

**Measurement of the  
polarization vector  
of the  $e^+$   
from the decay of polarized  $\mu^+$   
as a test  
of time reversal invariance**

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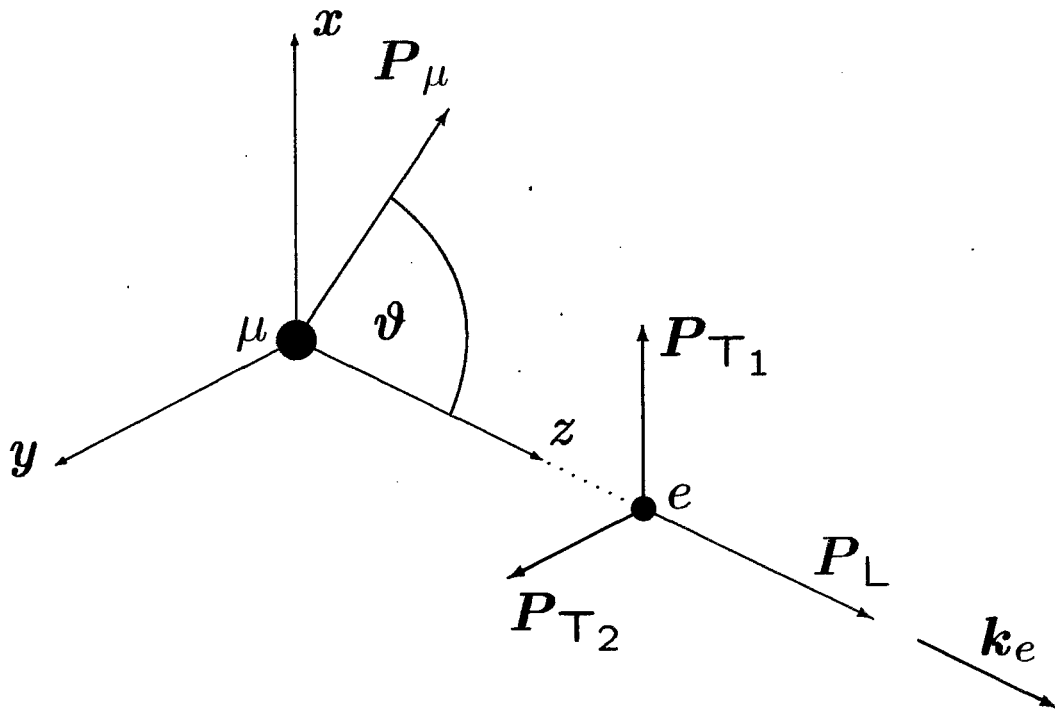
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$$P_{T1} = f_1(E, \vartheta, \eta, \eta'')$$

$$P_{T2} = f_2(E, \vartheta, \frac{\alpha}{A}, \frac{\beta}{A})$$

The standard model predicts:

$$\langle P_{T1} \rangle_E = 0.003$$

$$P_{T2} \equiv 0$$

A nonzero  $P_{T2}$  would signal violation of time reversal invariance. This is the only purely leptonic reaction for which TRI has been tested up to now.

### 3. Matrix Element

$$\mathcal{M} = \frac{4G_F}{\sqrt{2}} \sum_{\substack{\gamma=S,V,T \\ \varepsilon,\mu=R,L}} g_{\varepsilon\mu}^{\gamma} \langle \bar{e}_{\varepsilon} | \Gamma^{\gamma} | (\nu_e)_n \rangle \langle \bar{\nu}_m | \Gamma_{\gamma} | (\mu)_{\mu} \rangle$$

The index  $\gamma$  labels the type of interaction:

$$\begin{aligned} \Gamma^S &= 4\text{-scalar} \\ \Gamma^V &= 4\text{-vector} \\ \Gamma^T &= 4\text{-tensor} \end{aligned}$$

The indices  $\varepsilon, \mu$  indicate the chiralities of the spinors of the observed (charged) leptons. The chiralities  $n, m$  of the neutrinos are uniquely determined for given  $\gamma, \varepsilon$  and  $\mu$ .

The transverse polarization component  $P_{T_1}$  yields the low energy parameter  $\eta$  *without* the suppression factor  $m_e/m_{\mu}$  of  $\eta$  in the energy spectrum of the decay positron. These interference terms allow for sizeable effects.

$$\eta = \frac{1}{2} \text{Re} \{ g_{LL}^V g_{RR}^{S*} + g_{RR}^V g_{LL}^{S*} + g_{LR}^V (g_{RL}^{S*} + g_{RL}^{T*}) + g_{RL}^V (g_{LR}^{S*} + g_{LR}^{T*}) \}$$

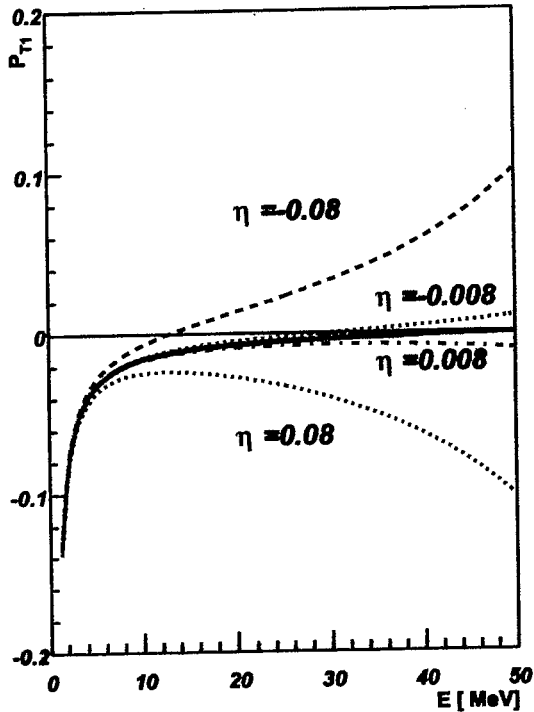
In the standard model

$$\begin{aligned} g_{LL}^V &= 1 \\ g_{\varepsilon\mu}^{\gamma} &= 0 \quad (\text{all other interactions}) \end{aligned}$$

Assuming 1 additional interaction and knowing that

$$g_{LL}^V \approx 1,$$

one deduces:



$$P_{T_1}(E_e) \rightarrow \eta \approx \frac{1}{2} \text{Re}\{g_{RR}^S\}$$

$$P_{T_2}(E_e) \rightarrow \frac{\beta'}{A} \approx \frac{1}{4} \text{Im}\{g_{RR}^S\}$$

Main scientific interests:

$P_{T_1}$ : Precise determination of Fermi coupling constant  $G_F$

$P_{T_2}$ : Test of time reversal invariance

## 4. Fermi coupling constant

Should be independent of masses and radiative corrections:

Universal coupling constant

$$G_F^2 = 192\pi^3 \cdot \frac{\hbar}{\tau_\mu} \cdot \frac{1}{m_\mu^5} \cdot \left\{ 1 + \frac{\alpha}{2\pi} \left( \pi^2 - \frac{25}{4} \right) \right\} \cdot \left\{ 1 - \frac{3}{5} \left( \frac{m_\mu}{m_W} \right)^2 \right\} \cdot \left\{ 1 - 4\eta \cdot \frac{m_e}{m_\mu} - 4\lambda \cdot \frac{m_{\nu\mu}}{m_\mu} + 8 \left( \frac{m_e}{m_\mu} \right)^2 + 8 \left( \frac{m_{\nu\mu}}{m_\mu} \right)^2 \right\}$$

New:

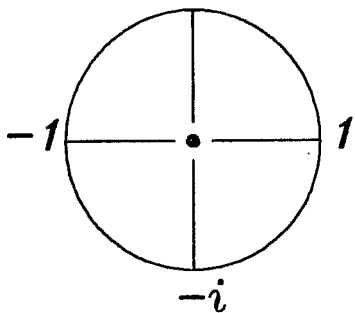
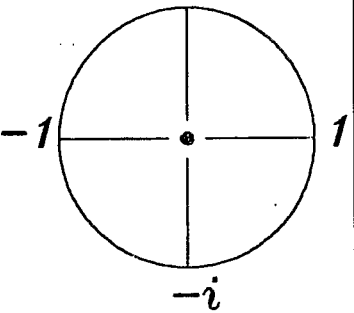
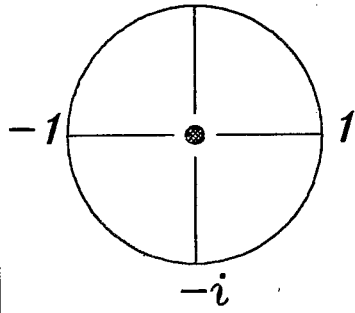
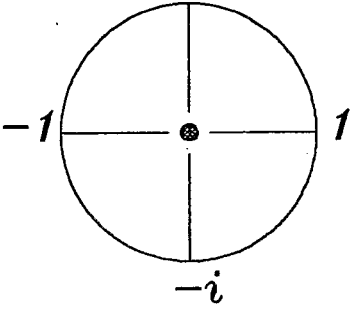
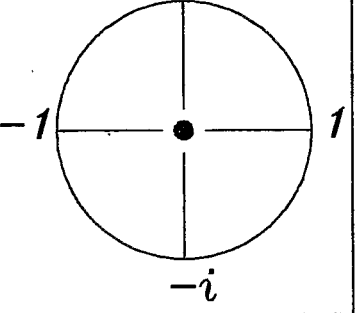
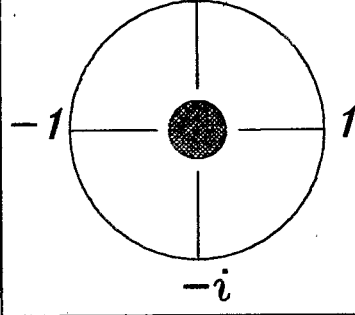
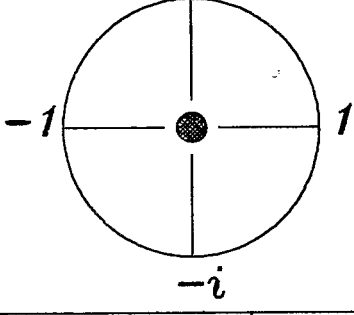
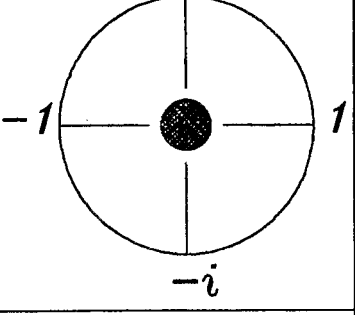
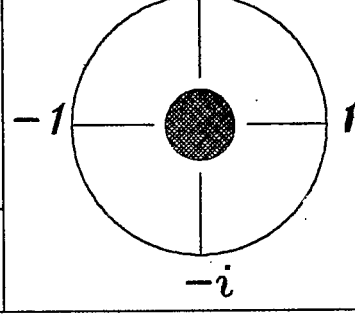
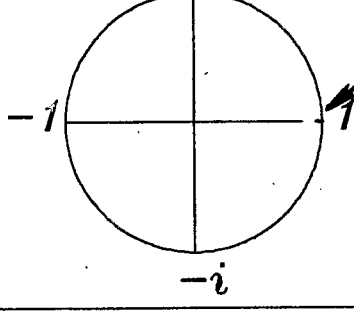

$$\lambda \approx \frac{1}{2} \text{Re} \{ g_{LL}^V \cdot g_{LR}^{V*} \}$$

In left-right symmetric models with mixing angle  $\zeta$ :

$$\lambda \approx \frac{1}{2} \zeta$$

Contribution from	$\frac{\Delta G_E}{G_F} [10^{-6}]$	
	$\mu^+ \rightarrow \bar{\nu}_\mu e^+ \nu_e$	$\tau^+ \rightarrow \bar{\nu}_\tau \mu^+ \nu_\mu$
$\Delta m_W$	0.0	1
$\Delta m_{\mu,\tau}$	0.2	421
$\Delta \tau$	9.1	2 070
$\Delta(\lambda m_{\bar{\nu}})$	70.0	22 500
$\Delta \eta$	193.0	6 500
$\Delta \Gamma_{\tau \rightarrow \mu}$	-	100 000

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$$

$g_{\epsilon\mu}^{\gamma_1}$	$S$	$V$	$T$
	$i$  $-i$	$i$ $V+A$  $-i$	
$e_R \mu_R$			
	$i$  $-i$	$i$  $-i$	$i$  $-i$
$e_L \mu_R$			
	$i$  $-i$	$i$  $-i$	$i$  $-i$
$e_R \mu_L$			
	$i$  $-i$	$i$ $V-A$  $-i$	
$e_L \mu_L$			



## 5. Experimental method

Measure the complete polarization vector of the decay positrons:

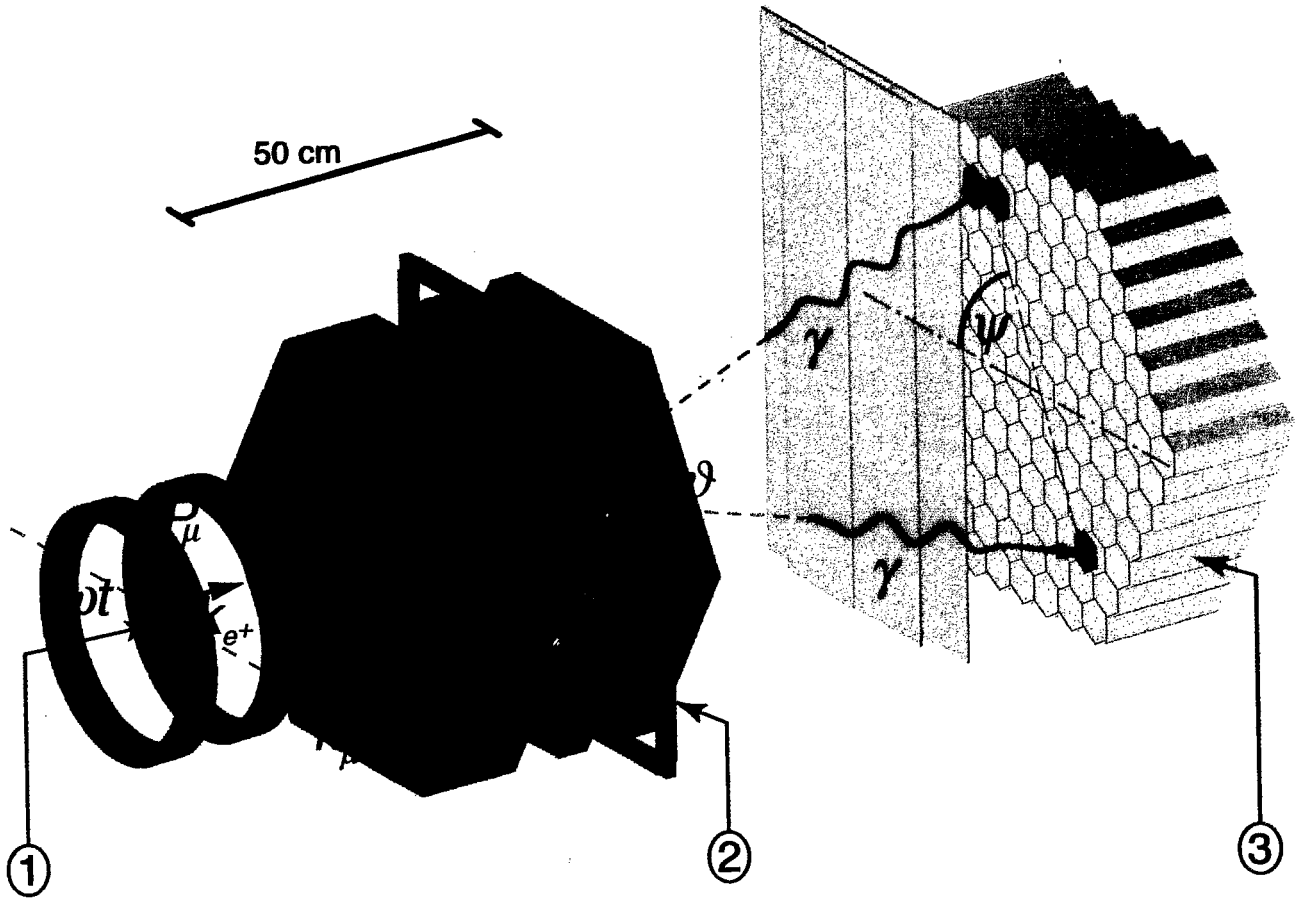
$$\mathbf{P}_{e^+} = \begin{pmatrix} P_{T1} \\ P_{T2} \\ P_L \end{pmatrix} \equiv \begin{pmatrix} P_T \cdot \cos \varphi \\ P_T \cdot \sin \varphi \\ P_L \end{pmatrix}$$

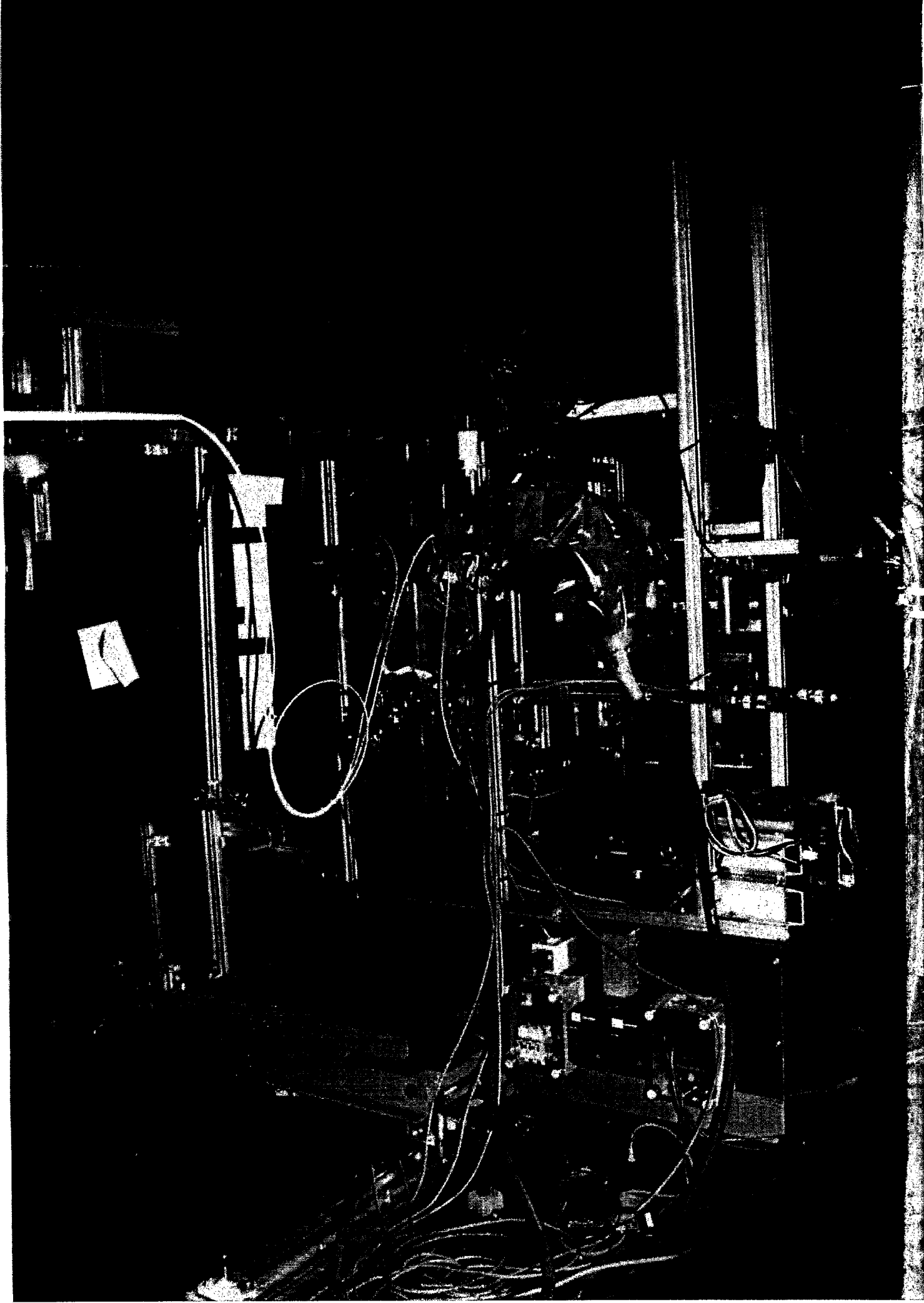
with 3 simultaneous and independent measurements:

Observable	Method
$P_T$	Time dependence of annihilation
$\varphi$	Remnant $\mu$ SR effect
$P_L$	Spatial dependence of annihilation

# 6. Experimental setup

- 6.1 Highly polarized  $\mu^+$  beam at  $\mu E1$  area of PSI: (91%)
- 6.2 Muon stop rate in Be target:  
(20 – 80)  $\times 10^6$  s<sup>-1</sup>
- 6.3 Precession in homogeneous  $B$  field;  
precession frequency = cyclotron frequency  
(50.8 MHz)
- 6.4 Burst width 3.9 ns (FWHM)  
 $\implies$  80% muon polarization in Be stop target
- 6.5 Positron tracking with drift chambers
- 6.6 Annihilation with polarized  $e^-$
- 6.7 Detection of annihilation quanta with 127 BGO crystals





# 7. Event reconstruction and data analysis

7.1  $e^+$  annihilation-in-flight as analysing reaction for transverse polarization

7.2 Muon decay asymmetry ( $\mu$ SR yields time zero)

7.3 Energy calibration with cosmic rays

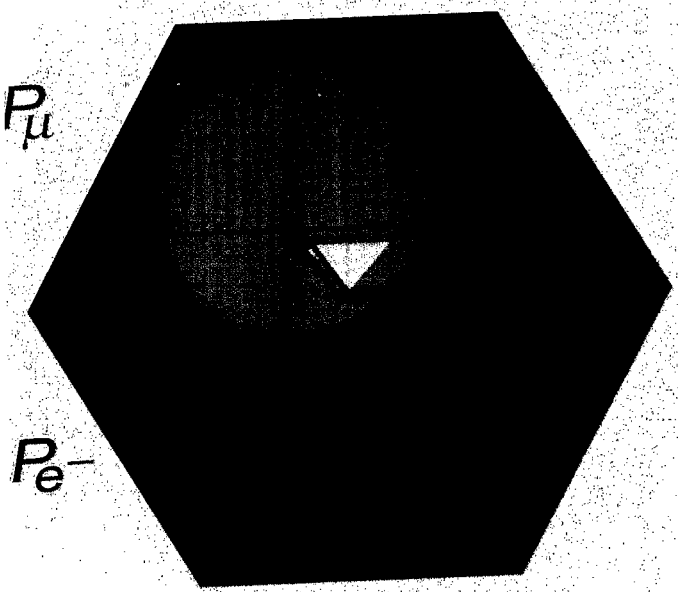
7.4 Cluster recognition

7.5 Background suppression

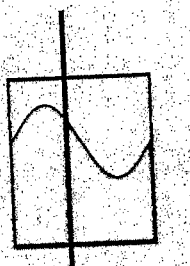
7.6 Analysing power of annihilation events

7.7 Energy distribution of transverse polarization

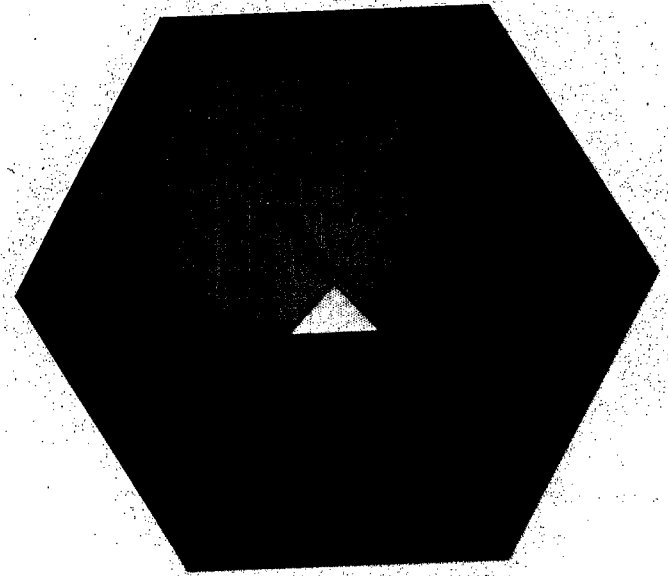
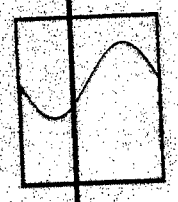
7.8 Longitudinal polarization



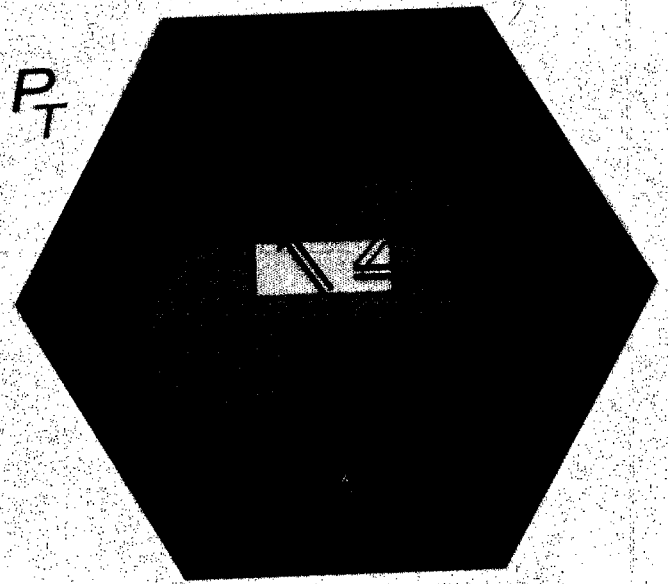
Selected Positrons •



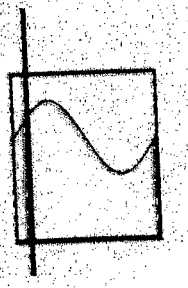
+



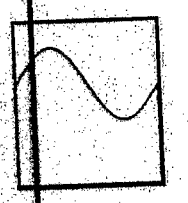
$\mu$ SR Effect



Selected Positrons •



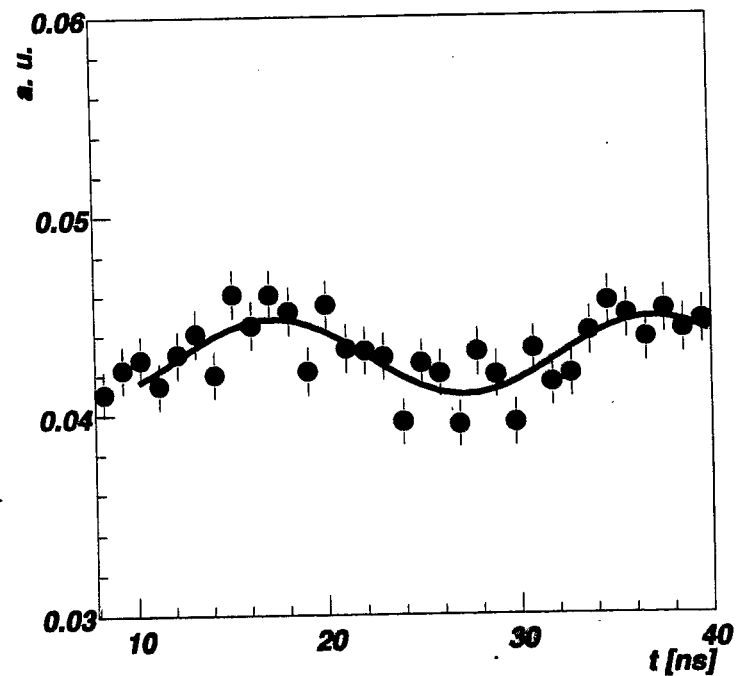
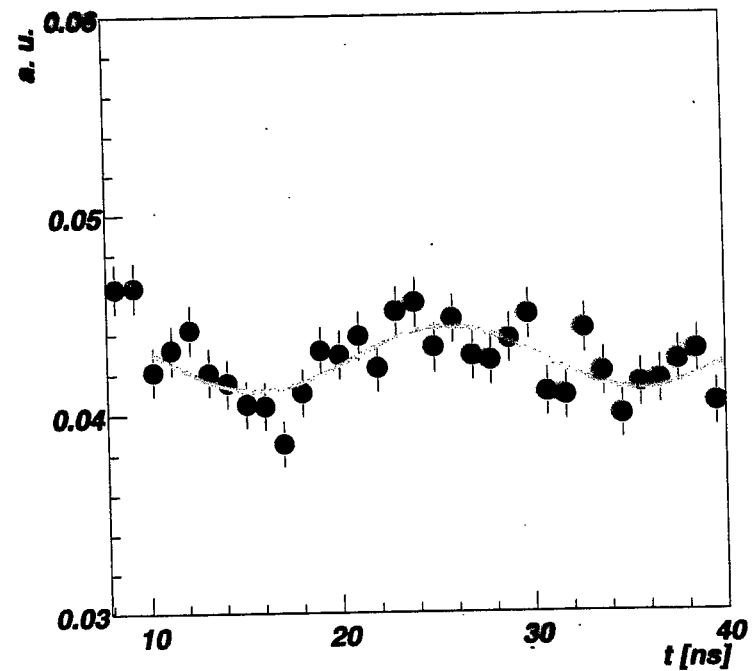
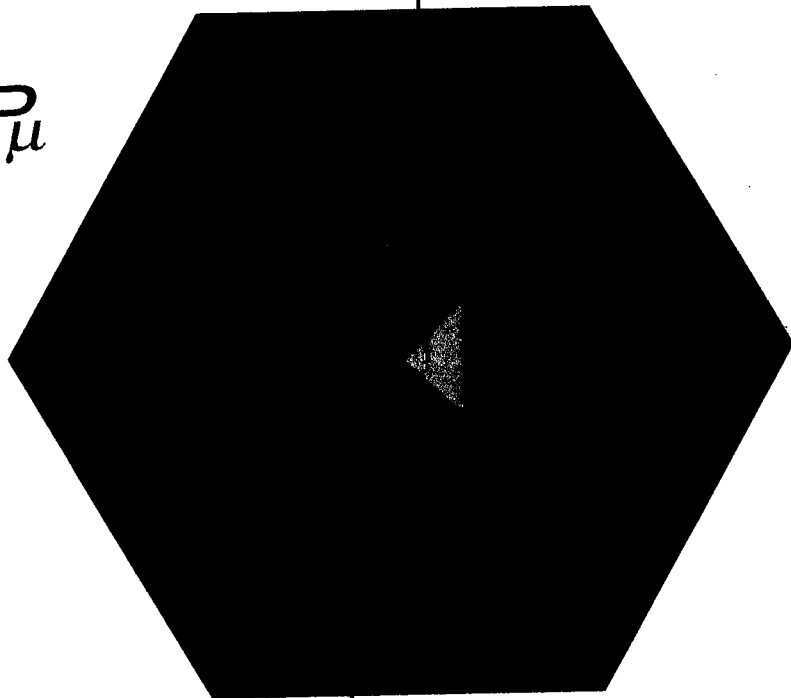
+



$P_T$  Effect

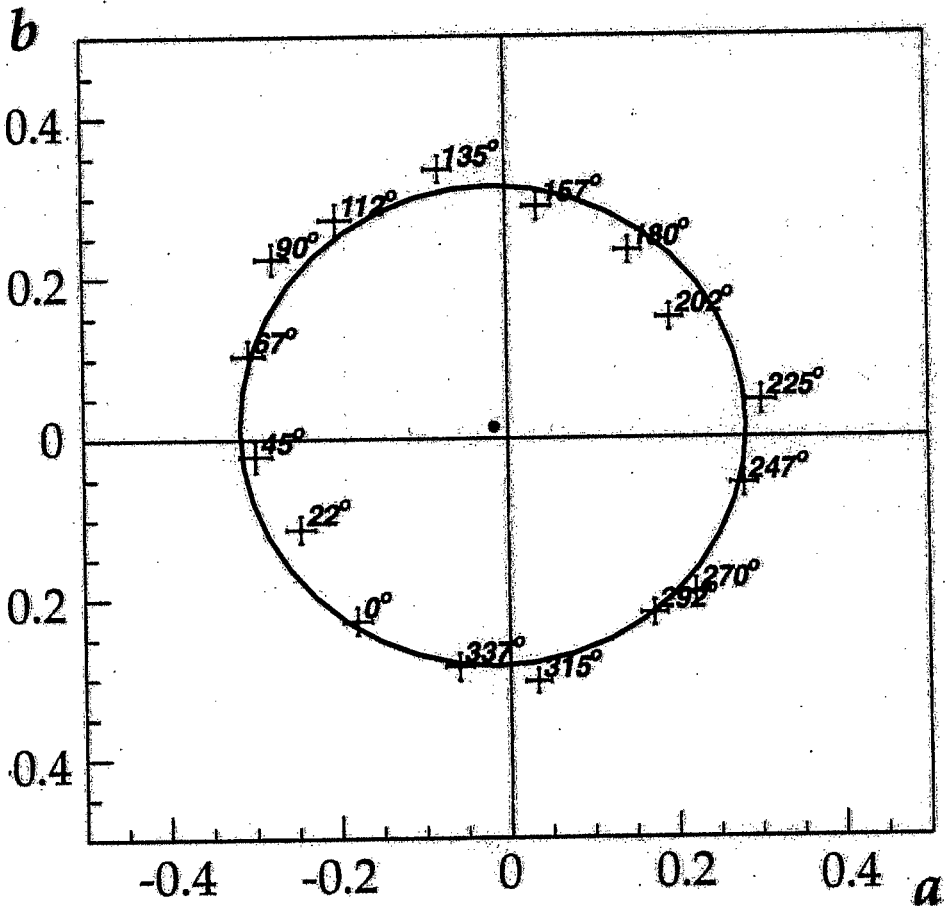


$P_{\mu}$



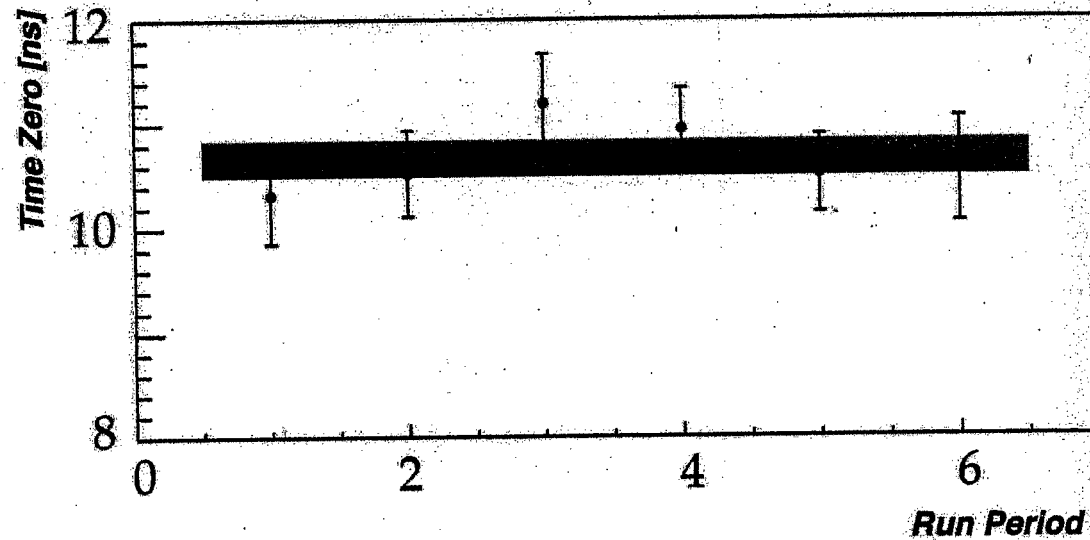
# Phases and Amplitudes for all azimuthal bins ( $\varphi$ )

$$N \sim (1 + a \cos \omega t + b \sin \omega t)$$



**Average  $\mu$ SR amplitude:  $0.297 \pm 0.004$ ,  
consistent with theory.**

## Stability of the Time-Zero during the whole run

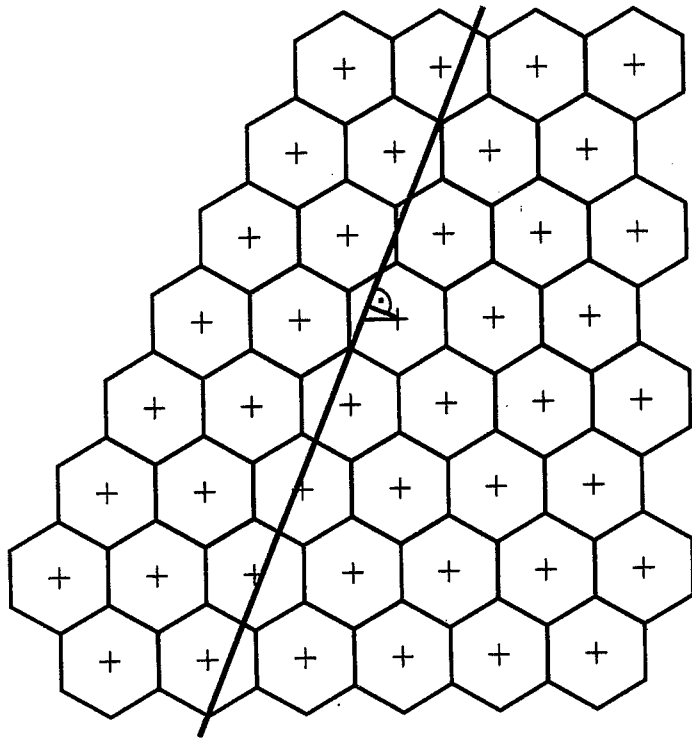


**Time-Zero :  $10.67 \pm 0.187$  ns gives the time,  
when the muon spin shows upward.**

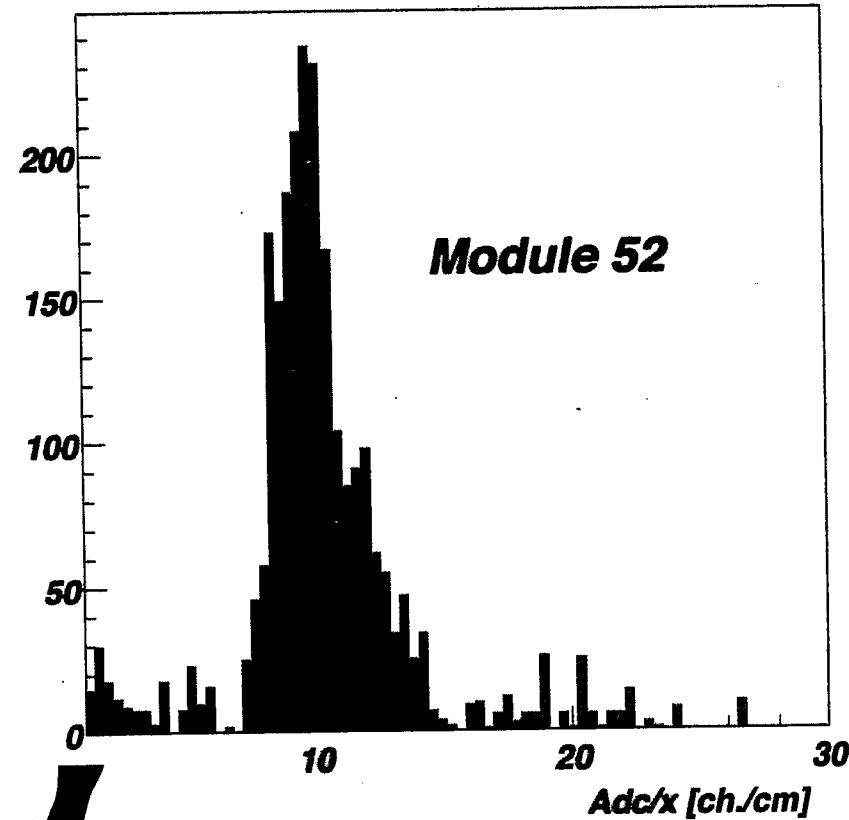


# Data is Calibrated with Cosmic Muons

Reconstructing the cosmic tracks gives the tracklength in BGOs



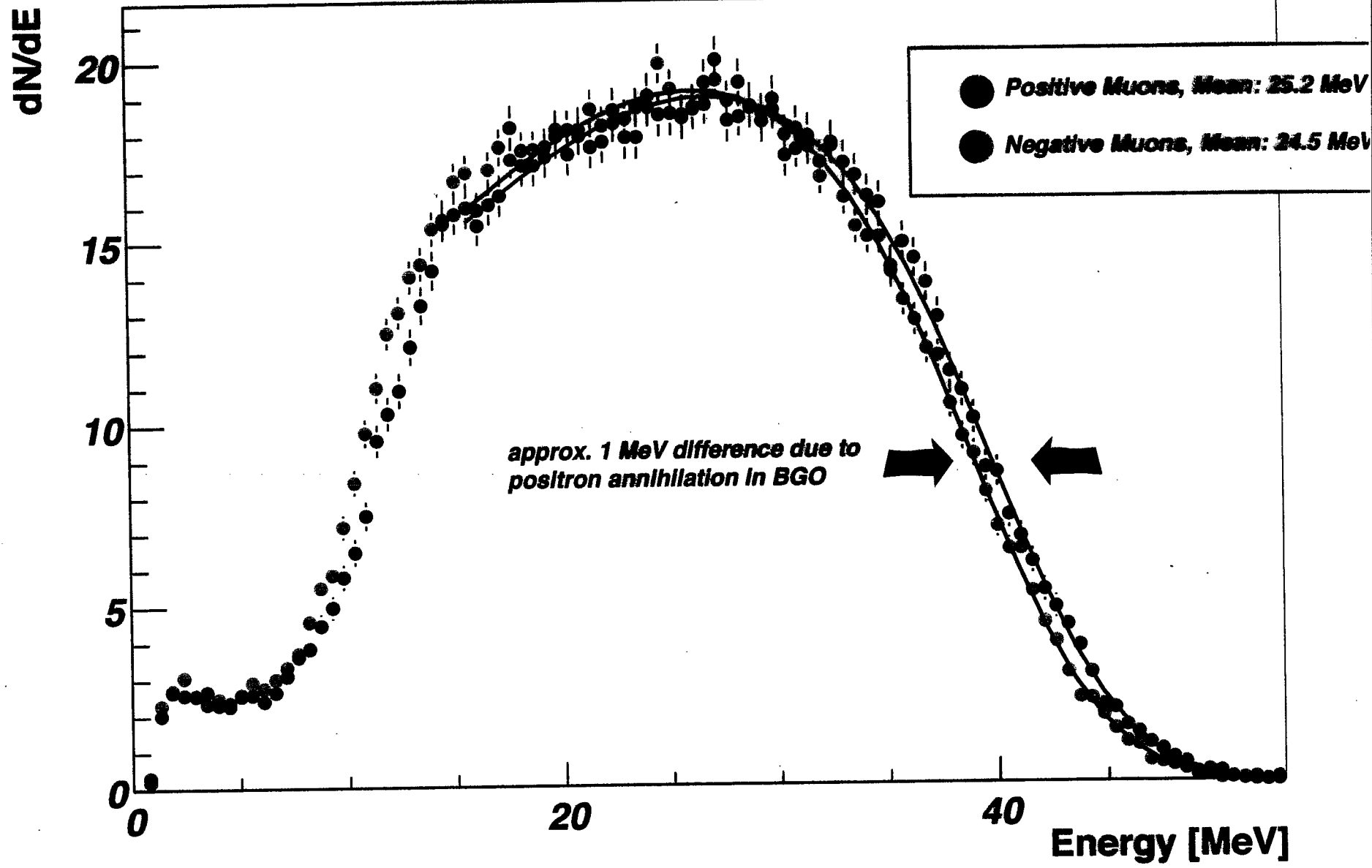
Histogram Adc/x for cosmic

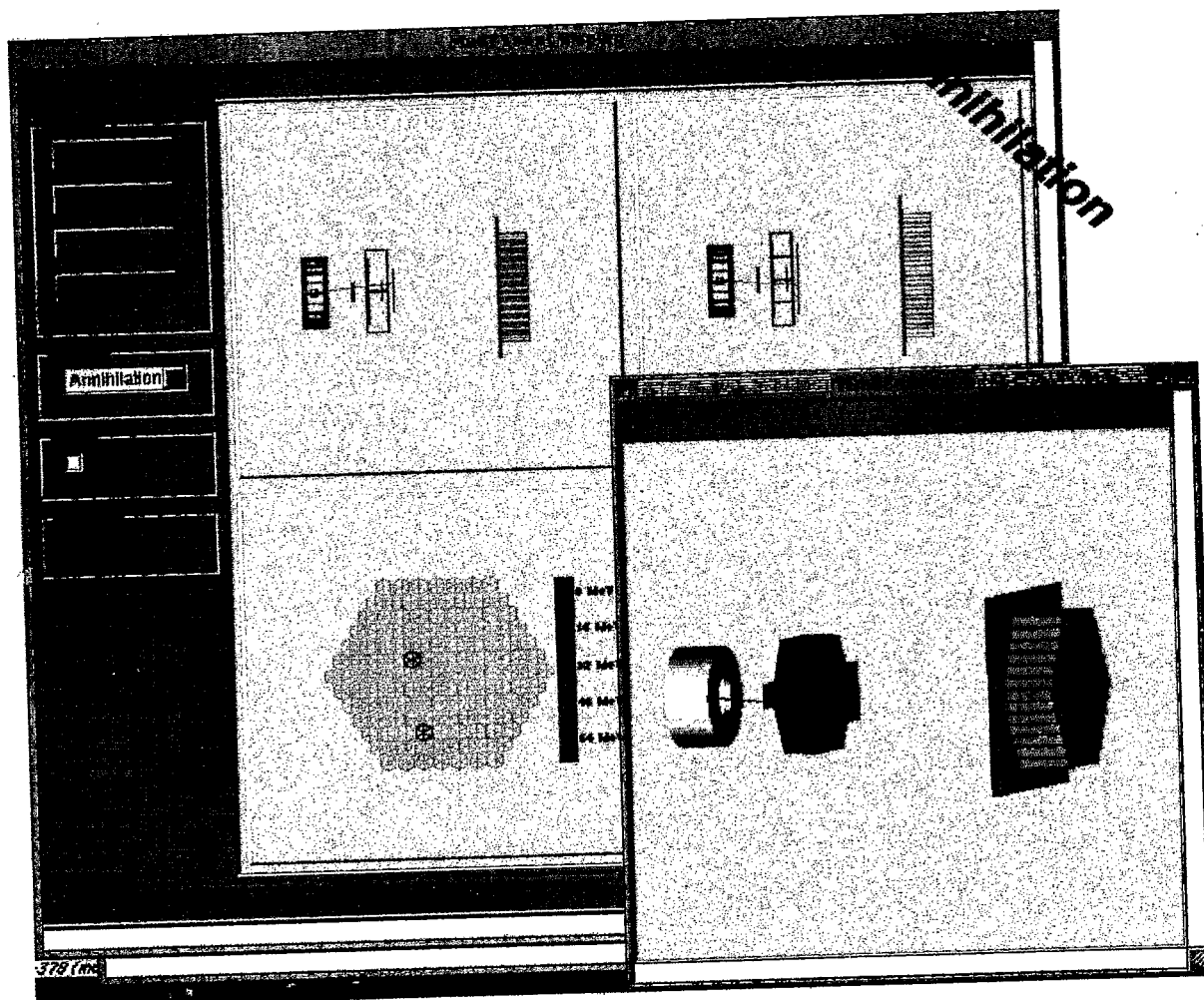
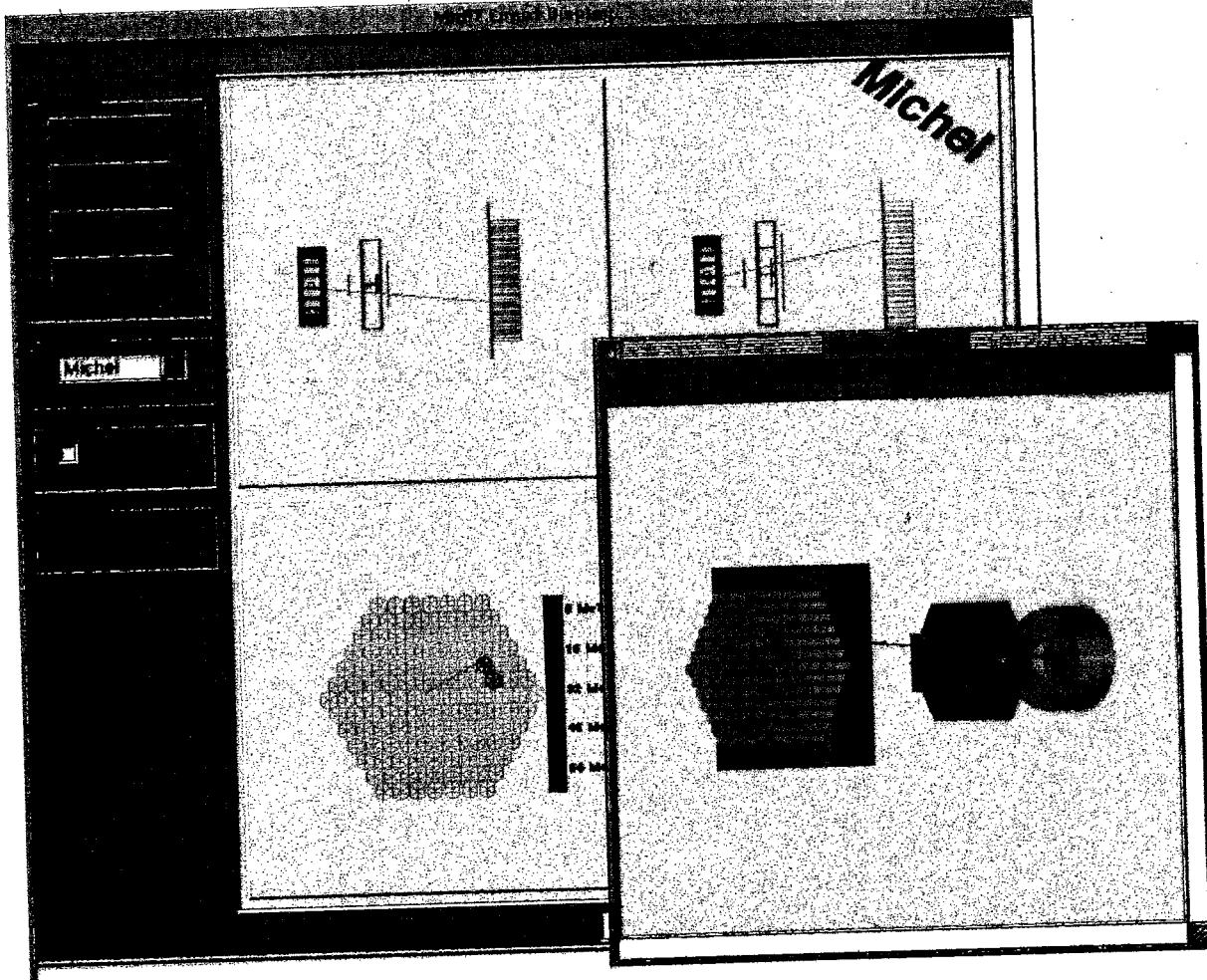


Monte Carlo

$$\frac{1}{E_i/x_i} \times \frac{ADC_i}{x_i} = \frac{ADC_i}{E_i} = c_i$$

Calibration constant for BGO number  $i$





# Triggering on Two Photons at a minimal Distance: The Cluster Recognition Unit

## How does it work?

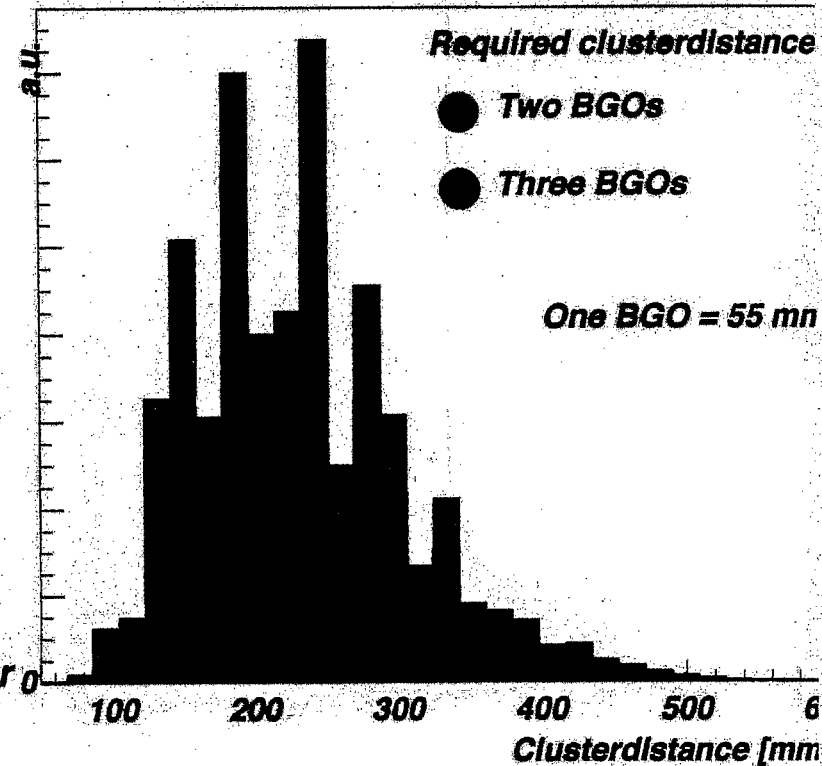
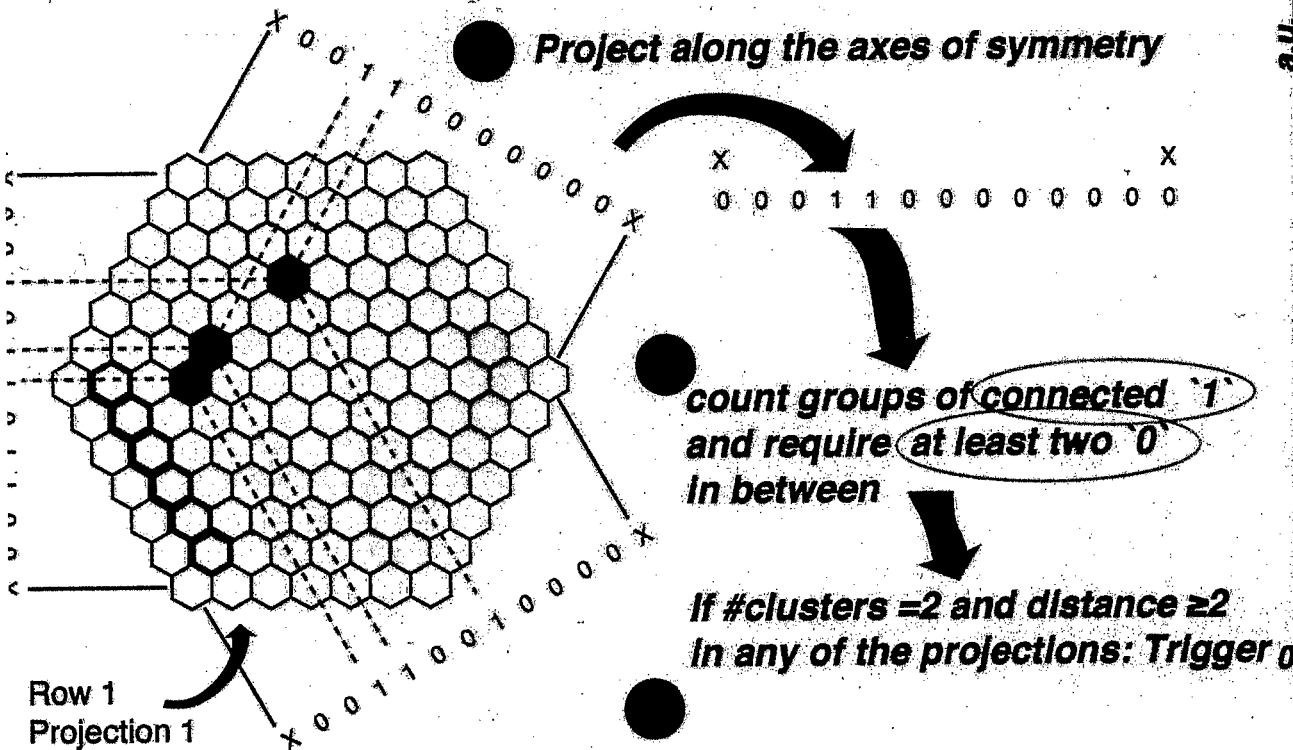
Kinematics require a minimal cluster distance

$$\cos \vartheta = 1 - m_e \frac{4}{E_{e^+}} \quad d = 2z \tan \frac{\vartheta}{2}$$

$$d_{min} = d(E_{e^+} = 50 \text{ MeV}) \approx 16 \text{ cm}$$

FPGA approach allows redefining trigger conditions 'on the fly'

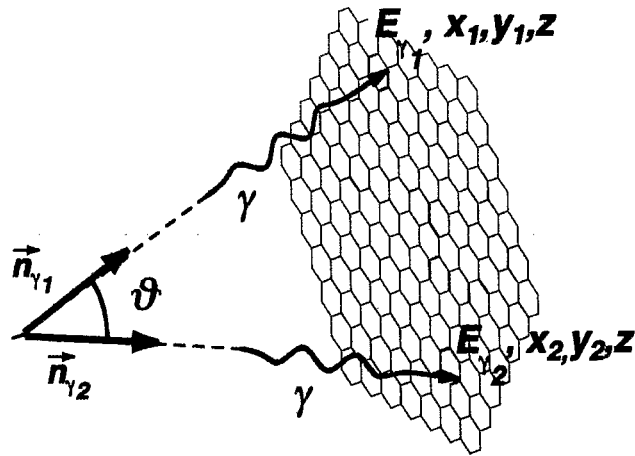
## Distance between two clusters



# Cut on the Kinematics to Extract the 'Good' Events

Calculate  $\vartheta$  in two different ways:

Sample from the last run



**Energy:**

$$\cos \vartheta = 1 - m_e \frac{E_{\gamma 1} + E_{\gamma 2}}{E_{\gamma 1} \cdot E_{\gamma 2}}$$

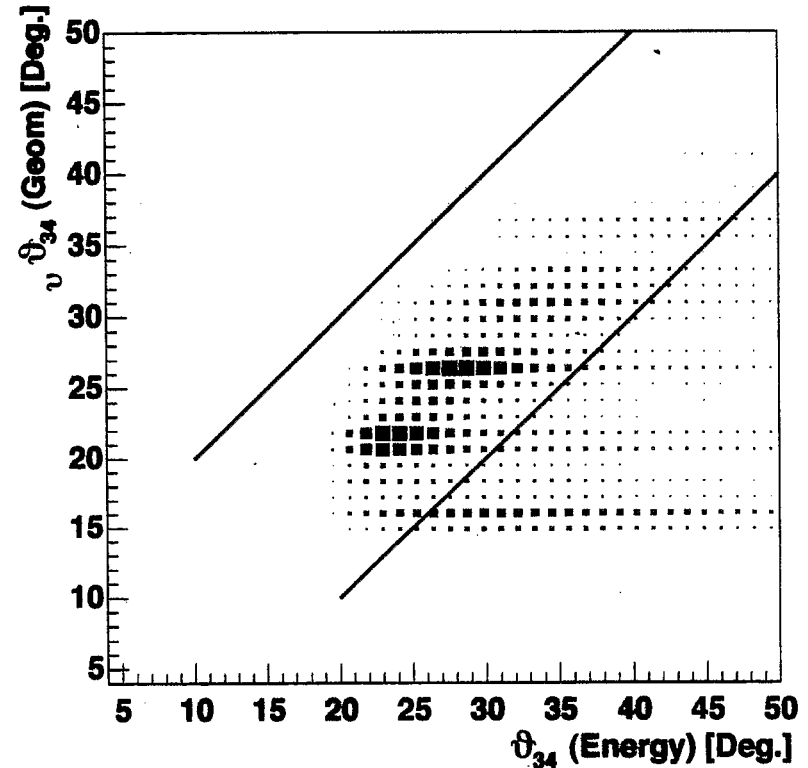
**Geometry**

$$\cos \vartheta = \vec{n}_{\gamma 1} \cdot \vec{n}_{\gamma 2}$$

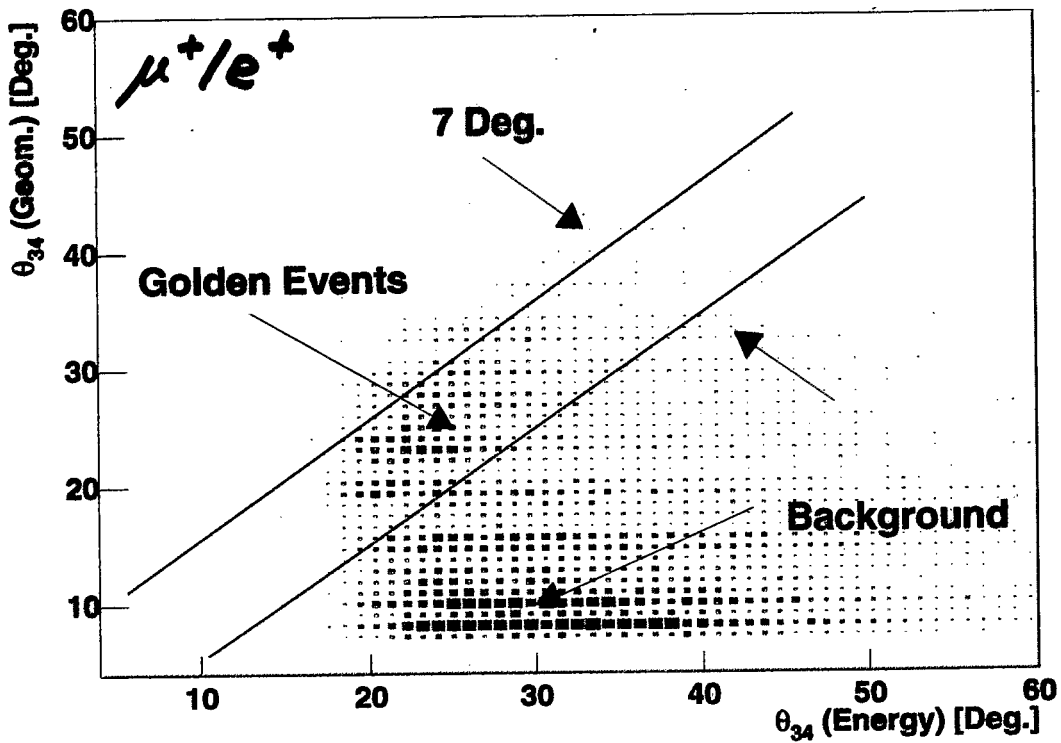


$$\vartheta^{\text{Geom}} = \vartheta^{\text{Energy}}$$

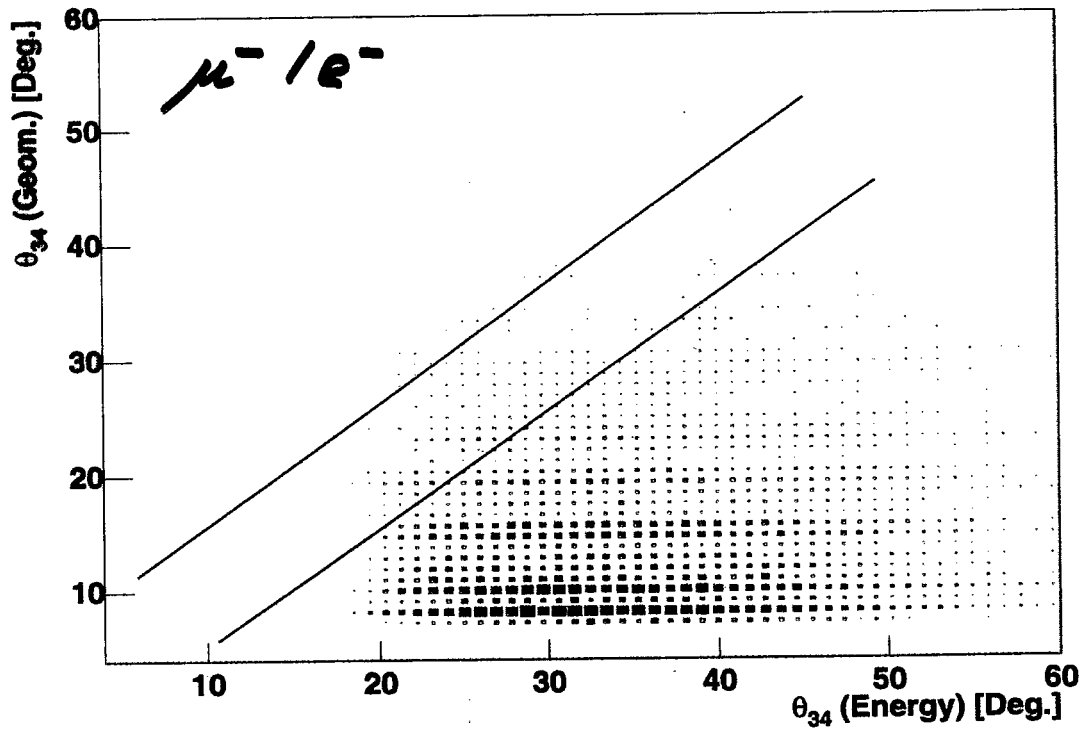
**for 'good' annihilations**



$\theta_{34}$  (Energy) vs.  $\theta_{34}$  (Geom.)

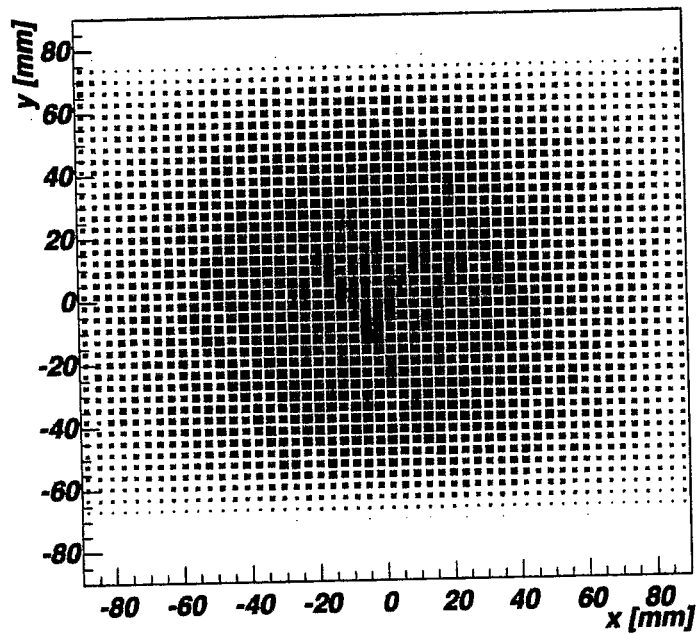


$\theta_{34}$  (Energy) vs.  $\theta_{34}$  (Geom.)



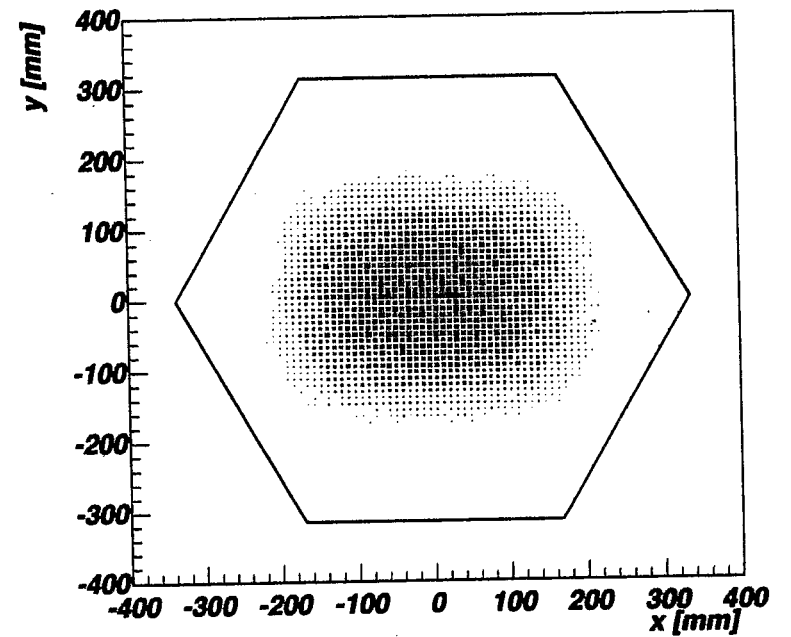
# Where do 'good' annihilations

come from,



reconstructed position  
of annihilations on the foil

and where do they go?

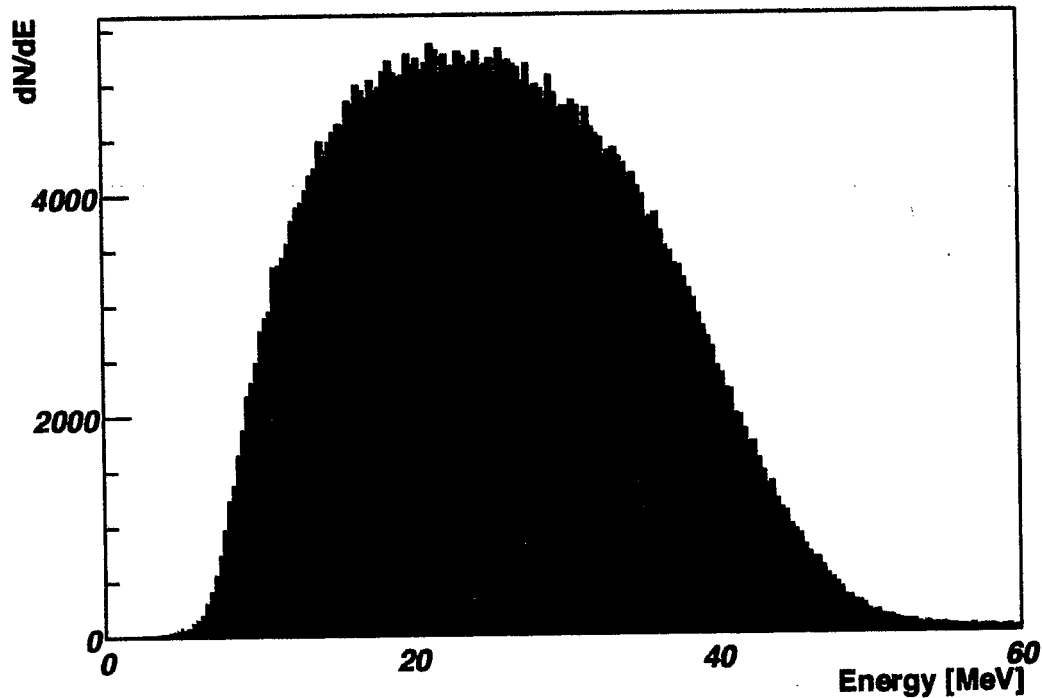


center of energy on the BGO-wall

# ***Analysis, Step Two: From Raw Data to `Good` Events***

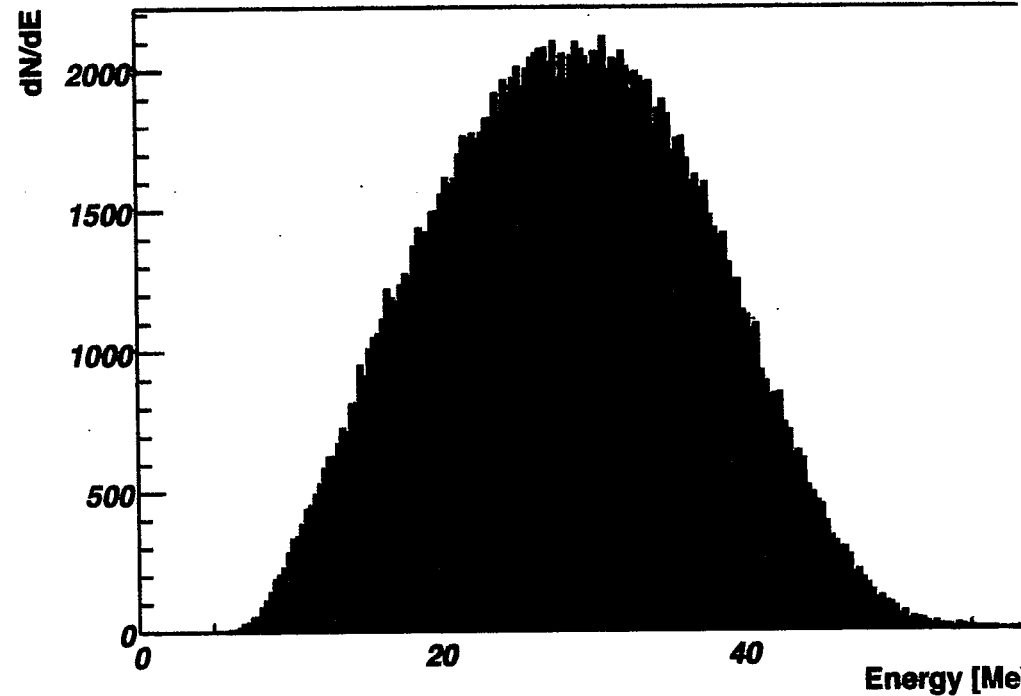
***Energy ( $E_{\text{tot}} = E_{\gamma_1} + E_{\gamma_2}$ ) spectra of annihilation events***

***After reconstruction ( $\approx 50\%$  Eff.)***



***charged track reconstructed,  
exactly two clusters in the BGO.***

***After selection ( $\approx 34\%$  Eff.)***



***event kinematics must be  
consistent with annihilation  
hypothesis.***

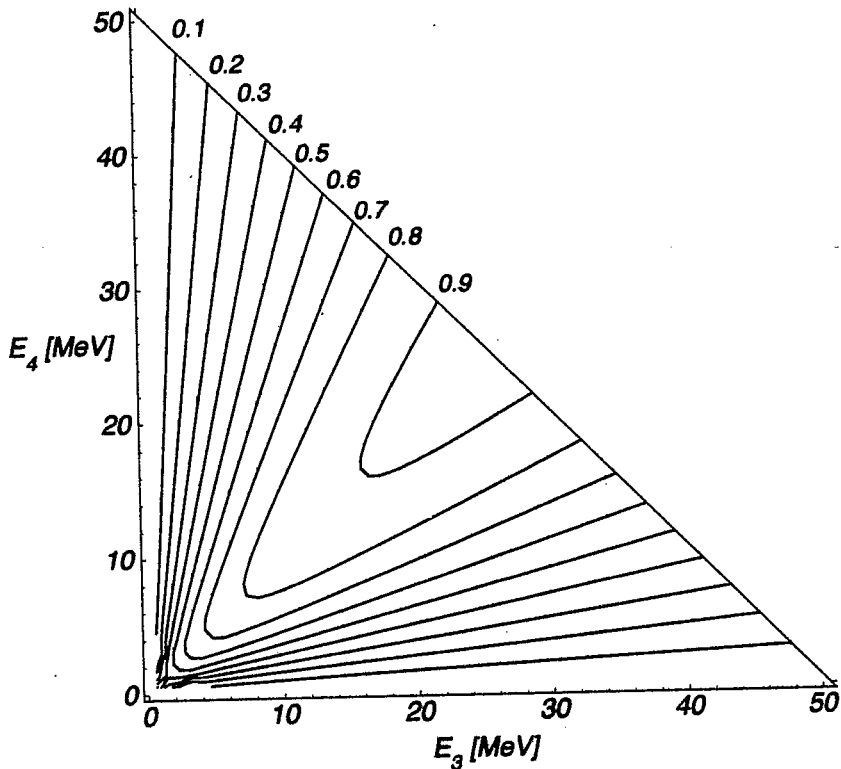
***After all cuts  $\approx 17\%$  `good` annihilation events remain.***



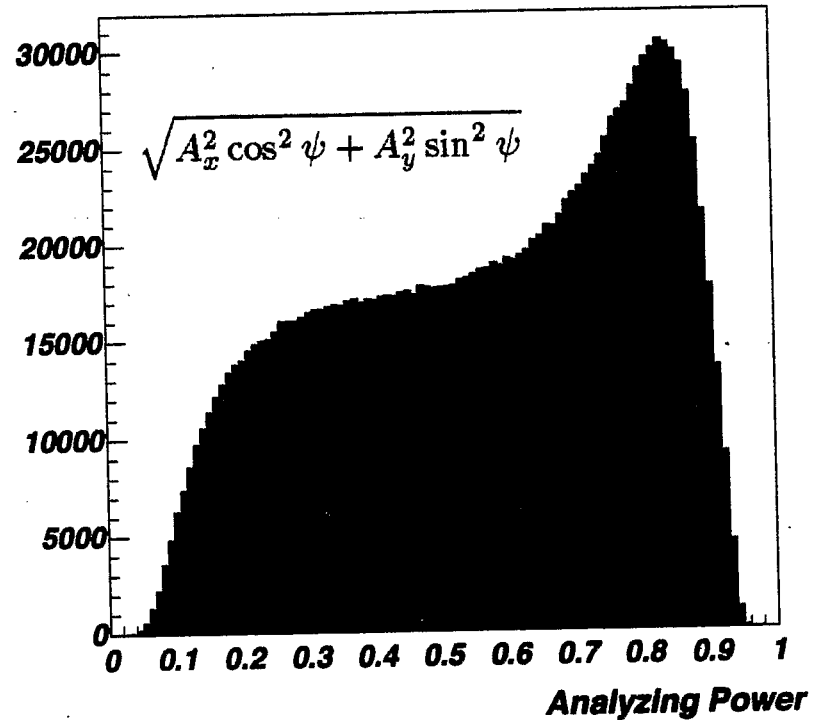
# Theoretical and Experimental Analyzing Power

The analyzing power is the amplitude of the expected oscillation

$$A = S \cdot \sqrt{P_1^2 + P_2^2} \cdot \sqrt{A_x^2 \cos^2 \psi + A_y^2 \sin^2 \psi}$$



Theoretical analyzing power,  
 $\Psi = 90^\circ$   $P_{e^-} = 100\%$ ,  $|P_T| = 1$

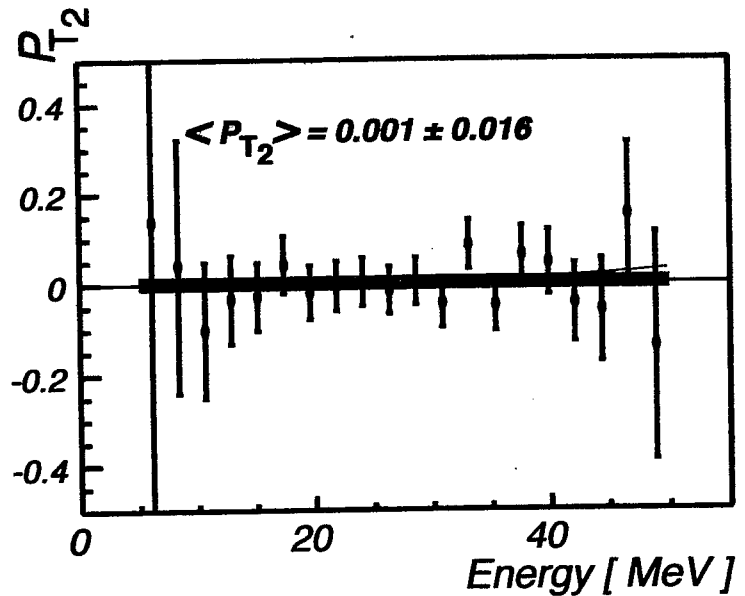
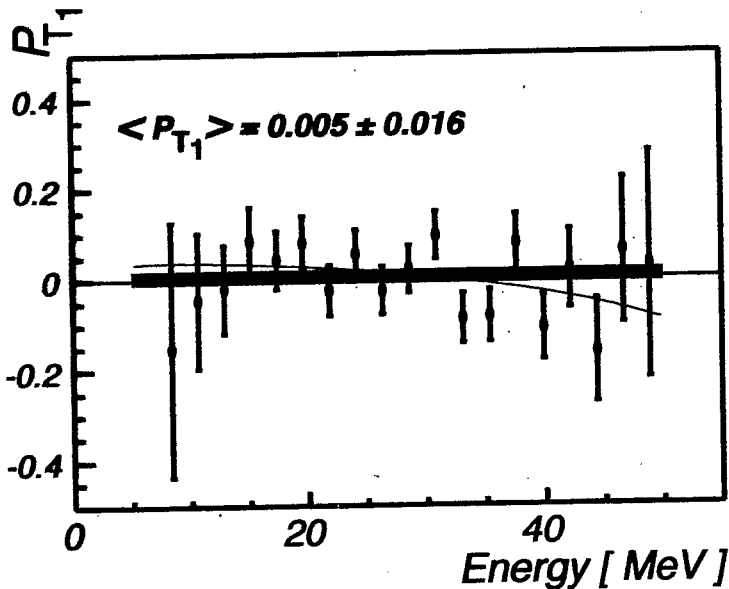


Analyzing Power of 'good' annihilation events.

$A_x$  and  $A_y$  are functions of the photon energies

## Transverse Polarization Components $P_{T1}$ and $P_{T2}$

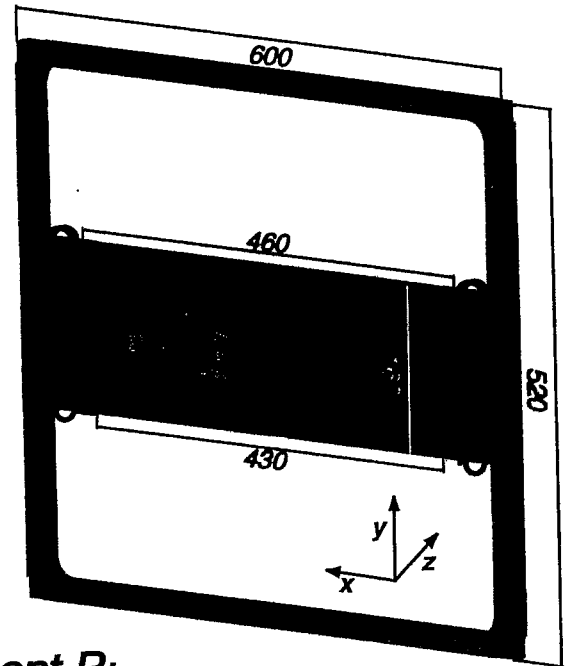
- time zero from  $\mu$ SR Effect (→ orientation of muon spin relative to  $P_1$  and  $P_2$ )
- rotation of transverse polarization components in the field of the spin precessing magnet ( MC )
- convolution with energy loss of the positron in the apparatus ( MC )
- sums of  $P_{T1}$  and  $P_{T2}$  for negative and positive foil polarization



# Measurement of the Longitudinal Polarization

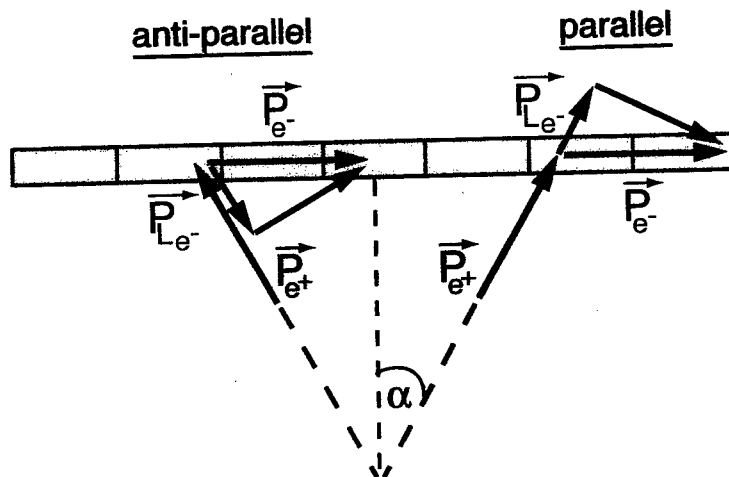
using information about position on magnetized Vacoflux foil (determined by tracks reconstructed from drift-chamber data) where annihilations take place

area on foil taken into account:  $140^2 \text{ mm}^2$   
 area divided into rectangular bins (ij),  
 17 bins in x- and y-direction, respectively



Tracks that do not hit the center of the foil 'see' a longitudinal component  $P_{Le^-}$  of the polarization of the electrons in the foil.

This  $P_{Le^-}$  can either be parallel or anti-parallel to the positron polarization :



# Longitudinal Polarization $P_L$ of the Positrons

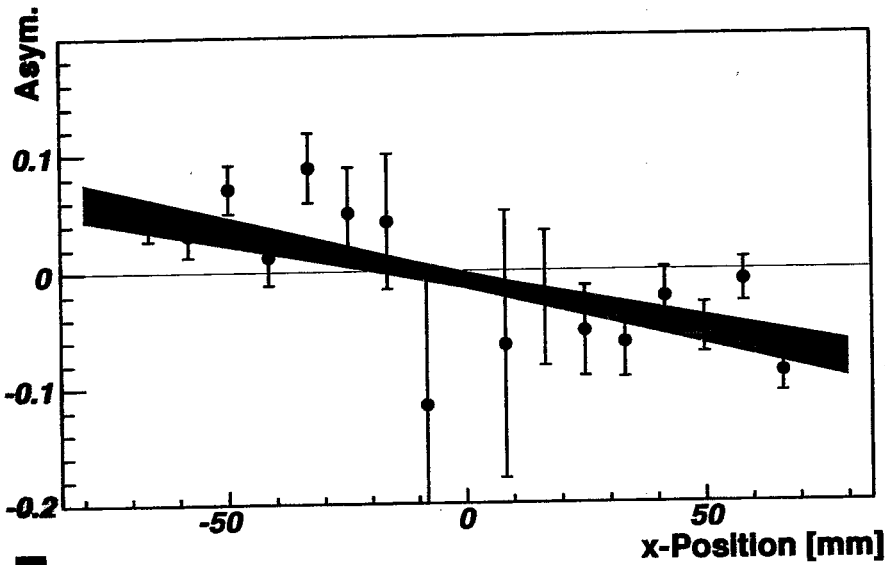
annihilation cross section depends on relative orientations of spins; it is larger if both spins are anti-parallel

Asymmetry: 
$$A_{ij} = \frac{\frac{n_{ij}^-}{N^-} - \frac{n_{ij}^+}{N^+}}{\frac{n_{ij}^-}{N^-} + \frac{n_{ij}^+}{N^+}}$$

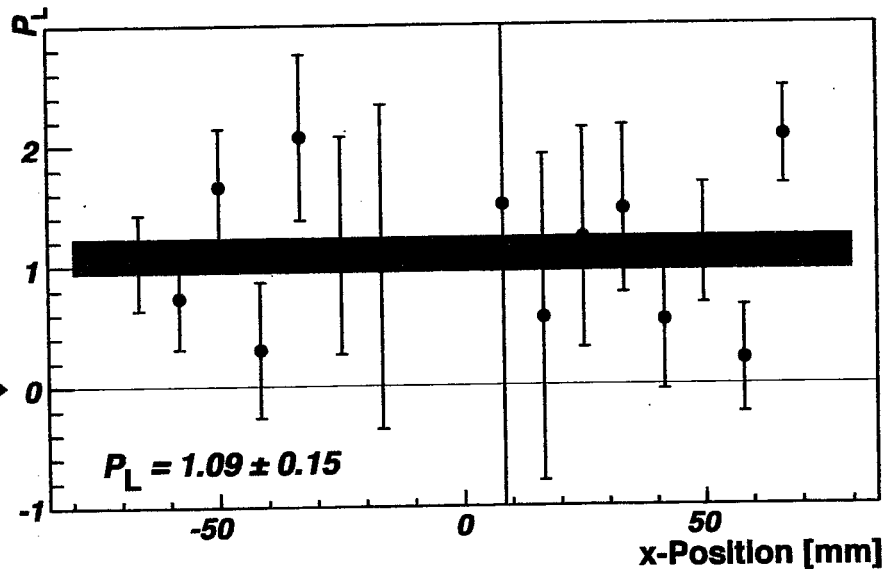
where  $n_{ij}^+$  : number of annihilations in bin  $ij$  for positive foil polarization

$N^+$  : total number of annihilations for positive polarization

$n_{ij}^-, N^-$  : same for negative polarization



- angle  $\alpha$
- elektron polarization in foil ( $P_{e^-} = 7.2\%$ )
- analysing power of 0.79
- background factor of 0.75 (backgr. ratio 25 %, mainly due to bremsstrahlung)

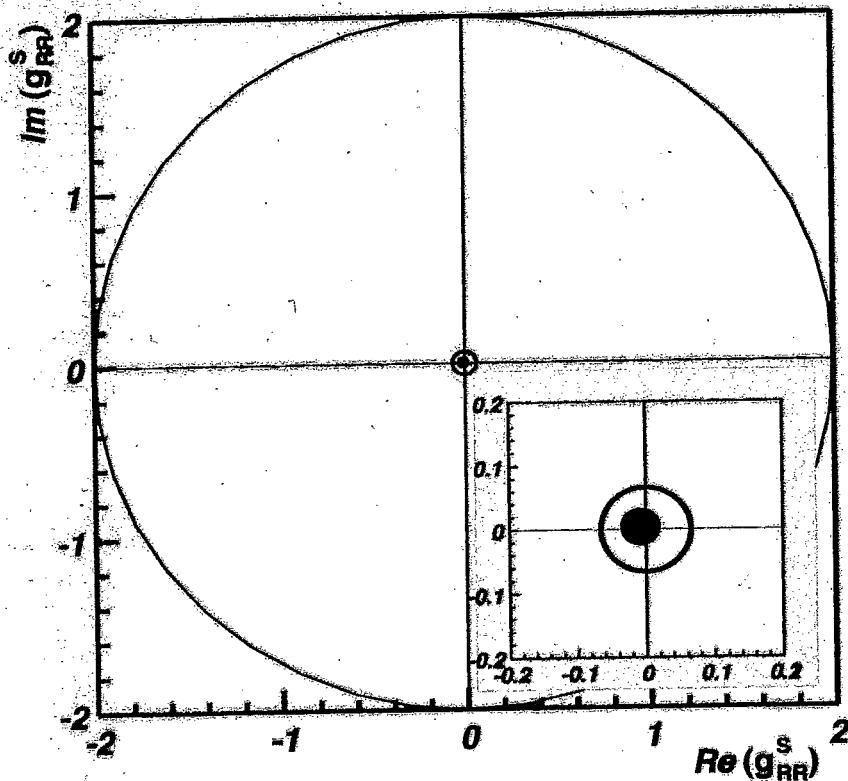


## 8. Preliminary Results

	General analysis	V - A + $g_{RR}^S$
$10^3 \times \langle P_{T_1} \rangle_E$	$5 \pm 16$	$5 \pm 16$
$10^3 \times \langle P_{T_2} \rangle_E$	$1 \pm 16$	$1 \pm 16$
$10^3 \times \eta$	$95 \pm 60$	$-4 \pm 14$
$10^3 \times \eta''$	$98 \pm 57$	$-10^3 \times \eta$
$10^3 \times \alpha'/A$	$-13 \pm 29$	0
$10^3 \times \beta'/A$	$8 \pm 16$	$1 \pm 7$
$10^3 \times \text{Re } g_{RR}^S$	—	$-8 \pm 28$
$10^3 \times \text{Im } g_{RR}^S$	—	$4 \pm 28$
$10^3 \times  g_{RR}^S $	—	$9 \pm 28$

## Implications from the Results

### 1. No Evidence for Additional Scalar Couplings in Muon Decay



Green circle: Result of a general analysis including all possible left- and righthanded scalar, vector and tensor couplings.

Red circle: Only one additional righthanded scalar coupling interferes with the lefthanded vector coupling in the SM.

$$\eta = -0.004 \pm 0.014$$

$$\frac{\beta'}{A} = 0.001 \pm 0.007$$

$$\text{Im}(g_{RR}^s) = 0.004 \pm 0.028$$

$$\text{Re}(g_{RR}^s) = -0.008 \pm 0.028$$

$$|g_{RR}^s| = 0.009 \pm 0.028$$

## 8. Outlook

Improve precision of previous experiment [1] by almost one order in magnitude to:

$$\Delta \langle P_{T_1} \rangle = 0.004$$

$$\Delta \langle P_{T_2} \rangle = 0.004$$

Assuming  $V - A$  and one additional coupling, this will reduce the limits for  $\eta$  and  $g_{RR}^S$  to

$$\Delta\eta = 0.004$$

$$\Delta \text{Re} \{g_{RR}^S\} = 0.009$$

$$\Delta \text{Im} \{g_{RR}^S\} = 0.009$$

[1] H. Burkard, F. Corriveau, J. Egger, W. Fetscher, H.-J. Gerber, K.F. Johnson, H. Kaspar, H.-J. Mahler, M. Salzmann, F. Scheck

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