

RIXS for 5d oxides in “atomic limit” Mott insulators

From Atomic Lego Blocks to
Frustrated Magnets and Chern Insulators

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IXS 2019, Stonybrook, 24 June 2019



UNIVERSITY OF
TORONTO



NSERC
CRSNG

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Oakridge National Lab and Ohio State Univ

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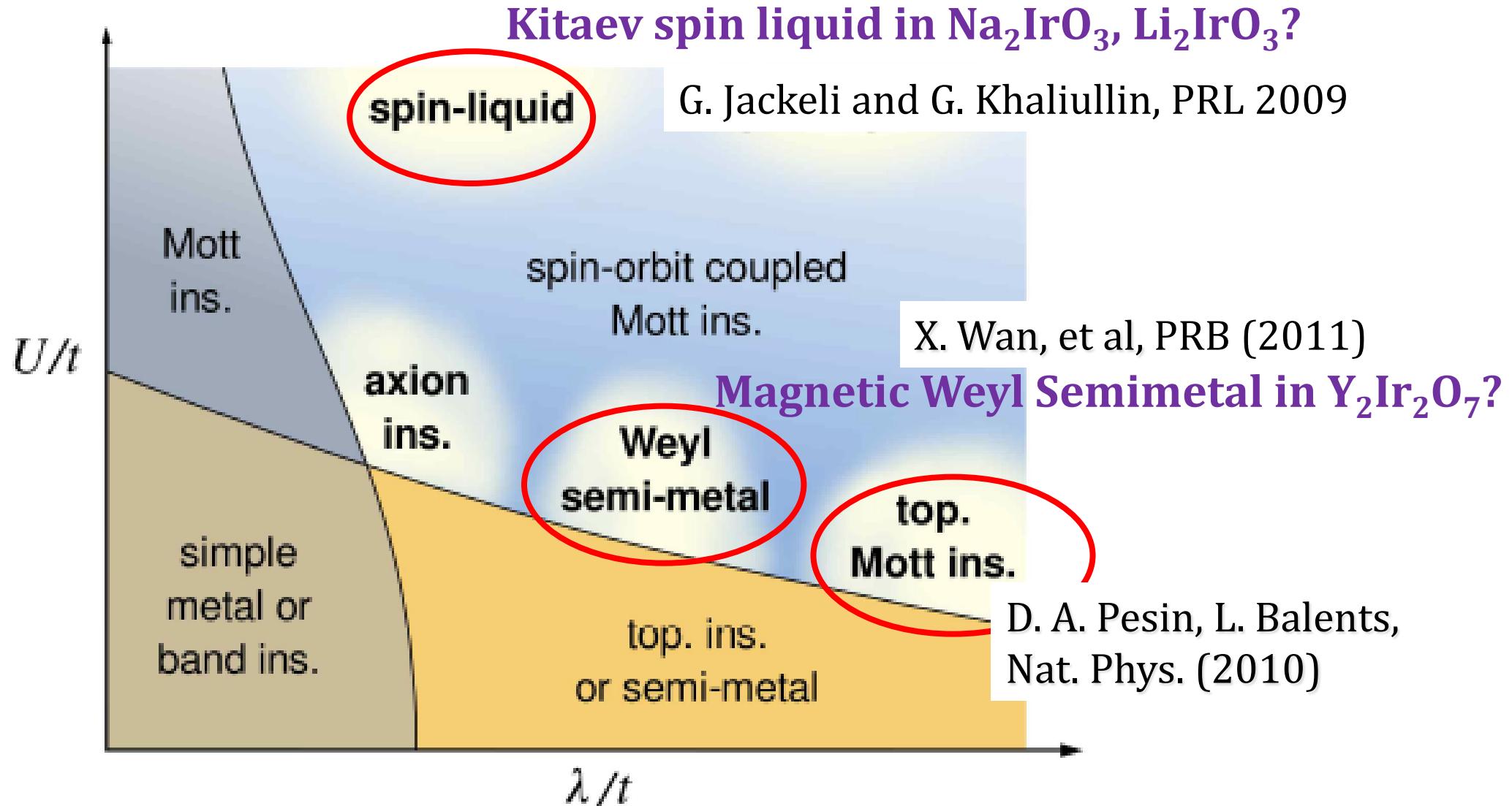
India

T. Saha-Dasgupta (DFT), A. Mukerjee,
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University of Cologne, IFW-Dresden, Stockholm

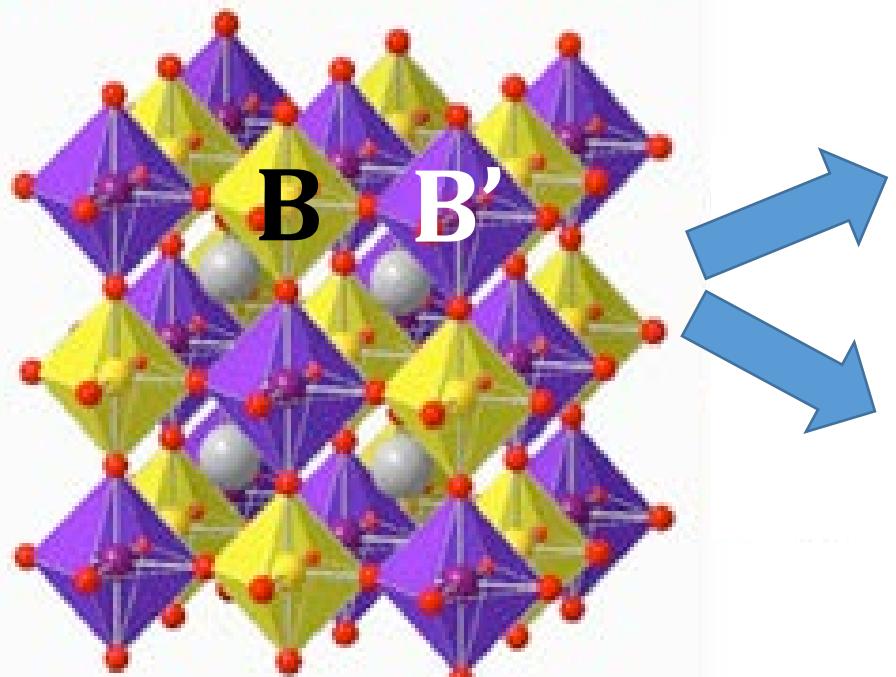
M. Gruninger, S. Trebst, T. Lorenz, D. Kiese, F. Lasse-Buessen, S. Streltsov
J. Attig, D. Khomskii, M. Revelli, M. Hermanns, J. van den Brink, et al

Interactions vs Spin-Orbit coupling



Ordered Double Perovskites

Double perovskite lattice



Mott insulators: Frustrated FCC Magnets

$Ba_2Y\text{Mo}O_6$, $Sr_2Y\text{Re}O_6$, $Ba_2Y\text{Ir}O_6$, $Ba_2Y\text{Ru}O_6$, ...

B sublattice: Inert closed shell atom

B' sublattice: Magnetic ion with SOC

T. Aharen et al, Phys. Rev. B 81, 224409 (2010)

J. P. Carlo, et al, Phys. Rev. B 88, 024418 (2013)

G. Cao, et al, Phys. Rev. B 87, 155136 (2013)

G. Cao, et al, Phys. Rev. Lett. 112, 056402 (2014)

High Tc Ferromagnetic Metals

$Sr_2\text{FeMo}O_6$, $Ba_2\text{FeRe}O_6$, $Pb_2\text{CrRe}O_6$, ...

B sublattice: 3d ion with strong magnetism

B' sublattice: 4d/5d ion with strong SOC

K.-I. Kobayashi, et al, Nature (1998) – Tokura group

W. Prellier, et al, J. Phys. Cond Matter (2000) – R. Greene group

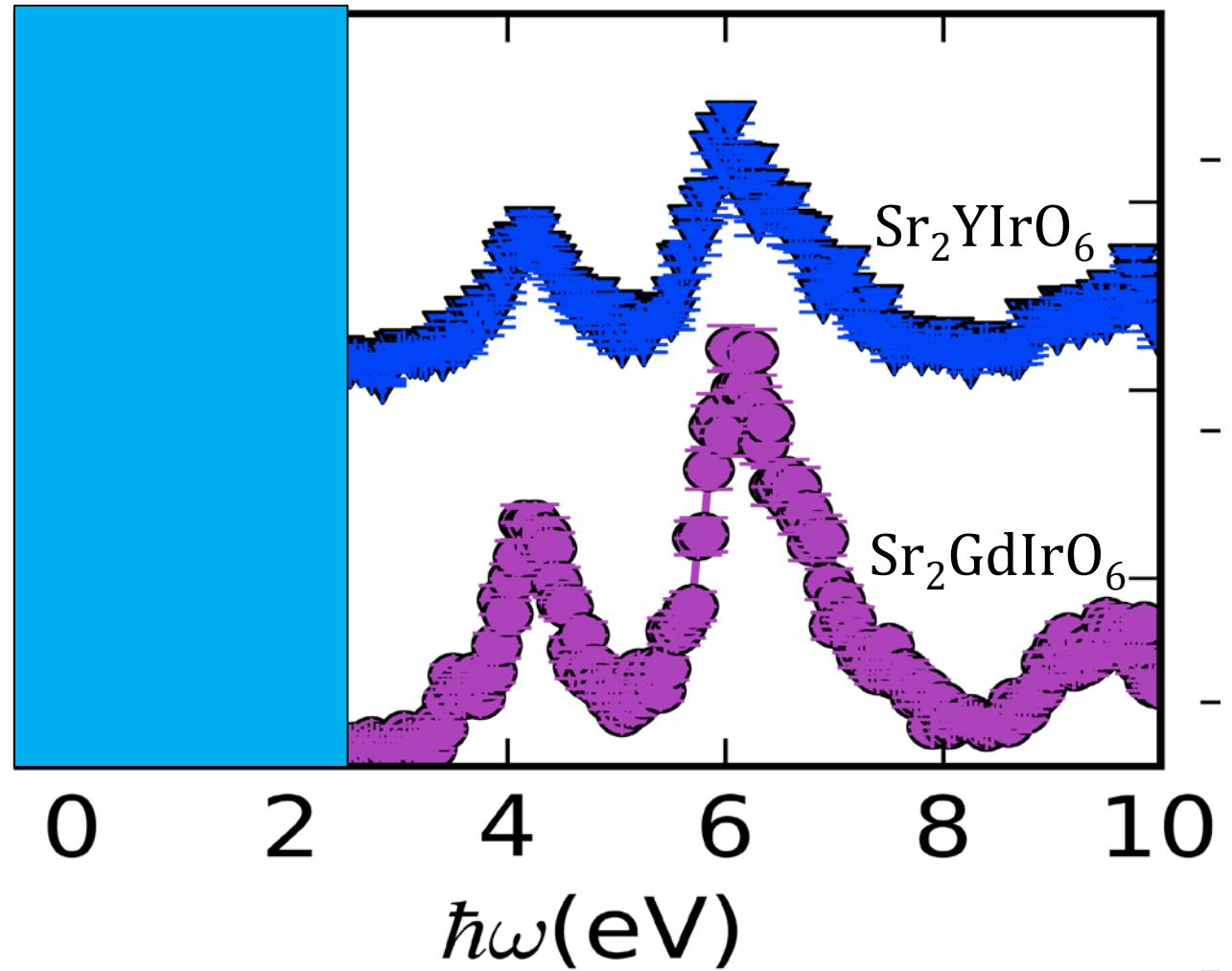
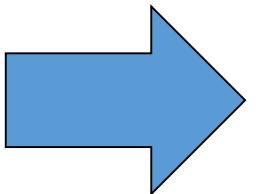
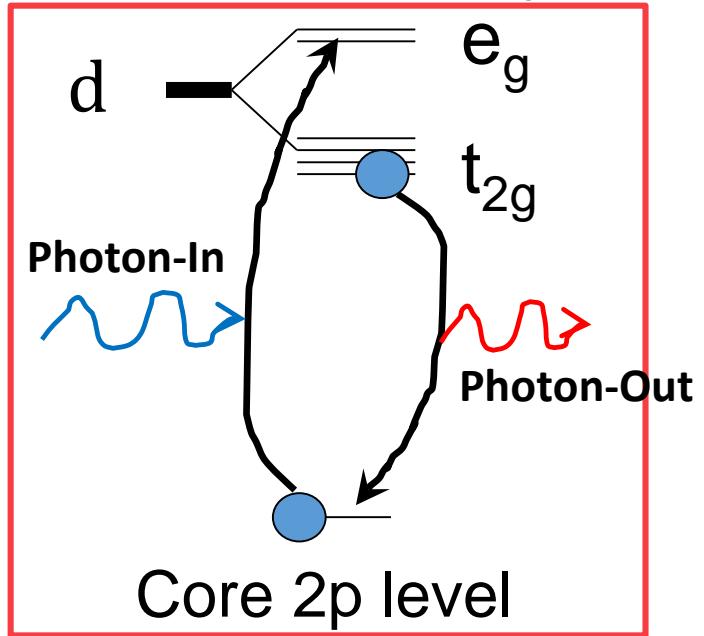
N. Trivedi, M. Randeria, P. Woodward groups - OSU

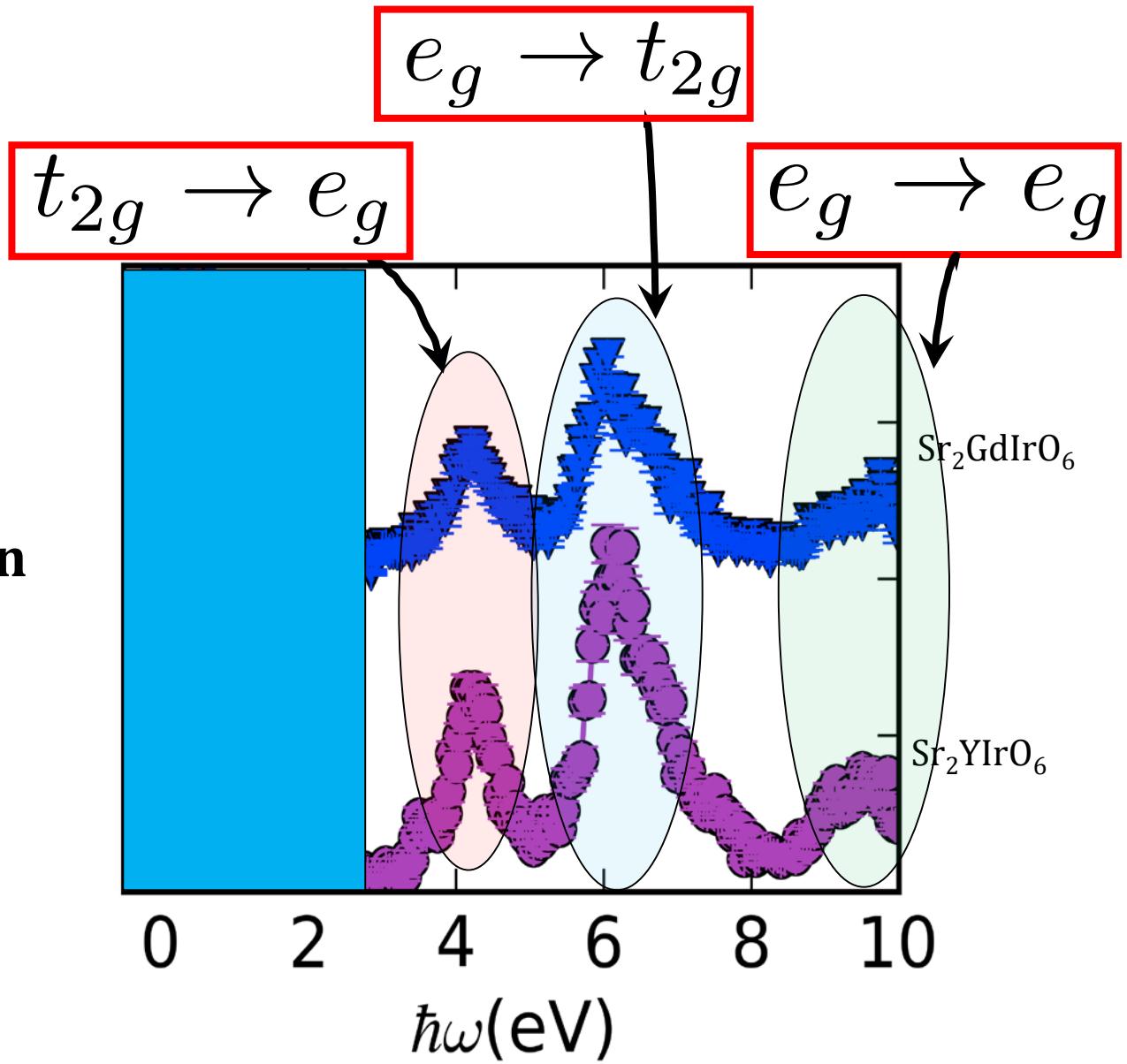
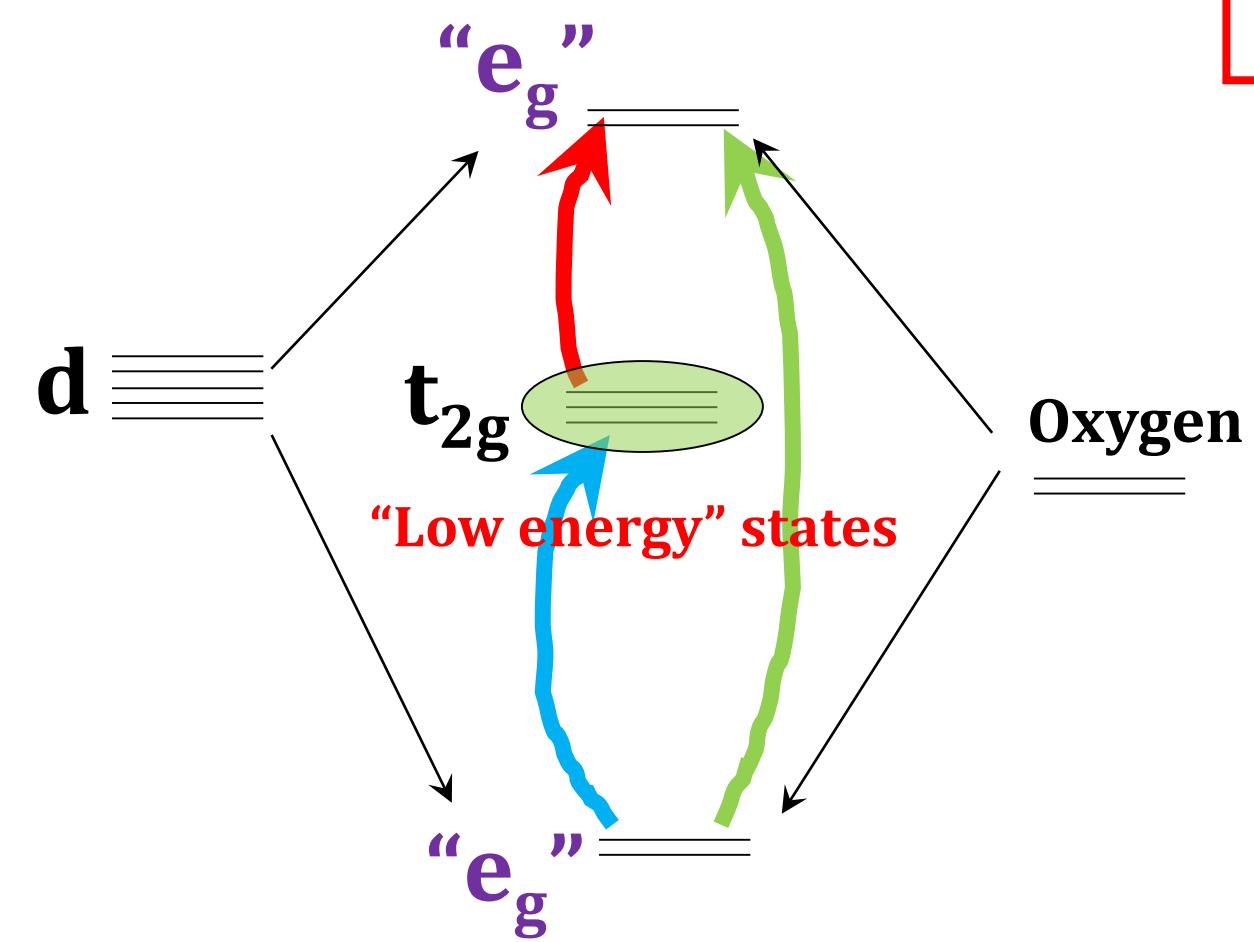
Outline

- **Lego Block: $B' O_6$ Octahedra**
- **Deconstructing the Lego Block: RIXS**
- **Lego Block Assembly I: Frustrated Mott Insulators**
- **Lego Block Assembly II: High Tc Chern Insulators**

“Lego Block”: Crystal Field Levels

Resonant Inelastic X-ray Scattering



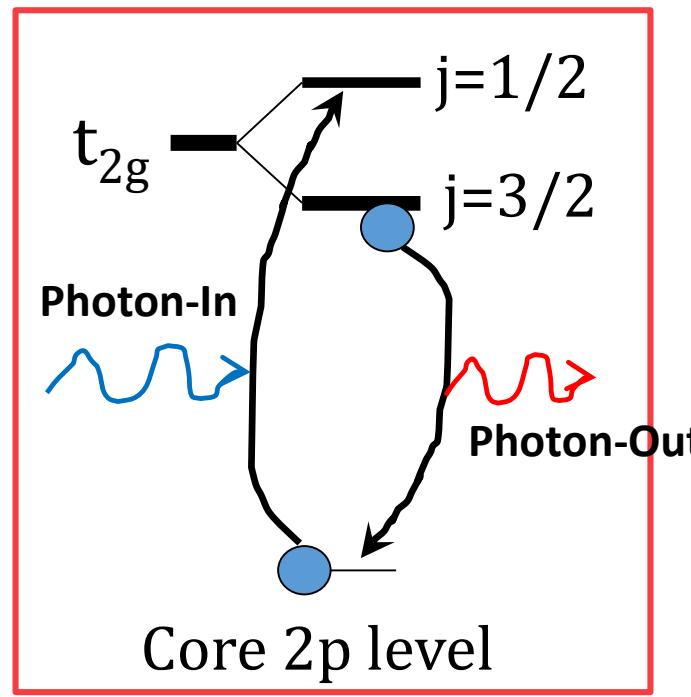


“Lego Block”: Spin-Orbit Coupling

Spin orbit coupling and interactions

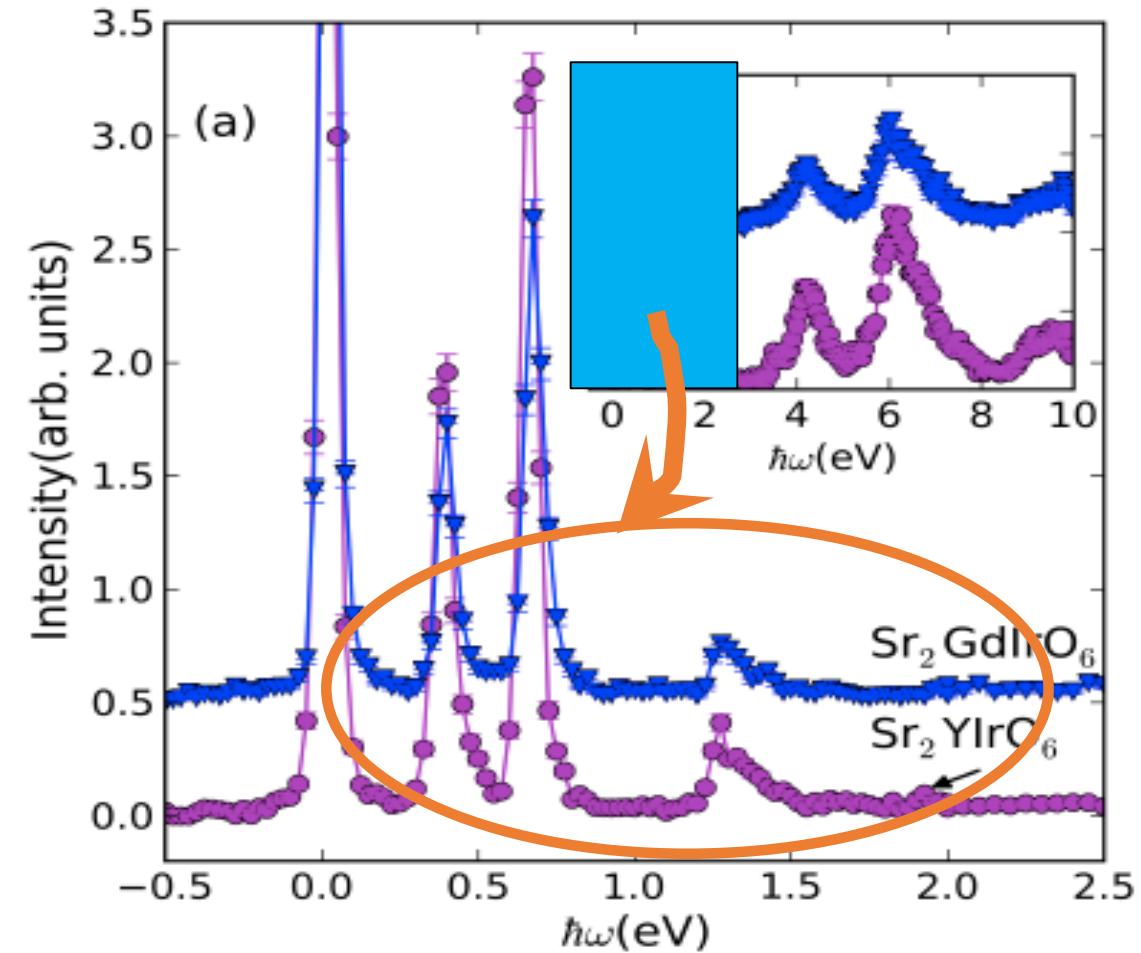
$$P_{t_{2g}} \vec{L} P_{t_{2g}} = -\vec{\ell} \quad (\ell = 1)$$

$$H_{\text{s.o.}} = -\lambda \vec{\ell} \cdot \vec{s}$$



Naive: Single peak at $E = 3 \lambda/2$
Works perfectly for d^5 iridates

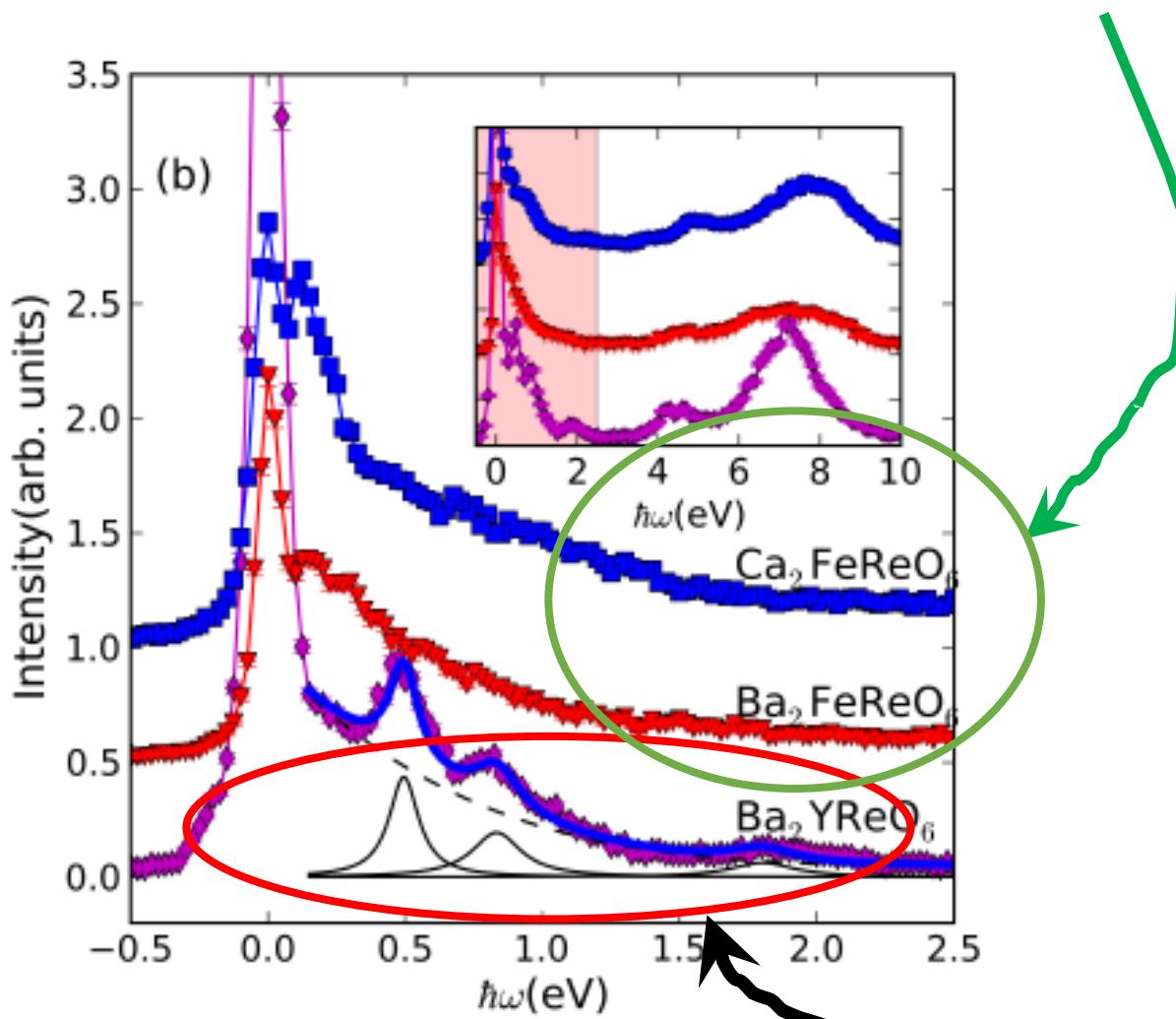
d^4 Iridates: Much richer structure



B. Yuan, et al, AP, Y. J. Kim, PRB (2017)

d² Rhenates

Metal/Semimetal
Nothing much visible

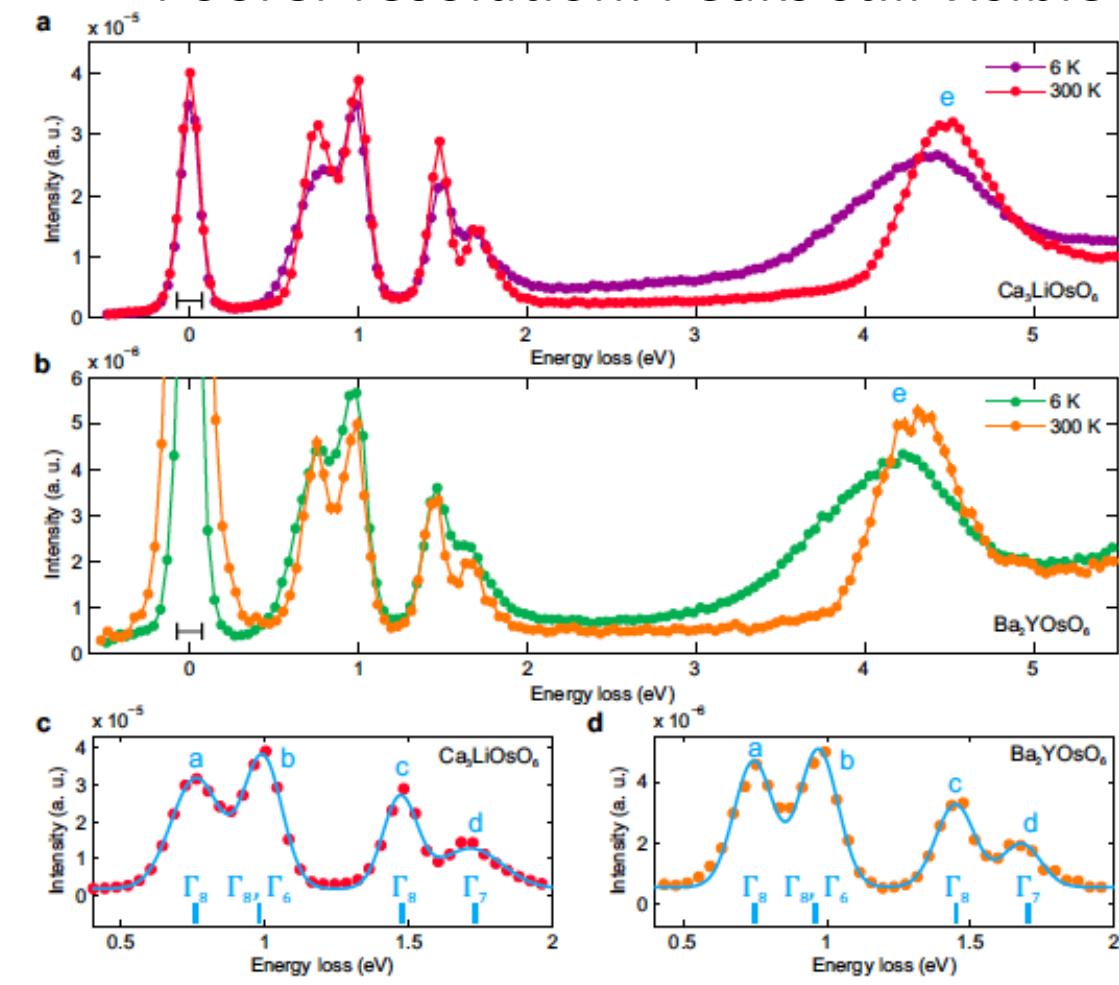


Bo Yuan, et al, PRB 2017

Insulators: Peaks

d³ Osmates

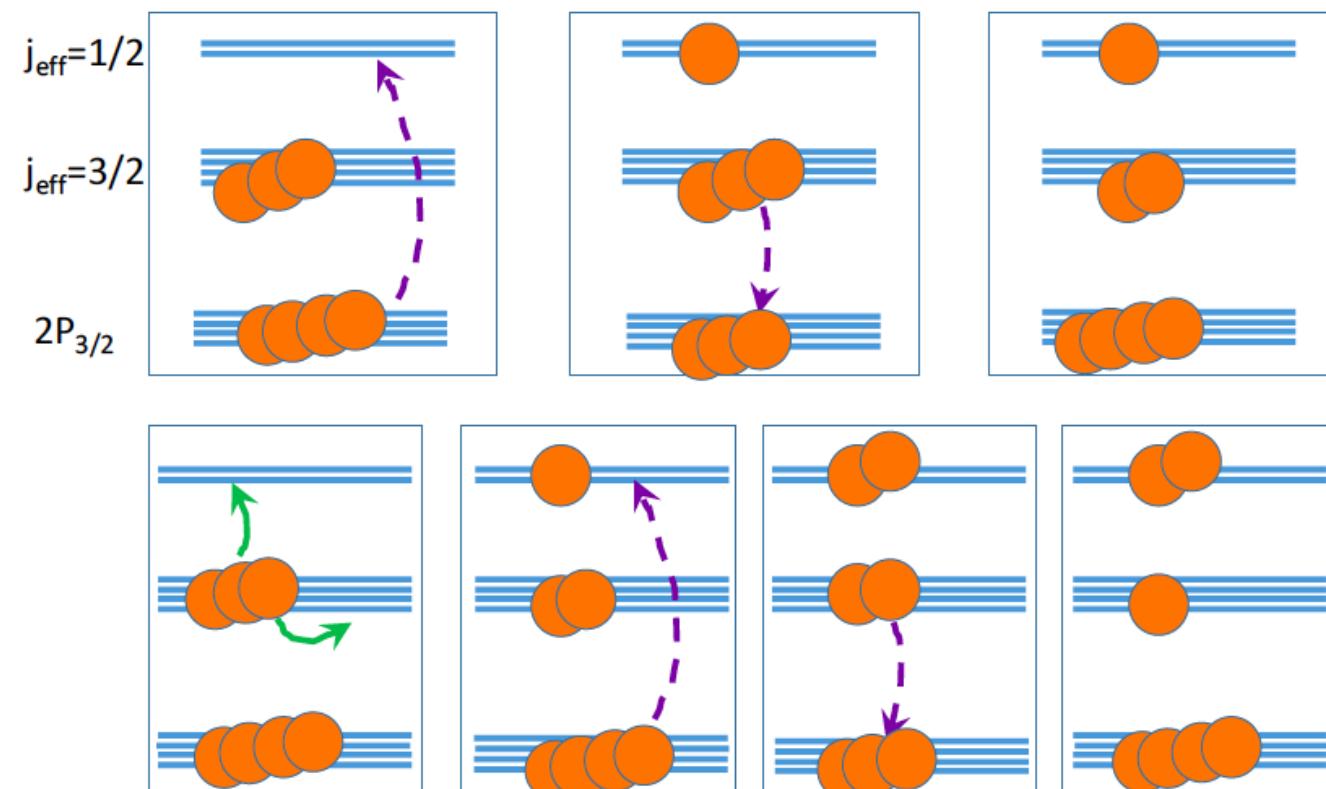
Poorer resolution: Peaks still visible



ORNL: A Taylor, et al, PRL 2017)

RIXS: Simple Picture

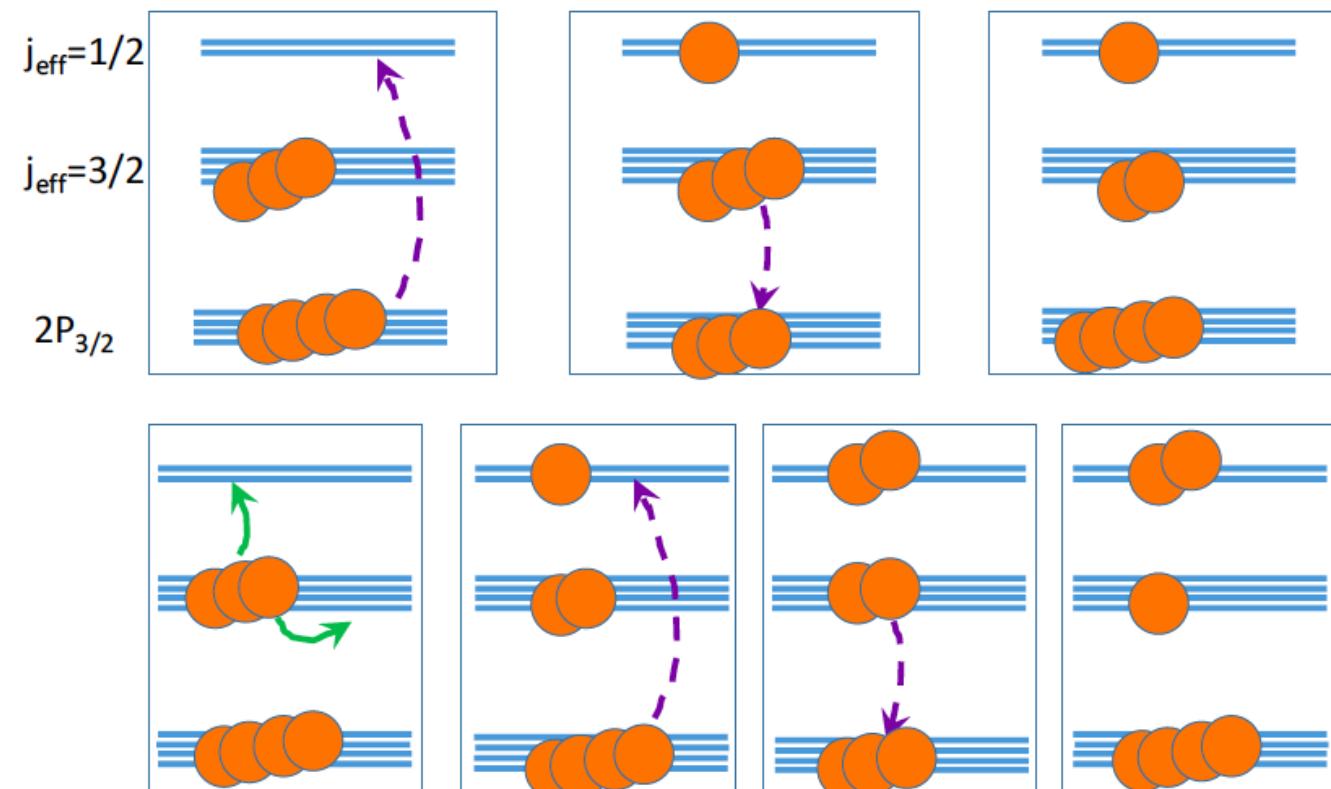
$$H_{\text{int}} = \frac{U}{2} : (\sum_{\ell} n_{\ell})^2 : -5 \frac{J_H}{2} \sum_{\ell < \ell'} n_{\ell} n_{\ell'} \\ - 2 J_H \sum_{\ell < \ell'} \vec{S}_{\ell} \cdot \vec{S}_{\ell'} + J_H \sum_{\ell \neq \ell'} d_{\ell \uparrow}^{\dagger} d_{\ell \downarrow}^{\dagger} d_{\ell' \downarrow} d_{\ell' \uparrow}$$



RIXS: Simple Picture

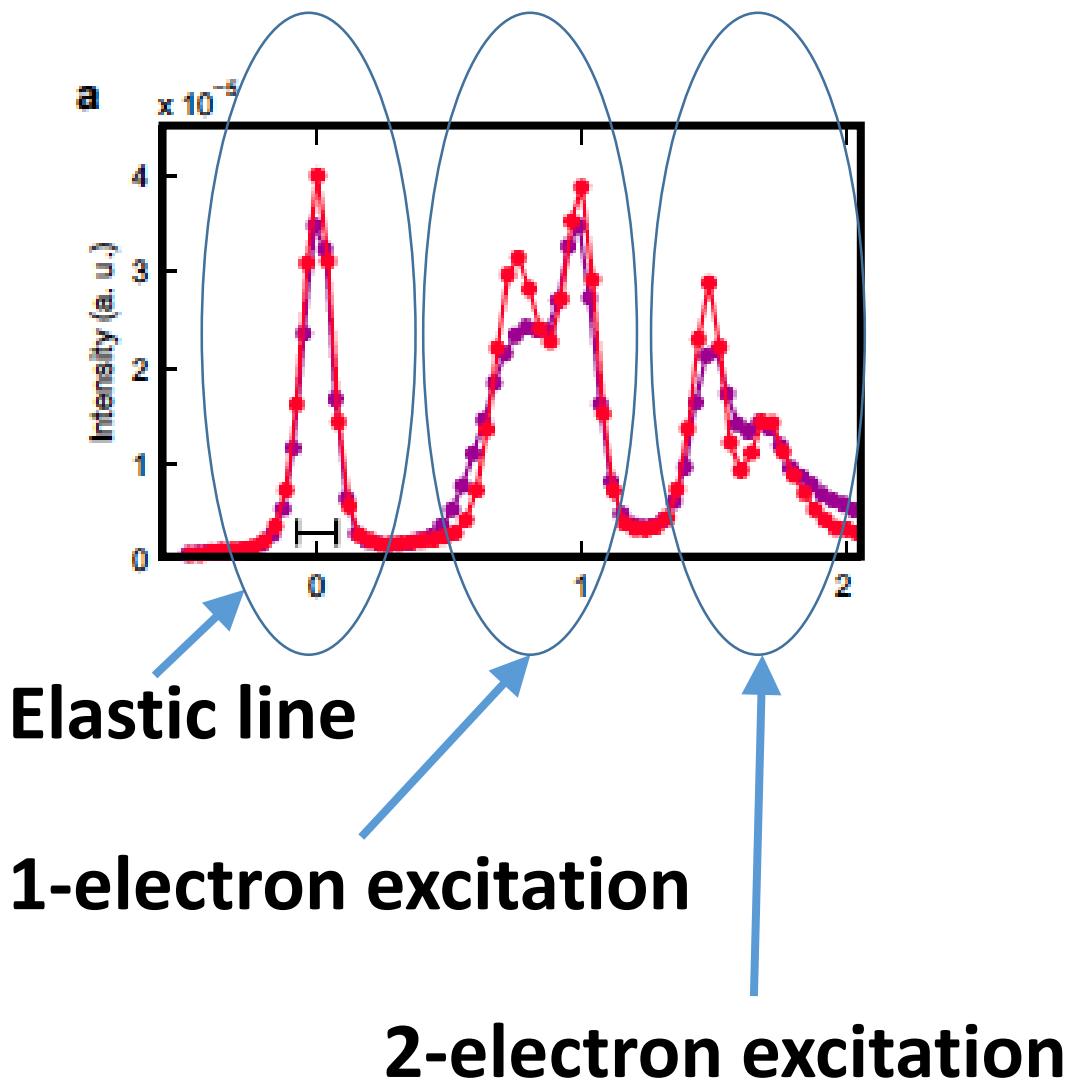
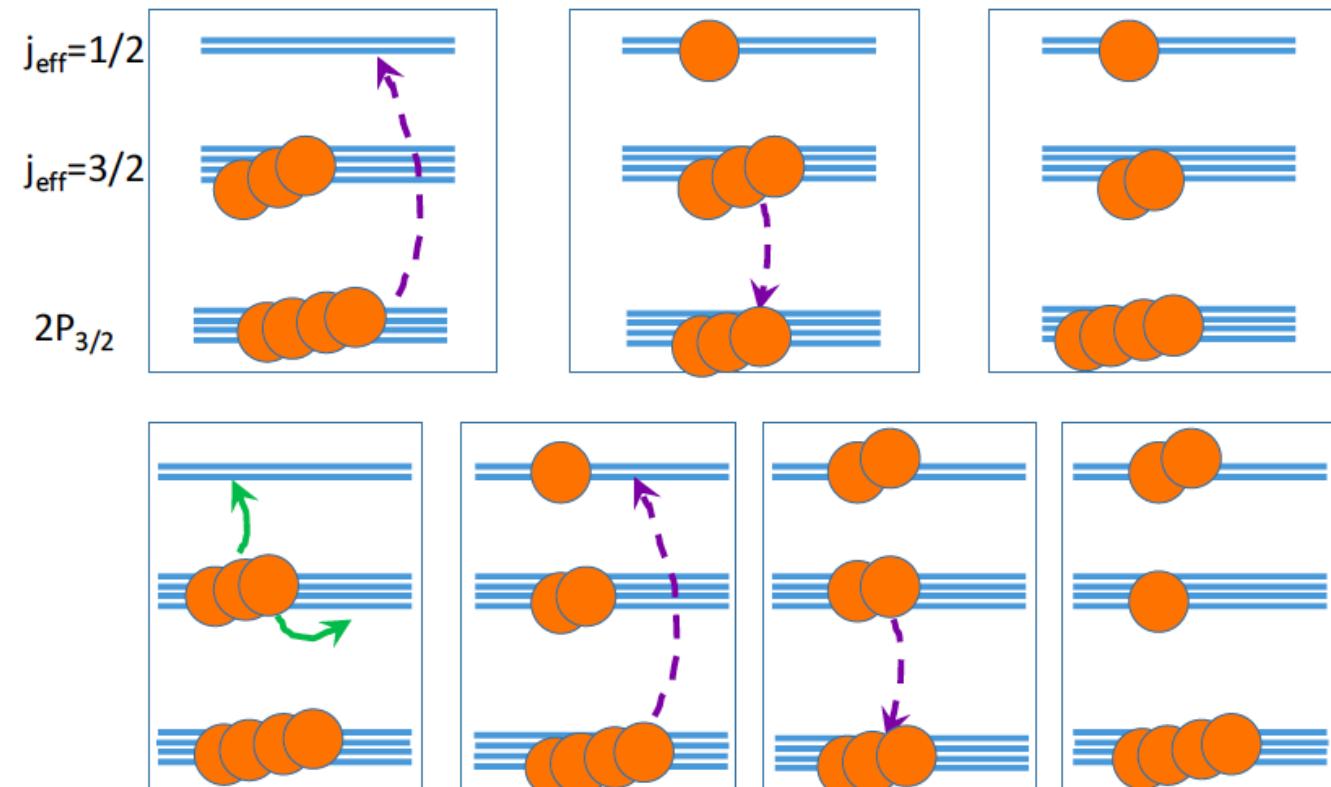
$$H_{\text{int}} = \frac{U}{2} : \left(\sum_{\ell} n_{\ell} \right)^2 : -5 \frac{J_H}{2} \sum_{\ell < \ell'} n_{\ell} n_{\ell'}$$

$$- 2 J_H \sum_{\ell < \ell'} \vec{S}_{\ell} \cdot \vec{S}_{\ell'} + J_H \sum_{\ell \neq \ell'} d_{\ell \uparrow}^{\dagger} d_{\ell \downarrow}^{\dagger} d_{\ell' \downarrow} d_{\ell' \uparrow}$$



RIXS: Simple Picture

$$H_{\text{int}} = \frac{U}{2} : \left(\sum_{\ell} n_{\ell} \right)^2 : -5 \frac{J_H}{2} \sum_{\ell < \ell'} n_{\ell} n_{\ell'} - 2 J_H \sum_{\ell < \ell'} \vec{S}_{\ell} \cdot \vec{S}_{\ell'} + J_H \sum_{\ell \neq \ell'} d_{\ell \uparrow}^{\dagger} d_{\ell \downarrow}^{\dagger} d_{\ell' \downarrow} d_{\ell' \uparrow}$$



RIXS Theory: Simplifying Assumptions

$$\frac{d^2\sigma}{d\Omega dE_i} \approx \left| \frac{1}{E_g - \bar{E}_n + E_i + i\frac{\bar{\Gamma}_n}{2}} \right|^2$$

$$\sum_f \left| \sum_n \langle f | T^\dagger | n \rangle \langle n | T | g \rangle \right|^2 \delta(E_g - E_f + E_i - E_o)$$

Average out denominator;
short core-hole lifetime

Focus on matrix elements, energies
Ignore core hole interactions

Dipole Approximation, Ignore Core-Hole interaction

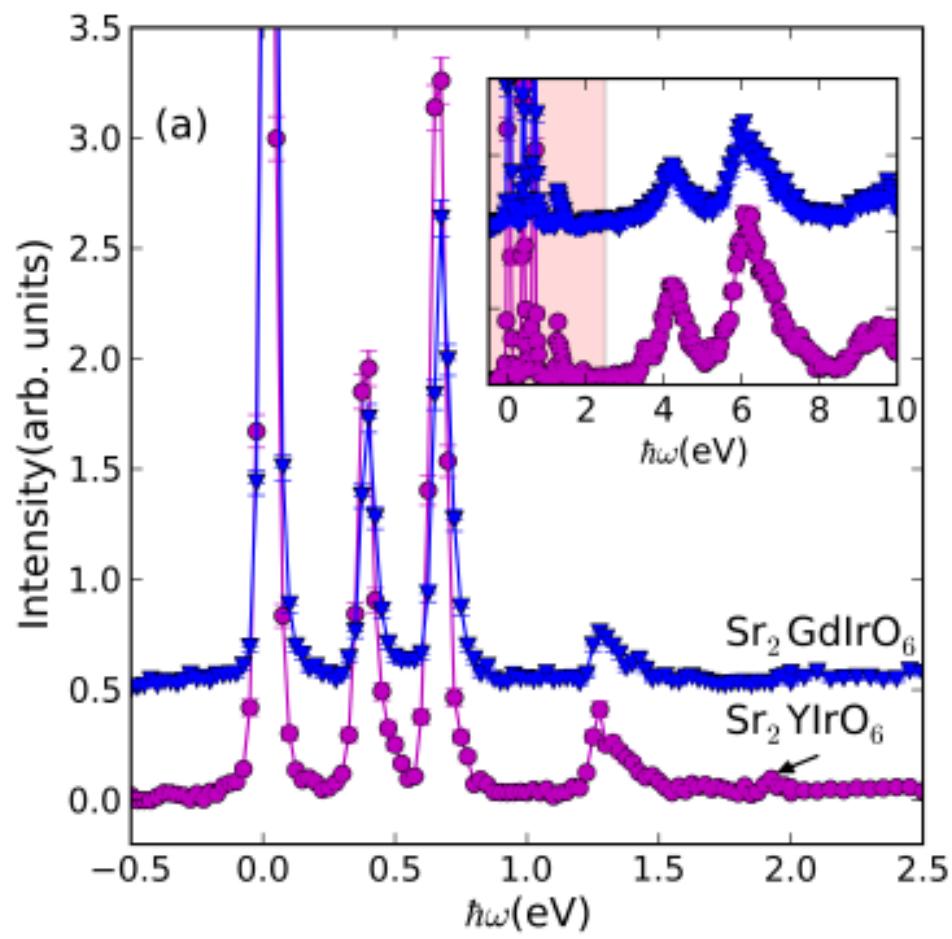
$$\langle n | T | g \rangle = \hat{\epsilon}_{in}^\alpha \langle n | p_{\beta\sigma}^\dagger d_{\alpha\beta\sigma}^\dagger | g \rangle$$

$$\langle f | T^\dagger | n \rangle = \hat{\epsilon}_{out}^\mu \langle f | d_{\mu\nu\sigma'}^\dagger p_{\nu\sigma'} | n \rangle$$

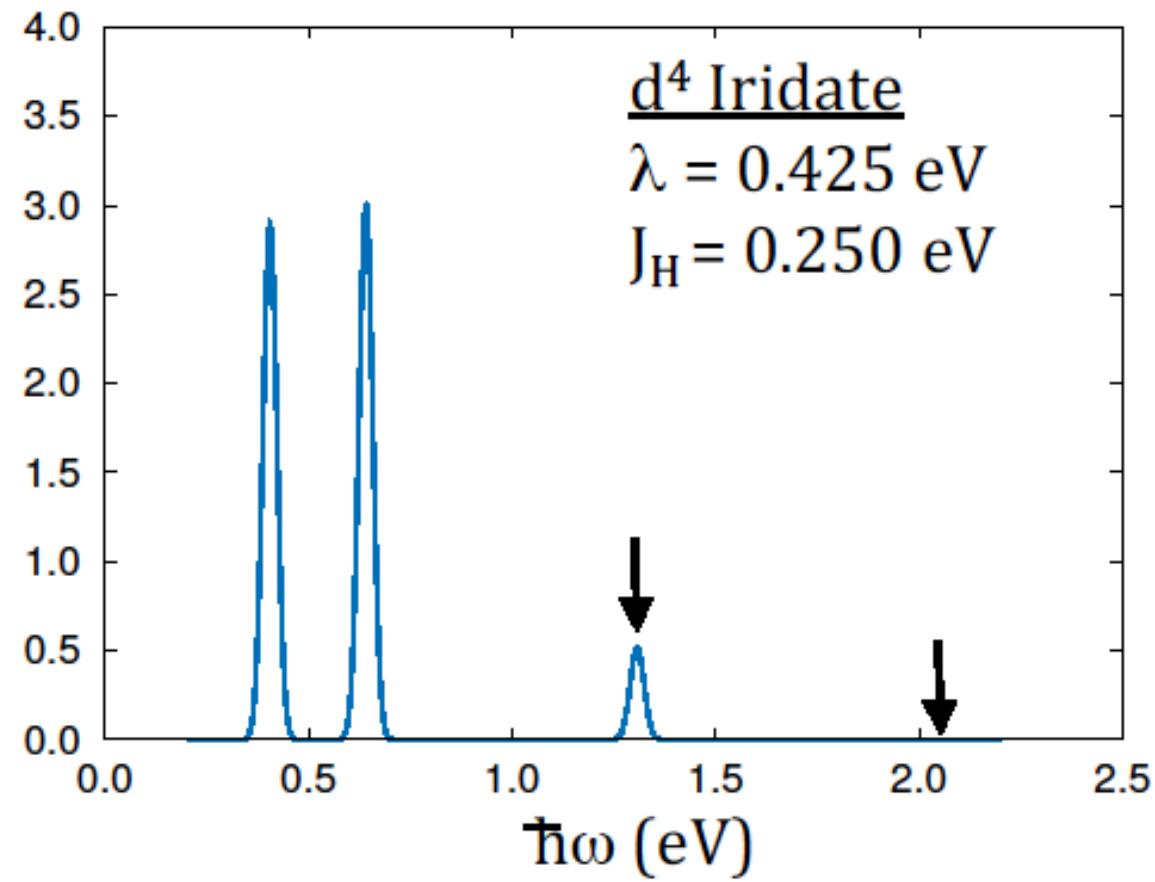
$$H_{int} = \frac{U}{2} : (\sum_\ell n_\ell)^2 : -5 \frac{J_H}{2} \sum_{\ell < \ell'} n_\ell n_{\ell'} \\ - 2 J_H \sum_{\ell < \ell'} \vec{S}_\ell \cdot \vec{S}_{\ell'} + J_H \sum_{\ell \neq \ell'} d_{\ell\uparrow}^\dagger d_{\ell\downarrow}^\dagger d_{\ell'\downarrow} d_{\ell'\uparrow}$$

Exact diagonalization results

d⁴ Iridates

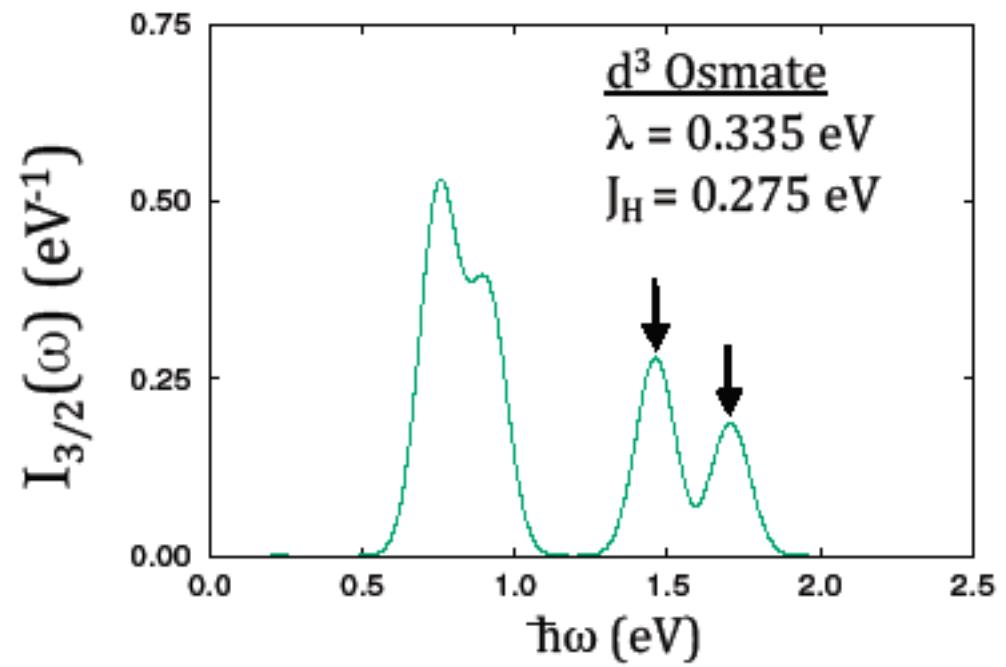
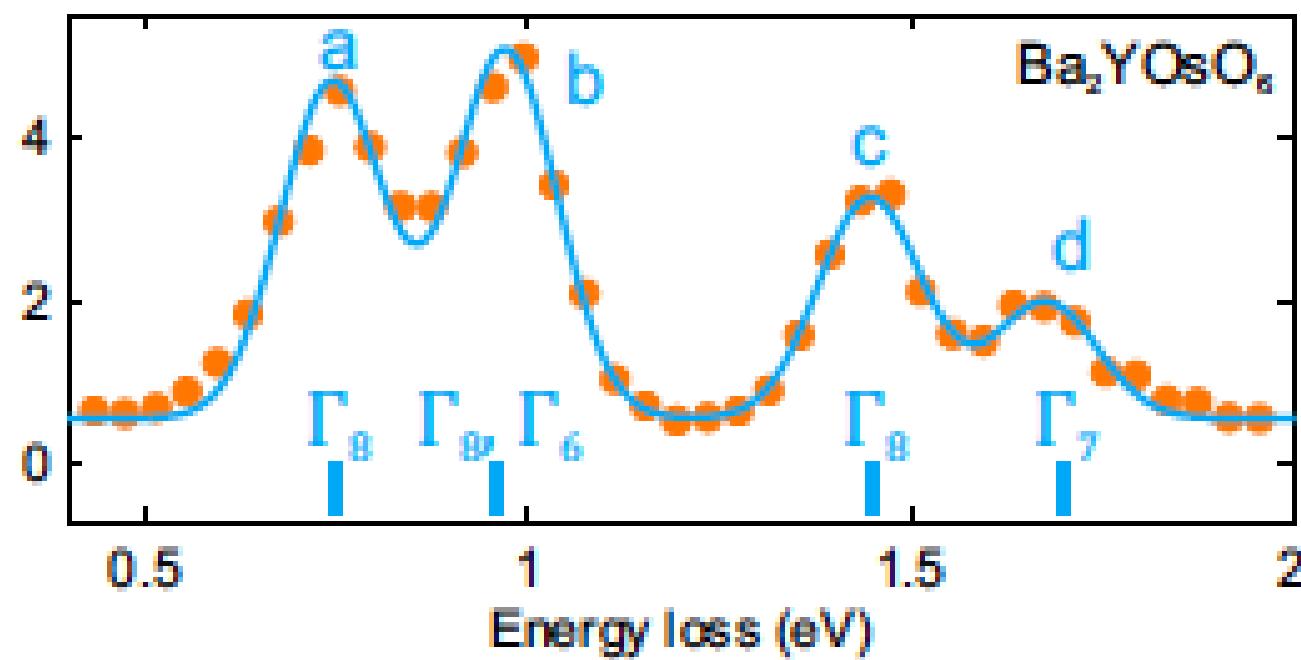


$$I_{3/2}(\omega) \text{ (eV}^{-1}\text{)}$$



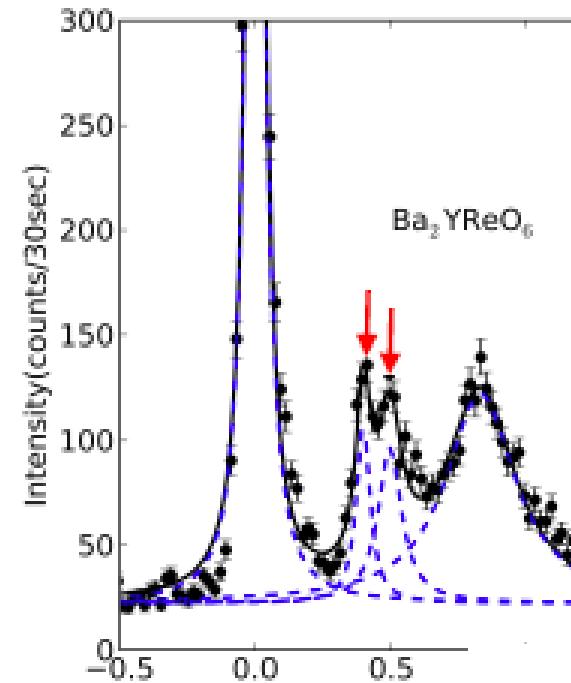
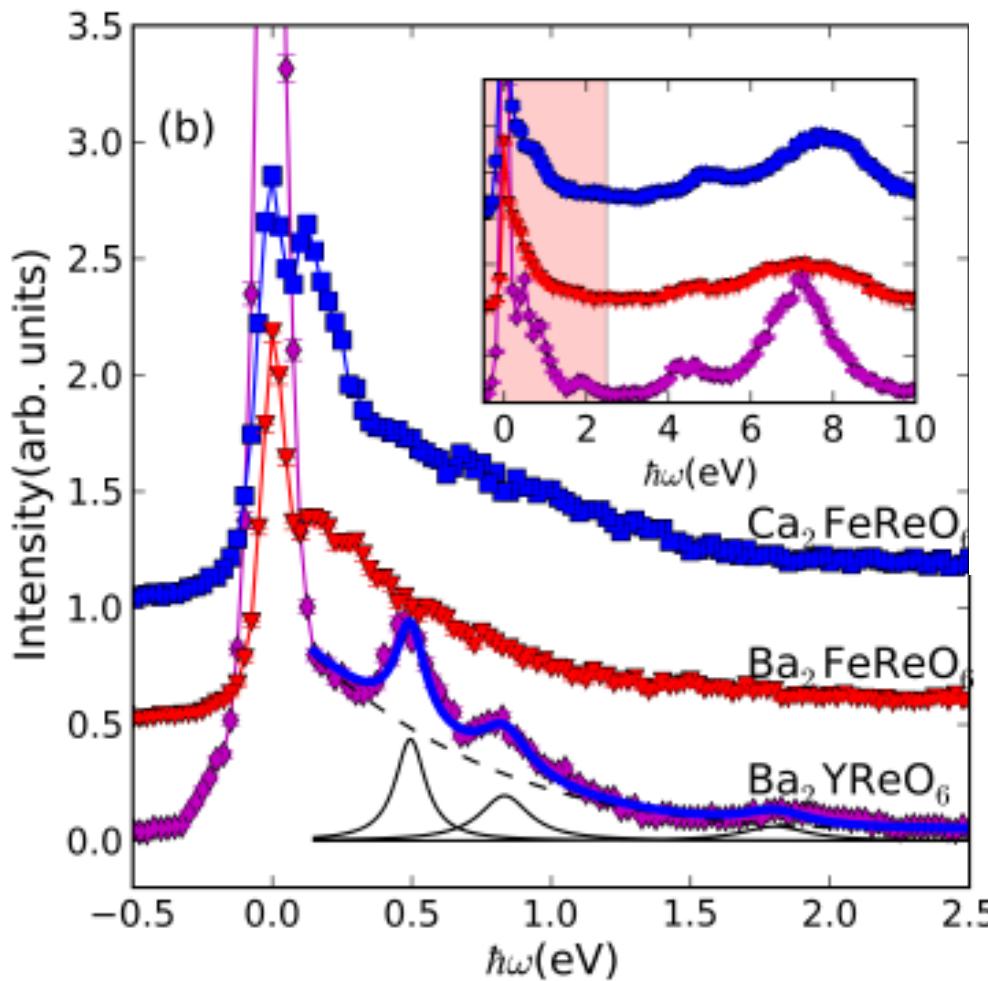
Exact diagonalization results

d³ Osmates

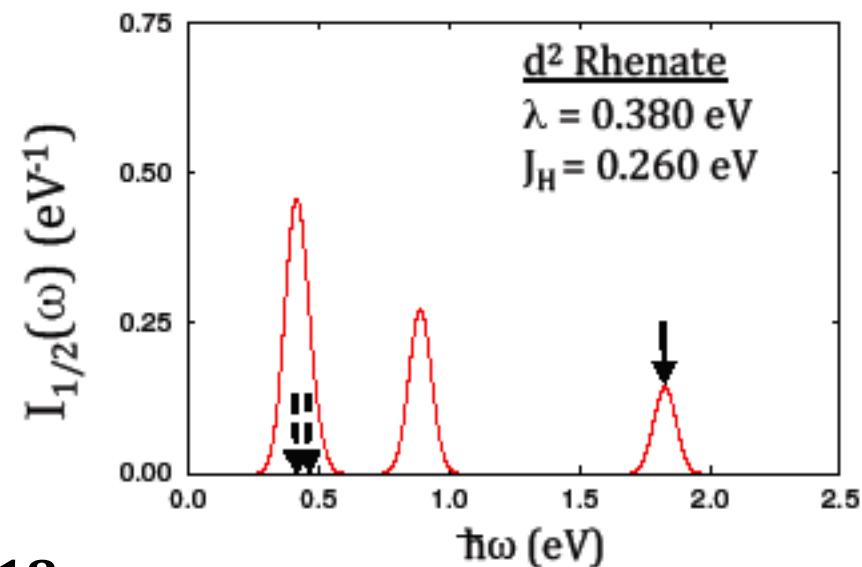


Exact diagonalization results

d² Rhenates

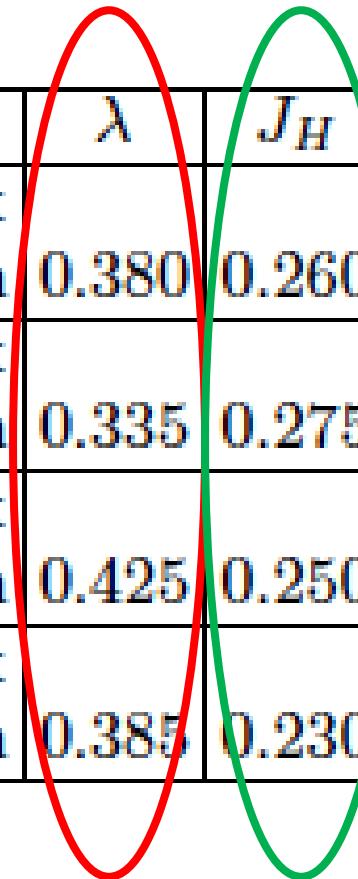


d² configuration has **only one** 2-electron excited state



“Lego Block” Hamiltonian

$$H = H_{\text{SOC}}[\lambda] + H_{\text{int}}$$



Material		λ	J_H	Peak 1	Peak 2	Peak 3	Peak 4
Ba_2YReO_6 (ref.55, this)	Ex			[0.40]	[0.50]	0.83	1.85
	Th	0.380	0.260	0.41	0.47	0.89	1.83
$\text{Ba}_2\text{YO}_{\text{Os}}\text{O}_6$ (ref.54)	Ex			0.745	0.971	1.447	1.68
	Th	0.335	0.275	0.75	0.91	1.46	1.71
Sr_2YIrO_6 (ref.55)	Ex			0.39	0.66	1.30	~ 2
	Th	0.425	0.250	0.41	0.64	1.31	2.06
Ba_2YIrO_6 (ref.56)	Ex			0.35	0.60	1.18	-
	Th	0.385	0.230	0.37	0.58	1.19	1.88

d³ (Osmates)

SOC still important!

d⁴ (Ba_2YIrO_6)

Magnetism?

Impurities?

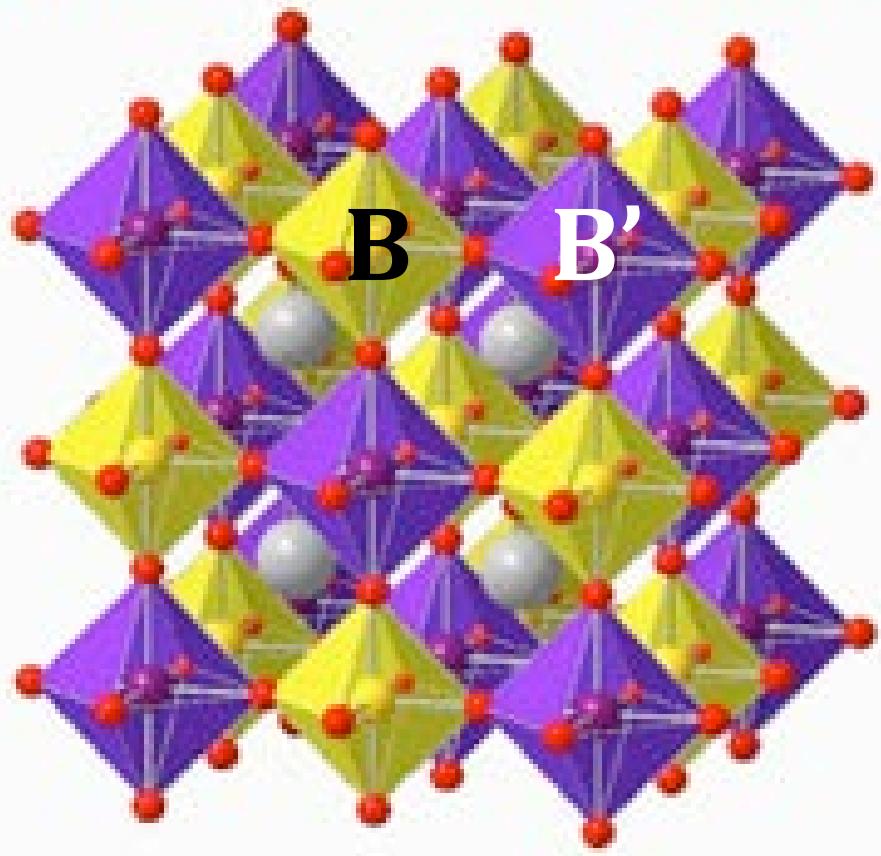
Cluster Excitons?

Material Variations

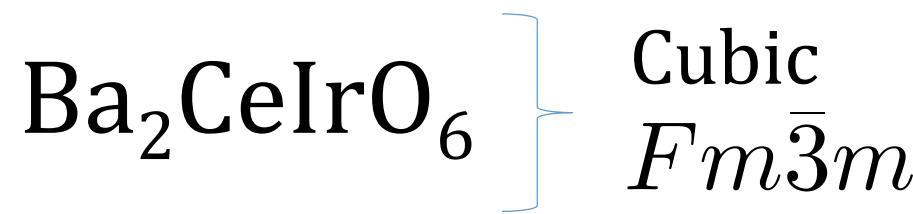
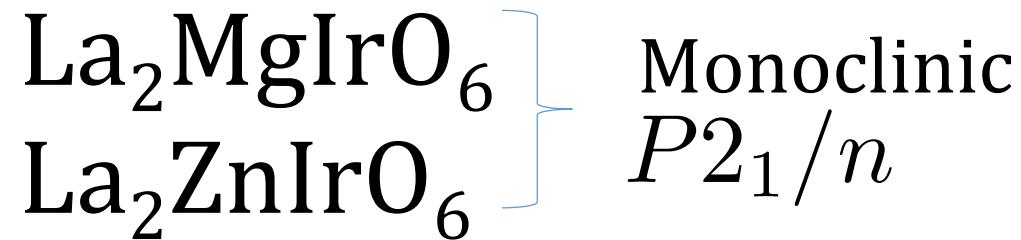
- Different Atomic Numbers
- Different Degrees of Oxygen Hybridization

Coupling the “Lego Blocks” Mott Insulators on the Frustrated FCC Lattice

Double perovskite lattice



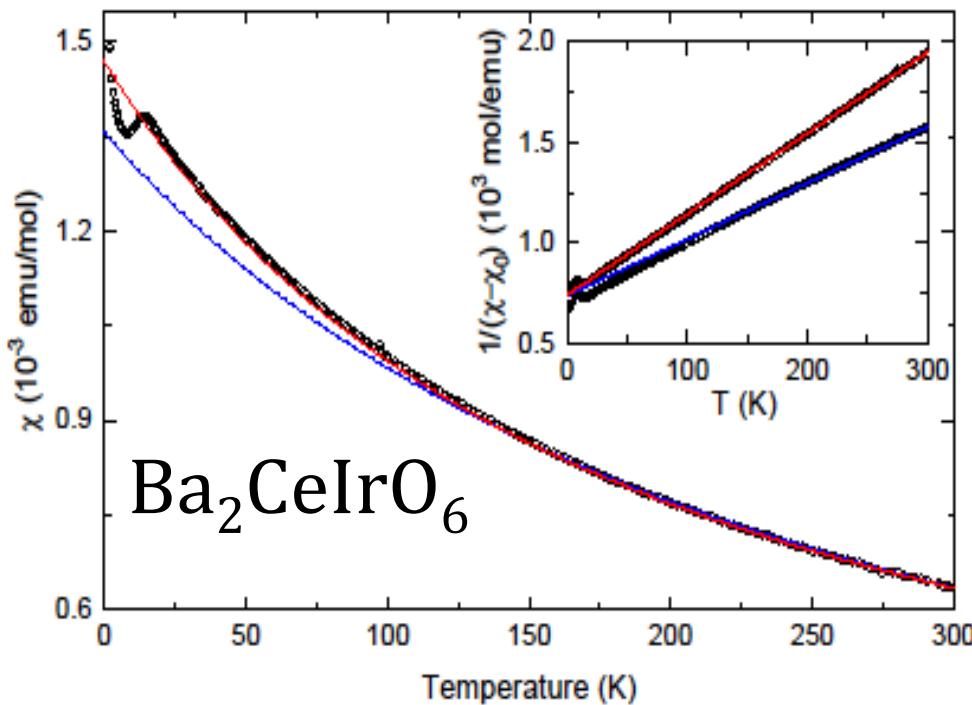
d⁵ iridates



Other fillings:

Chen, Pereira, Balents (2010); Chen, Balents (2011);
Svoboda, Randeria, Trivedi (2017, 2018); AP et al (2018)

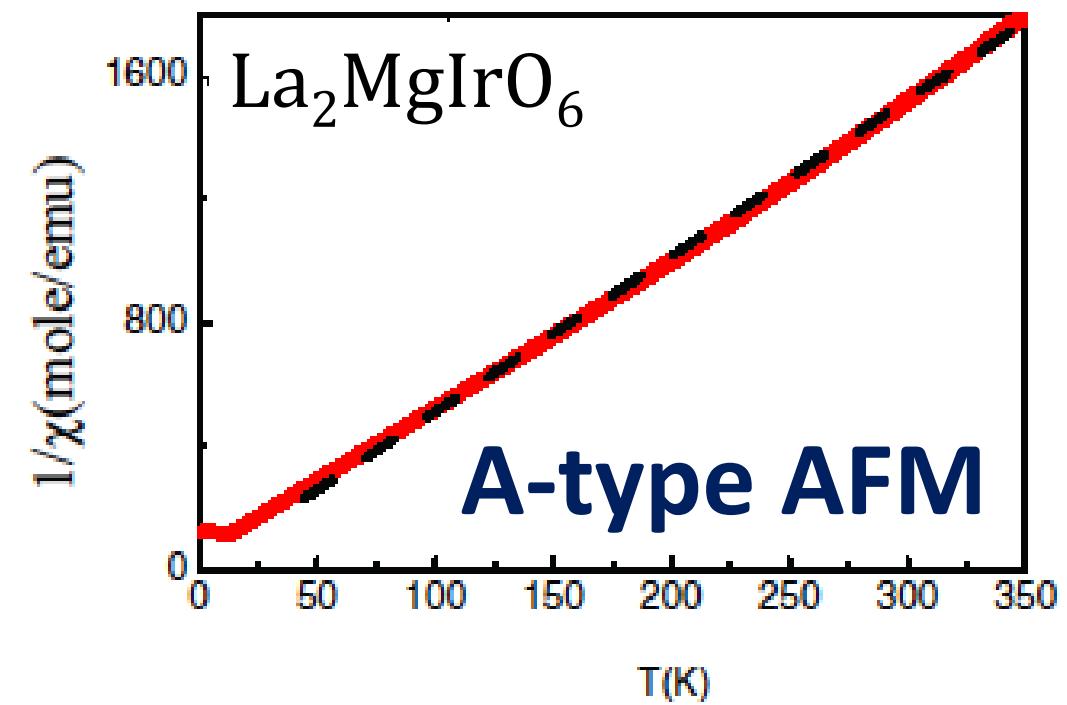
$j=1/2$ Mott Insulators on the FCC Lattice



$$\Theta_{\text{CW}} \approx -200K \quad T_N \approx 14K$$

A.Revelli, et al, arXiv:1901.06215 (Cologne)

A.Aczel, et al, arXiv:1901.08146 (ORNL/UTK); A-type AFM



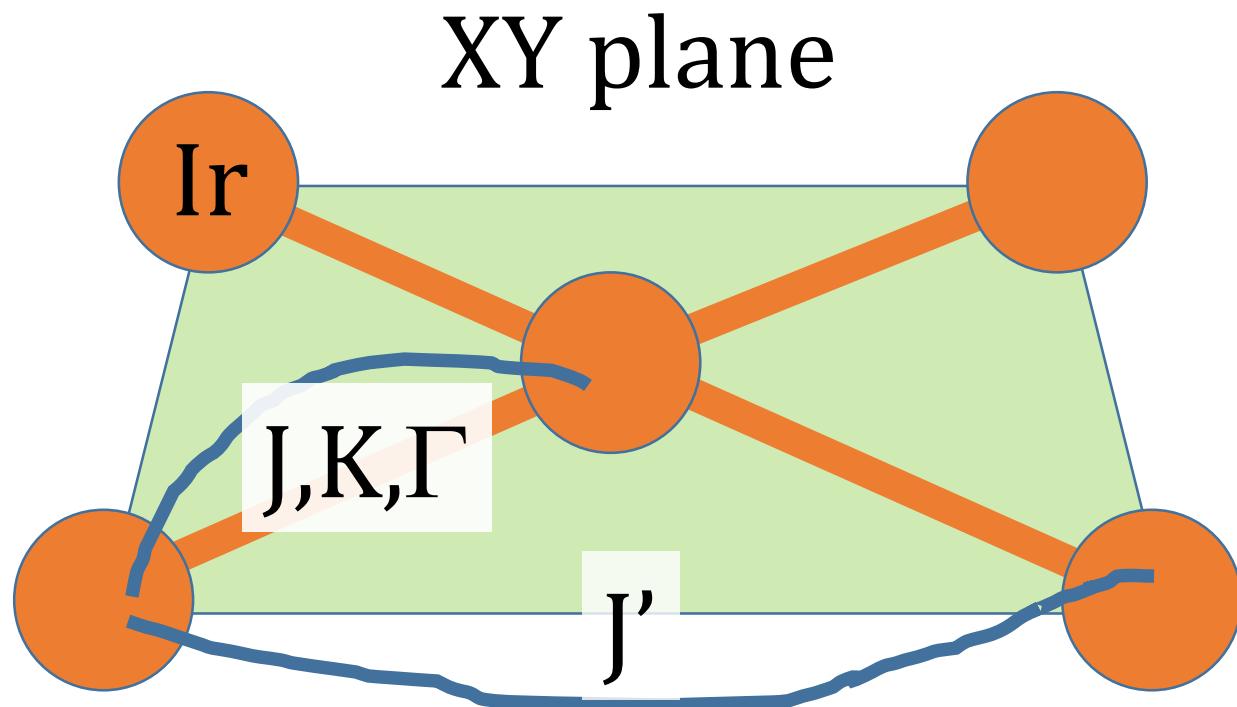
$$\Theta_{\text{CW}} \approx -24K \quad T_N \approx 12K$$

Guixin Cao, et al, PRB (2013) (ORNL/UTK)

A.Aczel, et al, PRB (2017); A.Cook, et al, PRB (2017)

- Strong Frustration for ideal FCC magnet
- Highly Tuned by Weak Distortion: Magnetoelastic

$j=1/2$ Mott Insulators on the Ideal FCC Lattice Theory: Exchange Couplings



Perturbative + DFT

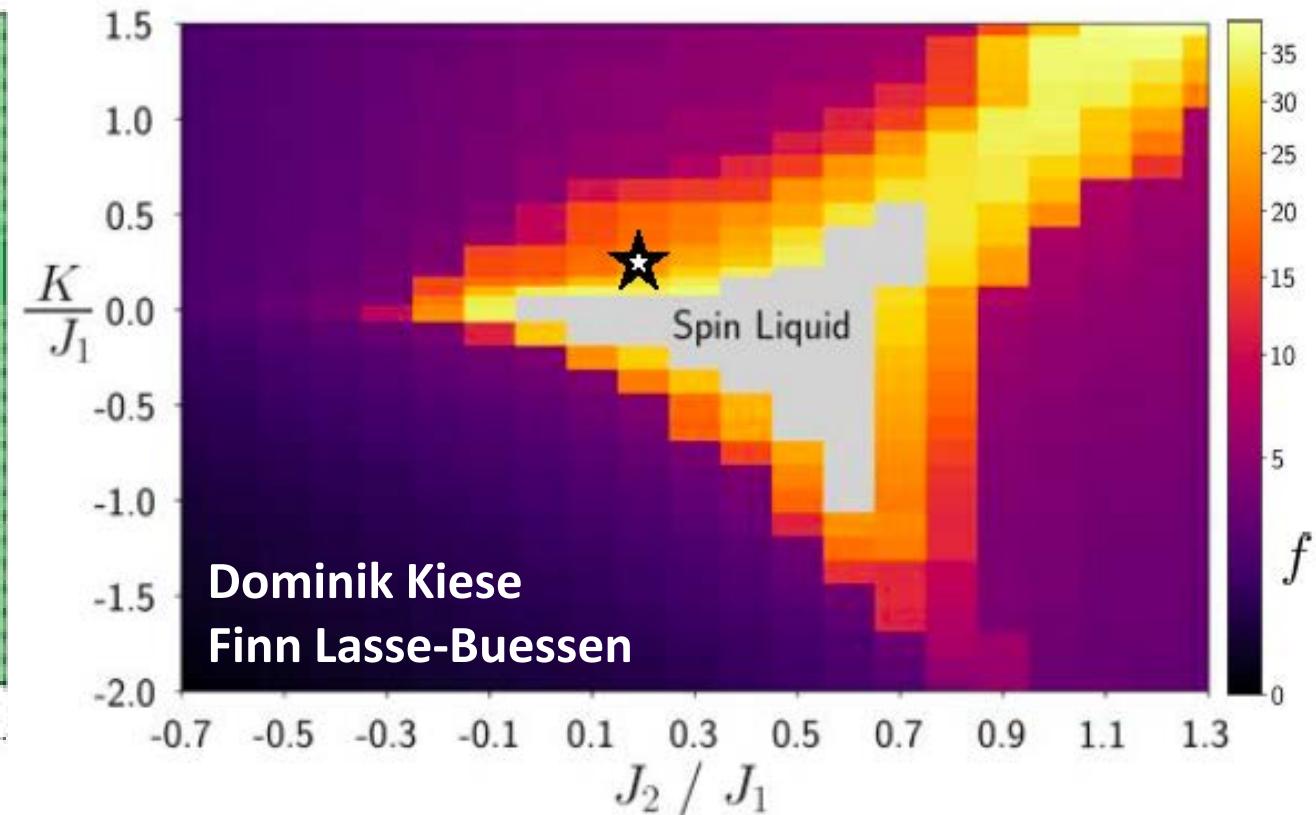
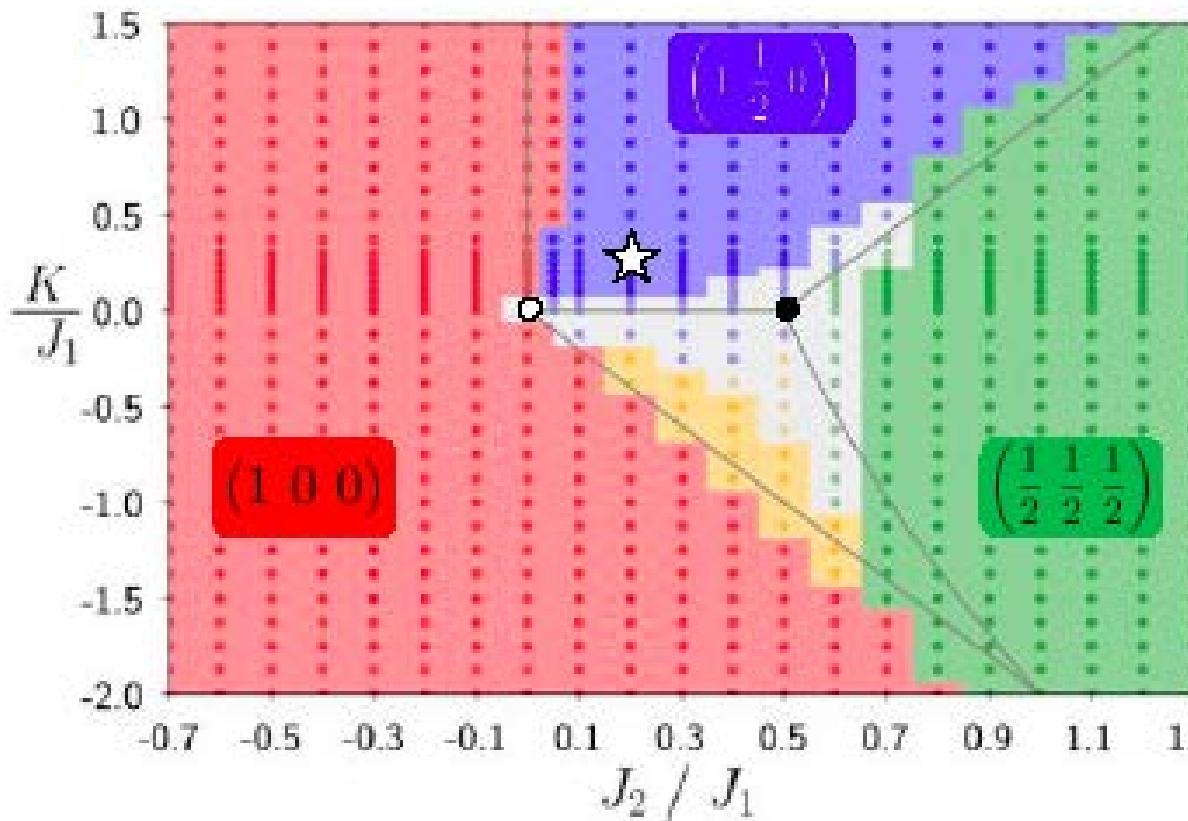
$K/J \approx 0.2$ AFM

$J'/J \approx 0.2$ nnbr

$\Gamma/J \approx 0$

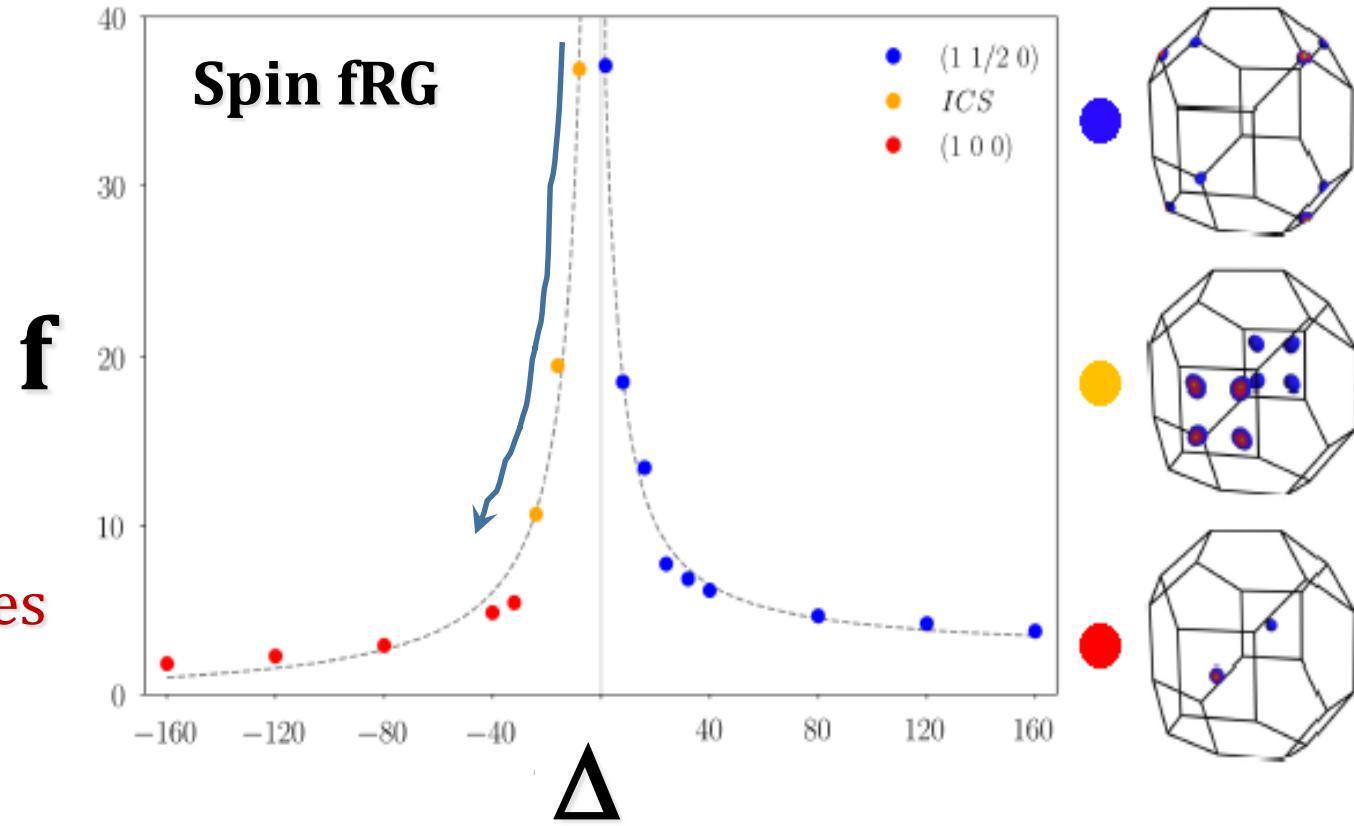
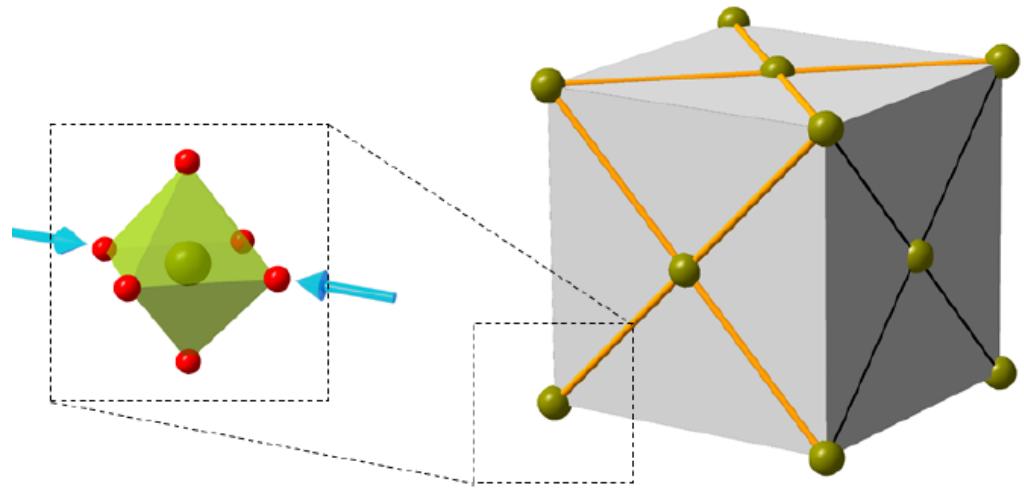
$$H_{ij} = J \vec{S}_i \cdot \vec{S}_j + K S_i^z S_j^z + \Gamma (S_i^x S_j^y + S_i^y S_j^x)$$

$j=1/2$ Mott Insulators on the Ideal FCC Lattice Theory: Spin Functional-RG



- Proximate to Quantum Spin Liquid (type of QSL unclear)
- Very large frustration parameter, $f \sim 30$ ($>>$ classical)
- Kitaev interaction **suppresses** frustration

Mott Insulators on the FCC Lattice Theory: Impact of Weak Distortion



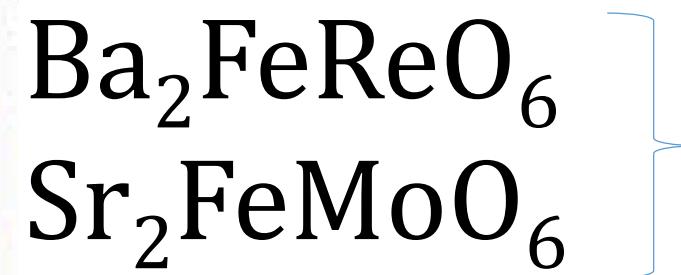
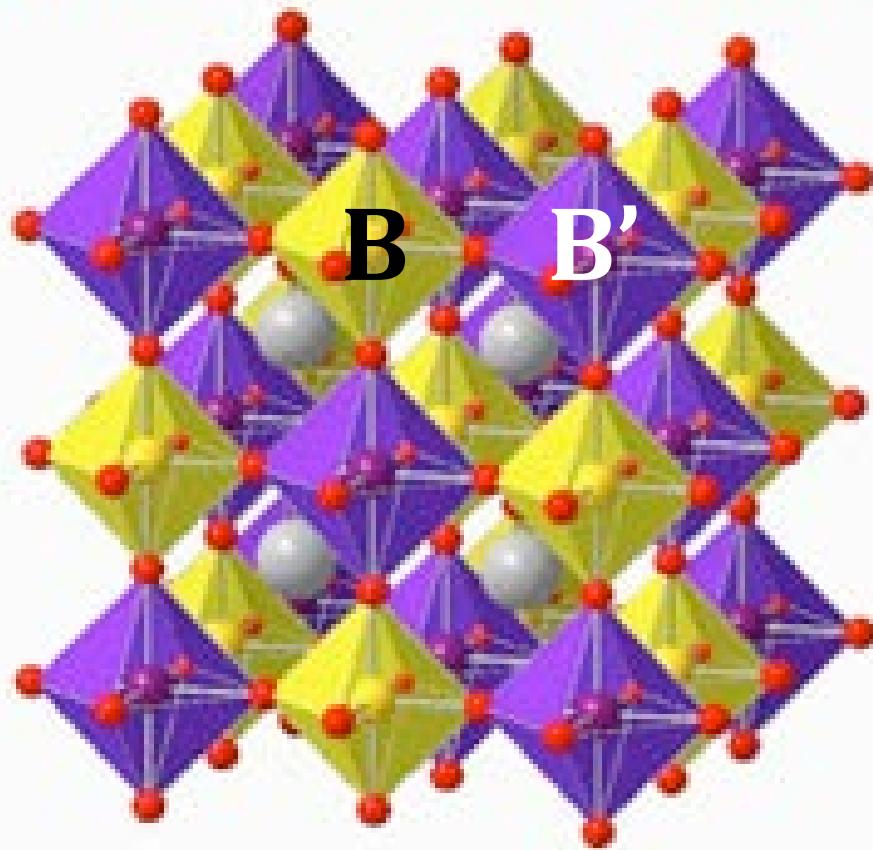
- Z-compression ($\Delta > 0$) strongly enhances XY-plane Heisenberg exchange
- Z-stretch ($\Delta < 0$) strongly suppresses XY-plane Heisenberg exchange

Strong magneto-elastic effect

$\text{La}_2\text{MgIrO}_6$: Global distortion, $f \sim 2$
 $\text{Ba}_2\text{CeIrO}_6$: Local distortions, $f \sim 14$

Coupling the “Lego Blocks” High Tc Chern Insulators in Ultrathin Films

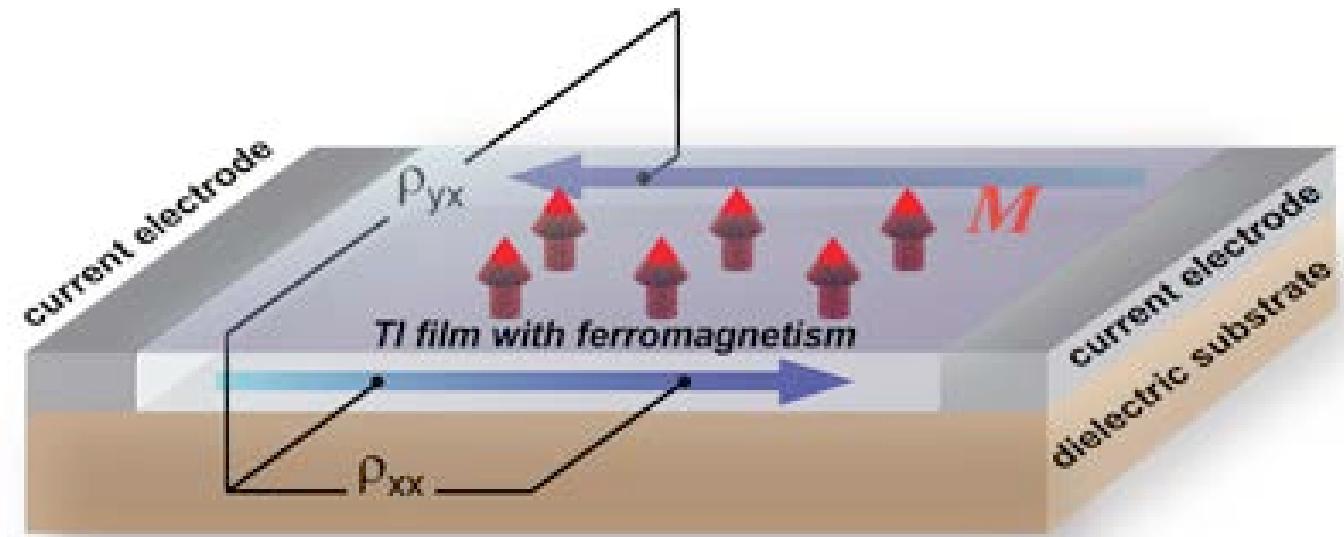
Double perovskite lattice



- Half-Metallic FMs
- Spin-Polarized
- High Bulk FM Tc

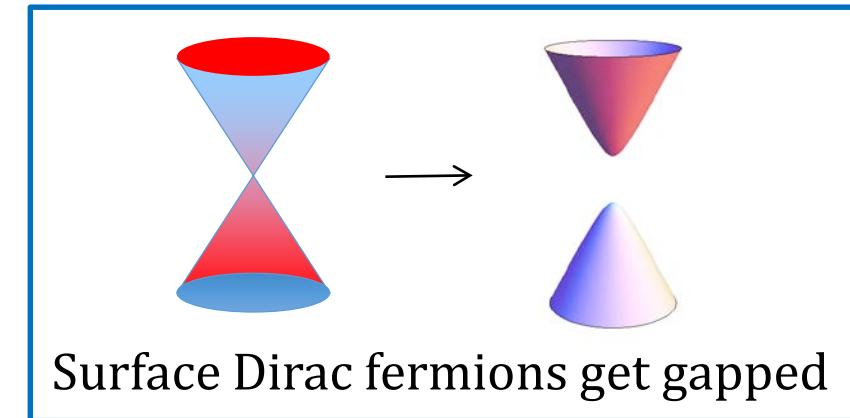
Quantum Anomalous Hall Effect

- Topological insulators: Surface Dirac Fermions
- Breaking time-reversal with dopant magnetization, no B-field!
- Spin orbit coupling is crucial: Gives a “**Lorentz force**”



R. Yu, et al (Science 2010)

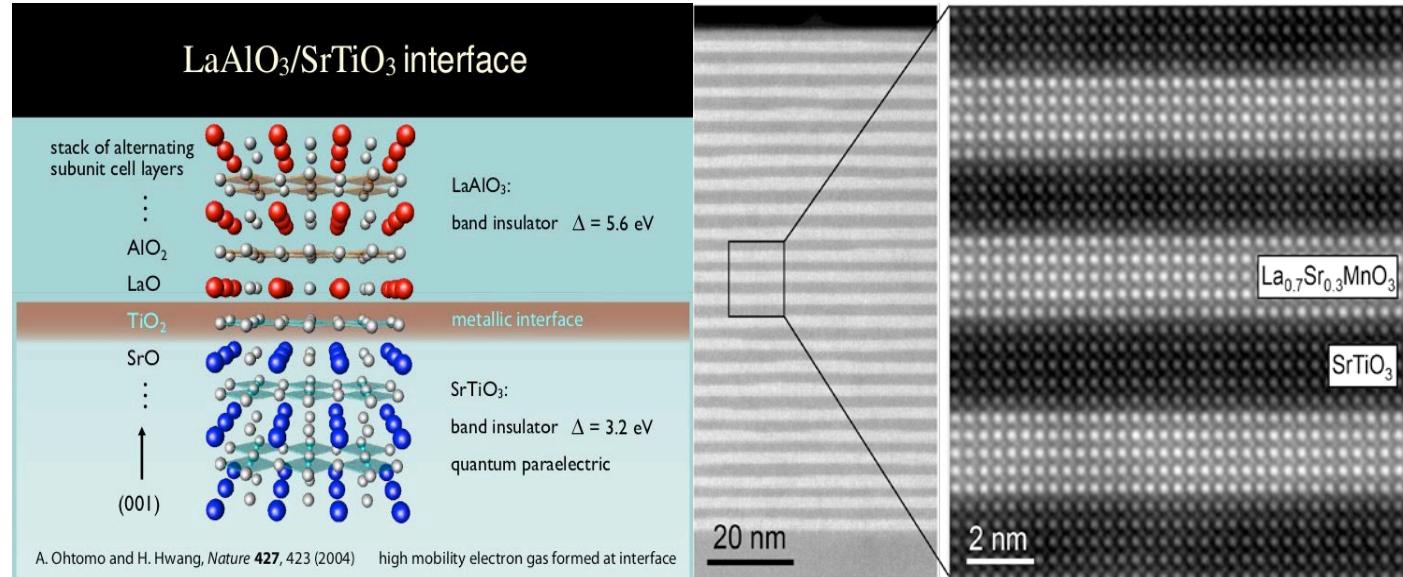
$(\text{Bi},\text{Sb})_2\text{Te}_3$ film doped with Cr or V



Ferromagnetic $T_c \sim 10\text{K}$
Hall quantization: $T \sim 25\text{ mK}$

C.Z. Chang et al, Science 2013 (Xue group, Tsinghua)
C.Z. Chang et al, arXiv (M. Chan + J. Moodera groups)
A. J. Bestwick et al, arXiv (Goldhaber-Gordon group)
A. Kandala, et al, arXiv (N. Samarth group)

Oxide heterostructures



A. Ohmoto, H. Hwang, Nature (2004)

**Boost temperature scale for
Quantum Anomalous Hall effect?**

Multicomponent Superlattices



Exchange bias in LaNiO₃-LaMnO₃ superlattices

Marta Gibert^{1*}, Pavlo Zubko¹, Raoul Scherwitzl¹, Jorge Íñiguez² and Jean-Marc Triscone¹

**Triscone Group (Geneva)
M. Gibert (Zurich)**



Local electronic and magnetic studies of an artificial La₂FeCrO₆ double perovskite
Benjamin Gray, Ho Nyung Lee, Jian Liu, J. Chakhalian, and J. W. Freeland

Citation: *Applied Physics Letters* 97, 013105 (2010); doi: 10.1063/1.3455323

Jak Chakhalian group (Rutgers)

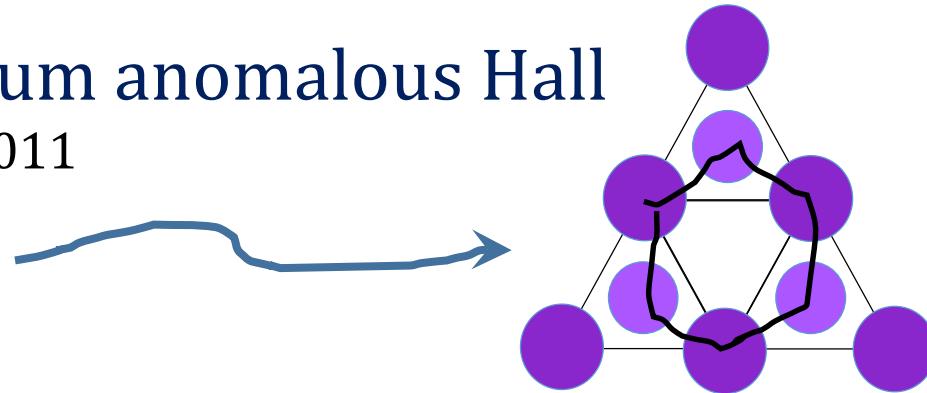
**H. Takagi group (APL 2015)
SrIrO₃/SrTiO₃**

**ORNL/UTK/ANL (H.N. Lee)
SrIrO₃/SrMnO₃ – large AHE
Nat Comm. 2016**

Topological Bands in (111) Heterostructures

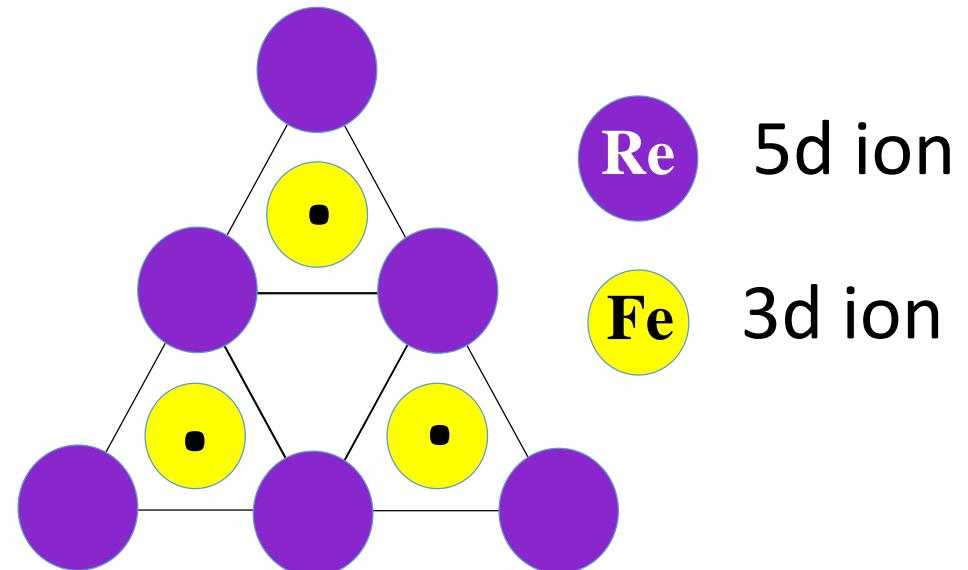
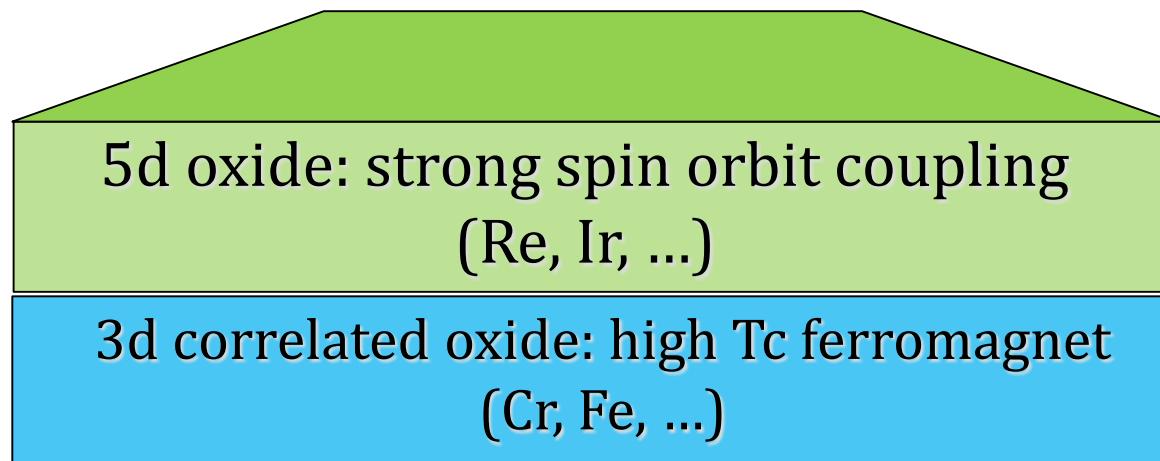
Transition metal oxide bilayers: Quantum anomalous Hall

- D. Xiao, Zhu, S. Okamoto, Y. Ran, Nagaosa, Nat Comm 2011
- G. Fiete group, PRB 2011, PRL 2013
- K. Held group, PRL 2017 (SrRuO_3 bilayers)



Multicomponent bilayers

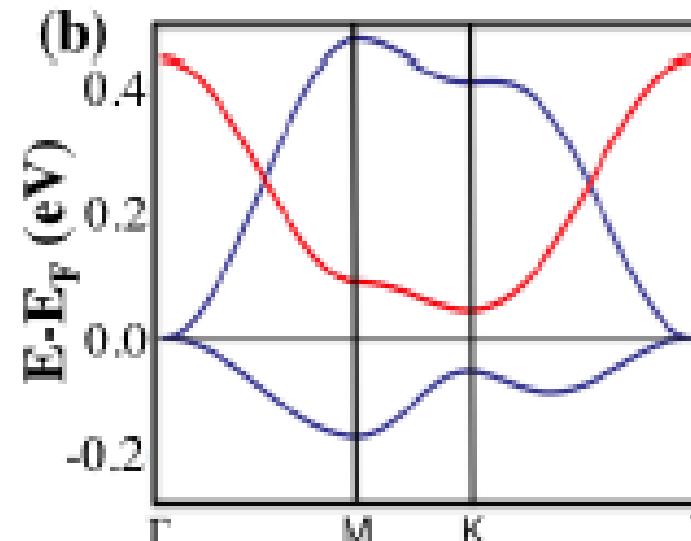
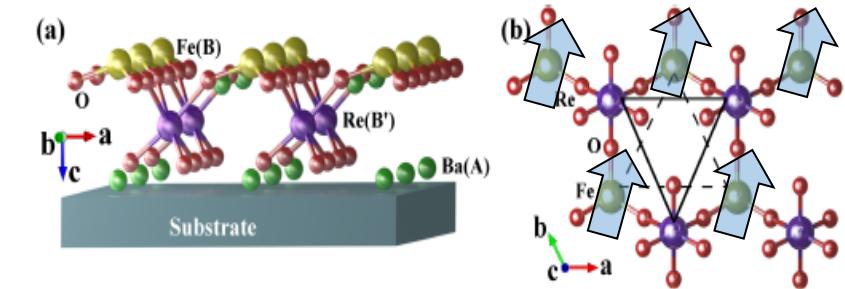
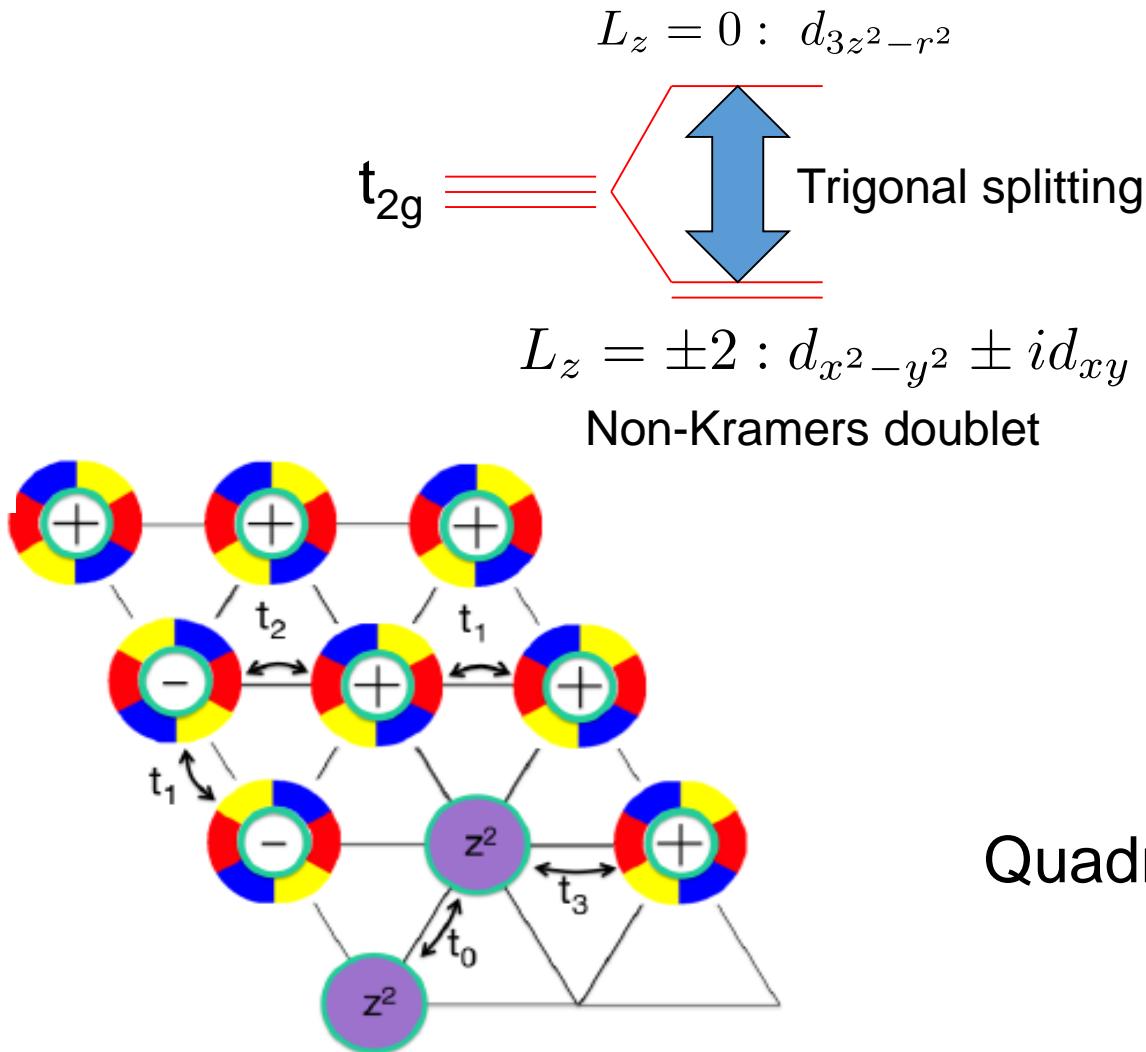
Predict high T_c QAH effect



A.M. Cook, AP (PRL 2014); K. Garrity, D. Vanderbilt, PRB (2014);
S. Baidya, U.Waghmare, AP, T.Saha-Dasgupta (PRB 2015+16)

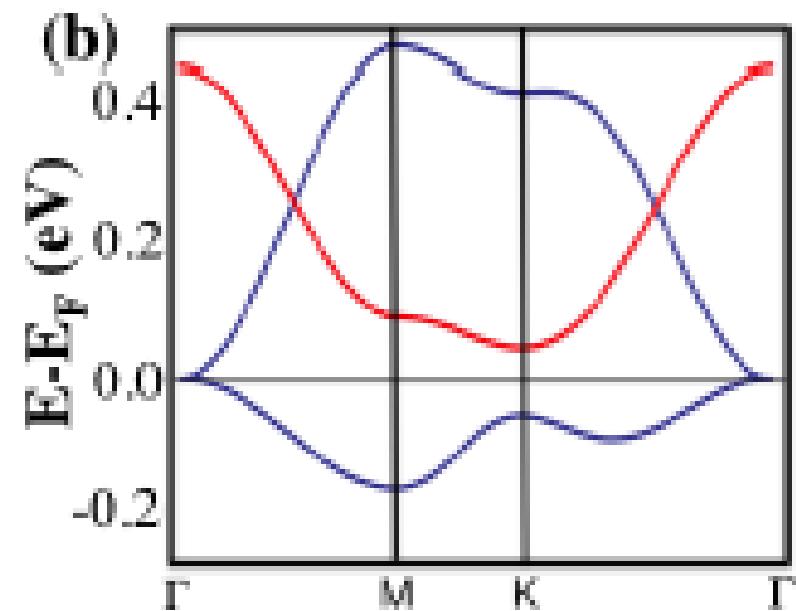
“Multicomponent” (111) bilayer

- Triangular half-metal

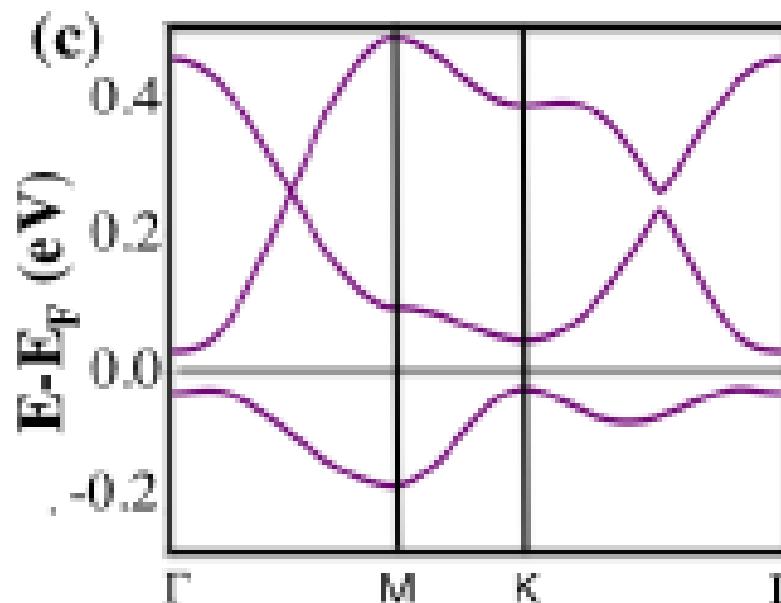


Quadratic band touching at Γ : C3 + “Time reversal”

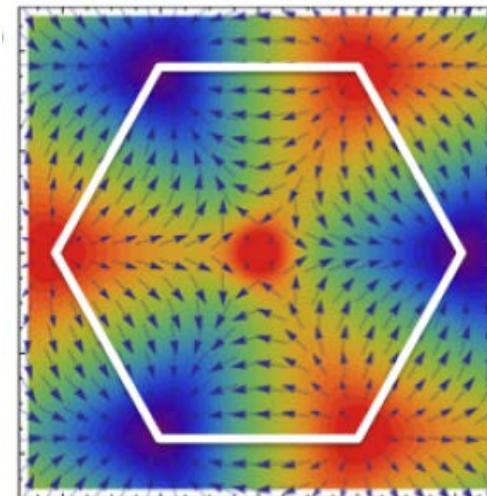
Large gap Chern insulator



SOC
→
 $\lambda L_z \langle S_z \rangle$



DFT: QAH gap ~ 100 meV
Monte Carlo: Ferromagnetic $T_c \sim 250$ K



k-space skyrmion

Summary

- **Double Perovskites: Lego-Blocks with Strong SOC**
- **RIXS: Can be used to extract SOC and Hund's Coupling**
- **Iridates, Rhenates, Osmates: Hund's coupling < SOC**
- **Lego Assembly: FCC Mott Insulators [Strong Magnetoelastic Effect]**
- **Lego Assembly: Chern Insulators [High Tc QAH Effect]**