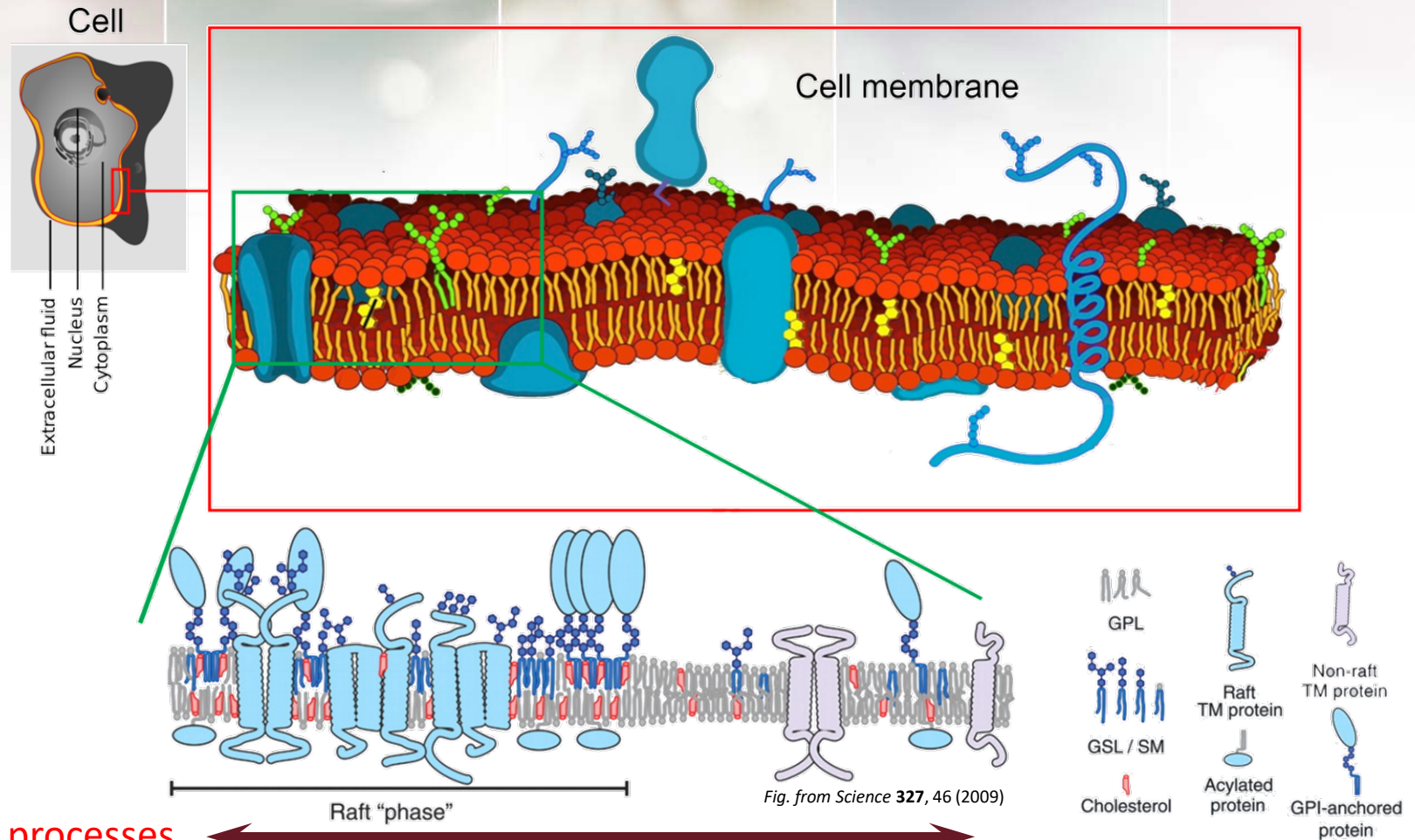


New insight into biophysics of lipid membranes with high resolution IXS

Mikhail Zhernenkov

06/27/2019

Fast and slow dynamics in a cell membrane



Dynamic processes in a cell membrane is virtually impossible to study due to its complexity

XPCS

Slow dynamics on mesoscale: seconds to μ s

Slow dynamics:

Lipid domain (rafts) formation and diffusion, anomalous subdiffusion

Ultrafast dynamics:

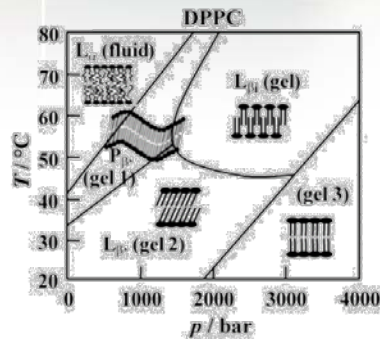
Protein and lipid diffusion, lateral and transmembrane transport

IXS

Fast dynamics on molecular scale: picoseconds

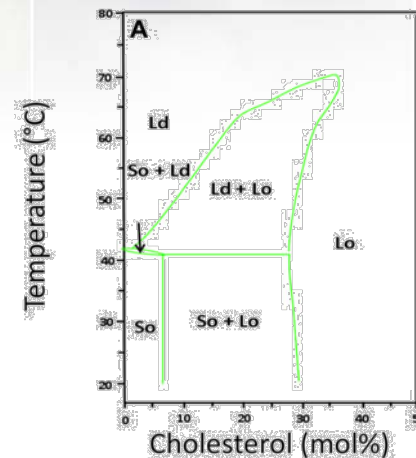
Model membranes as a test bed

Single lipid



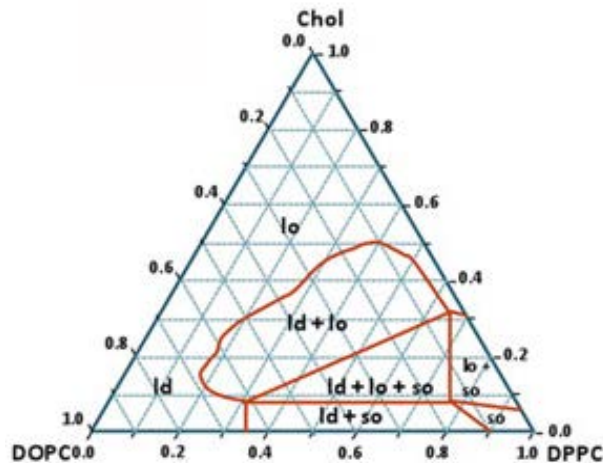
R. Winter, C. Jeworrek, Soft Matter, 2009, 5, 3157

Binary systems (DPPC-Chol)



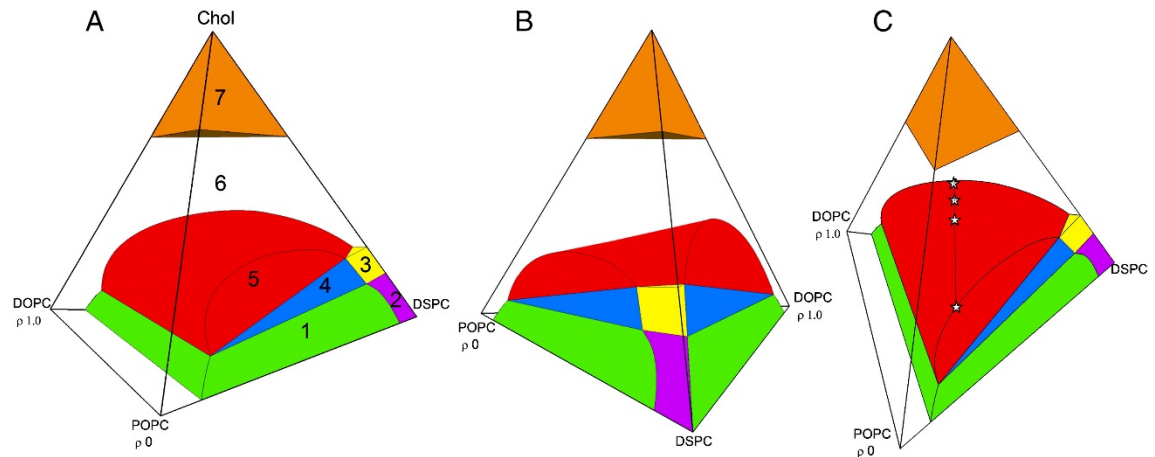
de Almeida&Joly, Front. Plant. Sci (2014)

Ternary systems



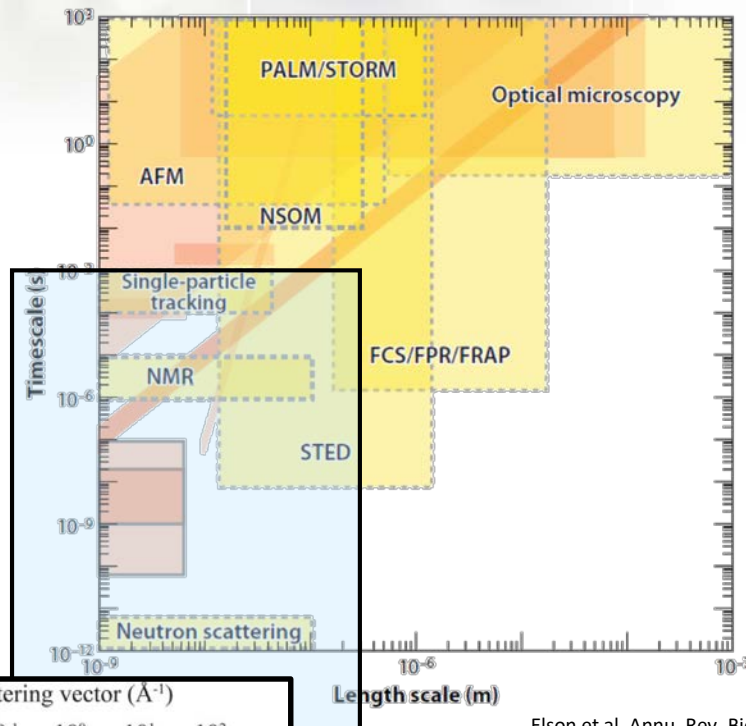
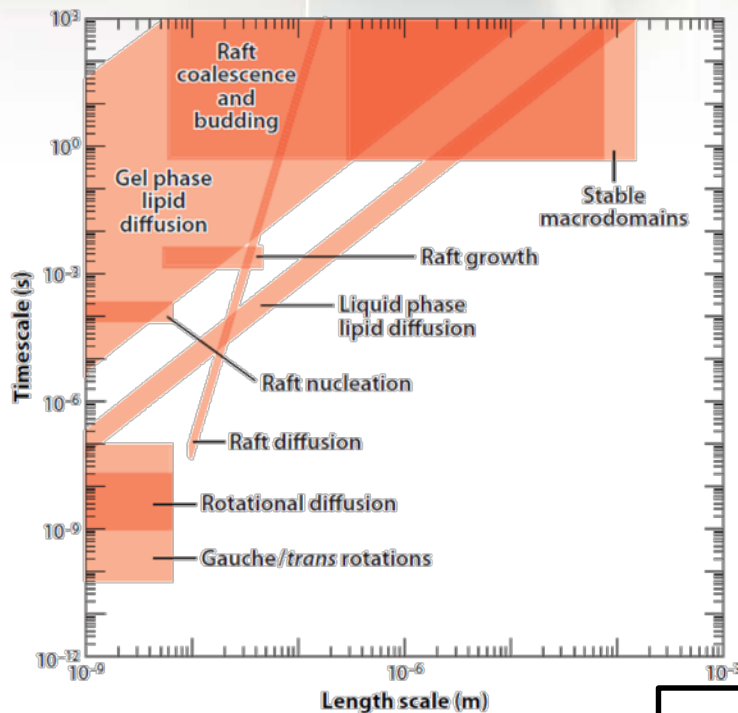
de Almeida&Joly, Front. Plant. Sci (2014)

Quaternary phase diagram



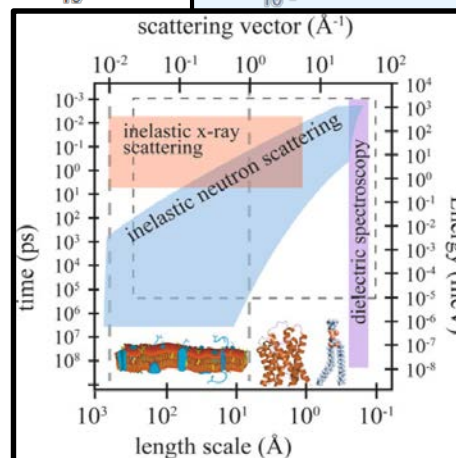
G. Feigenson, Cornell U

Model membranes as a test bed



Elson et al. Annu. Rev. Biophys. 2010. 39:207–26

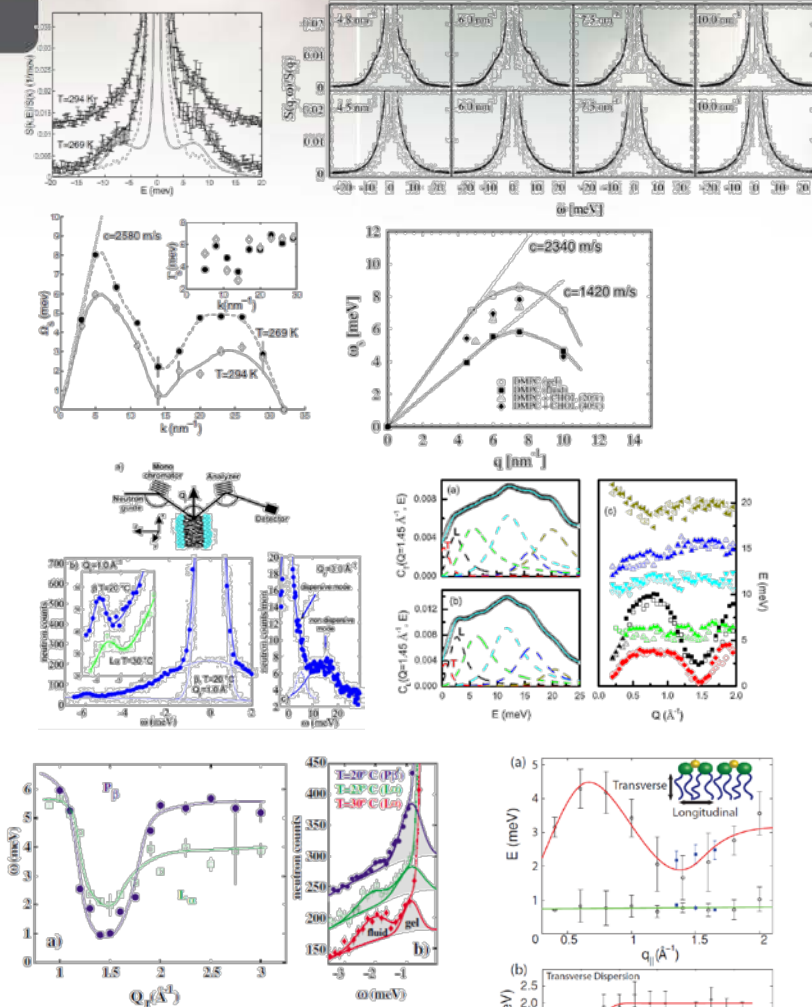
Predicted and observed phenomena in lipid layers span broad length and timescales



Experimental techniques cover most of “large” and relatively “slow” lipid behavior.

Still do not adequately determine “small” and “fast”

Previous efforts to study fast dynamics



➤ Different parts of lipid molecules contribute to different excitations, which all can, in principle, be probed by inelastic scattering

➤ Generalized three effective eigenmode (GTEE) theory to fit the IXS/INS data:

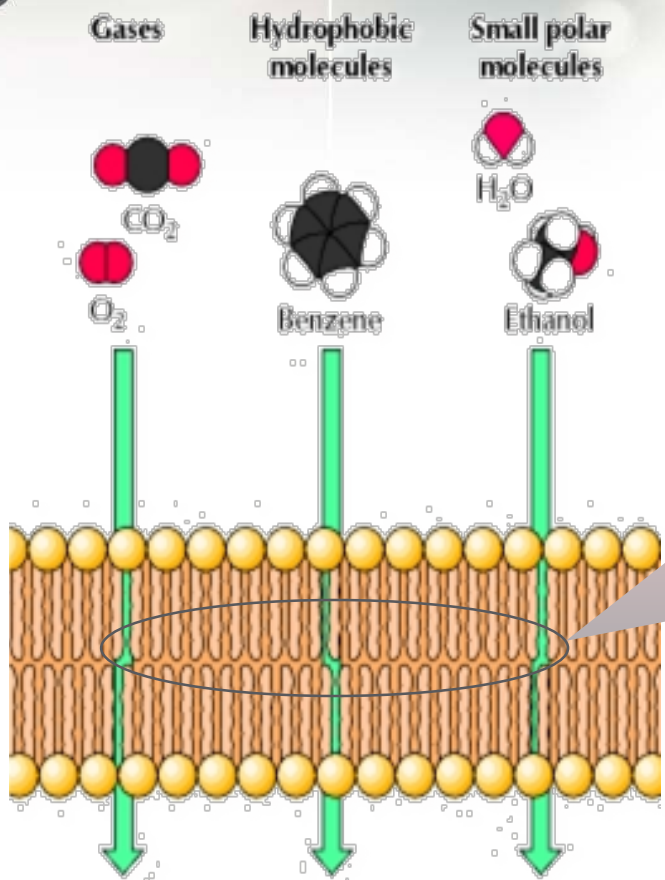
$$\frac{S(Q, \omega)}{S(Q)} = \frac{1}{\pi} \left[A_0 \frac{\Gamma_h}{\omega^2 + \Gamma_h^2} + A_s \left(\frac{\gamma_s + b(\omega + \omega_s)}{(\omega + \omega_s)^2 + \gamma_s^2} + \frac{\gamma_s - b(\omega - \omega_s)}{(\omega - \omega_s)^2 + \gamma_s^2} \right) \right]$$

➤ GTEE → treats the data in terms of simple liquids which is an oversimplification

Very few processes take place on ps time scale: trans-gauche isomerization, density fluctuations, ultra-fast diffusion

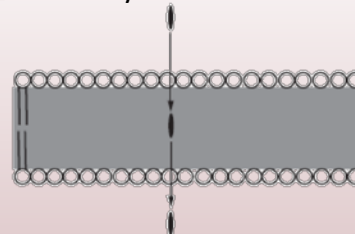
S.H. Chen et al. Phys. Rev. Lett. **86** (2001) 740
 M. Tarek et al. Phys. Rev. Lett. **87** (2001) 238101
 T. Weiss et al. Biophys. J. **84** (2003) 3767
 M. C. Rheinstadter et al. Phys. Rev. Lett. **93** (2004) 108107
 V. Conti Nibali et al. Phys. Rev. E **89** 050301(R) (2014)
 M. Kaye, et al., Phys. Rev. E **83** 050907(R) (2011)

Passive transport mechanisms

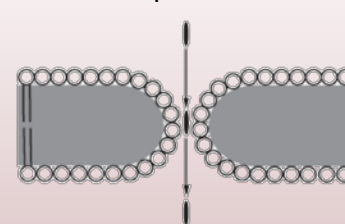


Possible transport mechanisms:

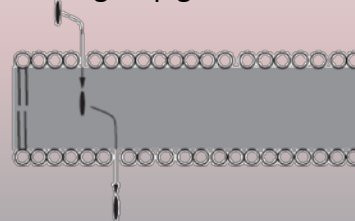
A Solubility-diffusion



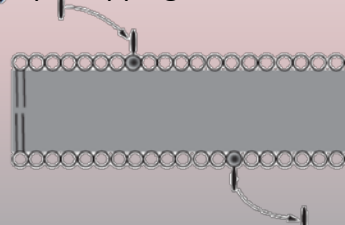
B Transient pores



C Head-group gates



D Lipid flipping-carriers



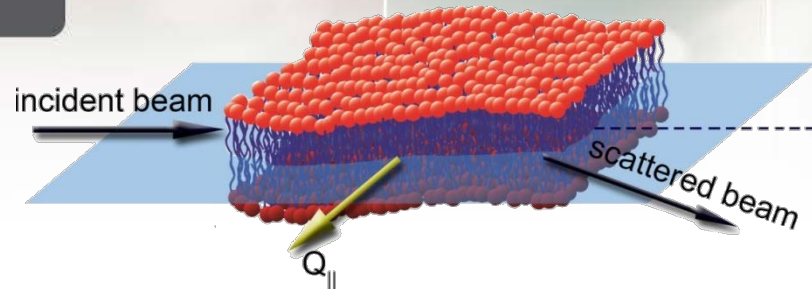
S.S. Mansy, Cold Spring Harb. Perspect. Biol. 2010; 2:a002188

IXS is perfect to study ultrafast **thermally-triggered** phonon-mediated transport mechanisms

Transient Pore Model
for unsaturated lipids

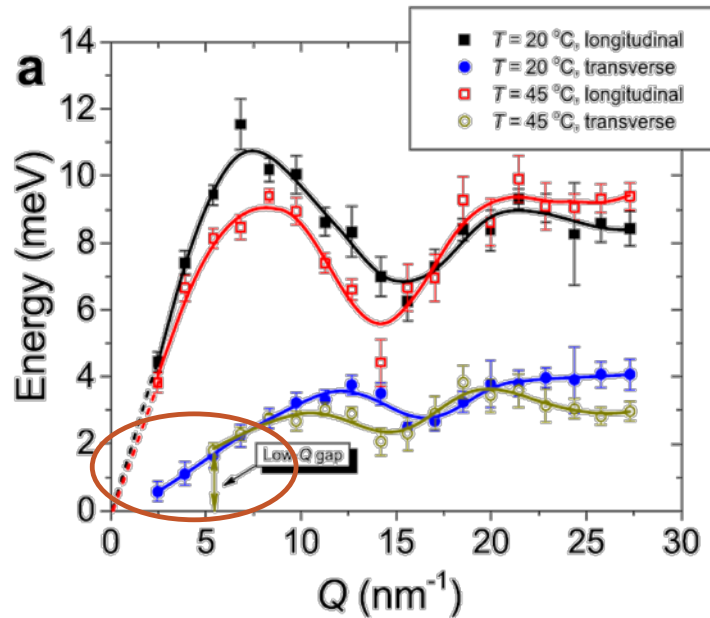
Packing Defect Model
for saturated lipids

IXS measurements: DPPC lipid

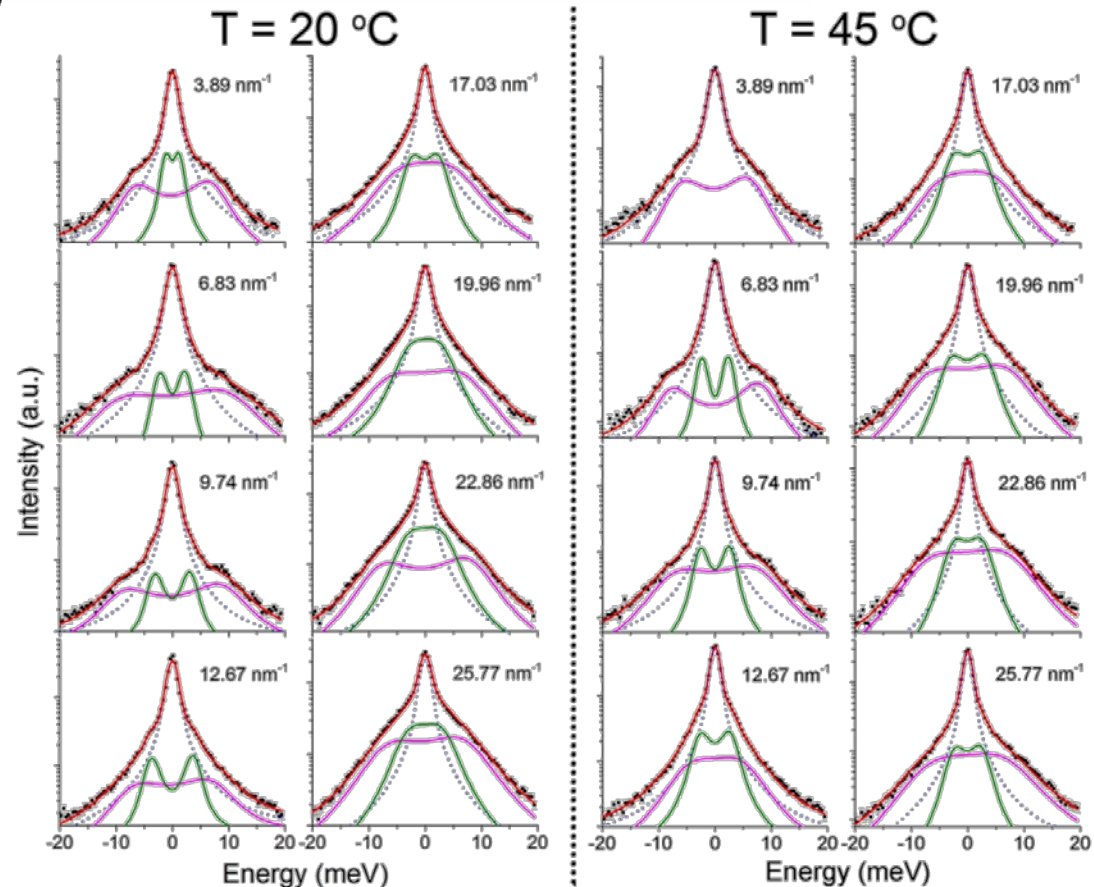


DPPC spin-coated on diamond substrate →
stack of thousands of lipid bilayers

- DPPC main transition temperature: 41 °C
- Measured at 20 °C and 45 °C; E = 21.78 KeV
- Saturated water vapor



M. Zhernenkov et al. *Nat. Commun.* **7**, 11575 (2016)



The discovery of the low-Q phononic gap in TA mode!

Low-Q phononic gap in TA mode

Low-Q phononic gap

The cut-off value Q_{gap}
(at which gap occurs)

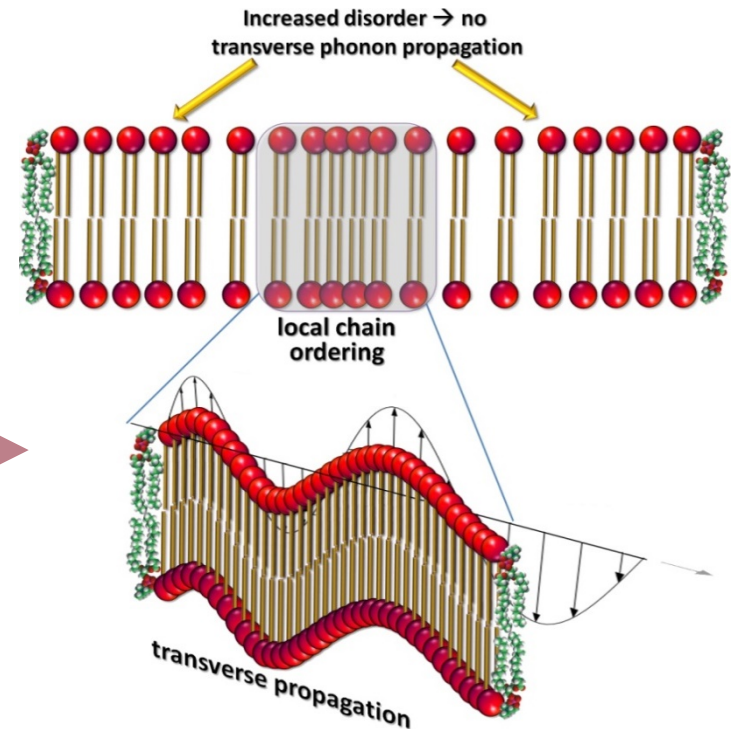
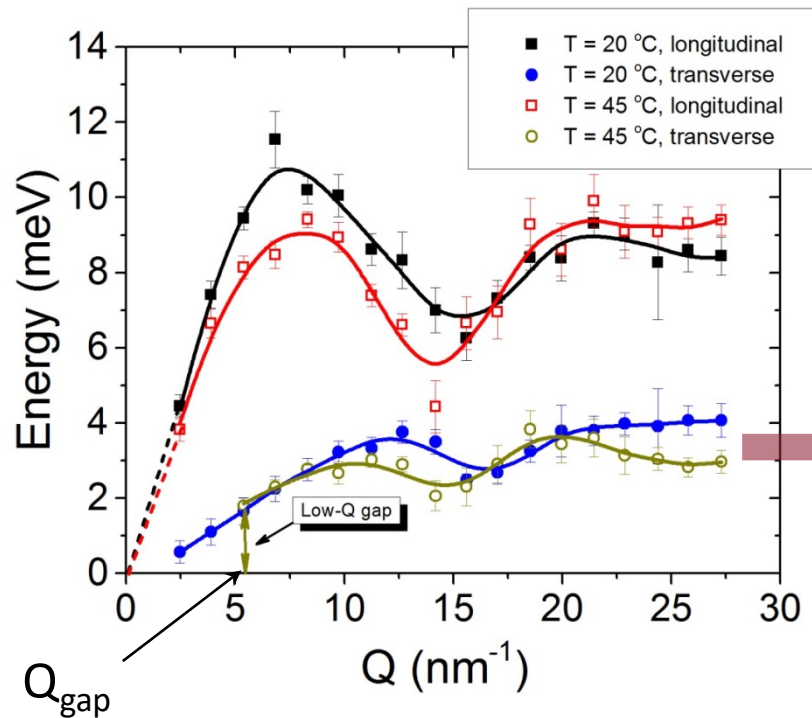
No phonon propagation
at $Q < Q_{\text{gap}}$

Characteristic length-scale

$d_{\text{gap}} = 2\pi/Q_{\text{gap}}$ of a local “order”
(no phonons at smaller $Q \rightarrow$ more disorder on long scale!)

D. Bolmatov, et al. *Ann. Phys.* **363**, 221-242 (2015)

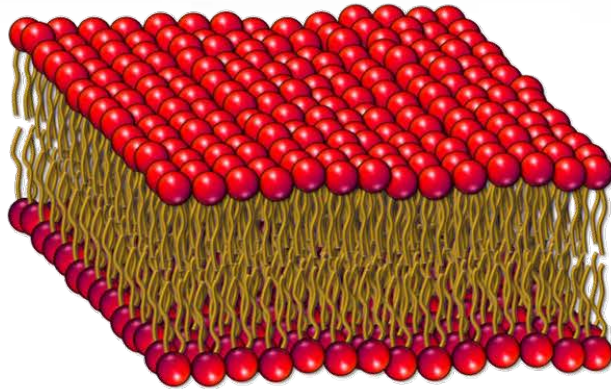
D. Bolmatov, et al. *J. Phys. Chem. Lett.* **6**, 3048-3053 (2015)



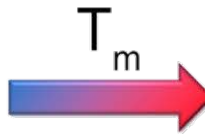
Phonon-mediated nm-scale clustering

NO phononic gap

Tight lipid packing on long length scale



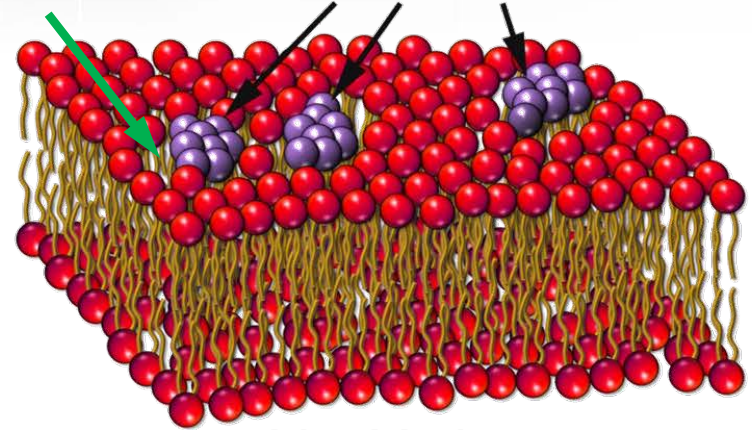
Gel phase



Low-Q phononic gap

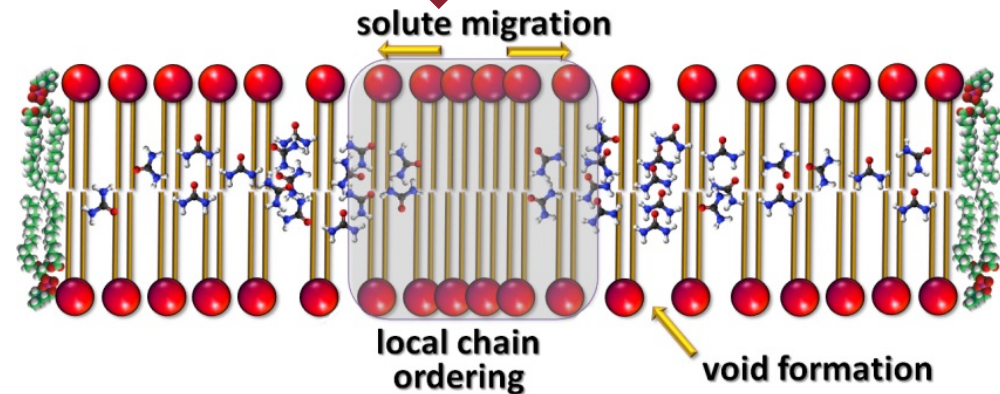
Rarified regions (voids)

Local DPPC clustering



Liquid phase

- the cluster size $d_{\text{gap}} = 2\pi/Q_{\text{gap}}$ is $\sim 1.1\text{-}1.6$ nm
- Lifetime: ~ 1 ps (agrees well with common assumptions – *a few ps*)



Phonon-mediated nm-scale clustering

Theory of solute diffusion through a membrane:

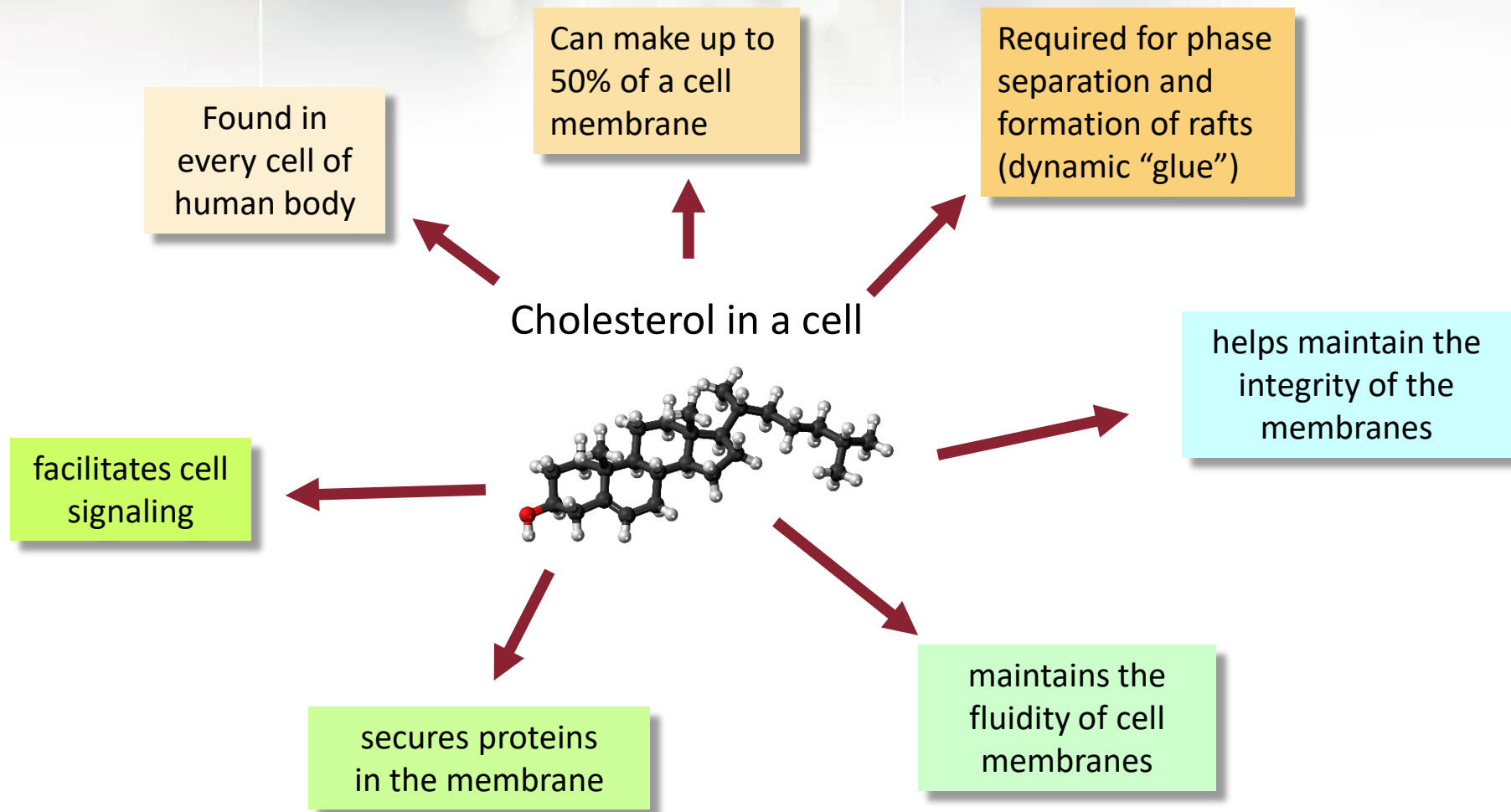
- ultra-fast “hopping”, or “rattling” between thermally-triggered voids
- partition coefficient strongly depends on the local chain ordering → solute exclusion within the region
- Potential formation of water fingers inside voids → proton translocation through membrane

Adv. Drug Deliv. Rev. **58**, 1357–1378 (2006)
J. Am. Chem. Soc. **117**, 4118–4129 (1995)
Cold Spring Harb. Perspect. Biol. (2010), 2, a002188

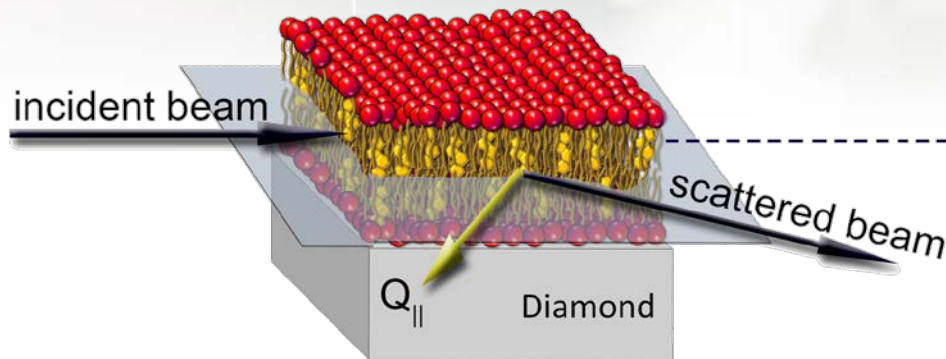
We observe:

- ✓ nm-scaled short-lived molecular clusters → local chain ordering, or density fluctuations
- ✓ Increased disorder beyond the cluster size → indication of the transient voids formation
- ✓ Size and the life time of the clusters agrees well with the theory prediction

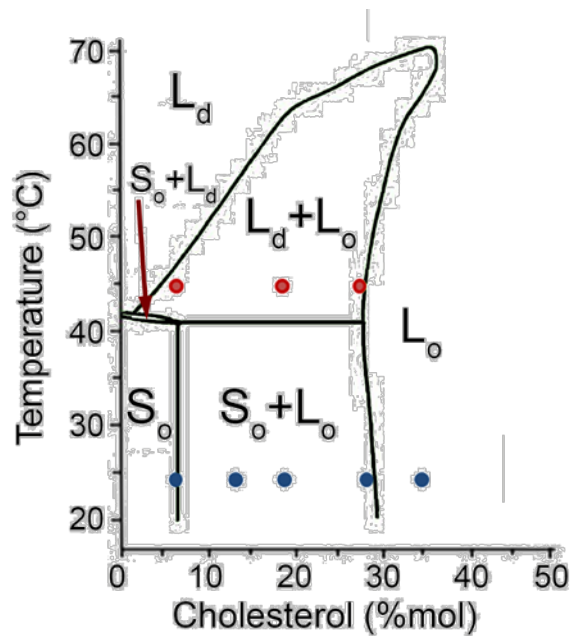
DPPC-Cholesterol binary mixtures



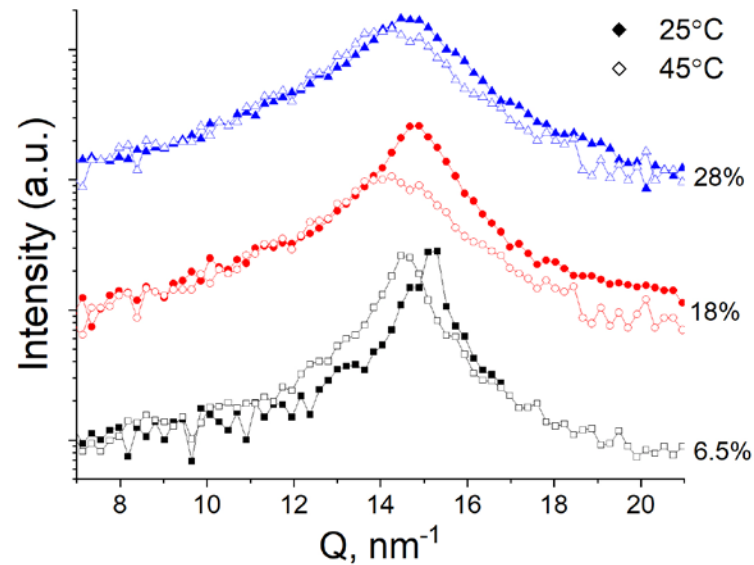
IXS measurements: DPPC-Chol



Top view of a lipid multilayer on diamond

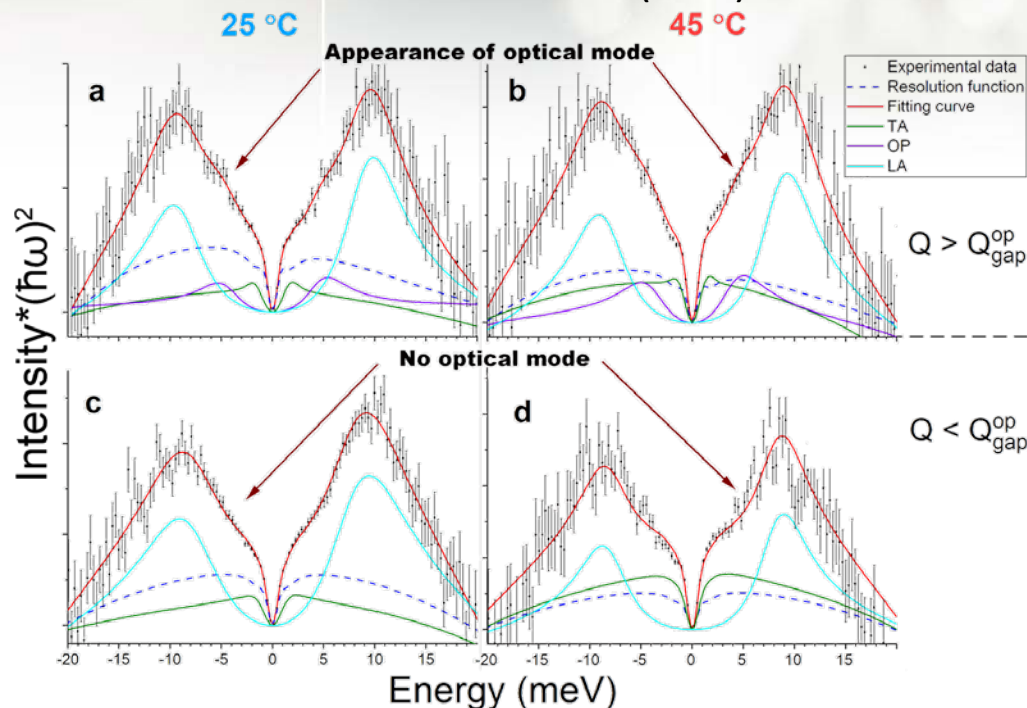


De Almeida & Joly, Front. Plant Sci. (2014)

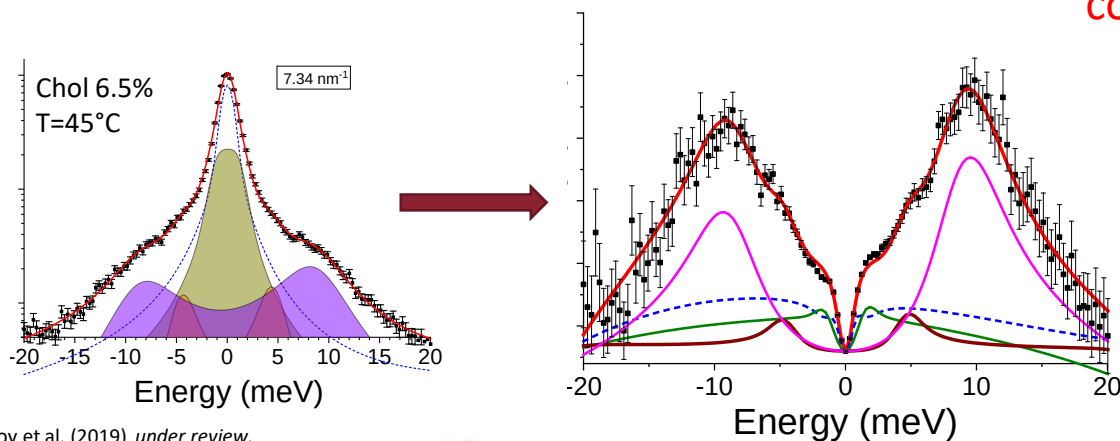


IXS phonon current spectra

IXS currents for Chol(28%)

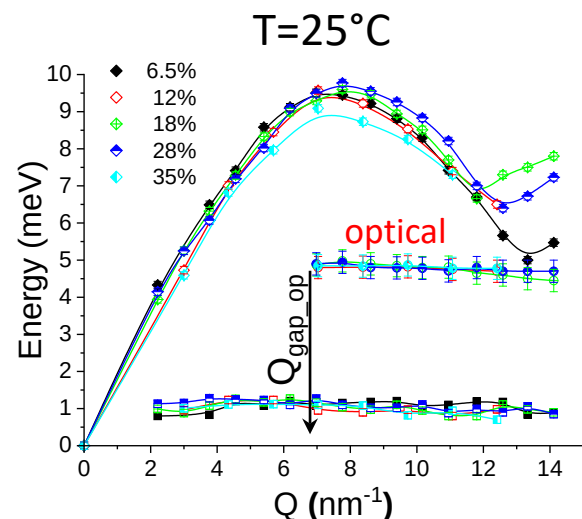
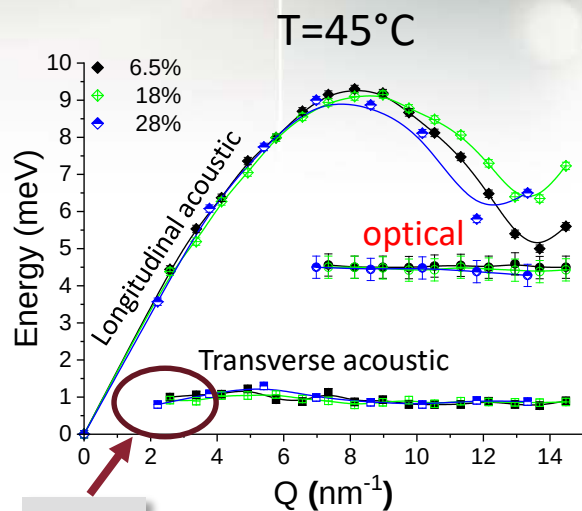


- Two or more “particles” per “unit cell” should generate the optical mode(s) due to out-of-phase movements
- Optical mode emerges at $Q_{\text{gap_op}} \geq 7 \text{ nm}^{-1}$ for all temperatures
- Transverse mode shows no $Q_{\text{gap_ta}}$ within measured Q-range (in contrast to pure DPPC!)



D. Soloviov et al. (2019) under review.

DPPC-Cholesterol binary mixtures

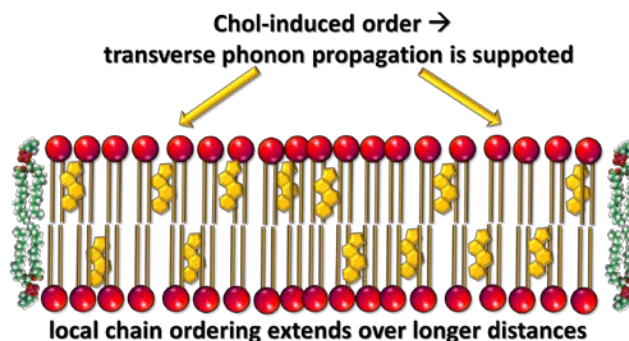


- Longitudinal acoustic (LA) → compression waves (sound) propagating in lipids, no Q_{gap} allowed
- Transverse acoustic (TA) → oscillatory movements of lipids due to shear restoring forces (short range)

No measurable Q_{gap} in TA mode

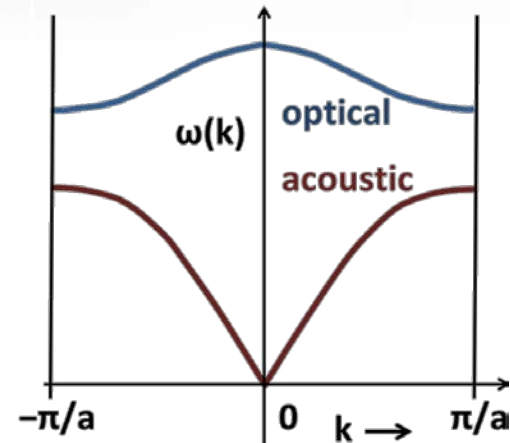
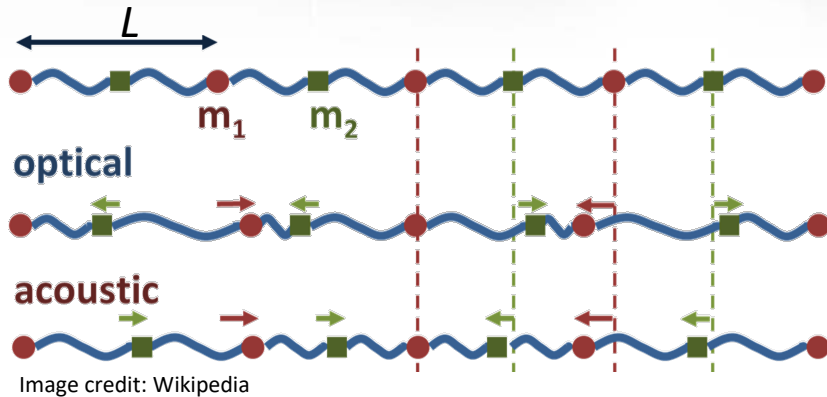
Cholesterol adds rigidity + local ordering

Reduced diffusion rates



Origin of phononic gap in optical mode

Optical and acoustic vibrations in linear diatomic chain



Optical mode frequency:

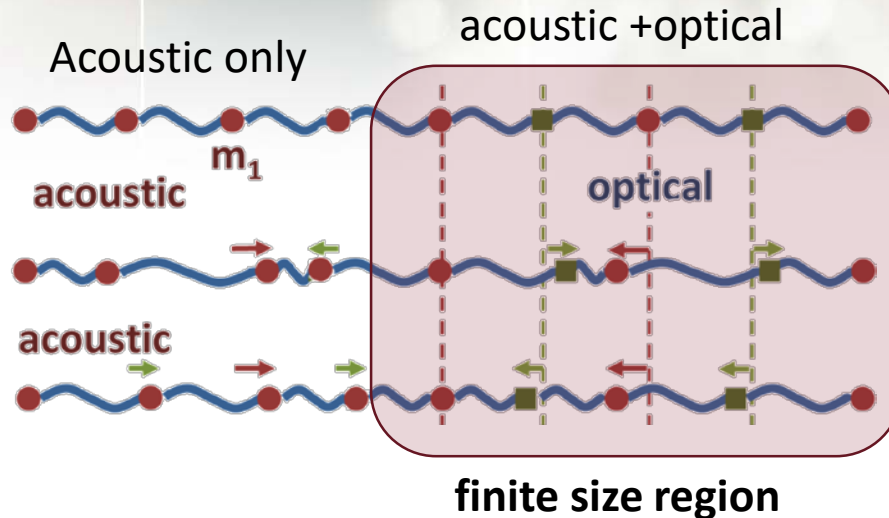
$$\nu^2 = \alpha \left(\frac{1}{M_1} + \frac{1}{M_2} \right) + \alpha \sqrt{\left(\frac{1}{M_1} + \frac{1}{M_2} \right)^2 - \frac{4 \sin^2(QL)}{M_1 M_2}}$$

at $Q \rightarrow 0$ $\nu^2 \sim 2\alpha \left(\frac{1}{M_1} + \frac{1}{M_2} \right)$

α - interparticle force constant

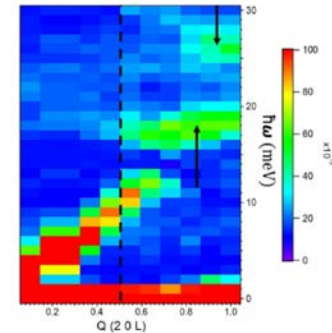
Under a continuous medium approximation ($Q \rightarrow 0$)
the optical mode cannot exhibit the phononic gap Q_{gap} !

Phononic gap in optical mode: finite size effect



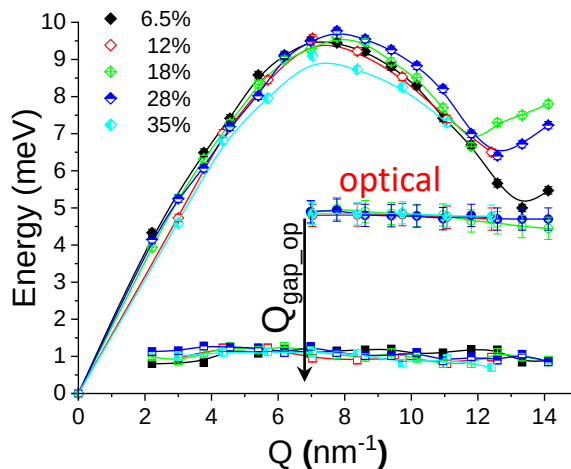
Recent example

Finite size effect on phonons of the quasi-one-dimensional system SrCuO_2 : doping breaks symmetry and suppresses optical mode at low Q



D. Bounoua et al. Physica B 536 (2018) 323–326

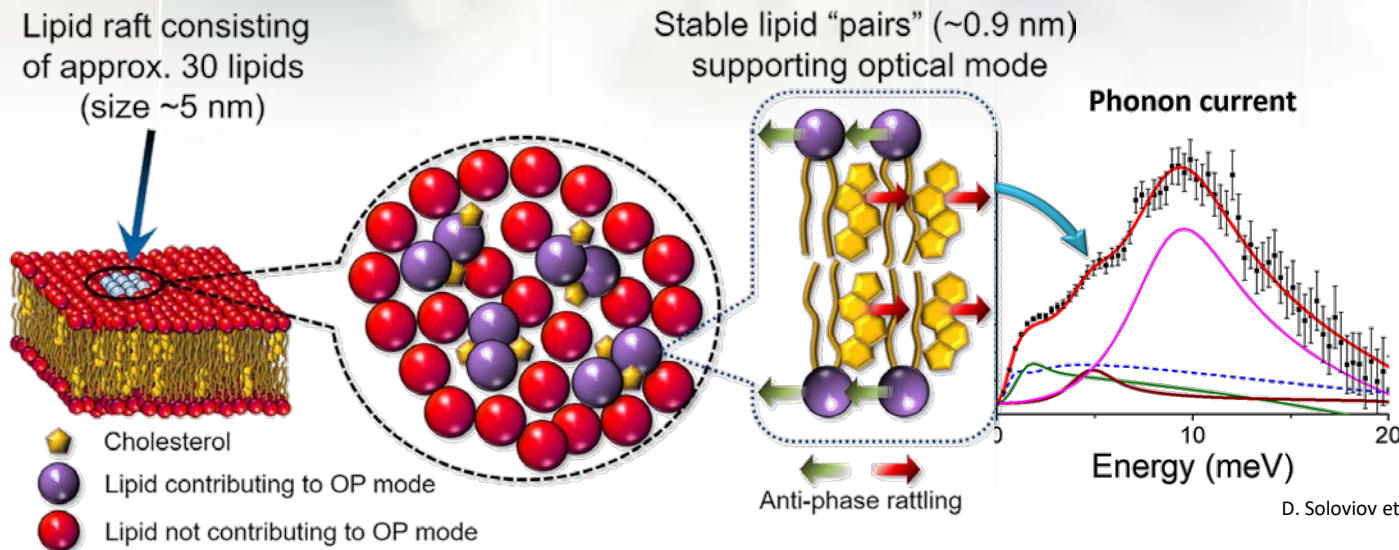
DPPC-Cholesterol case



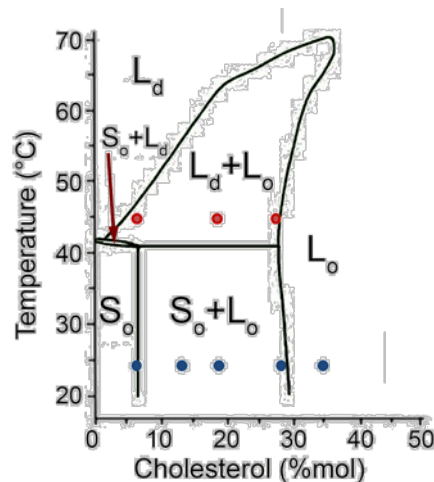
D. Soloviov et al. (2019) *under review*.

- Two lipid “particles” per “unit cell”
- $Q_{\text{gap_op}}$ in lipids $\rightarrow$ finite size regions throughout the sample supporting optical mode
- Size of the region: $2\pi/Q_{\text{gap_op}} \sim 9 \text{ \AA}$ or several lipids in diameter!

Phononic gap in optical mode: DPPC-Chol “pairs”



L_o phase → set of discrete patches with 2-3 coupled lipid pairs



De Almeida & Joly, Front. Plant Sci. (2014)

- OP mode (~5 meV) observed when L_o phase is present
- Optical phonon energy estimation for DPPC-Chol pair:
4-7.5 meV

- DPPC-Chol interaction energy (1-3.6 kT)*
- Simple Hooke's law
- $\nu^2 \sim 2\alpha(\frac{1}{M_1} + \frac{1}{M_2})$

*Heberle, F.; Feigenson, G. Cold Spring Harb. Perspect. Biol. (2011)
Zeno, W. et al. Langmuir (2016)

Ternary systems: DOPC/POPC-based

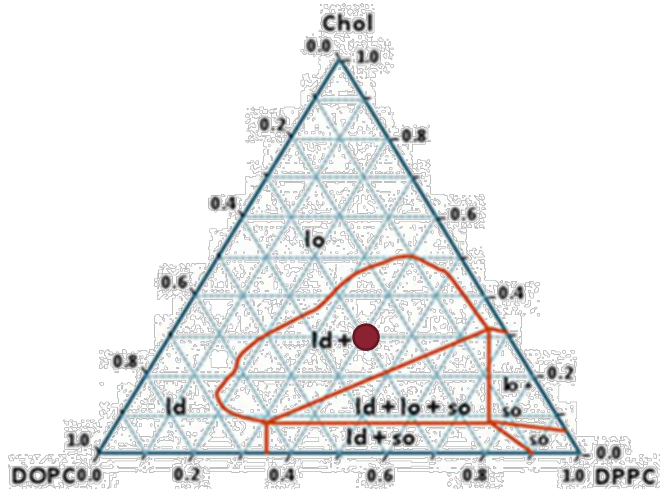
Ternary systems

3 lipid “particles”

3 combinations of “pairs”
(but not all “pairs” are equally favored!)

More optical modes allowed

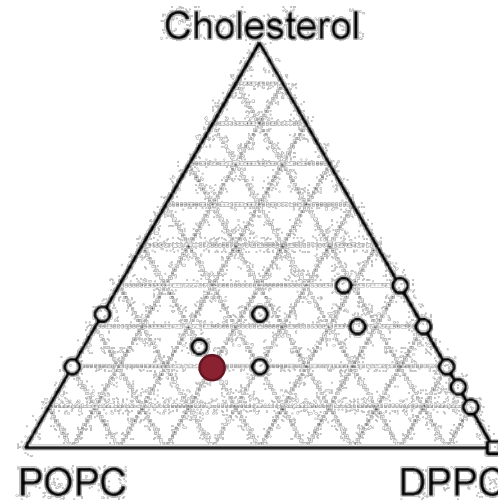
Expectation: in reality, optical modes should be mixed in soft matter



de Almeida&Joly, Front. Plant. Sci (2014)

DOPC(3):DPPC(4):Chol(3)

Lo+Ld

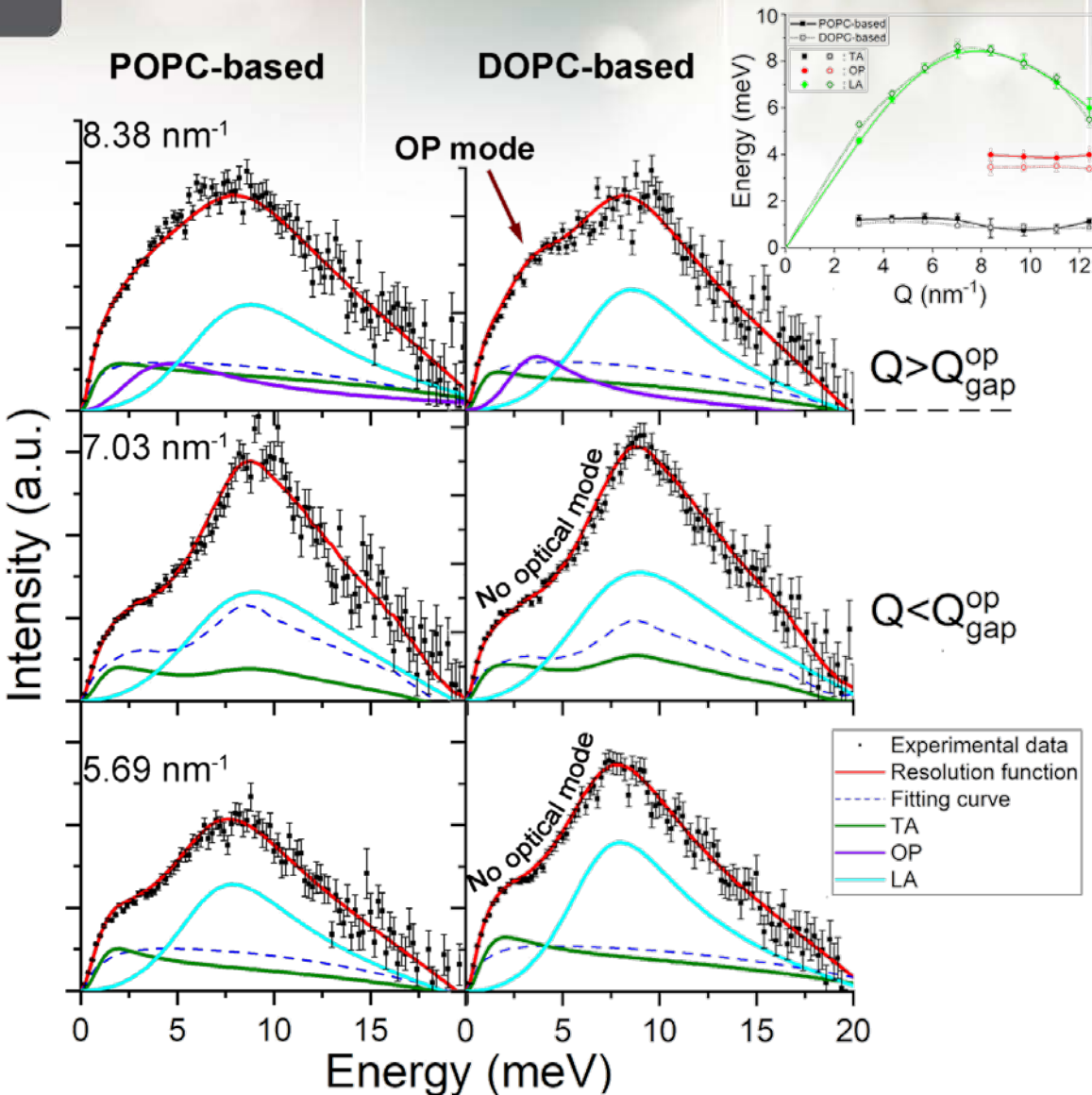


Veatch&Keller, Phys.Rev.Lett. (2005)

POPC(5):DPPC(3):Chol(2)

Coexisting liquid domains
are not previously observed

IXS: ternary systems at 37°C



D. Soloviov et al. (2019) *under review*.

- Optical mode emerges at $Q_{\text{gap_op}} \geq 8.3 \text{ nm}^{-1}$
- OP energy: 3.5-3.9 meV
- Size of the region: $2\pi/Q_{\text{gap_op}} \sim 7.5 \text{ \AA}$ or several lipids in diameter

- OP energy estimation for DOPC/POPC-DPPC pairs: 1.6 meV (unfavorable pairs)
- Unsaturated lipids reduce the size of discrete patches!

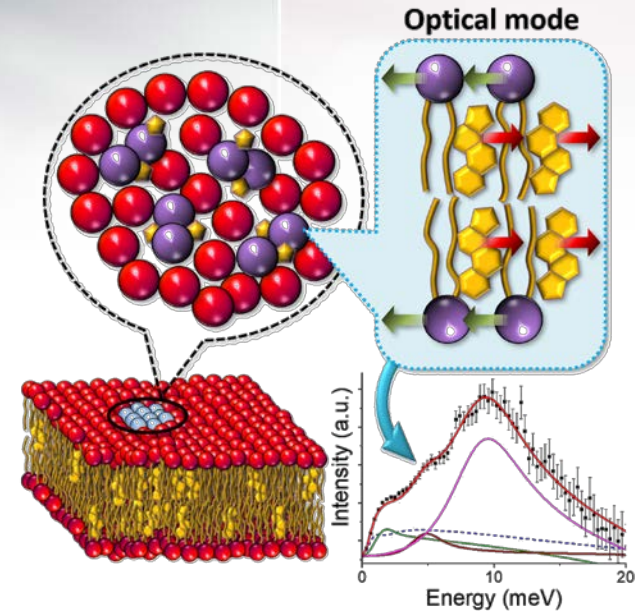


- No patches in Ld phase with excess of unsaturated lipids

Summary

Observations:

- Multicomponent phase separated membranes exhibit optical modes
- Optical modes are supported on a short length scale → indication of discrete nano-patches (*finite size effect*), or lipid “pairs”
(only groups of two+ lipids can support optical mode!)
- Optical mode only observed in phase coexistence, or when L_o phase is present



Hypothesis: in phase separated domains, a single domain with typical sizes of 10-200 nm on nanoscale consists of a set of mobile coupled discrete areas <1 nm in diameter. The large domains form when “pairs” transiently coalesce into larger entities.

- In multicomponent mixtures: different combination of “pairs” can couple due to the selective affinity between different lipid species

Acknowledgements



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THANK YOU!