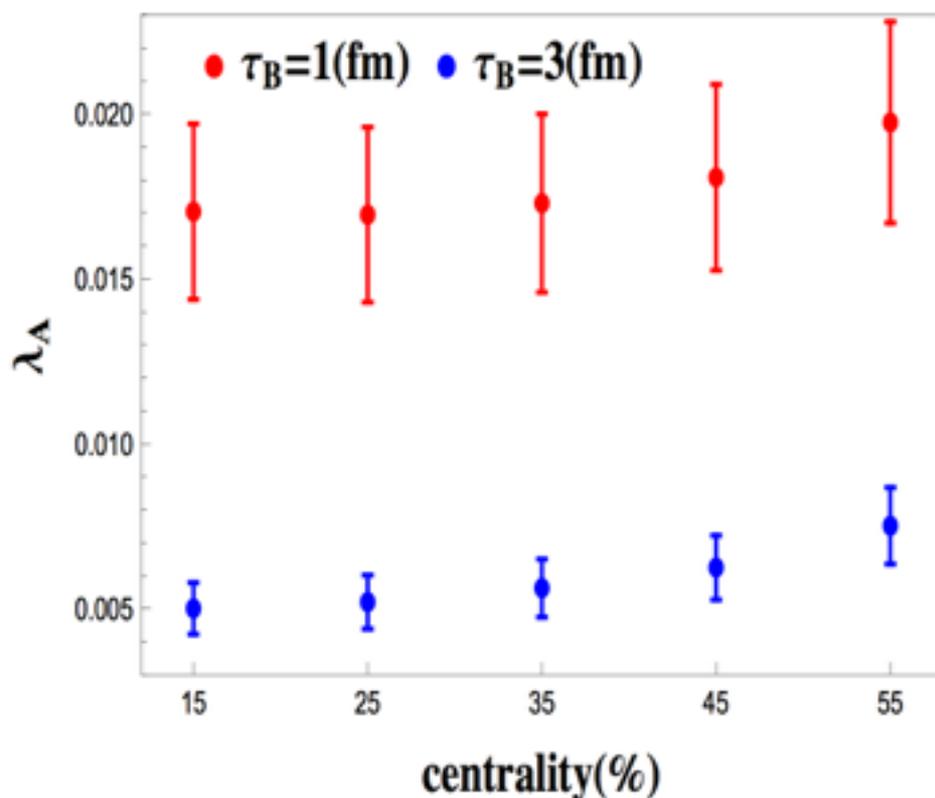
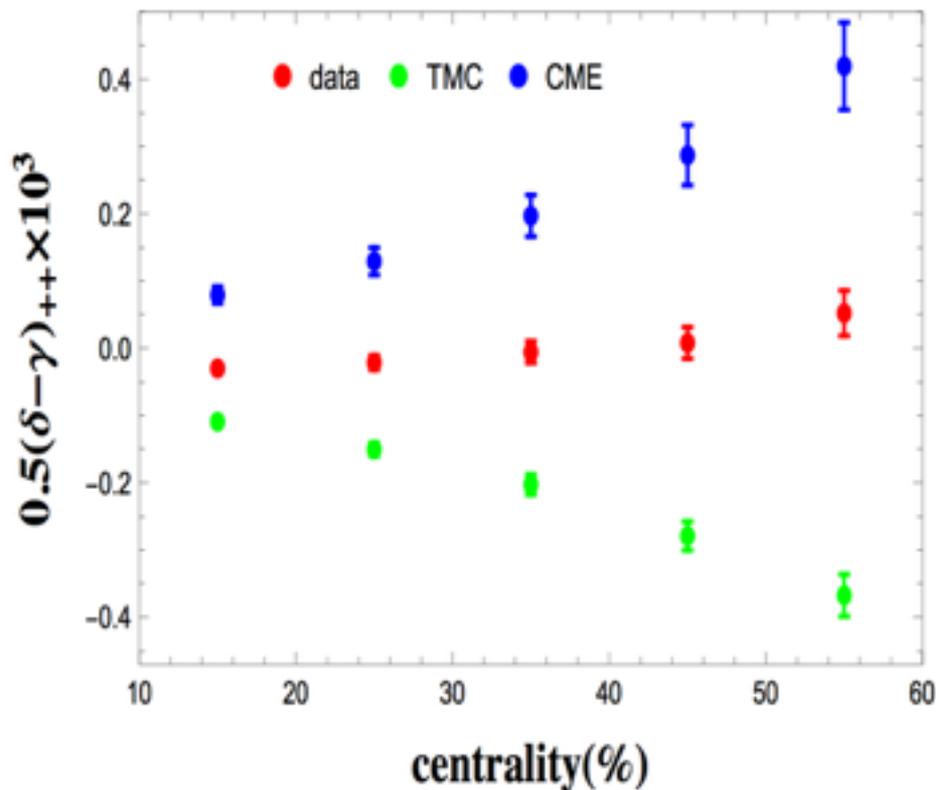
$$\delta_{\alpha\beta}^{\text{TMC}} \pm \gamma_{\alpha\beta}^{\text{TMC}} = \frac{[\langle p_{\perp} \rangle_{\alpha}(1 \pm \bar{v}_{2,\alpha})][\langle p_{\perp} \rangle_{\beta}(1 \pm \bar{v}_{2,\beta})]}{N_{\text{TMC}} \langle p_{\perp}^2 \rangle (1 \pm \bar{v}_2)}$$

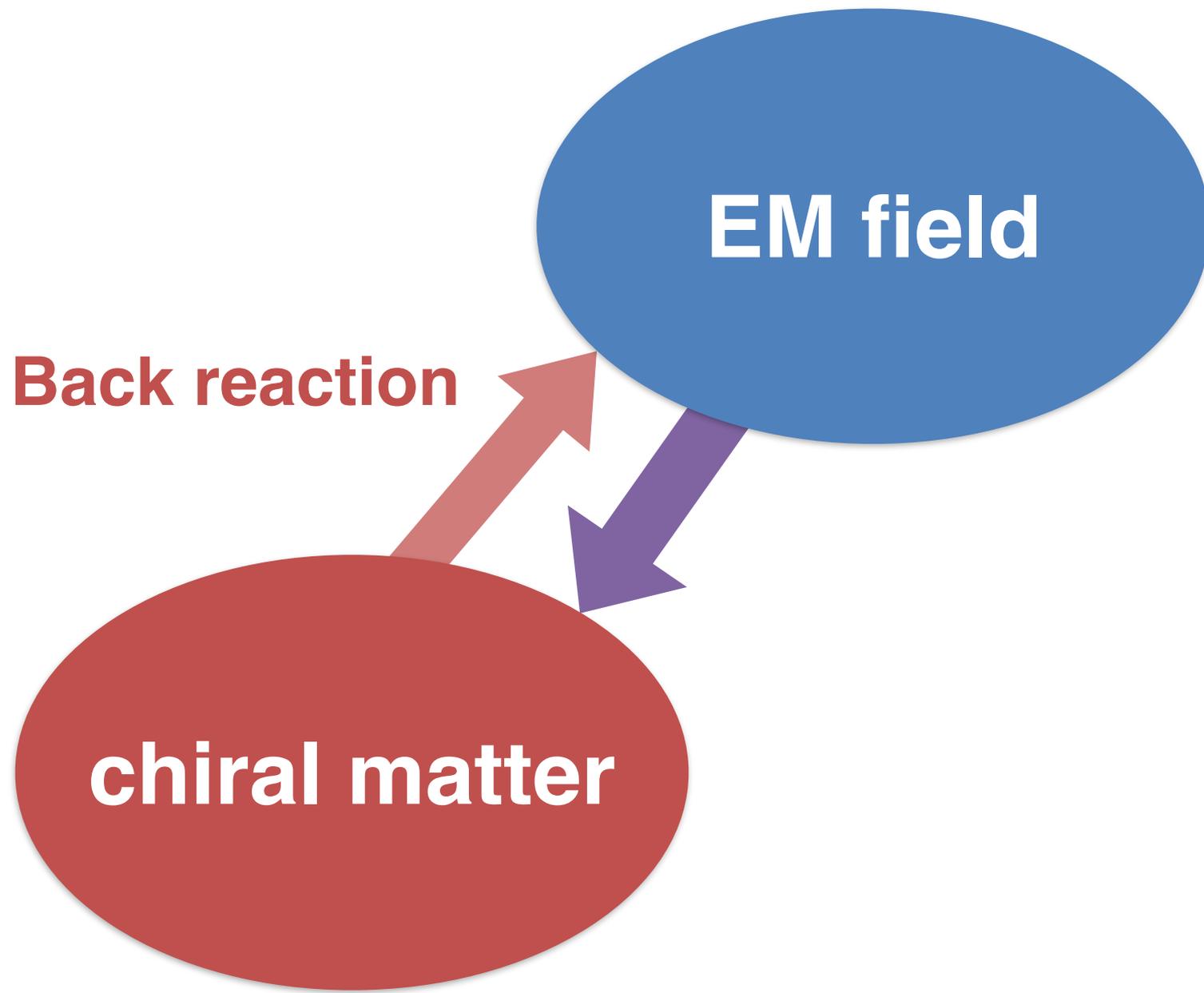
[Bzdak-Koch-Liao 1008.4919]

Hydrodynamics with chiral anomaly and charge separation in relativistic heavy ion collisions

Yi Yin^{1,*} and Jinfeng Liao^{2,3,†}

[1504.06906]





Magnetic & fermionic helicities

$$\partial_\mu \dot{j}_A^\mu = C_A \mathbf{E} \cdot \mathbf{B}$$



$$\frac{d}{dt} [\mathcal{H} + \mathcal{H}_F] = 0$$

$$\mathcal{H} = \int d^3x \mathbf{A} \cdot \mathbf{B}$$

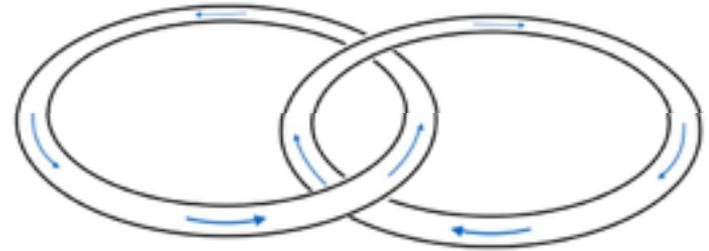
Magnetic helicity

$$\mathcal{H}_F = \frac{2}{C_A} \int d^3x n_A$$

Fermionic helicity

Magnetic helicity knows topology

$$\mathcal{H} = \int d^3x \mathbf{A} \cdot \mathbf{B}$$

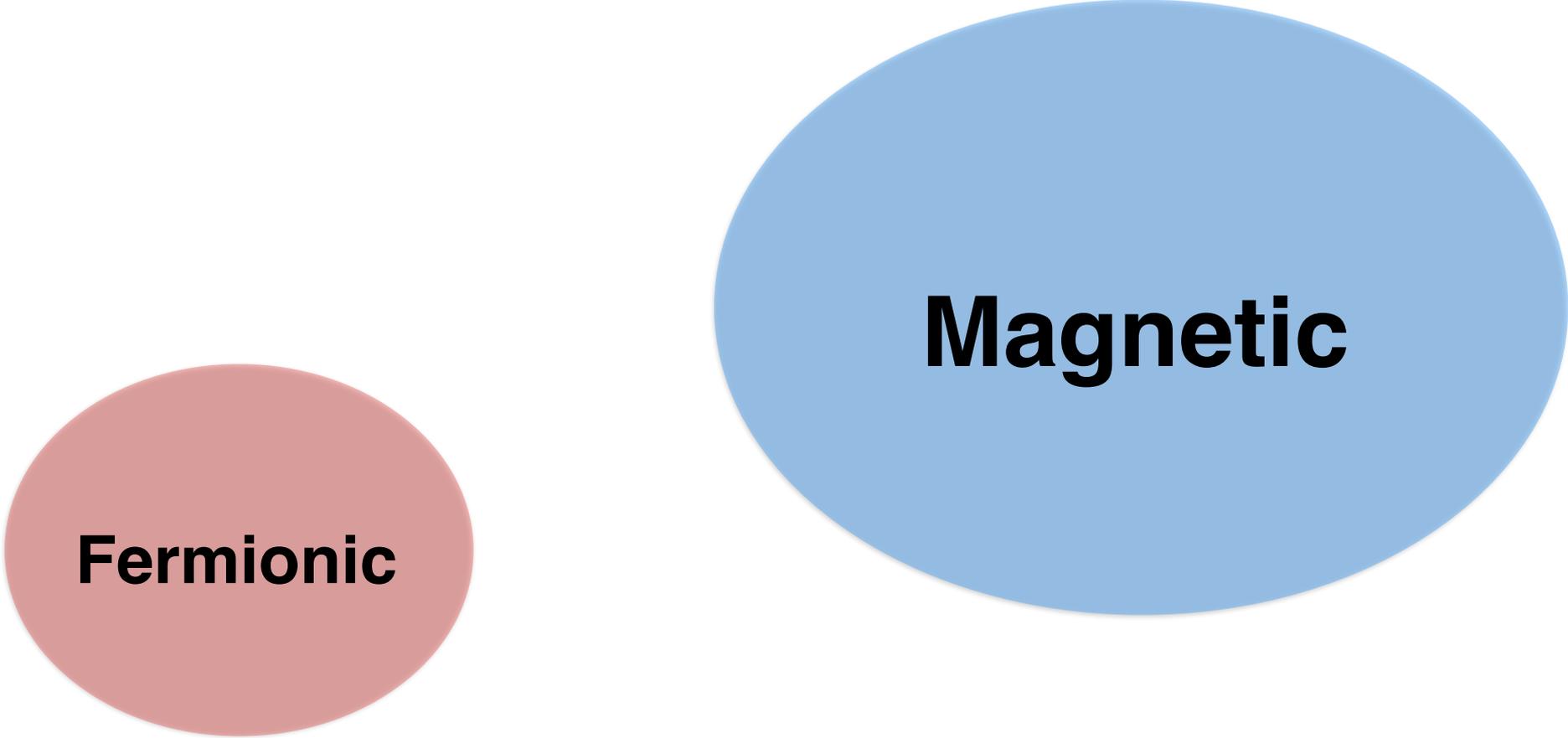


$$= \sum_i \mathcal{S}_i \varphi_i^2 + 2 \sum_{i,j} \mathcal{L}_{ij} \varphi_i \varphi_j$$

Self-linking number

Linking number

Helicity transfer

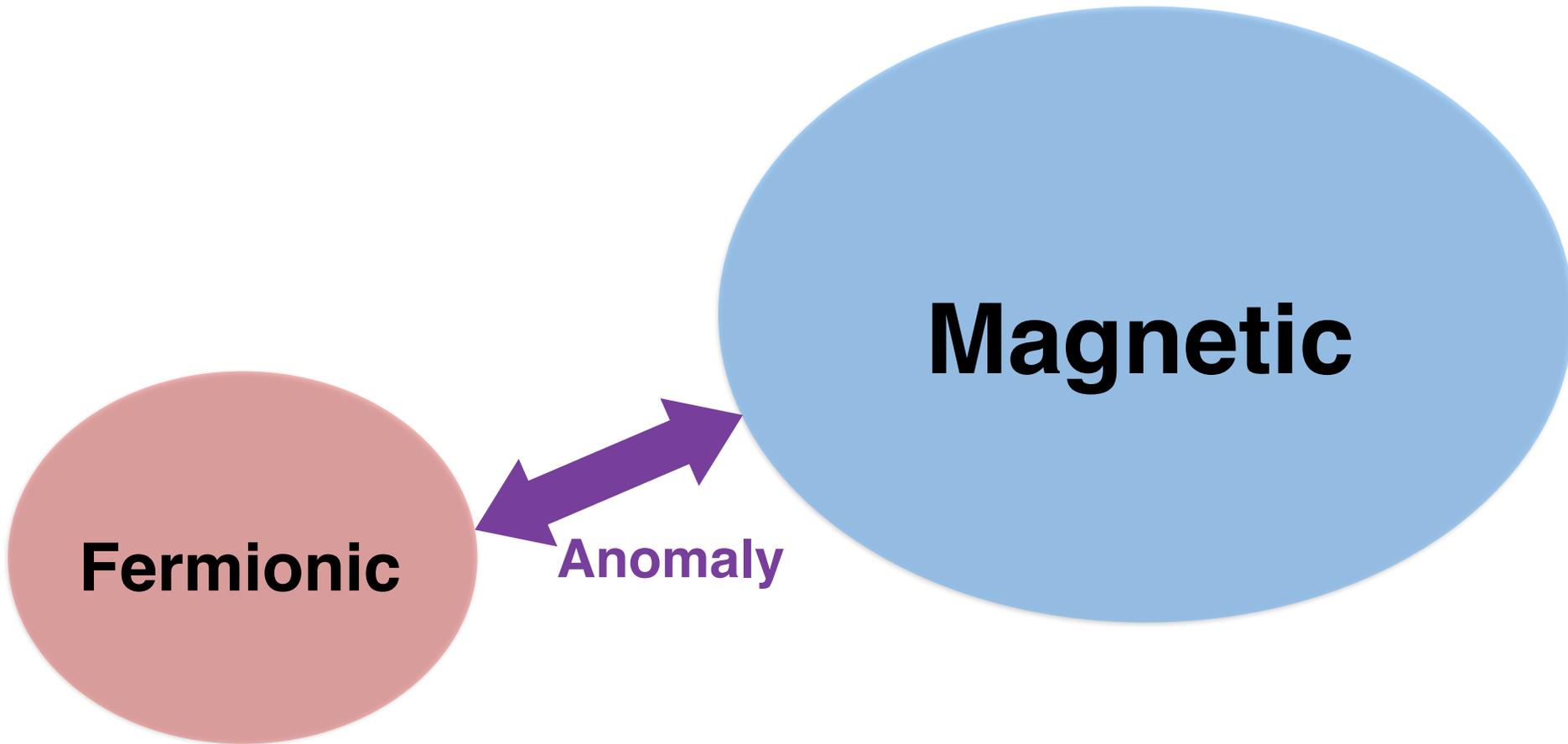


Fermionic

Magnetic

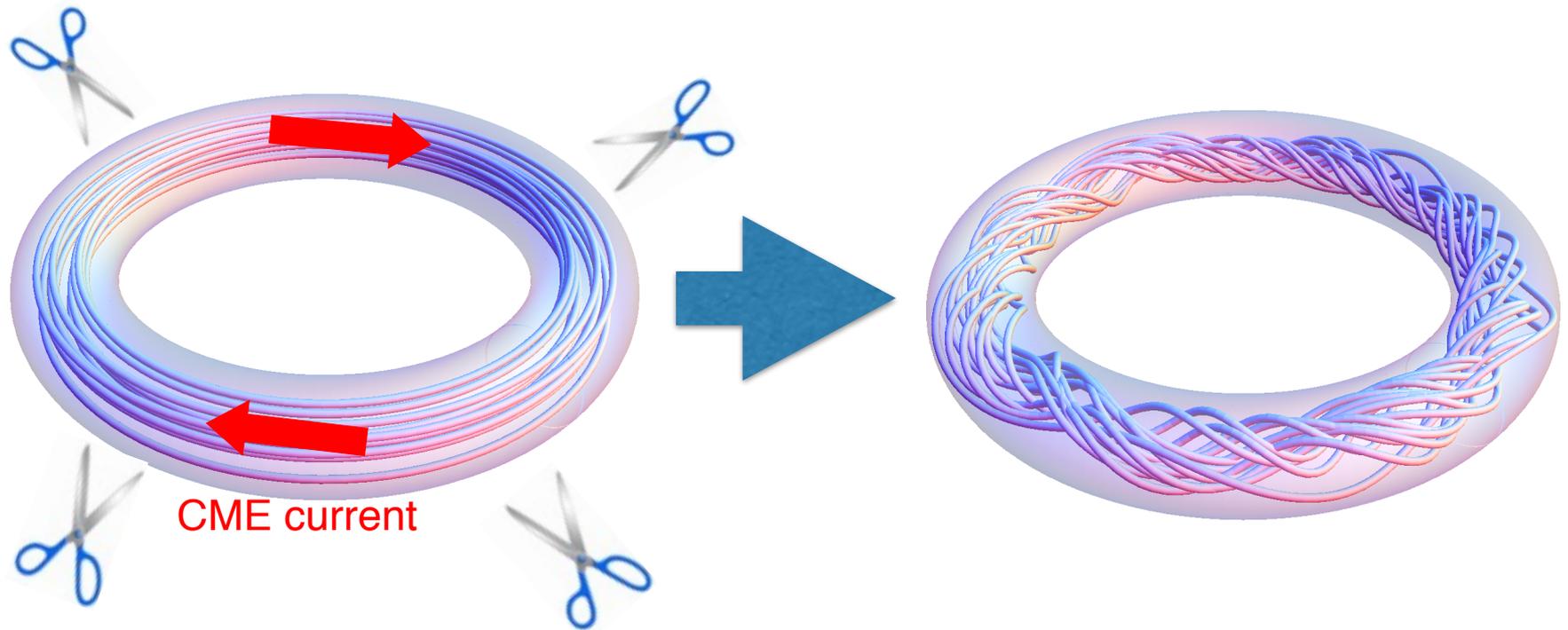
Without anomaly: two helicities are separately conserved

Helicity transfer



With anomaly: two helicities can be transferred
sum is conserved

Anomaly changes the topology of magnetic fields



What is the fate of the magnetic field?

Self-similar inverse cascade of magnetic fields

[Hirono-Khazeev-Yin 1509.07790]

Maxwell-Chern-Simons equations

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} + \mathbf{j}_{\text{EM}} \quad \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \cdot \mathbf{B} = \nabla \cdot \mathbf{E} = 0$$

$$\mathbf{j}_{\text{EM}} = \sigma \mathbf{E} + \boxed{\sigma_A \mathbf{B}}$$

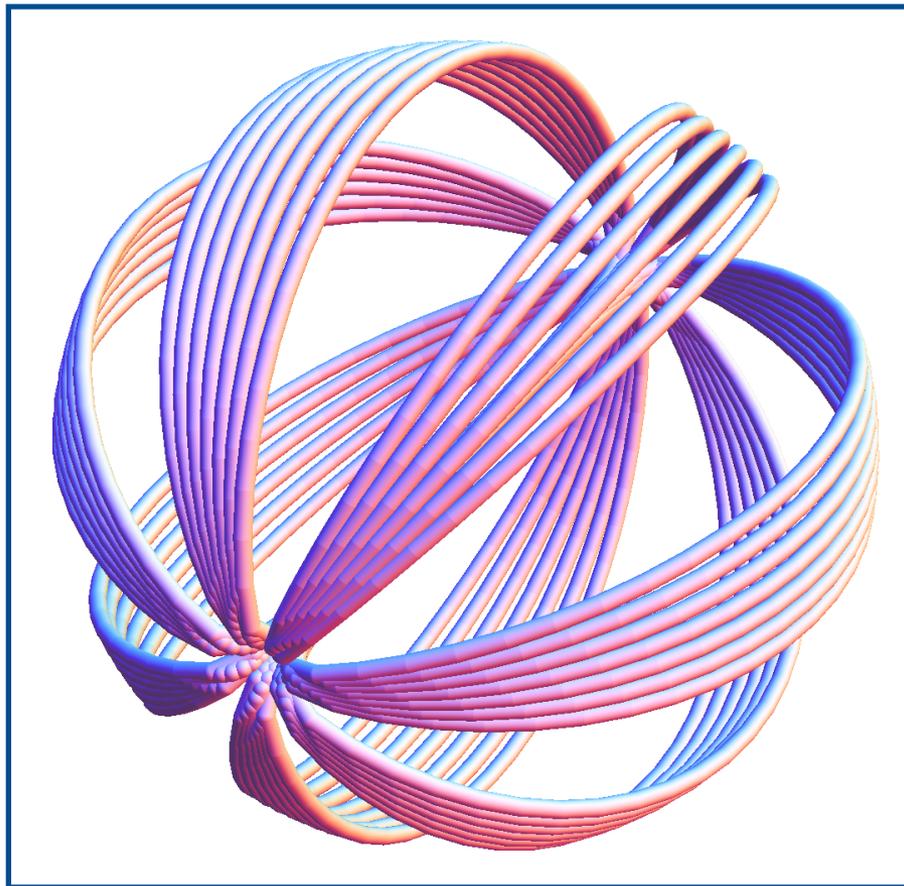
$$\sigma_A(t) = C_A \mu_A(t) = \frac{C_A n_A(t)}{\chi}$$

Axial charge is taken to be spatially constant

Chandrasekhar-Kendall states

$$\nabla \times \mathbf{W}_{lm}^{\pm}(\mathbf{x}; k) = \pm k \mathbf{W}_{lm}^{\pm}(\mathbf{x}; k)$$

Sign of helicity



Decomposition of the magnetic fields

$$\mathbf{B}(\mathbf{x}, t) = \sum_{l,m} \int_0^\infty \frac{dk}{\pi} k^2 [\alpha_{lm}^+(k, t) \mathbf{W}_{lm}^+(\mathbf{x}; k) + \alpha_{lm}^-(k, t) \mathbf{W}_{lm}^-(\mathbf{x}; k)]$$

$$h_m(t) = \int_0^\infty \frac{dk}{\pi} k [g_+(k, t) - g_-(k, t)]$$

Helicity spectral function

$$\mathcal{E}_m(t) = \int_0^\infty \frac{dk}{\pi} k^2 [g_+(k, t) + g_-(k, t)]$$

$$g_\pm(k, t) \equiv \sum_{l,m} |\alpha_{lm}^\pm(k, t)|^2$$

Energy cost to store a unit helicity in $\mathbf{W}^\pm(\mathbf{x}, k)$ is k

EOM for Fourier coeff.

$$\partial_t \alpha_{lm}^{\pm}(k, t) = \sigma^{-1} [-k^2 \pm \sigma_A(t)k] \alpha_{lm}^{\pm}(k, t)$$

If σ_A is positive,

- Negative helicity modes decays exponentially
- Positive helicity modes

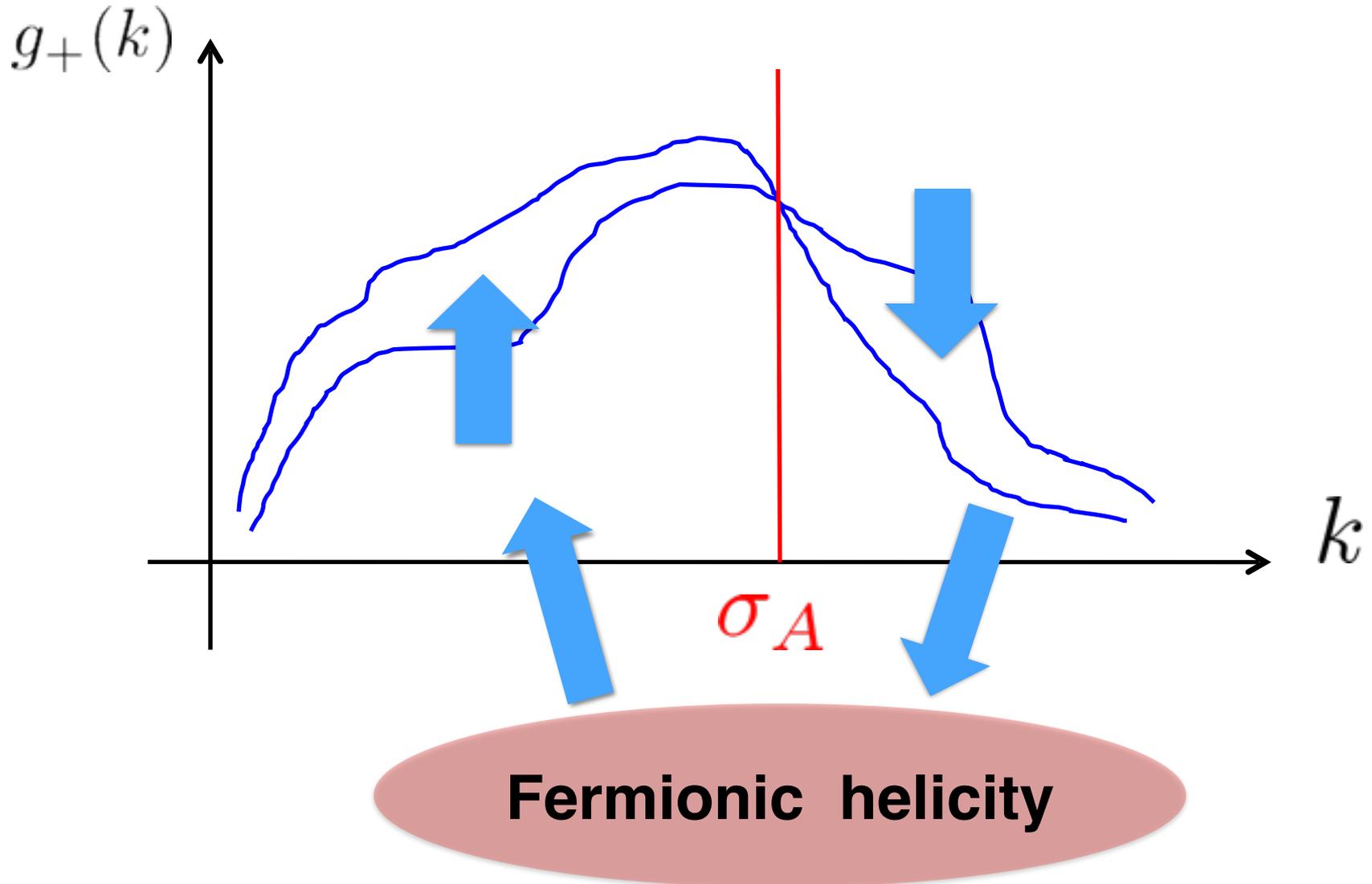
$$k > \sigma_A \quad \text{decay}$$

$$k < \sigma_A \quad \text{grow (chiral plasma instability)}$$

$\sigma_A(t)$: energy cost per unit fermionic helicity

$$h_F = \frac{1}{C_A} \int d^3x n_A \quad \sigma_A = C_A \mu_A = \frac{\mu_A}{1/C_A}$$

Magnetic helicity spectrum



We will see the fates of

$h_m(t)$: magnetic helicity

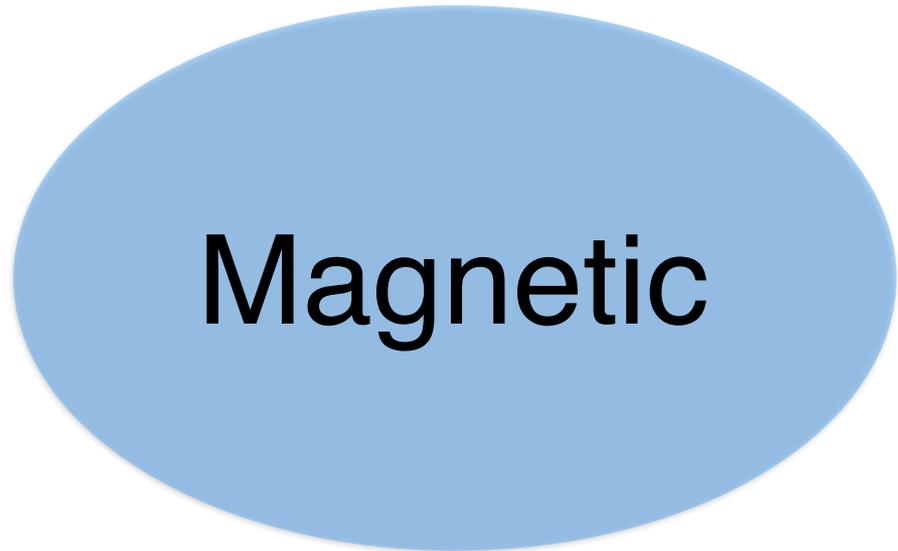
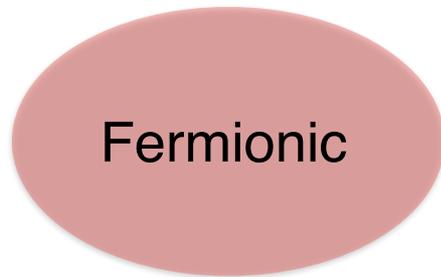
$\sigma_A(t)$: (prop. to) fermionic helicity

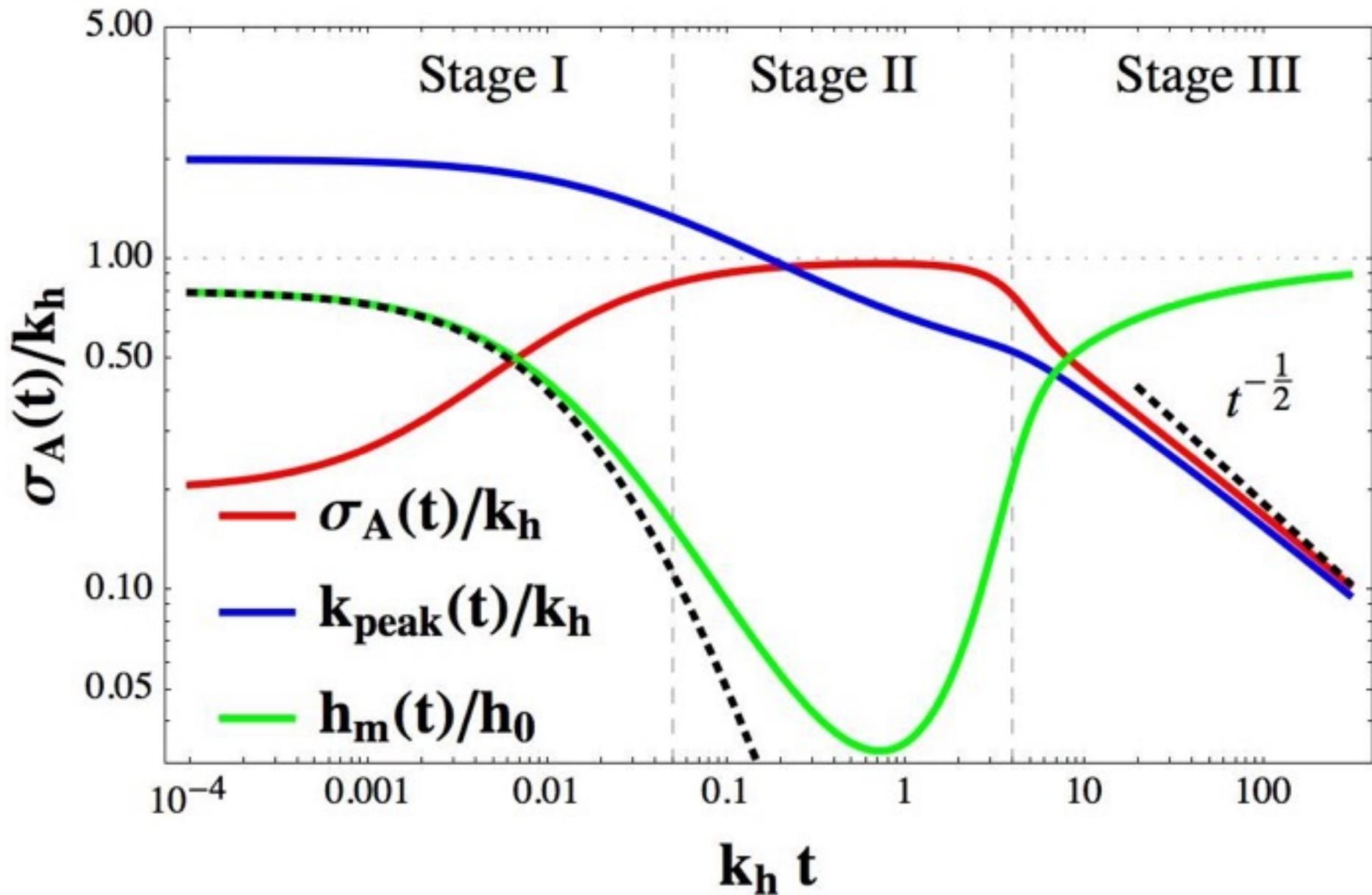
$k_{\text{peak}}(t)$: peak momentum of magnetic helicity spectrum

$$h_m(t) = \int_0^\infty \frac{dk}{\pi} k [g_+(k, t) - g_-(k, t)]$$

Initial helicity is dominated by magnetic helicity

$$h_m(t = 0) = 0.8h_0$$

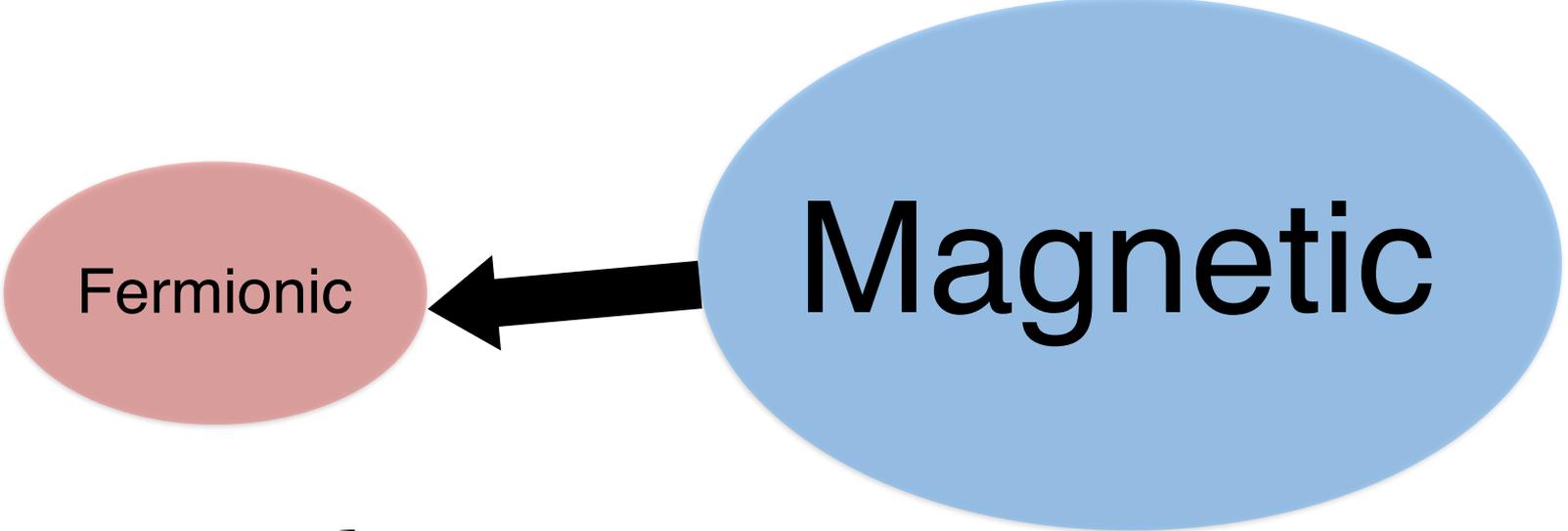




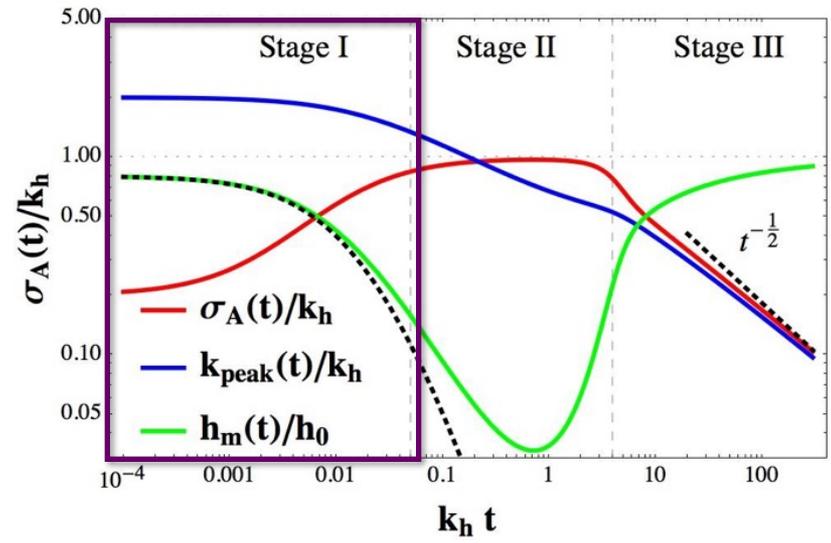
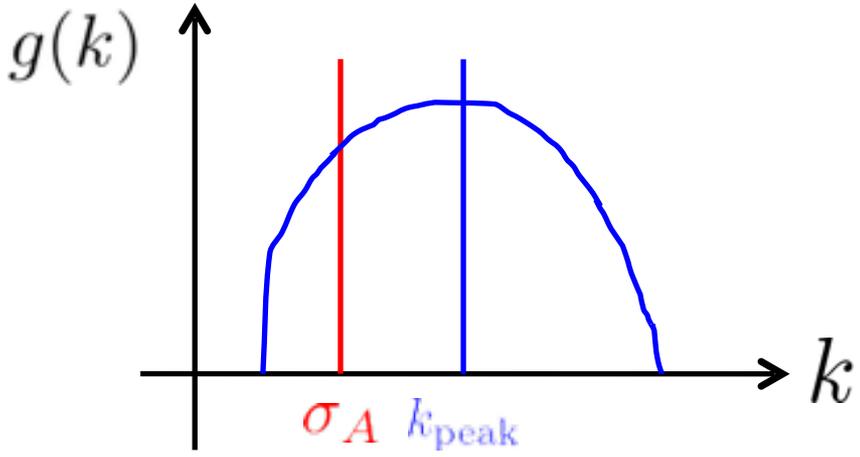
$$k_h \equiv \frac{C_A^2 h_0}{\chi V}$$

energy cost per unit fermionic helicity,
when all the helicity is stored in fermions

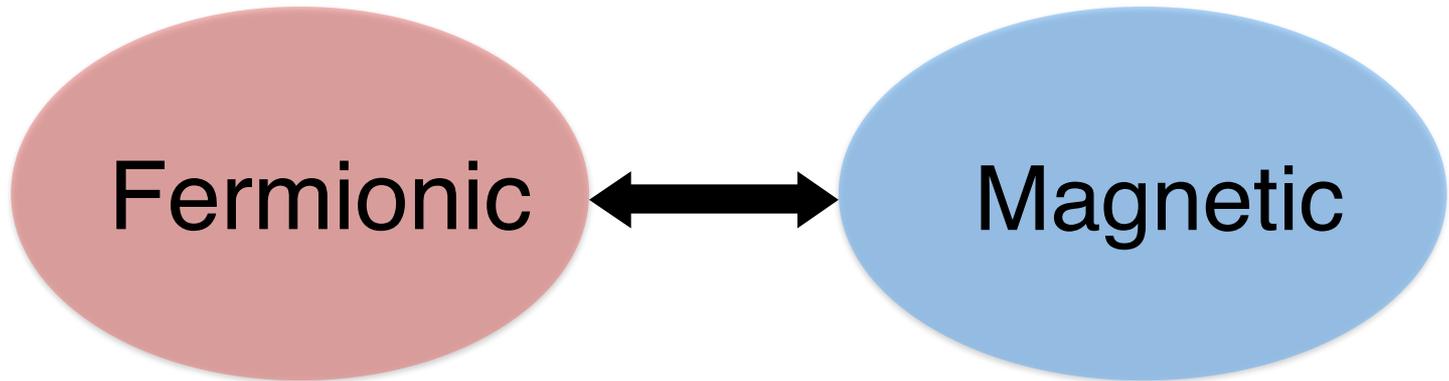
Stage I: helicity transferred to fermions



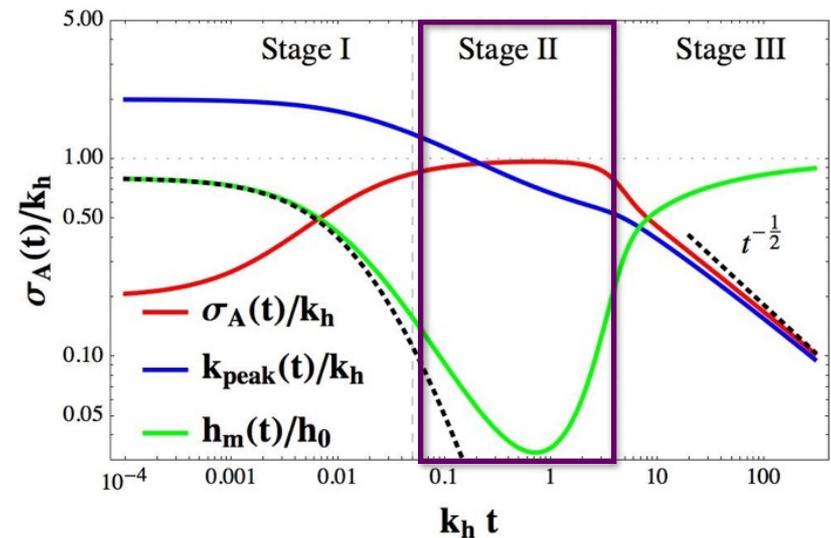
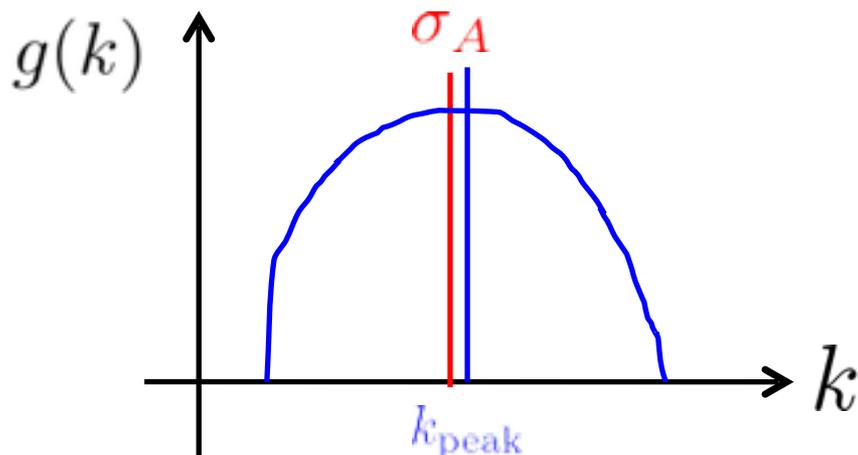
$$\sigma_A < k_{\text{peak}}$$



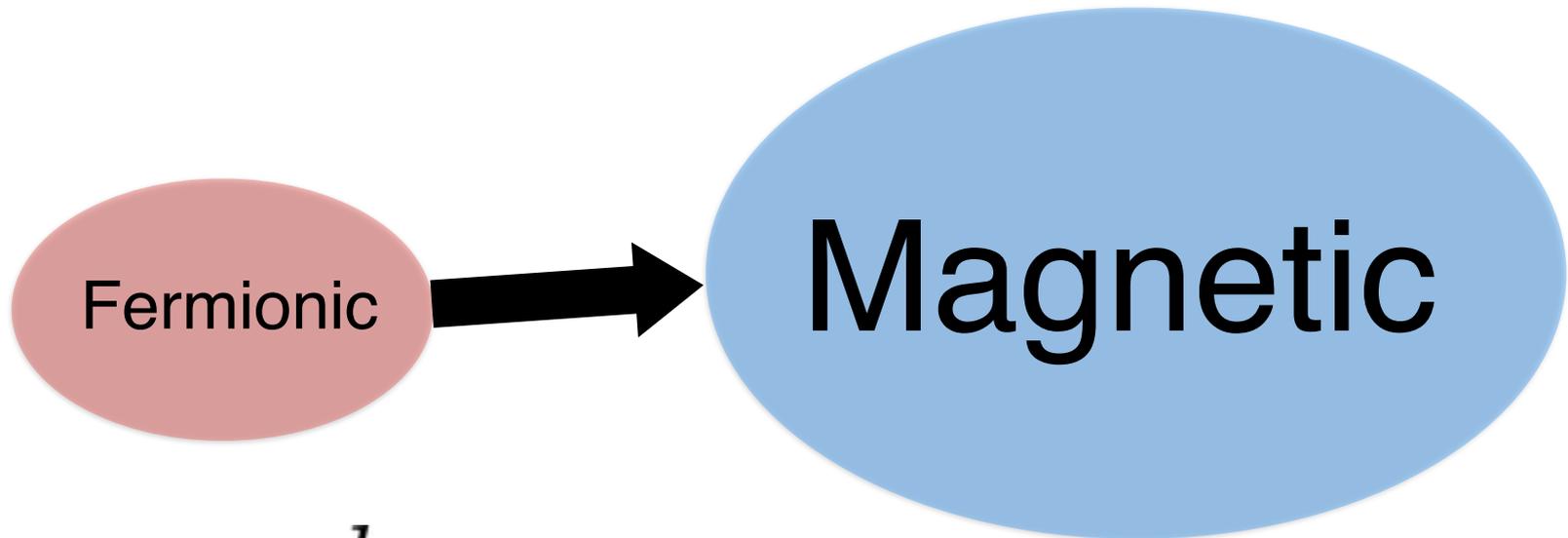
Stage II: intermediate stage



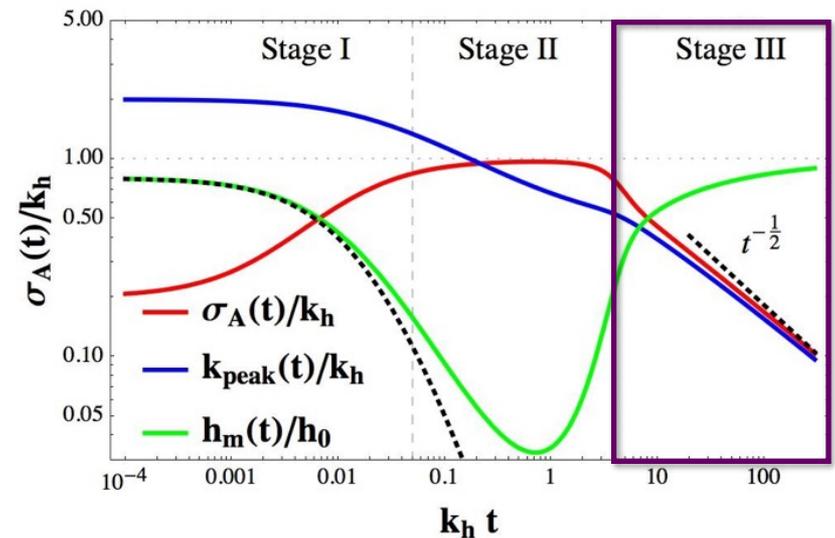
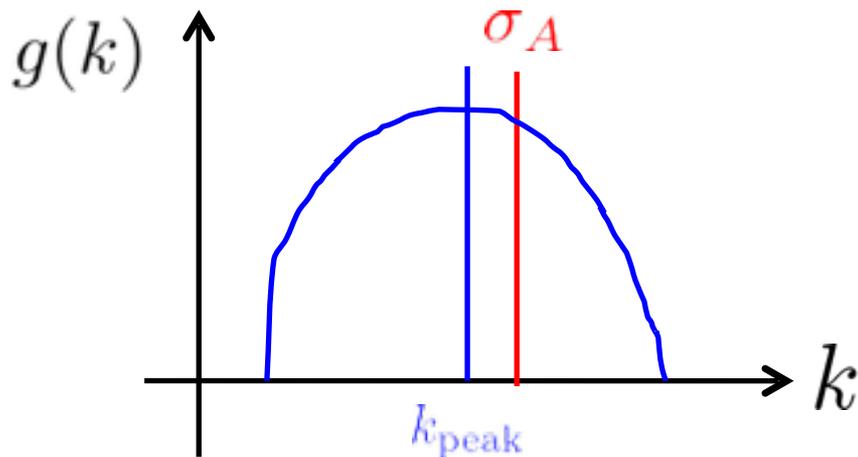
$$\sigma_A \approx k_{\text{peak}}$$



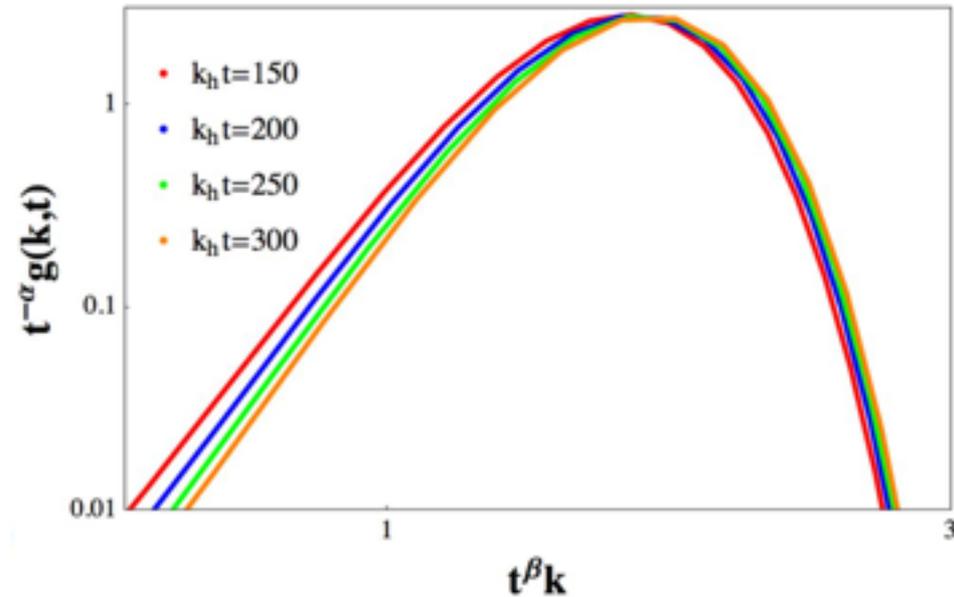
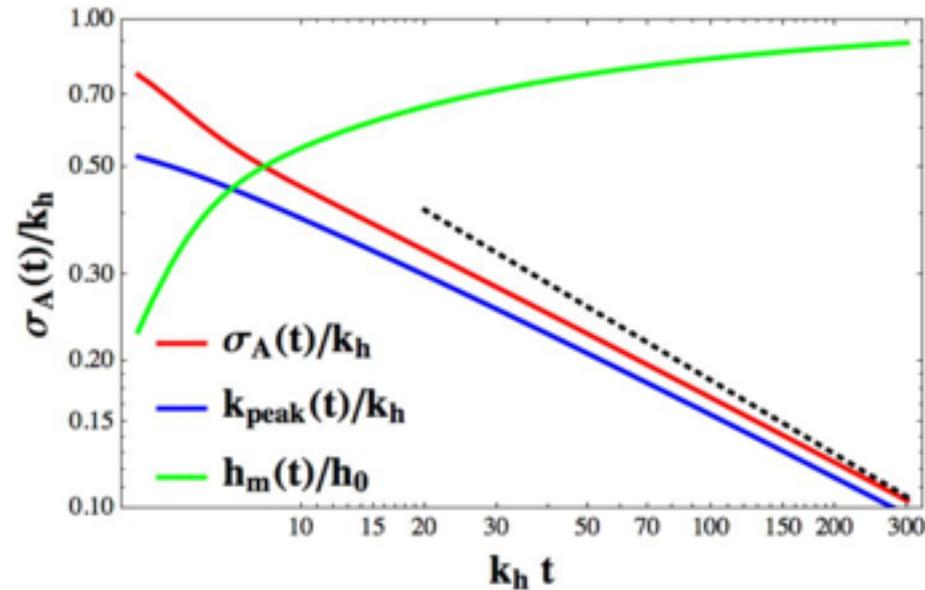
Stage III: self-similar evolution



$$\sigma_A > k_{\text{peak}}$$



Stage III: self-similar evolution



$$\sigma_A(t) \sim k_{\text{peak}}(t) \propto t^{-\beta}$$

$$g(k, t) \sim t^\alpha \tilde{g}(t^\beta k)$$

$$\alpha = 1, \quad \beta = 1/2$$

Stage III: self-similar evolution

$$g(k, t) \propto \exp \left[-2\sigma^{-1} (k - k_{\text{peak}})^2 t \right]$$
$$\rightarrow \delta(k - k_{\text{peak}}(t))$$

**The magnetic field approaches
a single CK state**

Self-similarity extended to chiral MHD

Scaling laws in chiral hydrodynamic turbulence

Naoki Yamamoto

[1603.08864]

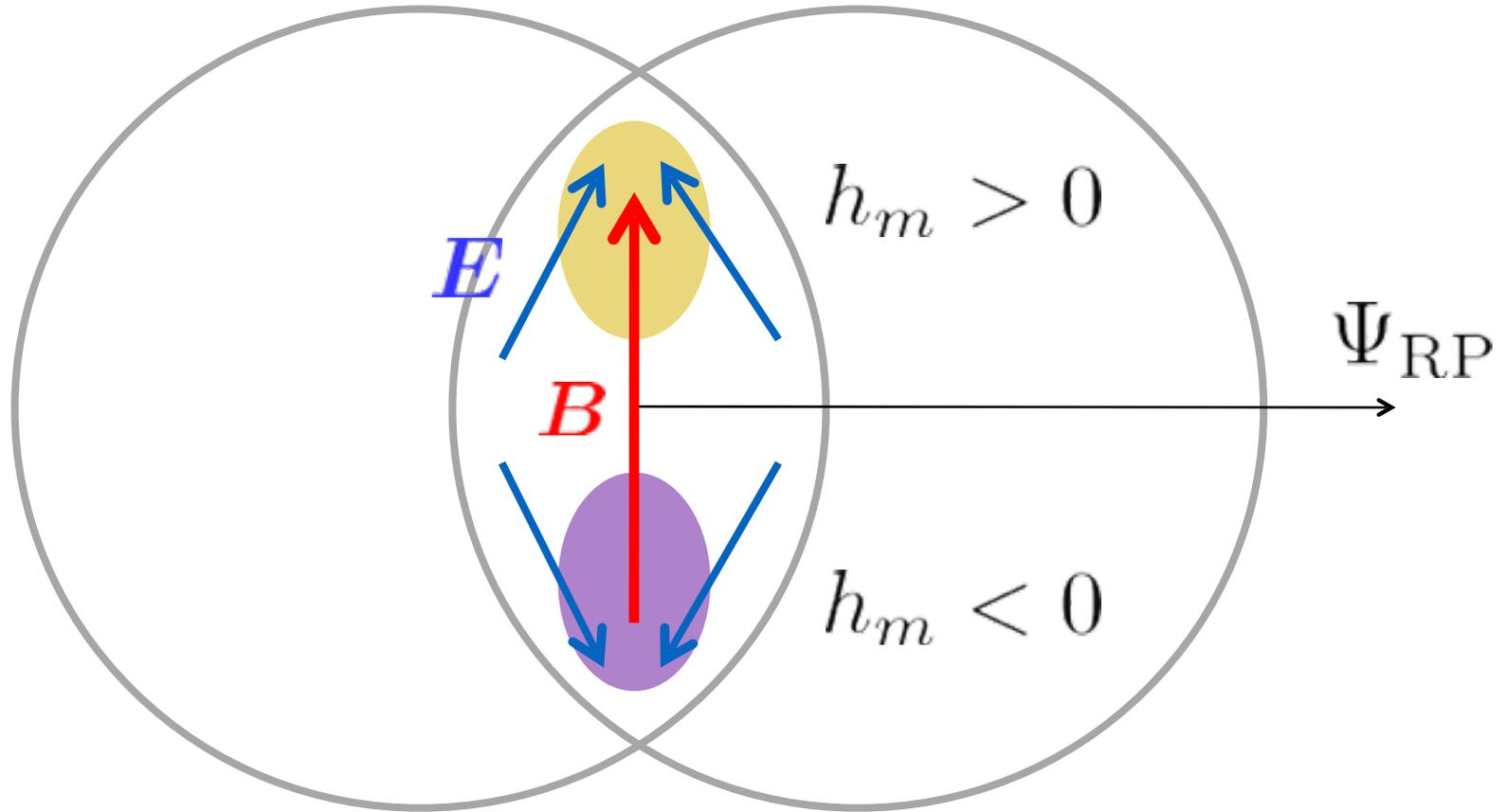
$$\bar{n}_5(t) = \bar{n}_5(t_s) \left(\frac{t_s}{t} \right)^{\frac{1}{2}}$$

$$\xi_B(t) = \xi_B(t_s) \left(\frac{t}{t_s} \right)^{\frac{1}{2}} \quad \text{magnetic correlation length}$$

$$\xi_v(t) = \xi_v(t_s) \left(\frac{t}{t_s} \right)^{\frac{1}{2}} \quad \text{kinetic correlation length}$$

**Realized in
heavy-ion collisions?**

Helical magnetic field in HIC

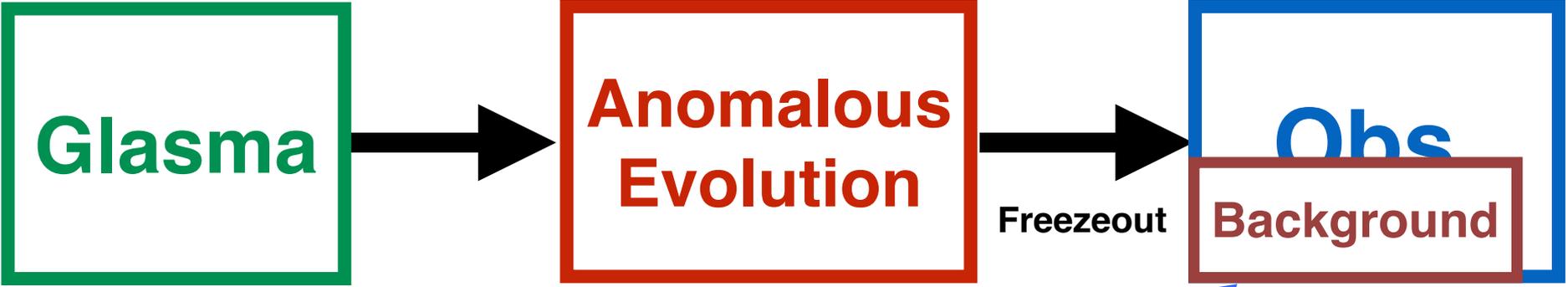


Condition for the inverse cascade to occur

$$eB = c_B m_\pi^2 \quad L \approx 10 \text{ [fm]}$$

$$c_B > 100 \quad \text{RHIC}$$

$$c_B > 26 \quad \text{LHC}$$



Outlook for future efforts

time
Magnetic
CME/CMW

Vorticity
CVE/CVW

Toward the detection of anomalous effects

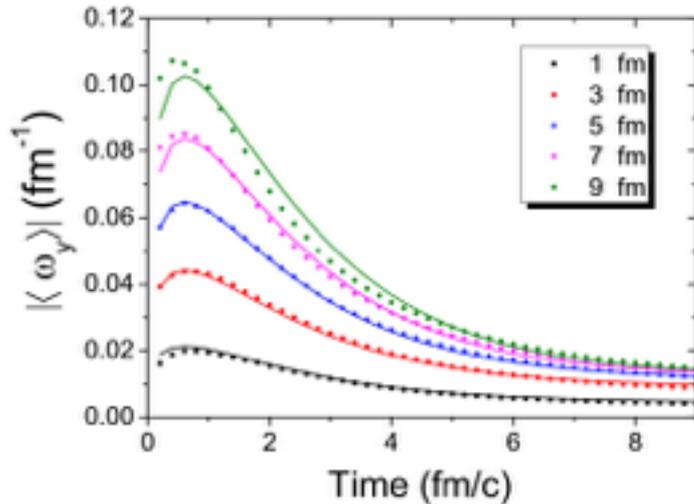
- **Challenges**

- Lifetime of B
- Initial axial charge density & CME current
- Chiral vortical effect
- Initial charge densities
- Modeling of background effects like LCC

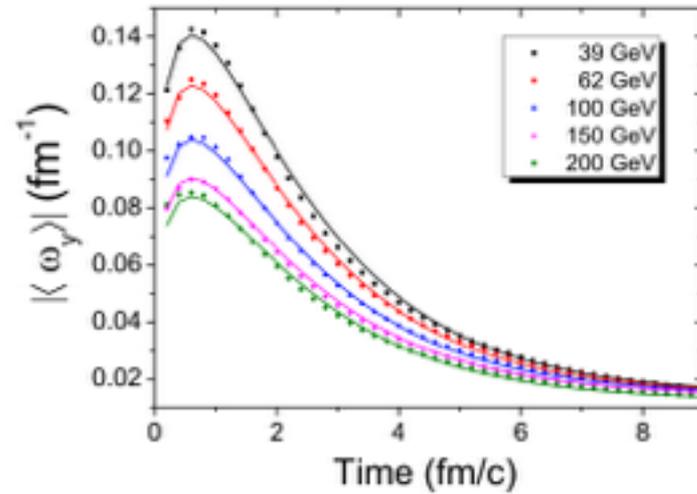
- **Chiral magnetohydrodynamics**

- Solve anomalous hydro + Maxwell equations consistently

Vorticity from HIJING/AMPT

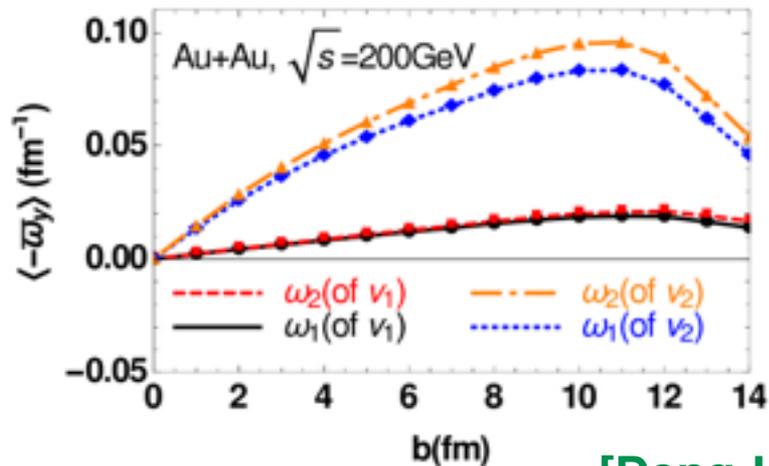


Impact parameter dep.

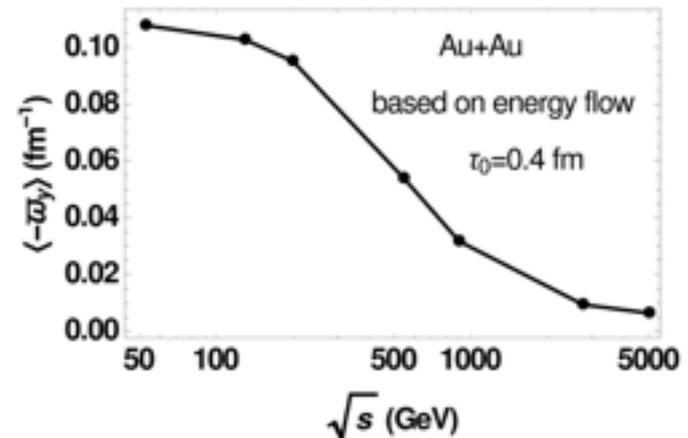


Collision energy dep.

[Jiang-Lin-Liao 1602.06580]



[Deng-Huang 1603.06117]



Chiral magnetic effect and anomalous transport from real-time lattice simulations

Niklas Müller,^{1,*} Sören Schlichting,^{2,†} and Sayantan Sharma^{2,‡}

[1606.00342]

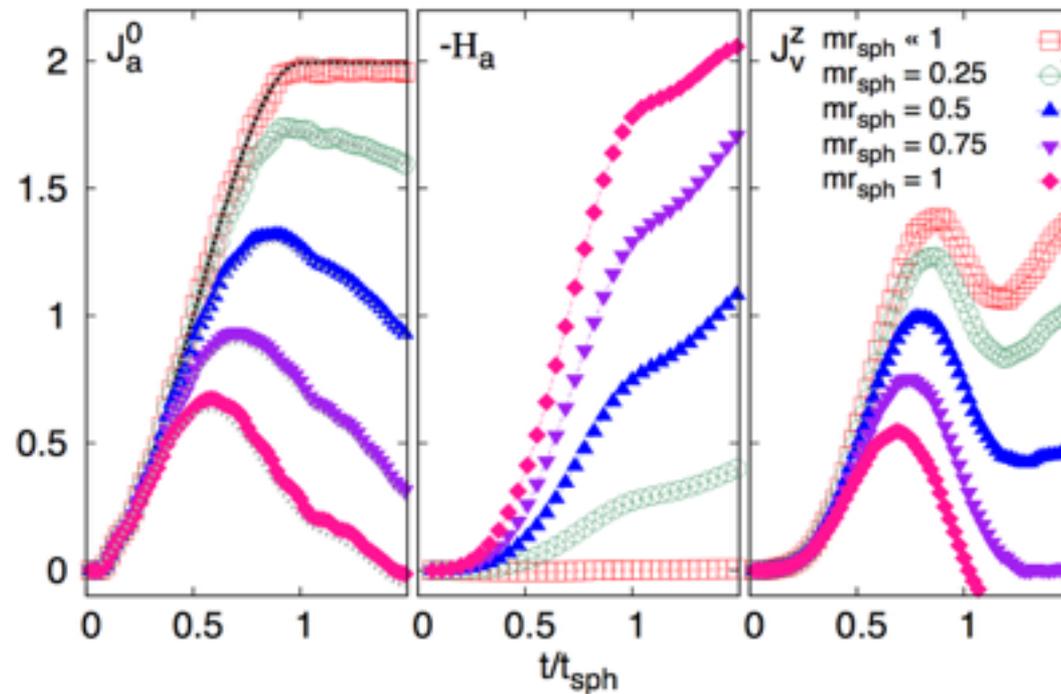


FIG. 3. Evolution of the axial charge J_a^0 (left), pseudoscalar condensate H_a (center) and vector current J_v^z (right) for different values of the fermion mass $mr_{\text{sph}} = 3 \cdot 10^{-3}, 0.25, 0.5, 0.75, 1.0$. Comparison with the gray lines in the left panel demonstrates that the axial anomaly relation (6) is satisfied in all cases.

Toward the detection of anomalous effects

- **Challenges**

- Lifetime of B
- Initial axial charge density & CME current
- Chiral vortical effect
- Initial charge densities
- Modeling of background effects like LCC

- **Chiral magnetohydrodynamics**

- Solve anomalous hydro + Maxwell equation consistently