

Three-Loop Contributions to $\Delta\rho$ and Δr for large m_t

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IV Weak Corrections:

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V Summary

I Introduction

central prediction of SM:

$$M_W = f(\underbrace{G_F, M_Z, \alpha}_{\text{Born}}; \underbrace{M_t, M_H, \dots}_{\text{radiative corrections}})$$

[similarly for couplings of fermions: $\sin^2 \theta_{\text{eff}}$]

$$M_W^2 \left(1 - \frac{M_W^2}{M_Z^2} \right) = \frac{\pi \alpha}{\sqrt{2} G_F} (1 + \Delta r)$$

$$\Delta r \text{ dominated by: } \Delta r = -\frac{c^2}{s^2} \Delta \rho + \Delta \alpha$$

$$\Delta \rho \sim (G_F M_t^2) + \dots$$

aim: experiment

δM_W [MeV]	$\delta \sin^2 \theta_{\text{eff}} \times 10^{-5}$	
33	17	status
12 – 16	20	Tevatron (run II), LHC
7	1.3	TESLA (GigaZ)

requires theory uncertainty better than 5 MeV!

correlation with δM_t :

$$\begin{aligned}\delta M_t = 5 \text{ GeV} &\Rightarrow \delta M_W = 32 \text{ MeV} \\ &\Rightarrow \delta \sin^2 \theta_{\text{eff}} = 16 \cdot 10^{-5}\end{aligned}$$

Tevatron (run II) 2.7 → 1.4 GeV

LHC 1.0 GeV

LC 0.2 GeV

remaining uncertainties:
SNOWMASS (Baur et al)

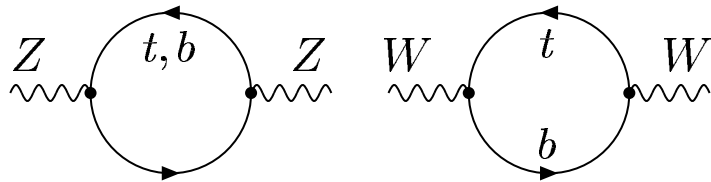
	order	sector	estimate	size ($\times 10^5$)	M_W [MeV]
2-loop {	α^2	fermionic	$N(\alpha/4\pi\hat{s}^2)^2$	8.7	complete [1]
	α^2	bosonic	$(\alpha/\pi\hat{s}^2)^2$	11.6	2.1
3-loop {	$\alpha\alpha_s^2$	top-bottom doublet	$N_C C_F C_A \alpha\alpha_s^2/4\pi^3\hat{s}^2$	4.7	complete [2]
	$\alpha\alpha_s^2$	light doublets	$2 N_C C_F C_A \alpha\alpha_s^2/4\pi^3\hat{s}^2$	9.4	complete
	$\alpha^3 m_t^6$	heavy top	$5.3 N_C^2 (\alpha m_t^2/4\pi\hat{s}^2 M_W^2)^3$	7.0	4.1
4-loop {	$\alpha^3 m_t^6$	heavy top	$3.3 N_C (\alpha m_t^2/4\pi\hat{s}^2 M_W^2)^3$	1.5	0.9
	$\alpha^2 \alpha_s m_t^4$	heavy top	$3.9 N_C C_F \alpha^2 \alpha_s m_t^4/16\pi^3\hat{s}^4 M_W^4$	7.8	4.5
	$\alpha\alpha_s^3 m_t^2$	heavy top	$N_C C_F C_A^2 \alpha\alpha_s^3 m_t^2/4\pi^4\hat{s}^2 M_W^2$	2.3	1.3
		total			7

complete
now

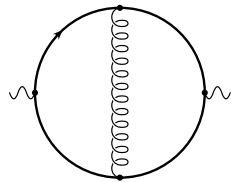
[1] (Freitas, Hollik, Walter, Weiglein)

[2] (Chetyrkin, JK, Steinhauser)

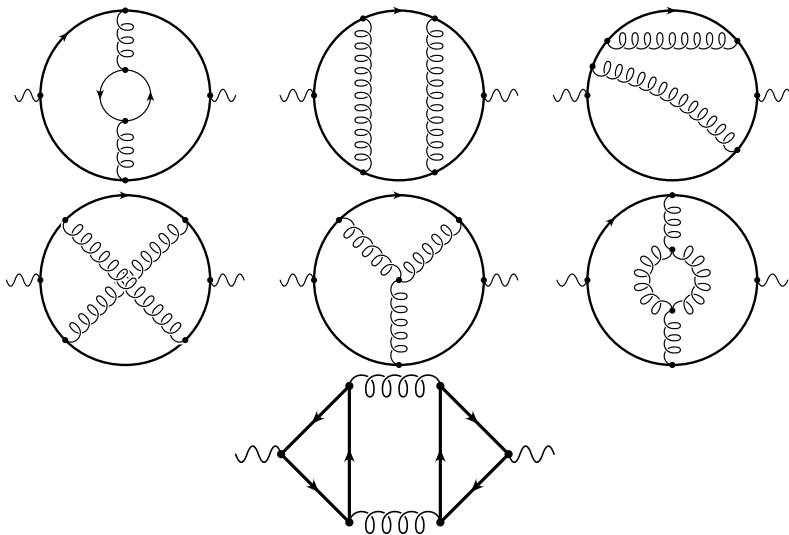
II Three-Loop QCD-Results



one-loop
(Veltman)



two-loop
(Djouadi; Kniehl, JK, Stuart)



three-loop
(Chetyrkin, JK, Steinhauser)

evaluation:

$$\Delta\rho = \frac{\Pi_{ZZ}(0)}{M_Z^2} - \frac{\Pi_{WW}(0)}{M_W^2}$$

$$\begin{aligned}\Delta\bar{r}_{tb} &\equiv \Delta r_{tb} - \tilde{\Pi}_{\gamma\gamma}(0) + \text{Re}\tilde{\Pi}_{\gamma\gamma}(M_Z^2) \\ &= \frac{c^2}{s^2} \left(\frac{\Pi_{WW}(0)}{M_W^2} - \frac{\Pi_{ZZ}(0)}{M_Z^2} \right) \\ &\quad + \frac{c^2}{s^2} \left(\Pi'_{WW}(0) - \Pi'_{ZZ}(0)|_t - \text{Re} \frac{\Pi_{ZZ}(M_Z^2)|_b}{M_Z^2} \right) \\ &\quad \quad \quad + \tilde{\Pi}_{\gamma\gamma}(0)|_t + \text{Re}\tilde{\Pi}_{\gamma\gamma}(M_Z^2)|_b - \Pi'_{WW}(0) \\ &\quad + \frac{c^2}{s^2} \left(\Pi''_{WW}(0) \frac{M_W^2}{2} - \Pi''_{ZZ}(0)|_t \frac{M_Z^2}{2} \right) + \tilde{\Pi}'_{\gamma\gamma}(0)|_t M_Z^2 - \Pi''_{WW}(0) \frac{M_W^2}{2} \\ &\quad + \dots\end{aligned}$$

(Chetyrkin, JK, Steinhauser; Avdaev, Fleischer, Michailov, Tarasov)

tadpoles! expansion in $\frac{M_W^2}{M_t^2}$!

Result:

δM_W in MeV	α_s^0	α_s^1	α_s^2
M_t^2	611.9	-61.3	-10.9
const.	136.6	-6.0	-2.6
$1/M_t^2$	-9.0	-1.0	-0.2

III Weak Corrections: Technicalities

approximation: $M_t^2 \gg M_W^2 \Rightarrow$ scalar bosons only

status: two-loop $\left\{ \begin{array}{l} \text{Barbieri, Beccaria, Ciafaloni, Curci, Vicere} \\ \text{Fleischer, Jegerlehner, Tarasov} \end{array} \right.$

$$\Delta\rho = X_t^2 f(M_t/M_H) \quad \text{with} \quad X_t \equiv \frac{G_F M_t^2}{8\sqrt{2}\pi^2} = \frac{g_{\text{Yukawa}}^2}{16\pi^2}$$

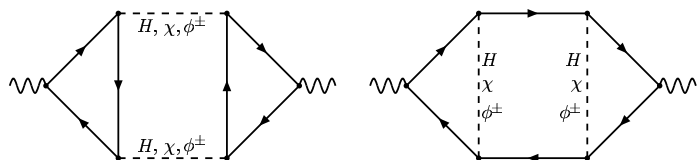
status: three-loop

$M_H = 0$ (van der Bij, Chetyrkin, Faisst, Jikia, Seidensticker)

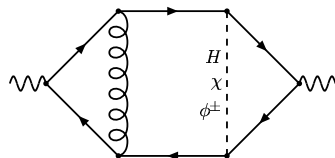
X_t^3 and $\alpha_s X_t^2$ tiny corrections

$M_H \neq 0$ (Faisst, JK, Seidensticker, Veretin)

requires: 3-loop tadpoles ($\Pi_{WW}(0)$, $\Pi_{ZZ}(0)$)

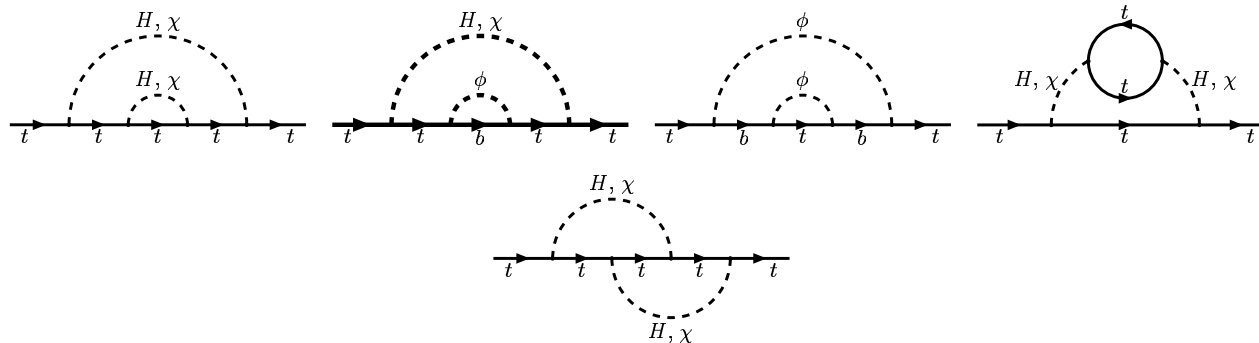


$$X_t^3 f(M_t/M_H)$$



$$\alpha_s X_t^2 f(M_t/M_H)$$

and two-loop on-shell diagrams: $M_t \iff m_t(\overline{\text{MS}})$



in general two (!) mass scales, three loops

special cases:

$M_H = 0$: $\checkmark$

$M_H \gg M_t$: hard mass expansion in $(M_t^2/M_H^2)^n$ mod. log

$M_H = M_t$: one scale

M_H in neighbourhood of M_t :

Taylor expansion: $\delta = (M_H - M_t)/M_t$

$\Rightarrow$ reduction to one-scale two- or three-loop integrals

Algebraic programs

- generating diagrams: **QGRAF** (Nogueira)
- hard-mass expansion: **EXP** (Seidensticker)
- Taylor-expansion / algebraic evaluation of tadpoles:
MATAD (Steinhauser)
- massless propagators:
MINCER (Larin, Tkachov, Vermaseren)
- on-shell diagrams:
TARCER (Mertig, Scharf),
ONSHELL2 (Fleischer, Kalmykov)

IV Results

define: $X_t = \frac{G_F M_t^2}{8\sqrt{2}\pi^2} \approx 3.2 \cdot 10^{-3}$ and $\frac{\alpha_s}{\pi} \approx 3.5 \cdot 10^{-2}$

• $M_H = 0$

$$\begin{aligned} \Delta\rho^{(3\text{ Loop})} = & X_t^3 n_c \left(68 + 96\zeta(2) \log 2 + 6\zeta(2) - 612\zeta(3) \right. \\ & \left. + 729S_2 + 324\zeta(4) - 72B_4 + 36D_3 \right) \\ & + X_t^3 n_c^2 \left(-\frac{6572}{15} - \frac{4374}{5}S_2 + \frac{1472}{15}\zeta(2) + 440\zeta(3) \right) \\ & + X_t^2 \frac{\alpha_s}{\pi} n_c C_F \left(\frac{185}{3} + \frac{729}{4}S_2 - 48\zeta(2) \log 2 \right. \\ & \left. - \frac{151}{6}\zeta(2) + 29\zeta(3) - 24\zeta(4) + 12B_4 \right) \end{aligned}$$

Numerically:

$$\begin{aligned} \Delta\rho^{(3\text{ Loop})} = & 3 X_t - 2.2176 X_t^2 + 249.74 X_t^3 \\ & - 8.5797 X_t \frac{\alpha_s}{\pi} + 2.9394 X_t^2 \frac{\alpha_s}{\pi} \Leftarrow \text{accidentally small} \end{aligned}$$

- $M_H \approx M_t$: Taylor expansion in $\delta = \frac{M_H - M_t}{M_t}$

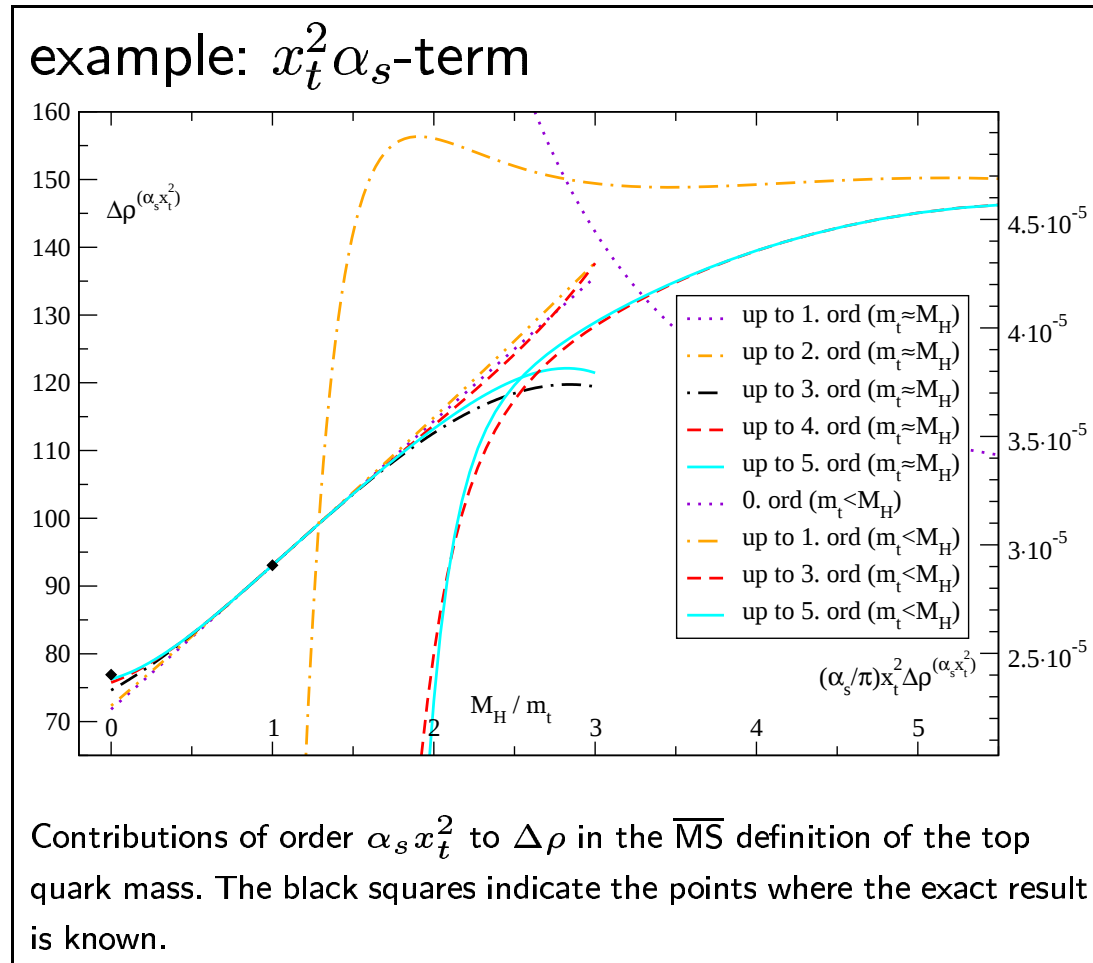
$$\Delta\rho^{(3\text{ Loop})} = X_t^3 (95.92 - 111.98\delta + 8.099\delta^2 + 9.36\delta^3 + \dots) \\ + X_t^2 \frac{\alpha_s}{\pi} (157.295 + 112.00\delta - 24.73\delta^2 + 7.39\delta^3 + \dots)$$

- $M_H \gg M_t$: asymptotic expansion in $Y = \frac{4M_t^2}{M_H^2}$

$$\Delta\rho^{(3\text{ Loop})} = X_t^3 \left(Y^{-1} (-3.17 - 83.25 \log Y) \right. \\ + 1 (-189.93 - 231.48 \log Y - 142.06 \log^2 Y + 2.75 \log^3 Y) \\ + Y (-332.34 + 77.71 \log Y - 68.67 \log^2 Y + 51.79 \log^3 Y) \\ \left. + Y^2 (227.55 - 510.55 \log Y + 87.77 \log^2 Y + 6.41 \log^3 Y) \right) \\ + X_t^2 \frac{\alpha_s}{\pi} \left(1 (79.73 - 47.77 \log Y + 42.07 \log^2 Y + 9.00 \log^3 Y) \right. \\ + Y (225.16 - 179.74 \log Y + 70.22 \log^2 Y - 19.22 \log^3 Y) \\ \left. + Y^2 (-76.07 + 25.33 \log Y - 9.17 \log^2 Y - 5.57 \log^3 Y) \right) + \dots$$

two steps:

- $\Delta\rho$ in $\overline{\text{MS}}$ -scheme

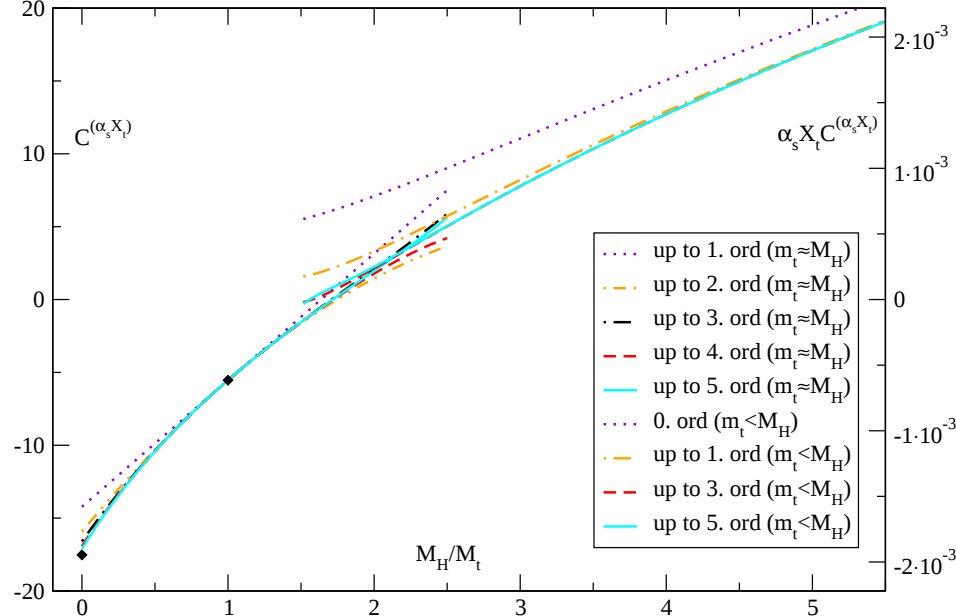


Expansion around $M_H/M_t = 1$ approaches the result at $M_H = 0$.

The two expansions agree for $M_H/m_t \approx 2$.

- transition from $\overline{\text{MS}}$ to on-shell top quark mass

$$\frac{m_t}{M_t} = 1 + \frac{\alpha_s}{\pi} C(\alpha_s) + X_t C(X_t) + \left(\frac{\alpha_s}{\pi}\right)^2 C(\alpha_s^2) + \frac{\alpha_s}{\pi} X_t C(\alpha_s X_t) + X_t^2 C(X_t^2) + \dots$$



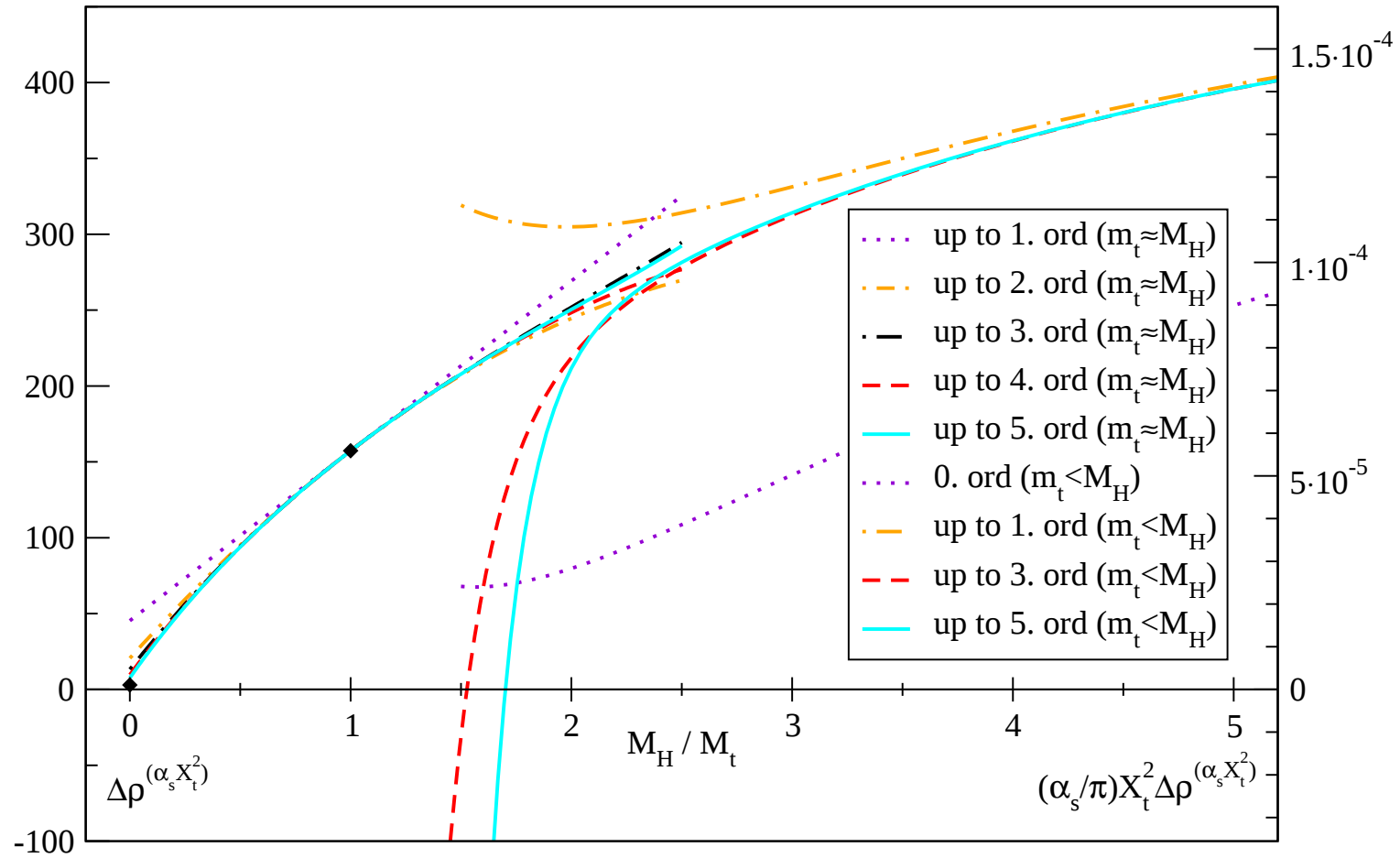
Contributions of order $\alpha_s x_t$ to the relation between the $\overline{\text{MS}}$ and on-shell top quark mass. The black squares indicate the points where the exact result is known.

smooth behaviour
in overlapping regions

corresponds to a shift
between M_t and m_t by
nearly 200 MeV!

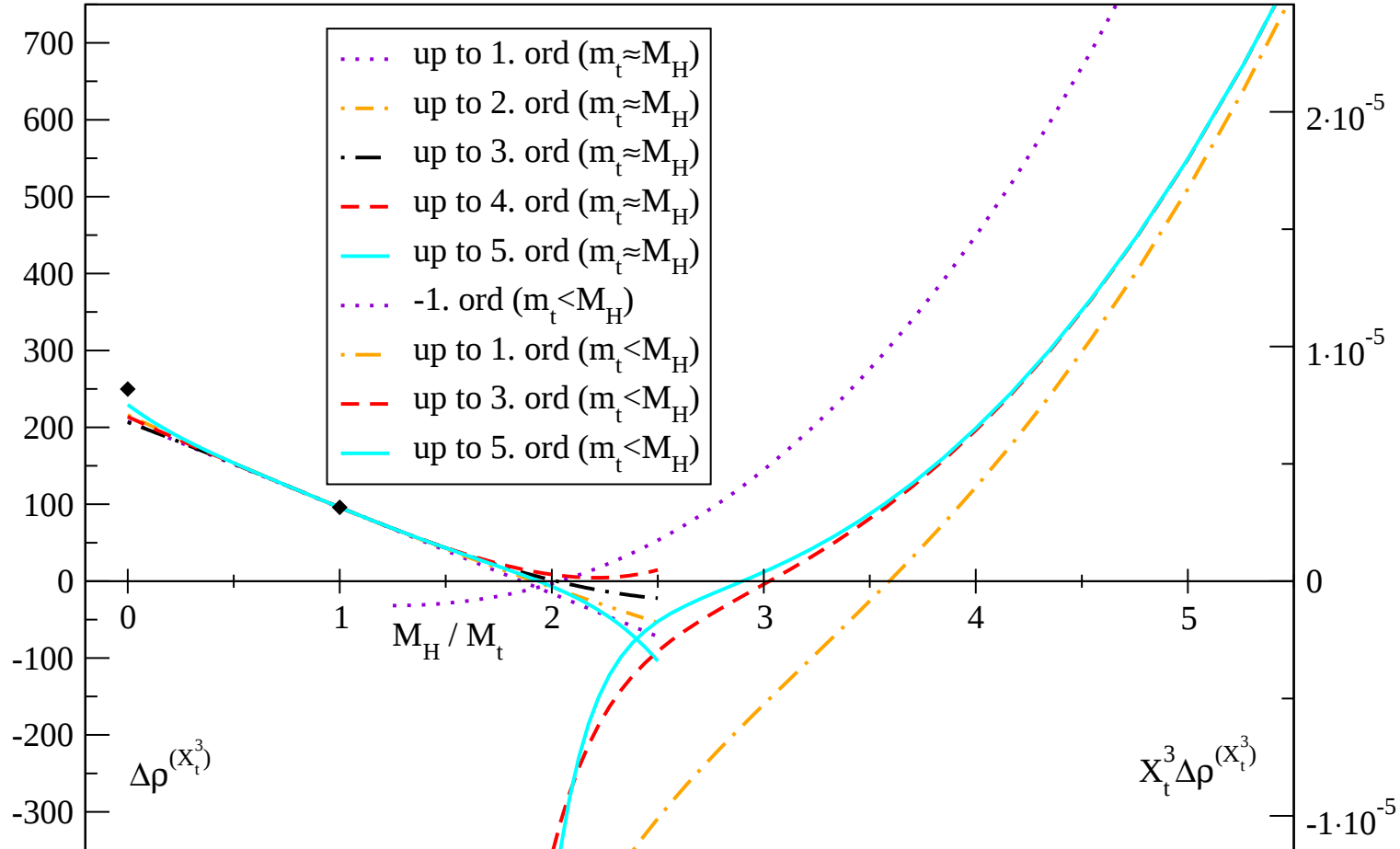
relevance for threshold
measurement!

on-shell result

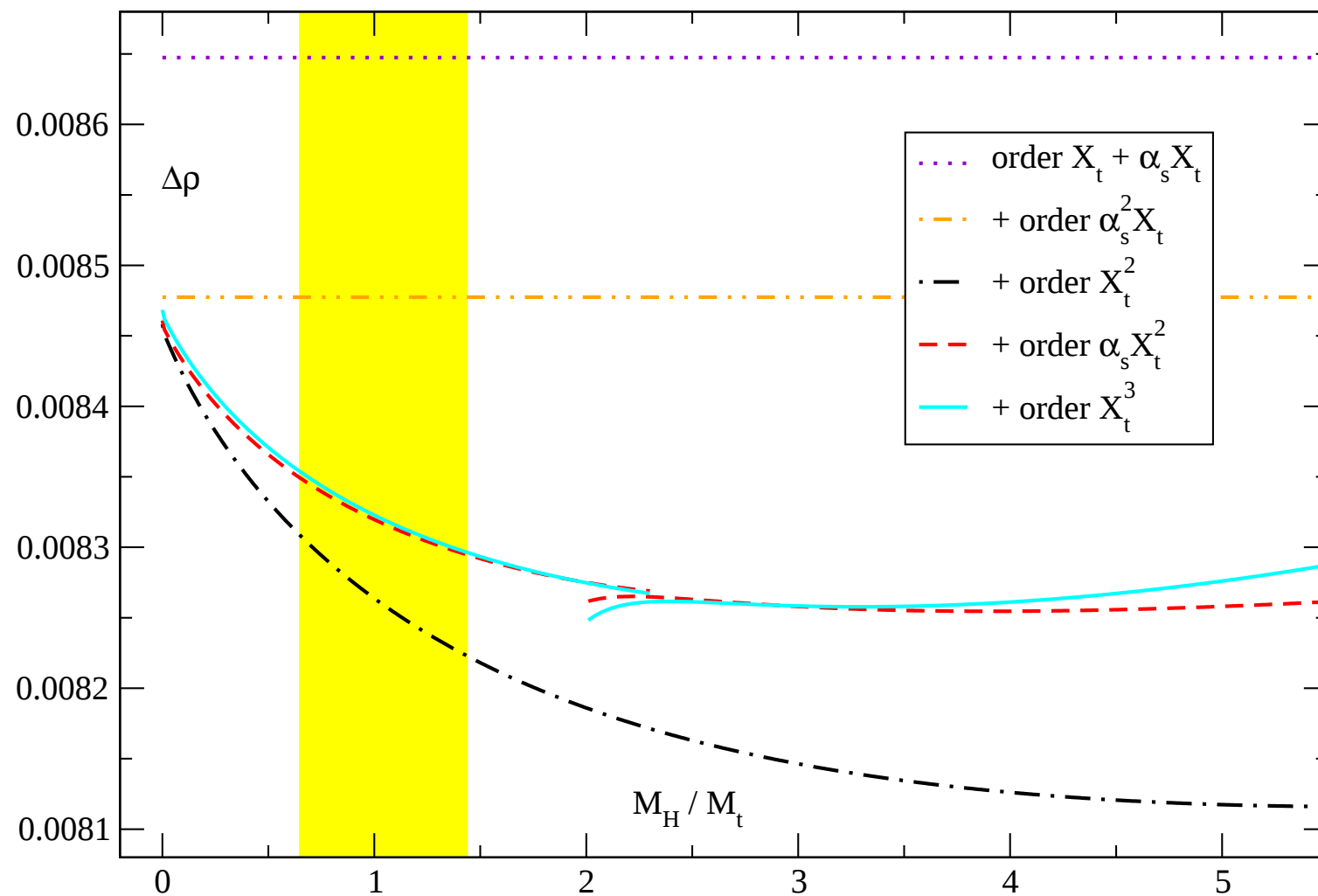


Contributions of order $\alpha_s X_t^2$ to $\Delta\rho$ in the on-shell definition of the top quark mass. The black squares indicate the points where the exact result is known.

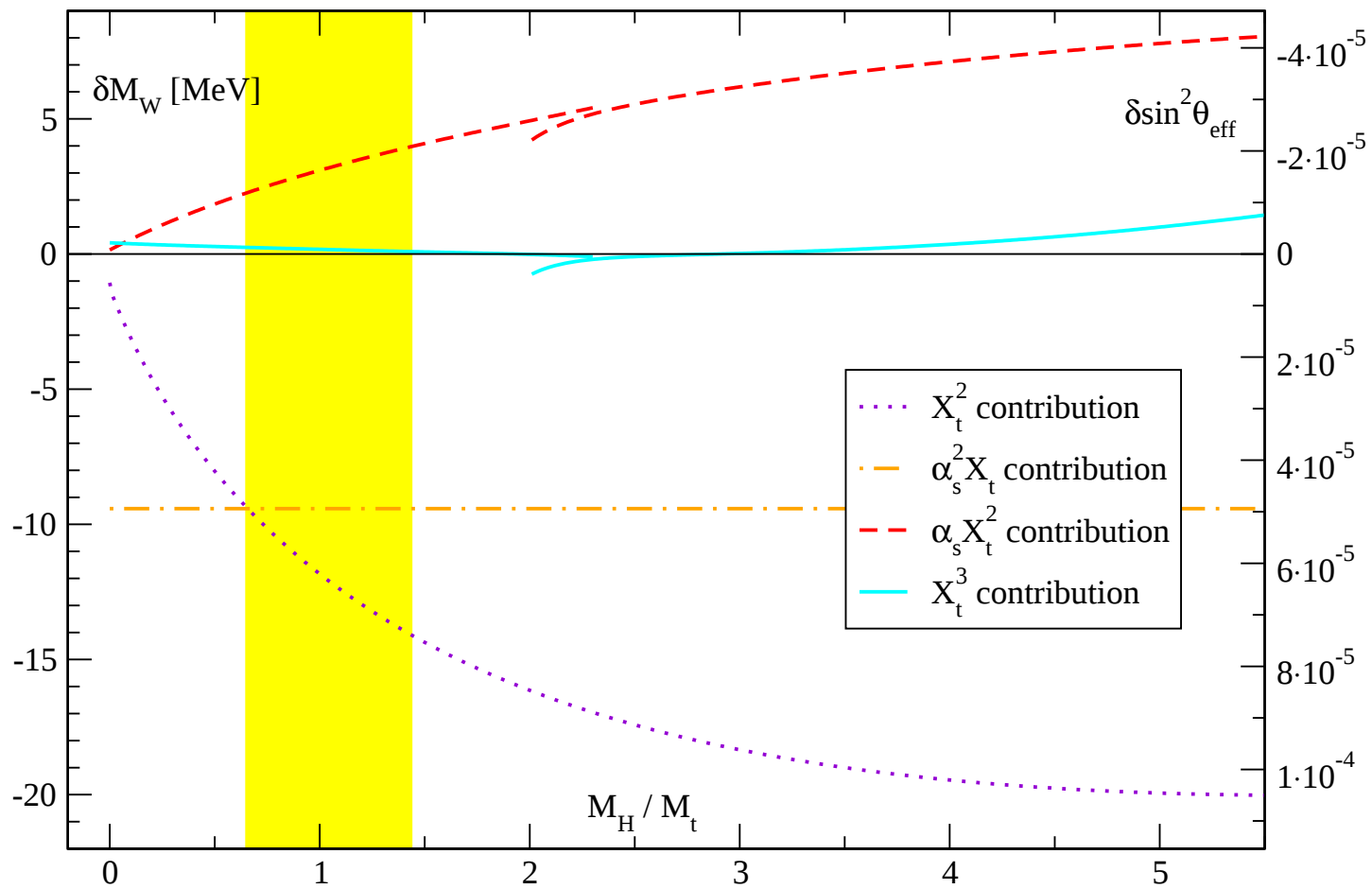
similarly for purely weak corrections



Contributions of order X_t^3 to $\Delta\rho$ in the on-shell definition of the top quark mass. The black squares indicate the points where the exact result is known.



Effect of several contributions to $\Delta\rho$ in the on-shell definition of the top quark mass.



δM_W between 2 and 4 MeV.

$\delta \sin^2 \theta_{\text{eff}}$ between 1 and $2 \cdot 10^{-5}$

V Summary

- dominant electroweak and mixed three-loop corrections are under control
- $\alpha_s X_t^2$ -term is comparable to GigaZ precision
- for fixed M_W , M_t upward shift of M_H by 5 GeV
- for fixed M_W , $M_H \rightarrow$ shift of M_t by 500 MeV
- corrections comparable to non-enhanced 2-loop effects in Δr

$\alpha_s X_t^2 \sim$ fermionic terms

$X_t^3 \sim$ bosonic terms

- essential for LC, about a factor 2–4 below LHC