

LDA+DMFT approach to calculation of RIXS spectra

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Collaborators

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LDA+DMFT approach to core-level spectroscopy

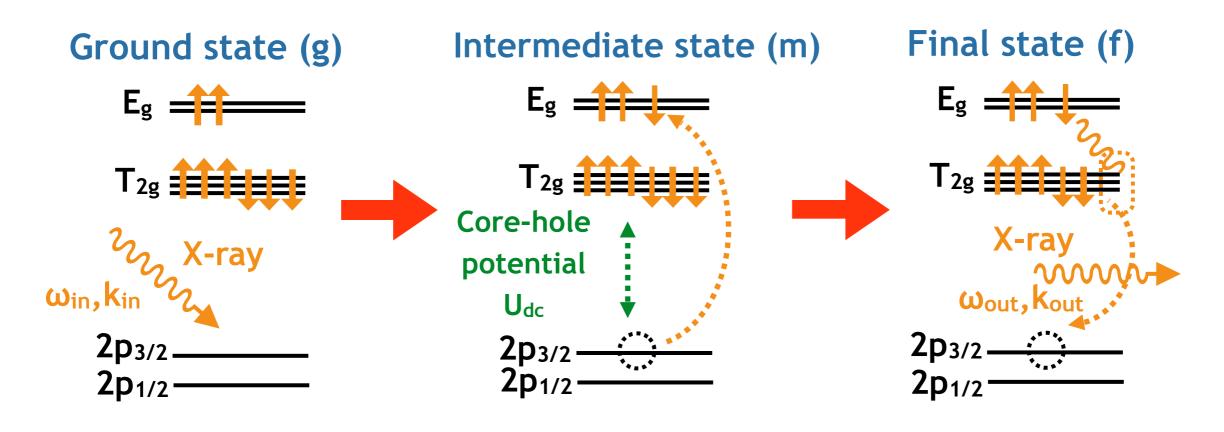
A. Hariki, T. Uozumi and JK, Phys. Rev. B **96**, 045111 (2017)

LDA+DMFT study of L-edge RIXS in high-valence cuprates

A. Hariki, M. Winder and JK, Phys. Rev. Lett. 122, 126403 (2018)

Resonant inelastic X-ray scattering (RIXS)

Resonant inelastic X-ray scattering: Powerful probe with broad sensitivity to **two-particle excitations** of charge, orbital, spin and lattice degrees of freedom



$$\begin{split} F_{\rm RIXS}(\omega_{\rm in},\omega_{\rm out}) = & \sum_{f} \Big| \sum_{m} \frac{\langle f|V_{E}|m\rangle\langle m|V_{I}|g\rangle}{\omega_{\rm in} + E_{g} - E_{m} + i\Gamma} \Big|^{2} \delta(\omega_{\rm in} + E_{g} - \omega_{\rm out} - E_{f}) \\ = & -\frac{1}{\pi} {\rm Im} \ \langle g|R^{\dagger}(\textbf{\textit{q}},\omega_{\rm in}) \ \frac{1}{\omega_{\rm loss} + E_{g} - H + i\delta} \ \frac{R(\textbf{\textit{q}},\omega_{\rm in})|g\rangle}{\Delta v_{\rm loss} + E_{g} - H + i\delta} \\ R(\textbf{\textit{q}},\omega_{\rm in}) = V_{E} \frac{1}{\omega_{\rm in} + E_{g} - H + i\Gamma} V_{I} \end{split}$$

H with core-hole

M. W. Haverkort, PRL 105, 167404 (2010)

Theoretical modeling of RIXS in correlated materials

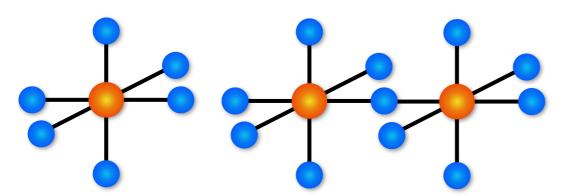
1. Atomic limit

Valence-valence and core-valence interaction

no lattice (=no dispersion)

cluster

multicluster



ab-initio cluster model (and beyond)

M. W. Haverkort et al. PRB **85**, 165113 (2012)

J. Luder et al. PRB **96**, 245131 (2017)

Multisite extension

- computationally expensive
- more tuning parameters ...

Recent development with DMRG etc.

A. Nocera et al. Sci. Rep. 8, 11080 (2018)

T. Toyama et al., PRB **92**, 014515 (2015)

2. Noninteracting limit

Material-specific band structure

approx. electron-hole interaction

e.g. Bethe-Salpeter approach

J. Vinson et al. PRB 83, 115106 (2011)

K. Gilmore et al. Comput. Phys. Commun. 197, 109 (2015)

T. Nomura. J. Phys. Soc. Jpn. 84, 094704 (2015)

Crude approximation on many-body effect in both ground and excited states

Merge advantages in the two approaches together?

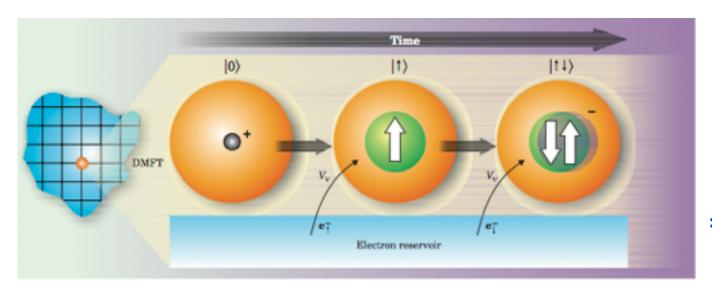
= ultimate goal of theoretical modeling

LDA+DMFT approach to RIXS in correlated materials

Local density approximation (LDA) + Dynamical mean-field theory (DMFT)

successful ab-initio many-body scheme in describing electronic properties of correlated materials

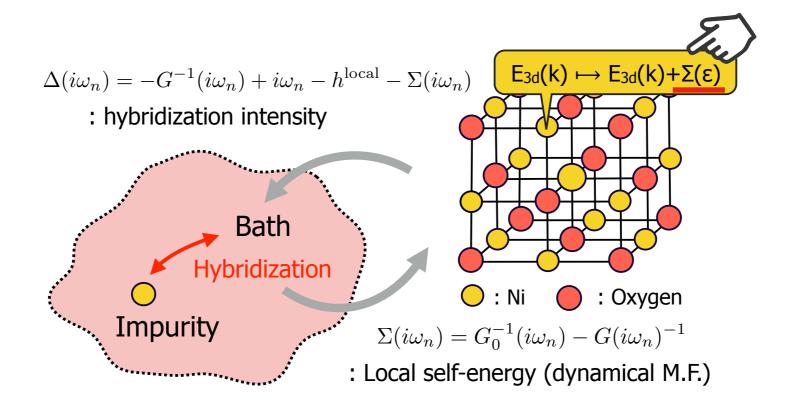
Ref. A. Georges et al., RMP **68**, 13 (1996) / G. Kotliar et al., RMP **78**, 865 (2006)



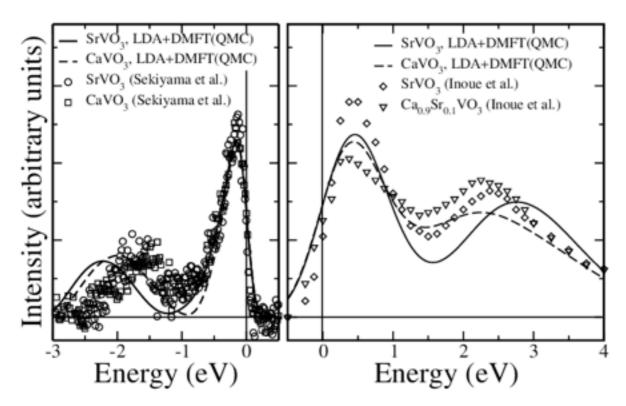
Material specific bands: LDA (multi-band Hubbard model)

Local electron correlation: DMFT

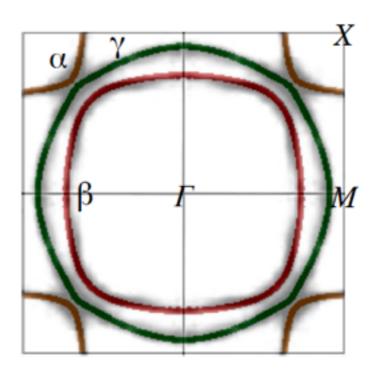
= Many-body dynamics in realistic materials



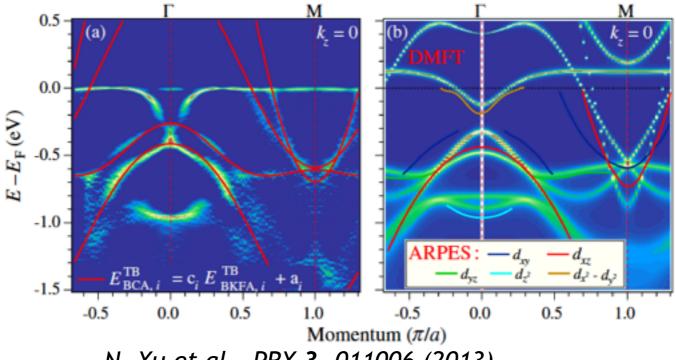
LDA+DMFT: 1P spectra



I. Nekrasov et al., PRB 72, 155106 (2005)

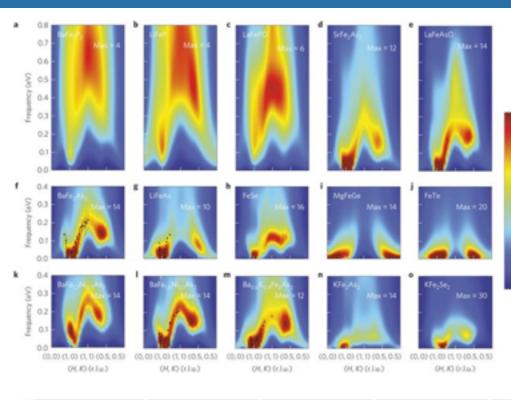


G. Zhang et al., PRL 116, 106402 (2016)



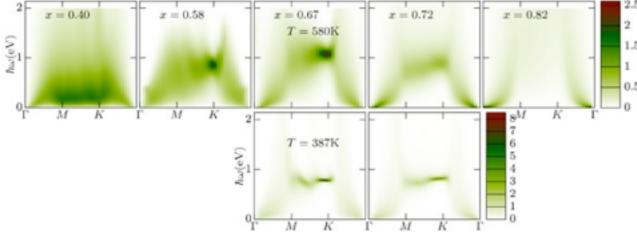
N. Xu et al., PRX 3, 011006 (2013)

LDA+DMFT: 2P spectra (limited)



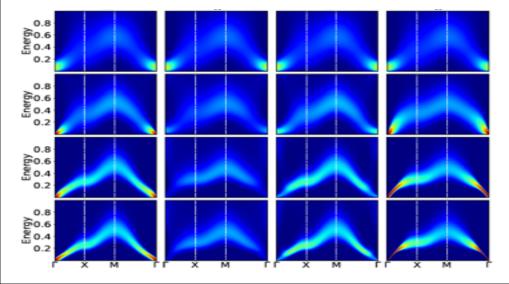
iron pnictides multi-orbital, modified RPA treatment

Z. P. Yin, K. Haule and G. Kotliar, Nat. Phys. 10, 845 (2014)



Na_xCoO₂ single-orbital, full DMFT treatment

L. Boehnke and F. Lechermann, pss **221**, 1267 (2014)



2-orbital Hubbard model, spontaneous symmetry breaking, full DMFT treatment

D. Geffroy et al. PRL **122**, 127601 (2019)

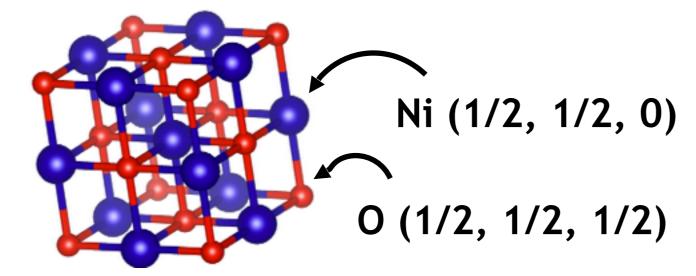
LDA+DMFT approach to RIXS in correlated materials

Hybridization function $V(\varepsilon)$

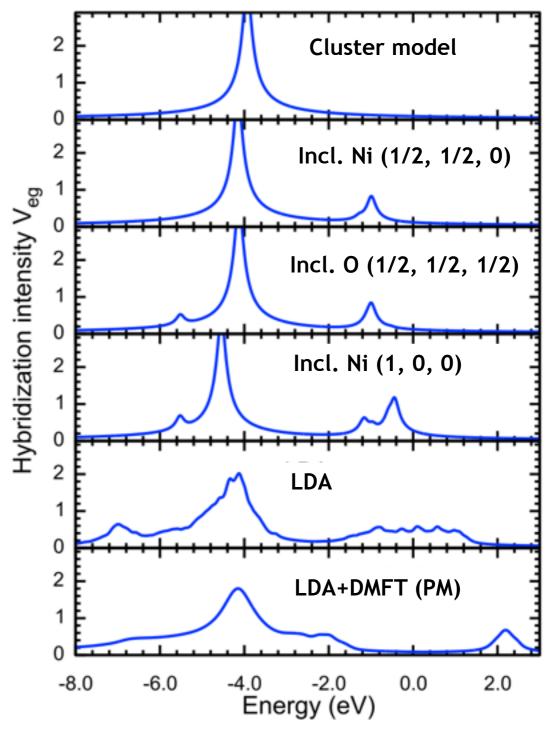
$$\begin{split} V_{\gamma}^{2}(\varepsilon) &= \sum_{i} |\langle \psi_{\gamma,i} | \hat{V}_{\text{imp-host}} | d_{\gamma} \rangle|^{2} \delta(\varepsilon - \varepsilon_{\gamma,i}) \\ &= -\frac{1}{\pi} \operatorname{Im} \langle d_{\gamma} | \hat{V}_{\text{imp-host}} \frac{1}{\varepsilon - H_{\text{host}}(\varepsilon)} \hat{V}_{\text{host-imp}} | d_{\gamma} \rangle \\ & \qquad \qquad \Delta(\varepsilon) \end{split}$$

$$\Delta(arepsilon) = arepsilon - H_{
m imp} - \Sigma(arepsilon) - G^{-1}(arepsilon)$$

DMFT self-energy



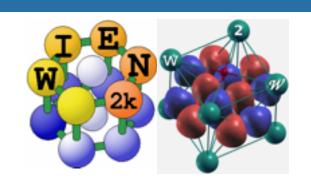
Hybridization function of Ni Eg orbital



LDA+DMFT approach: Computational details

LDA+DMFT:

LDA, tight-binding model: Wien2K + Wien2wannier + Wannier90



DMFT self-consistency: Continuous-time quantum Monte Carlo solver

RIXS calculation: Configuration interaction method

Fine bath discretization (~25 levels per spin, orbital) Efficient basis truncation (G.S theory + α)

3 (2) holes (electrons) below (above) Fermi energy

Lanczos diagonalization + Conjugate Gradient method

(with shift & seed switching)

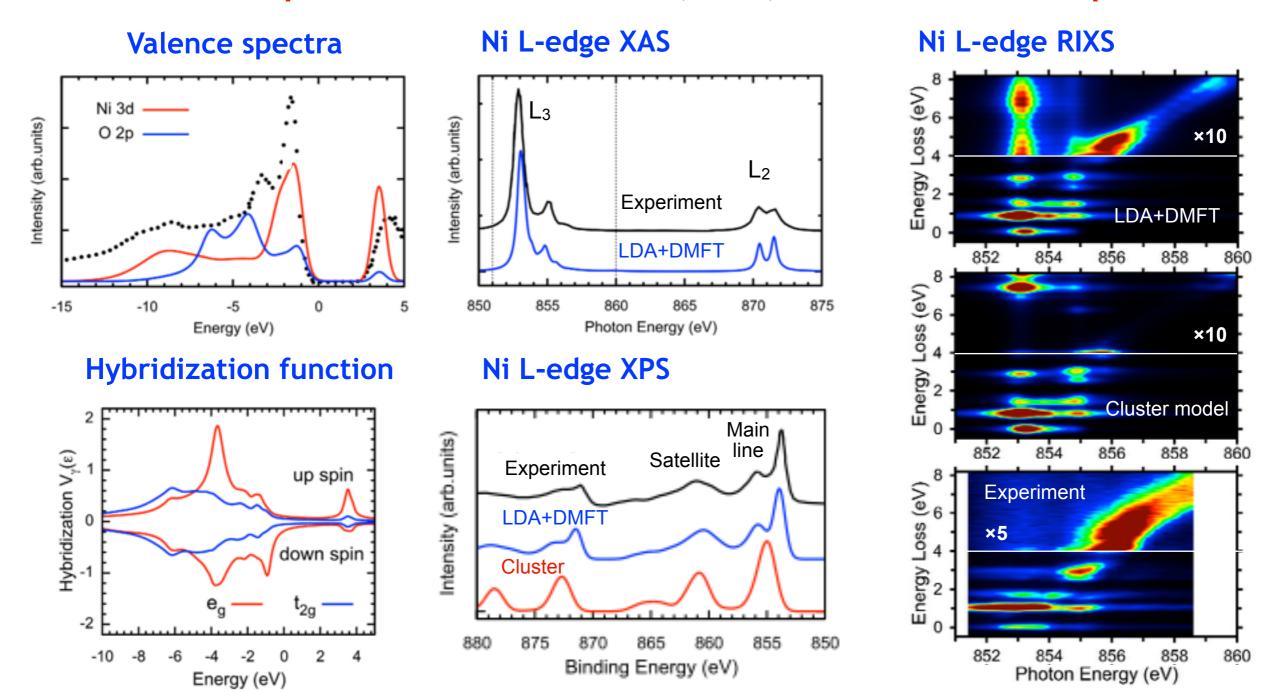
|g>: Lanczos method

$$\begin{split} F_{\rm RIXS}(\omega_{\rm out},\omega_{\rm in}) &= \sum_{f} \left| \sum_{m} \frac{\langle f|T_{\rm e}|m\rangle\langle m|T_{\rm i}|g\rangle}{\omega_{\rm in} + E_g - E_m + i\Gamma_{\rm L}} \right|^2 \times \delta(\omega_{\rm in} + E_g - \omega_{\rm out} - E_f) \\ &= \sum_{f} \left| \langle f|T_{\rm e} \frac{1}{\omega_{\rm in} + E_g - H_{\rm imp} + i\Gamma_{\rm L}} T_{\rm i}|g\rangle \right|^2 \times \delta(\omega_{\rm in} + E_g - \omega_{\rm out} - E_f) \end{split}$$

Conjugated Gradient method (shift technique for ω_{in})

Benchmark: Ni L-edge RIXS in NiO

A unified description of core-level RIXS, XAS, XPS and valence spectra



Good agreement with the experimental RIXS data

- dd excitations (0~4eV) : multiplet, crystal field
- CT excitations (4~8eV) : local and nonlocal CT
- unbound excitations (4eV~): fluorescence-like feature

Experimental data

Valence XPS: G. A. Sawatzky et al. PRL 53, 2339 (1984)

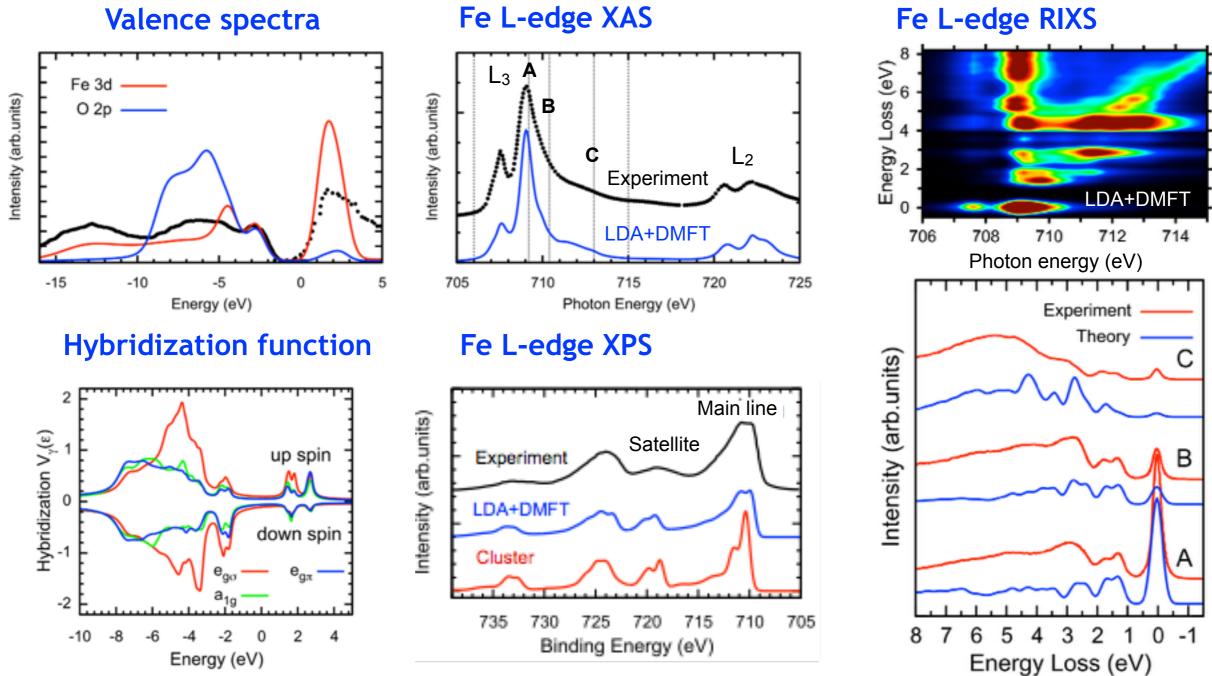
L-edge XAS: D. Alders et al. PRB **57**, 11623 (1998)

L-edge XPS: M. Taguchi et al. PRL 100, 206401 (2008)

L-edge RIXS: G. Ghiringhelli et al. PRL 102, 027401 (2009)

Benchmark: Fe L-edge RIXS in Fe₂O₃

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R. J. Lad et al. PRB 39, 13478 (1898)

L-edge XAS: M. L. Yang et al. PRB 80, 014508 (2017)

L-edge XPS: M. Miedema et al. JESRP 203, 8 (2015)

L-edge RIXS: J. Miyawaki et al. 96, 214420 (2017)

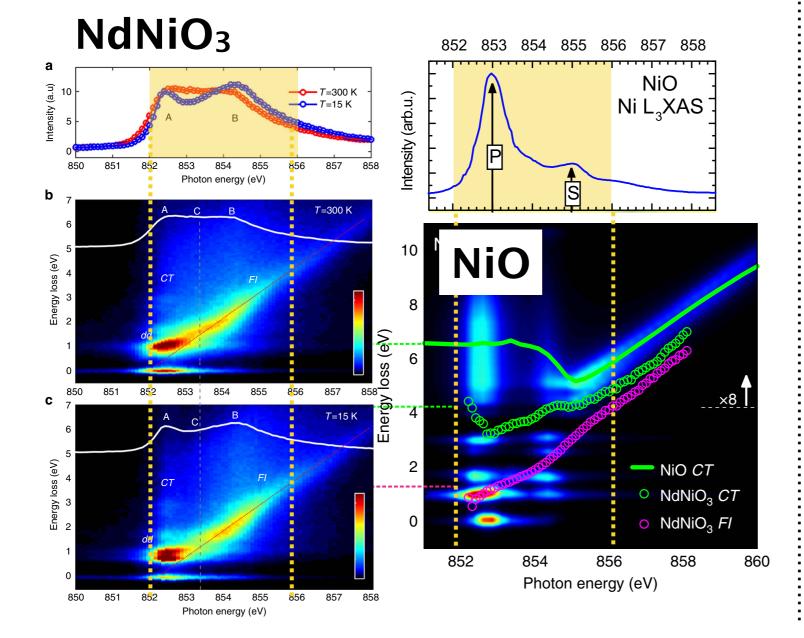
RIXS and XAS in high-valence transition-metal oxides

Received 13 Mar 2016 | Accepted 25 Aug 2016 | Published 11 Oct 2016

DOI: 10.1038/ncomms13017

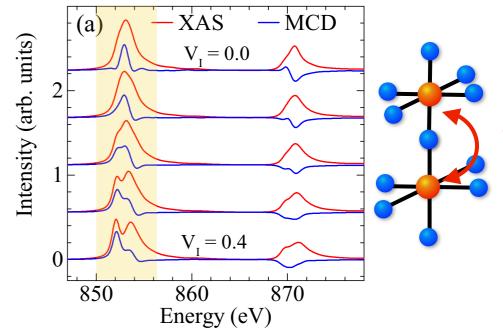
Ground-state oxygen holes and the metal-insulator transition in the negative charge-transfer rare-earth nickelates

Valentina Bisogni^{1,2}, Sara Catalano³, Robert J. Green^{4,5}, Marta Gibert³, Raoul Scherwitzl³, Yaobo Huang^{1,6}, Vladimir N. Strocov¹, Pavlo Zubko^{3,7}, Shadi Balandeh⁴, Jean-Marc Triscone³, George Sawatzky^{4,5} & Thorsten Schmitt¹

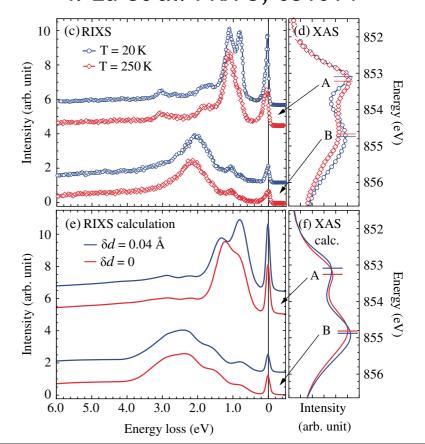


Double-cluster studies

R. J. Green et al. PRB **94**, 195127



Y. Lu et al. PRX 8, 031014

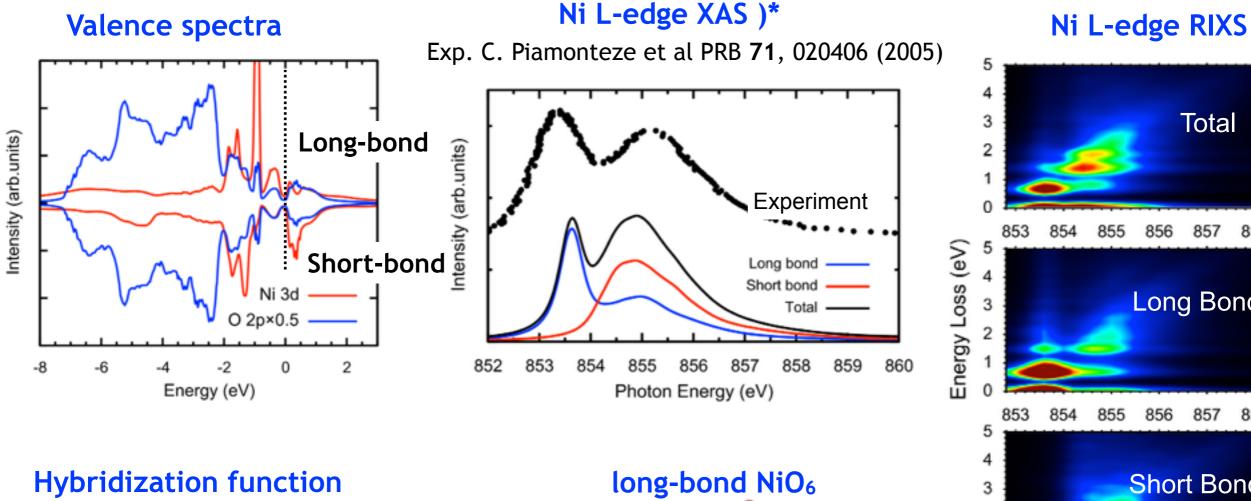


Ni L₃-RIXS in LuNiO₃: LDA+DMFT

Long-bond

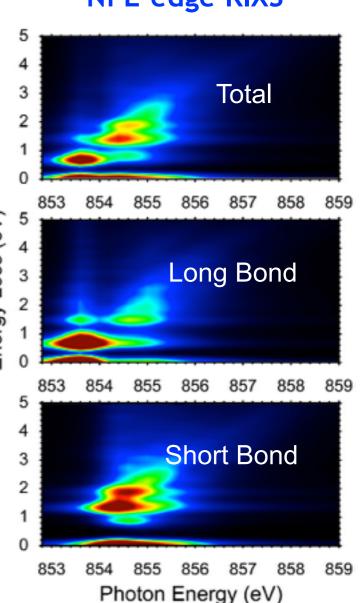
2

Energy (eV)



Short-bond

NiO₆

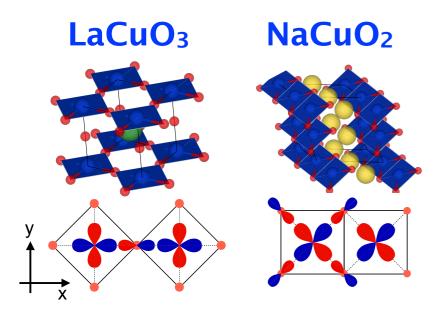


Hybridization -1 **Short-bond**

)* core-level shift (LDA) ~0.4eV (not included in these results)

Cu L₃-RIXS in LaCuO₃ and NaCuO₂





RNiO₃ V. Bisogni et al. (2016) а Intensity (a.u) - T=300 K → T=15 K 8<mark>5</mark>2 854 8<mark>5</mark>6 857 Photon energy (eV) Energy loss (eV)

Photon energy (eV)

856

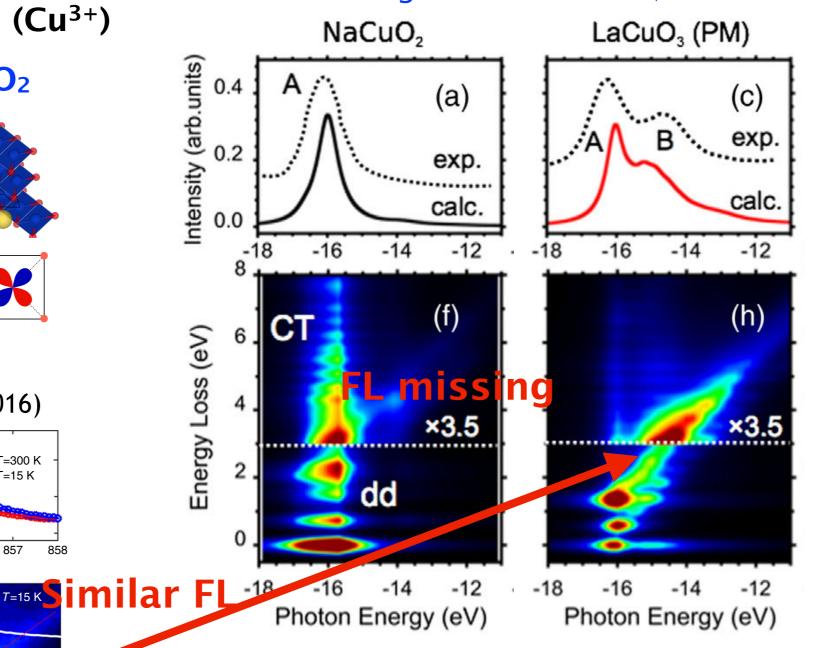
857

2

850

851

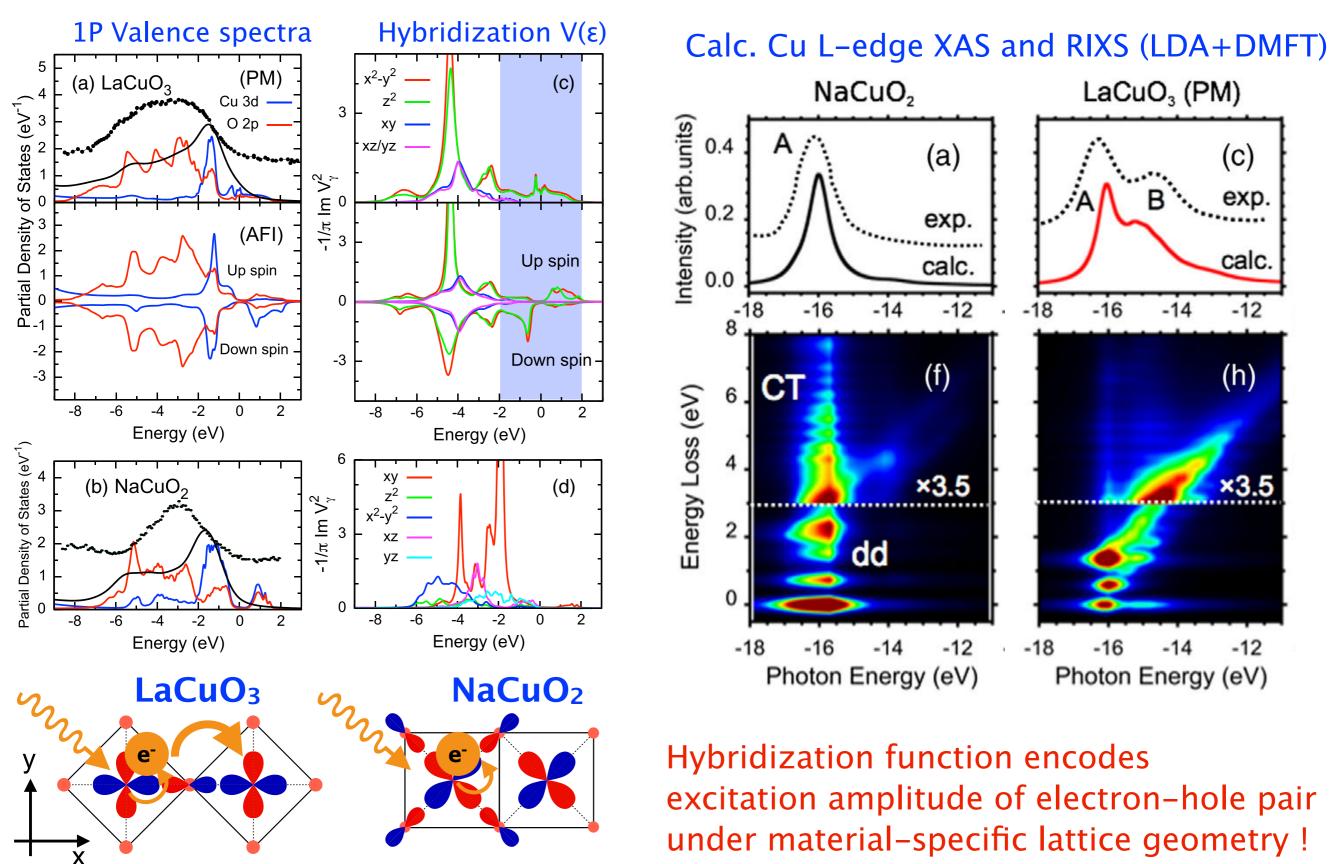
Calc. Cu L-edge XAS and RIXS (LDA+DMFT)



A. Hariki, M. Winder, and JK, PRL 121, 126403 (2018)

What determines the excitation rule of the unbound electron-hole continuum?

Cu L₃-RIXS in LaCuO₃ and NaCuO₂



Bound

Propagate!

Hybridization function encodes excitation amplitude of electron-hole pair under material-specific lattice geometry!

(c)

exp.

calc.

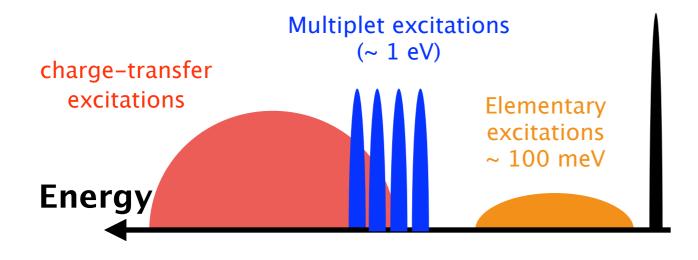
-12

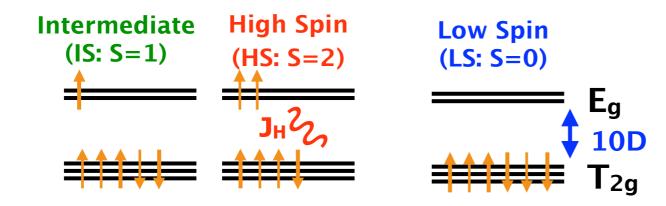
(h)

×3.5

Excitation spectrum in correlated materials

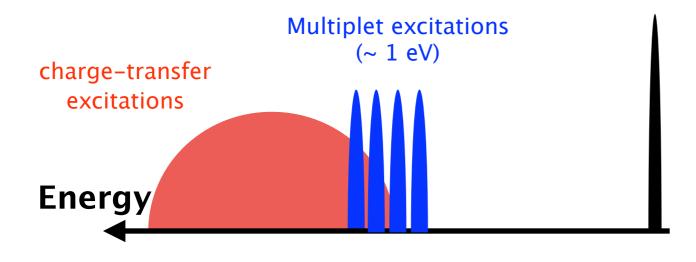
Excitonic physics in LaCoO₃ (Co³⁺: d⁶)

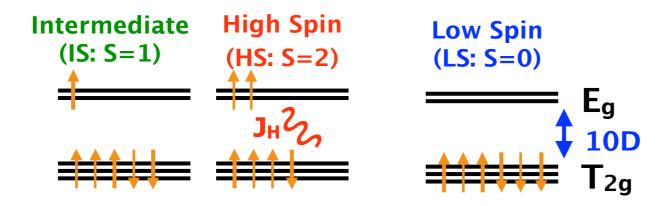




Excitation spectrum in correlated materials

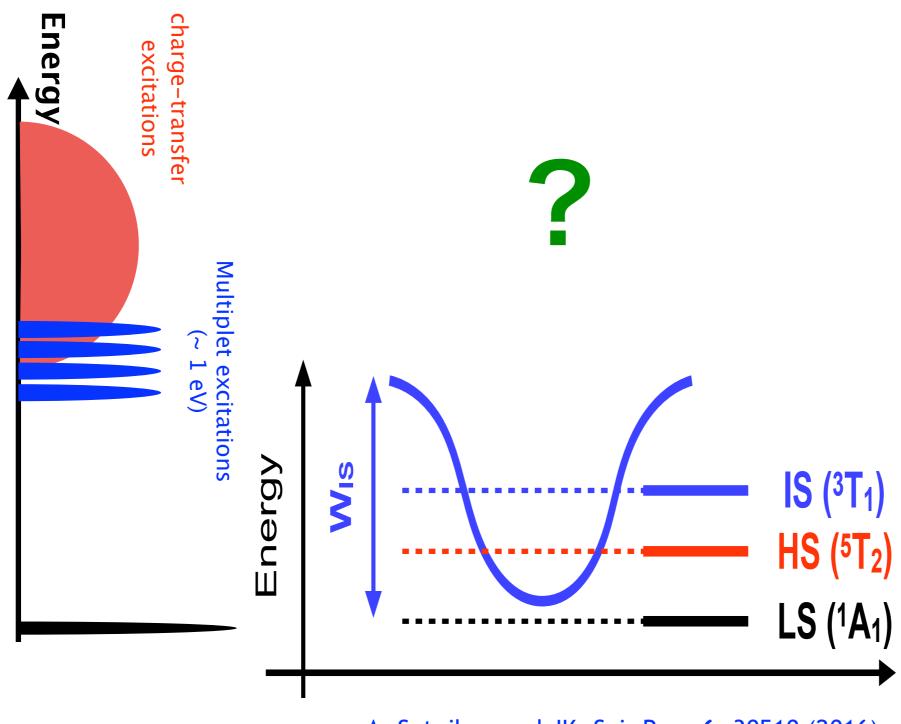
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Excitation spectrum in correlated materials

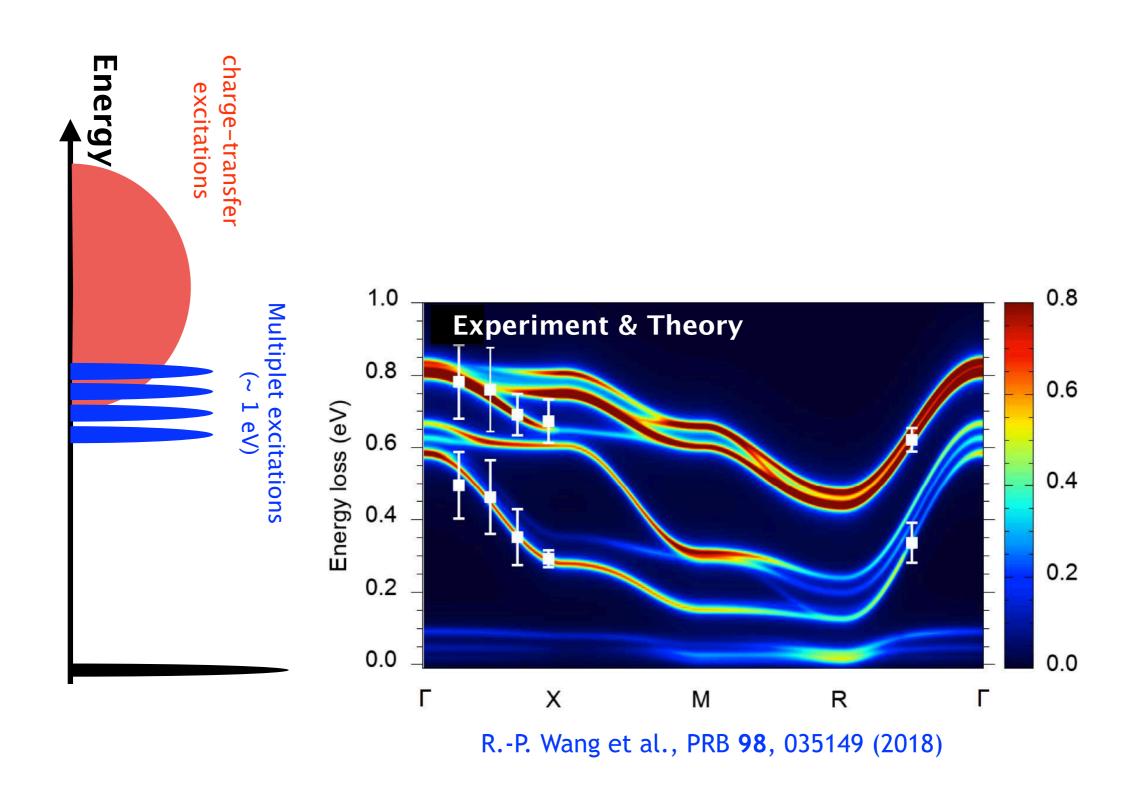
Excitonic physics in LaCoO₃ (Co³⁺: d⁶)



A. Sotnikov and JK, Sci. Rep. 6, 30510 (2016)

Excitation spectrum in correlated materials

Excitonic physics in LaCoO₃ (Co³⁺: d⁶)



Conclusion and future perspective

We developed computational method based on LDA+DMFT for calculation of core-level RIXS, XPS, XAS spectra in correlated materials

Anderson impurity model description with DMFT hybridization in RIXS analysis for correlated materials

Advantages:

- description of both bound (dd) excitation and unbound electron-hole pairs
- 2. low computational cost, applicable to wide range of correlated materials including covalent or high-valance materials

Future:

k-dependence in RIXS spectra (within DMFT or + beyond)

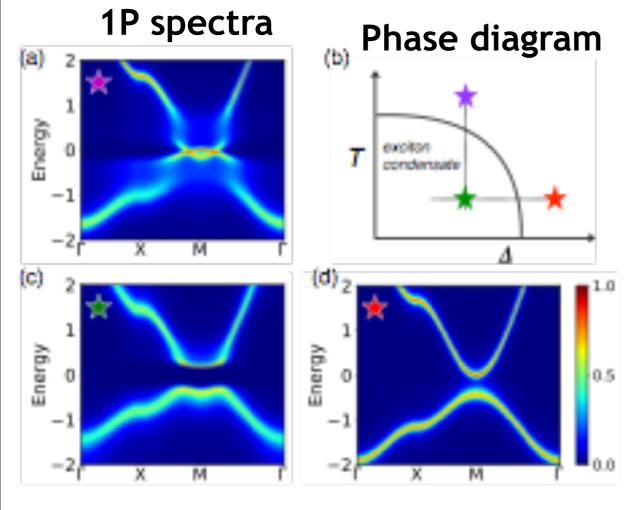
Conclusion and fut Comments on 2P function within DMFT?

Future challenge:
Recover k-dependence in RIXS spectra (within DMFT or +

beyond)

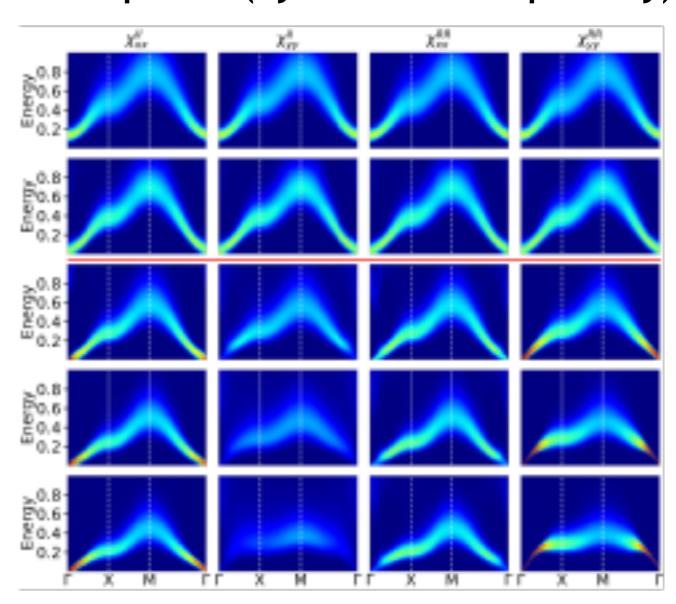
2-band Hubbard model near excitonic-condensation

instability



D. Geffroy et al. PRL 122, 127601 (2019)

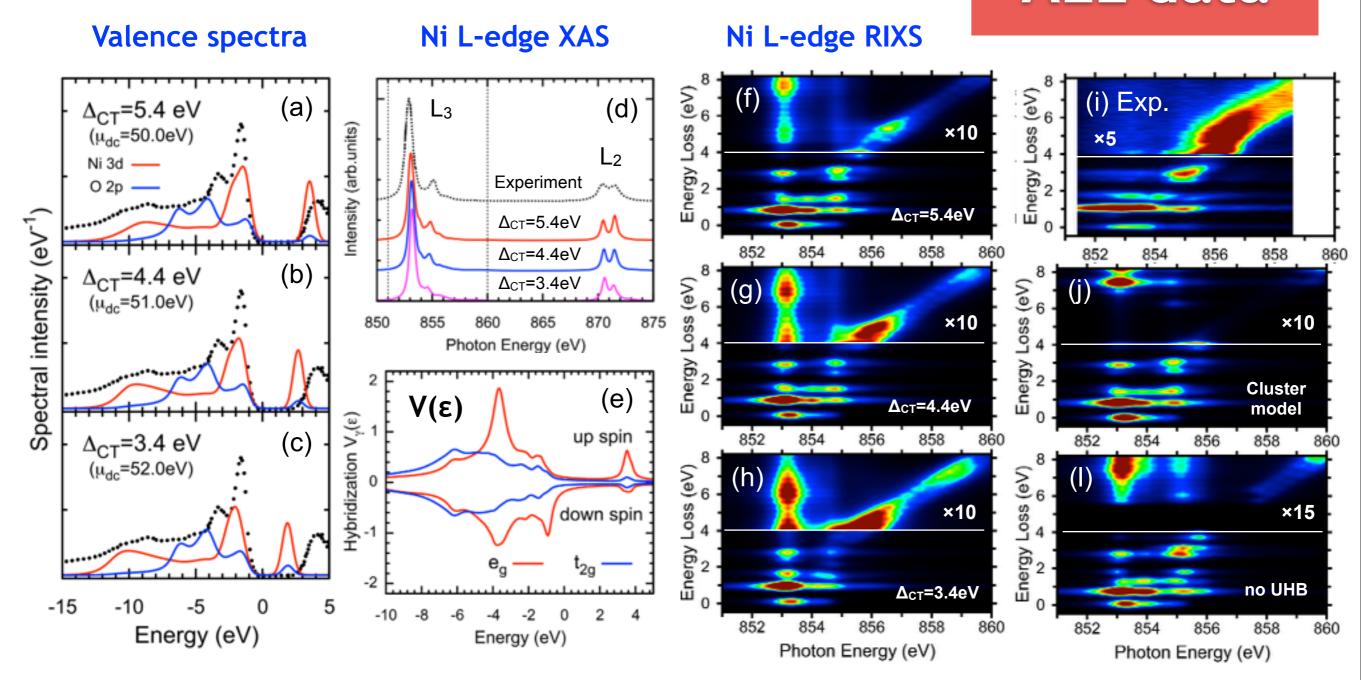
2P spectra (dynamical susceptibility)



$$\chi_{\eta\eta}^{OO}(\mathbf{k},i\nu_n)\!=\!\sum_{\mathbf{R}}\int_0^{\beta}\!d\tau e^{i(\nu_n\tau+\mathbf{k}\cdot\mathbf{R})}\langle O_{\mathbf{i}+\mathbf{R}}^{\eta}(\tau)O_{\mathbf{i}}^{\eta}(0)\rangle\!-\!\langle O^{\eta}\rangle^2.$$

Benchmark: Ni L-edge RIXS in NiO

ALL data



Good agreement with the experimental data

- localized excitations (0~4eV): multiplet, crystal field
- CT excitations (4-8eV): local and nonlocal CT
- unbound excitations (4eV~: fluorescence-like feature)L-edge RIXS: G. Ghiringhelli et al. PRL 102, 027401 (2009)

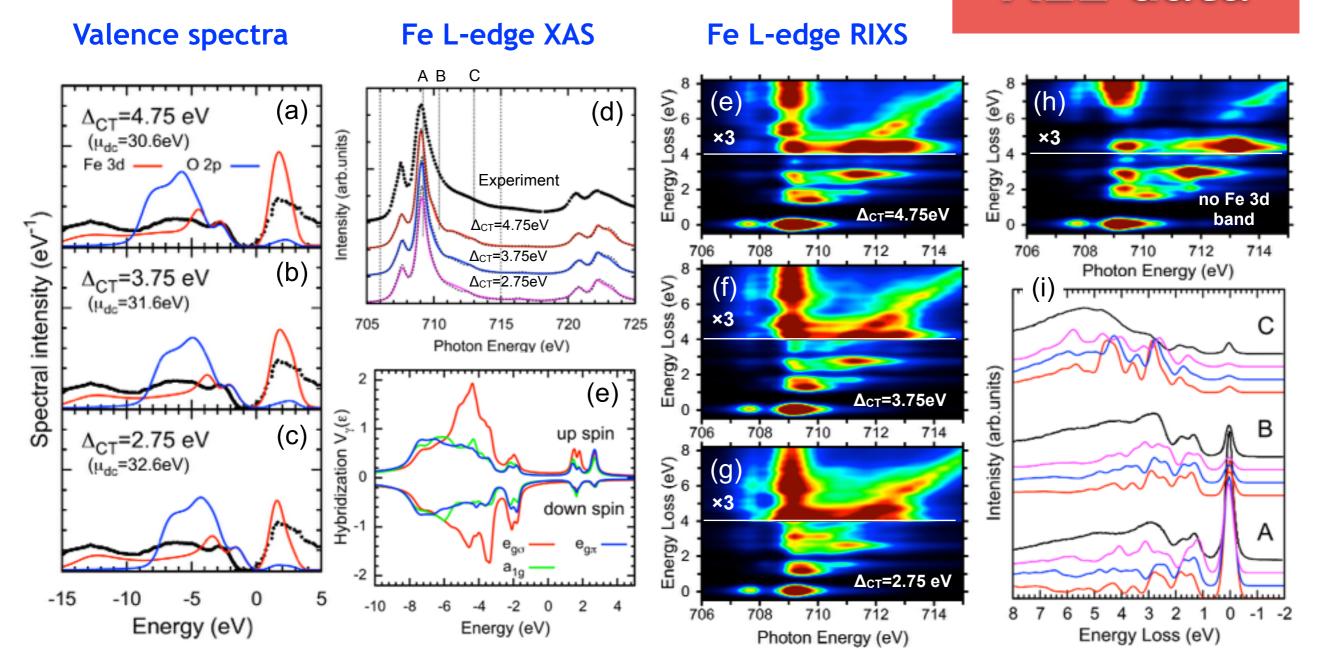
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R. J. Lad et al. PRB 39, 13478 (1898)

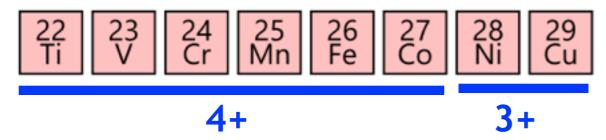
L-edge XAS: M. L. Yang et al. PRB 80, 014508 (2017)

L-edge RIXS: J. Miyawaki et al. 96, 214420 (2017)

High-valence transition-metal ox

for Prague only

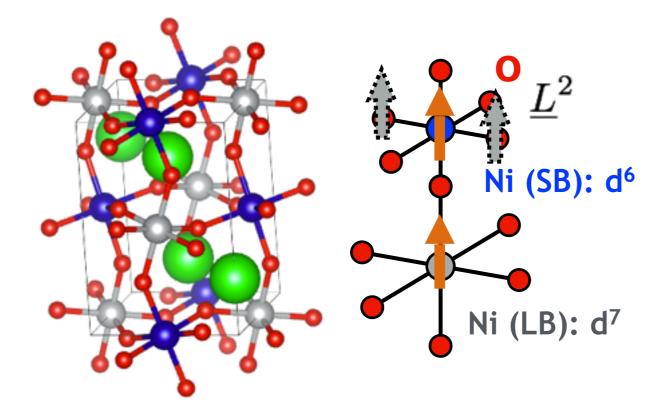
High-valence 3d transition-metal?



Small charge-transfer energy Δ

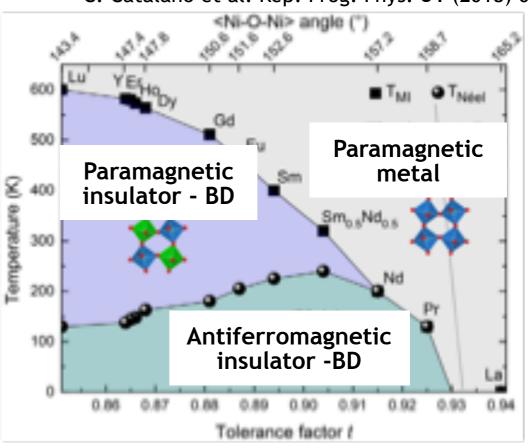
$$d^n \leftrightarrow d^{n+1}\underline{L}$$
 ($\underline{\mathsf{L}}$: Ligand hole)

Oxygen 2p states play a role in low-energy physics



Rare-earth Nickelates: RNiO₃

C. Catalano et al. Rep. Prog. Phys. 81 (2018) 046501



Charge-disproportionation

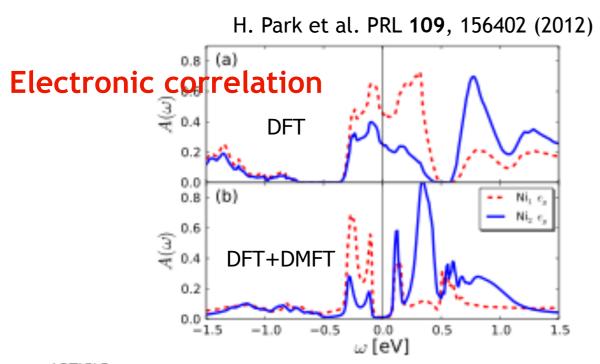
$$d^7 + d^7 \rightarrow d^6 + d^8$$
 T_{29} T_{29} T_{29}

Short Ni-O bond Long Ni-O bond

High-valence transition-metal ox

for Prague only

Electronic + Structural + Ligand

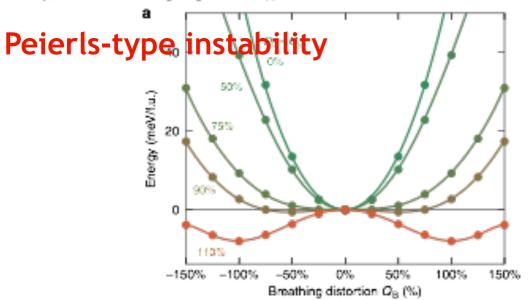


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DOI: 10.1038/s41467-017-01811-x

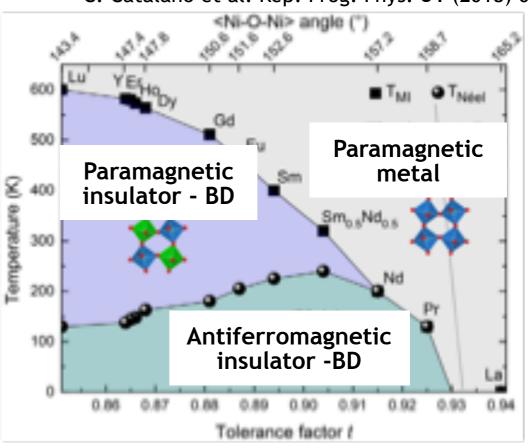
Structurally triggered metal-insulator transition in rare-earth nickelates Nat.Commun. 7, 13017 (2016)

Alain Mercy¹, Jordan Bieder^{1,2}, Jorge Íñiguez³ & Philippe Ghosez¹



Rare-earth Nickelates: RNiO₃

C. Catalano et al. Rep. Prog. Phys. 81 (2018) 046501



Charge-disproportionation

$$d^7 + d^7 \rightarrow d^6 + d^8$$
 $d^8 \underline{L} + d^8 \underline{L}$
 $d^8 \underline{L}^2 + d^8$
 $d^8 \underline{L}^2 + d^8$
 $d^8 \underline{L}^3 + d^8 \underline{L}^4$
 $d^8 \underline{L}^4 + d^8$

Short Ni-O bond Long Ni-O bond