

NYSDS

**October 23
2020**

Many-Body
Quantum
Systems as a
Challenge for
Machine Learning

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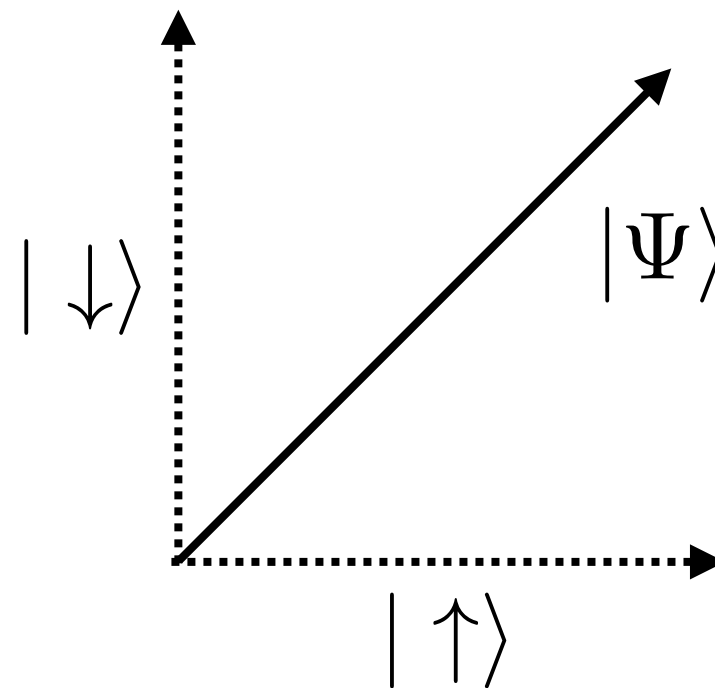
EPFL

A Major
Problem In
Quantum
Physics

The State of a Two-Level System

The state of a spin/qubit/etc is a complex-valued vector (the wave function)

$$|\Psi\rangle = c_{\uparrow}|\uparrow\rangle + c_{\downarrow}|\downarrow\rangle$$



Probability of Observing a Given State

$$P(\uparrow) = |c_{\uparrow}|^2$$

$$P(\downarrow) = |c_{\downarrow}|^2$$

A quantum spin can be found in either up or down state with a given probability

The State of a Many Body System

The wave function is a vector
in a huge (2^N)
Complex Vector Space

$$|\Psi\rangle = c_{\uparrow\uparrow\uparrow\dots} |\uparrow\uparrow\uparrow\dots\rangle + c_{\downarrow\uparrow\uparrow\dots} |\downarrow\uparrow\uparrow\dots\rangle + \dots + c_{\downarrow\downarrow\downarrow\dots} |\downarrow\downarrow\downarrow\dots\rangle$$

• Complex-Valued Coefficients

Schroedinger's Equation

$$\mathcal{H} |\Psi\rangle = E |\Psi\rangle$$

Eigenvalue Problem
for given Hamiltonian
(sparse matrix)

Quantum Many-Body Problem

[3000 BCE]

Papyrus



10 Qubits

[1455]

Book



15 Qubits

[1973]

**IBM
3340**



23 Qubits

[1993]

**IBM
3390**



35 Qubits

[2002]

**Earth
Simulator**



46 Qubits

[2019]

Summit



54 Qubits

Is this complexity “truly”
exponential for physical systems?

Corners of Hilbert space

Hilbert Space

Physical
States

Hilbert Space

Our Best Hope: Variational Formulation

$$E(\mathbf{W}) = \frac{\langle \Psi(\mathbf{W}) | \mathcal{H} | \Psi(\mathbf{W}) \rangle}{\langle \Psi(\mathbf{W}) | \Psi(\mathbf{W}) \rangle} \geq E_0$$

Rayleigh Quotient

Exact Ground-State Energy

Two ML-Inspired Approaches In This Talk

Classical Variational States

Require a CPU/GPU/TPU...

Quantum Variational States

Require a QPU...

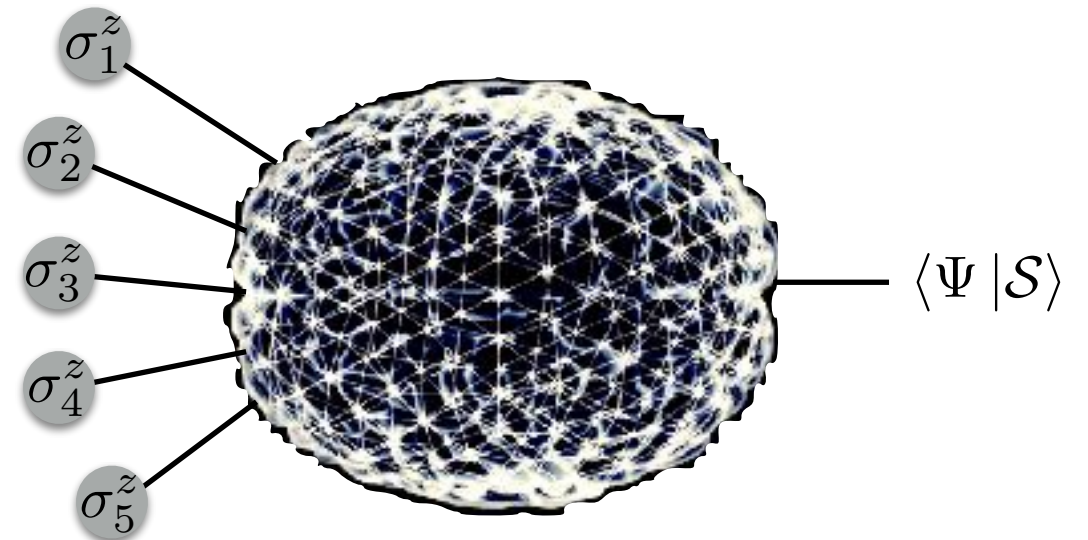
Classical Variational States

Issue

Can We Have A Sufficiently Flexible And Computationally Efficient Classical Ansatz?

Neural-Network Quantum States

Carleo, and Troyer
Science 355, 602 (2017)



$$C \sigma_1^z \sigma_2^z \dots \sigma_5^z$$

Many-Body Amplitudes

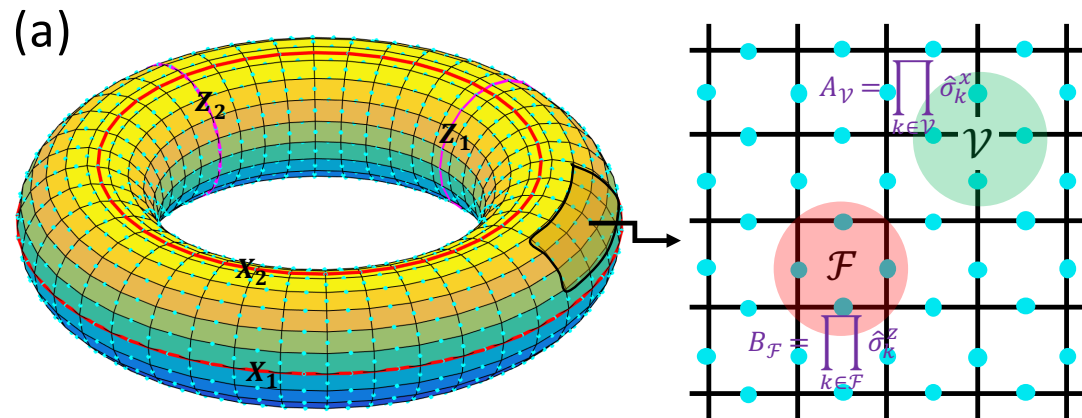
$$\Psi(s_1 \dots s_N) = g^{(L)} \circ W^{(L)} \dots g^{(2)} \circ W^{(2)} g^{(1)} \circ W^{(1)} \mathbf{s}$$

Quantum Numbers

Nonlinear Activation Functions Applied Component-Wise

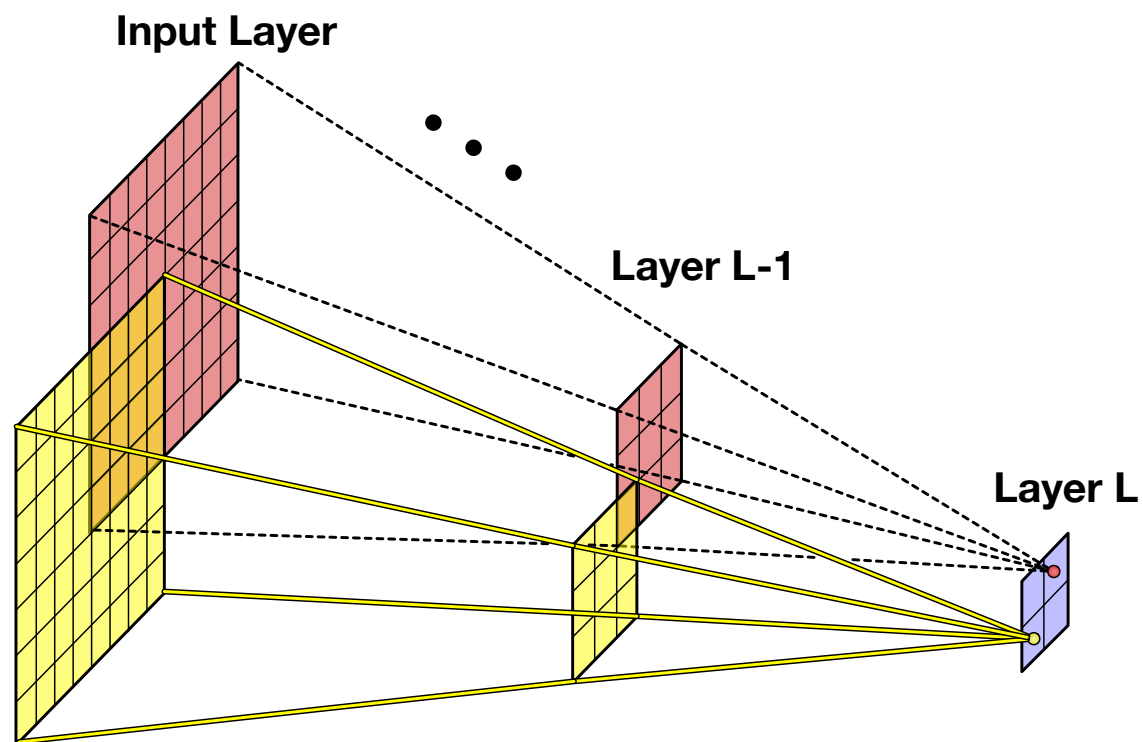
Variational Parameters

Some Properties



Compact Representations of
Almost all known **efficient**
Variational States (Jastrow,
Laughlin, MPS, etc)

For a review: see Carleo et al.
Reviews of Modern Physics 91, 045002 (2020)



Efficient Volume-Law States
With Shallow or (better) Deep
Convolutional Networks

Levine, Sharir, Cohen, and Shashua
Physical Review Letters 122, 065301 (2019)

Ground-State Search

Variational Sampling

Expectation Minimization

$$\mathcal{L}(\mathbf{W}) = \frac{\sum_s |\Psi(s, W)|^2 E_{\text{loc}}(s, W)}{\sum_s |\Psi(s, W)|^2}$$

McMillan
Phys. Rev. 138,
A442 (1965)

Loss Function

No "Dataset" here, closer in
spirit to reinforcement learning

Energy Gradient

$$\nabla_k E = 2 \left(\langle \mathcal{O}_k^* E_{\text{loc}} \rangle - \langle \mathcal{O}_k^* \rangle \langle E_{\text{loc}} \rangle \right)$$

$$E_{\text{loc}}(\sigma_1 \dots \sigma_N) = \sum_{\sigma'_1, \dots, \sigma'_N} \frac{\Psi(\sigma'_1 \dots \sigma'_N)}{\Psi(\sigma_1, \dots, \sigma_N)} \mathcal{H}_{\sigma, \sigma'}$$

$$\mathcal{O}_k(\sigma_1 \dots \sigma_N) = \frac{1}{\Psi(\sigma_1 \dots \sigma_N)} \frac{\partial \Psi(\sigma_1 \dots \sigma_N)}{\partial p_k}$$

$$\langle \dots \rangle = \frac{\sum_{\sigma_1 \sigma_2 \dots \sigma_N} |\Psi(\sigma_1, \sigma_2 \dots \sigma_N)|^2 \dots}{\sum_{\sigma_1 \sigma_2 \dots \sigma_N} |\Psi(\sigma_1, \sigma_2 \dots \sigma_N)|^2}$$

Computationally Tractable States

“Computationally Tractable” States

Van den Nest
arXiv:0911.1624 (2009)

Definition 1 *An n -qubit state $|\psi\rangle$ is called ‘computationally tractable’ (CT) if the following conditions hold:*

- (a) it is possible to sample in $\text{poly}(n)$ time with classical means from the probability distribution $\text{Prob}(x) = |\langle x|\psi\rangle|^2$ on the set of n -bit strings x , and*
- (b) upon input of any bit string x , the coefficient $\langle x|\psi\rangle$ can be computed in $\text{poly}(n)$ time on a classical computer.*

Theorem 3 *Let $|\psi\rangle$ and $|\varphi\rangle$ be CT n -qubit states and let A be an efficiently computable sparse (not necessarily unitary) n -qubit operation with $\|A\| \leq 1$. Then there exists an efficient classical algorithm to approximate $\langle \varphi|A|\psi\rangle$ with polynomial accuracy.*

Corollary 1 *Let $|\psi\rangle$ be an n -qubit CT state and let O be a d -local observable with $d = O(\log n)$ and $\|O\| \leq 1$. Then there exists an efficient classical algorithm to estimate $\langle \psi|O|\psi\rangle$ with polynomial accuracy.*

Examples

Matrix Product States
Are Computationally
Tractable

PEPS are **not**
Computationally
Tractable

Generic neural deep
quantum states are
not computationally
Tractable

NQS: MCMC Breaks Tractability

$$\langle \dots \rangle = \frac{\sum_{\sigma_1 \sigma_2 \dots \sigma_N} |\Psi(\sigma_1, \sigma_2 \dots \sigma_N)|^2 \dots}{\sum_{\sigma_1 \sigma_2 \dots \sigma_N} |\Psi(\sigma_1, \sigma_2 \dots \sigma_N)|^2}$$

Historically,
MCMC is used to
sample from
this
probability

For deep networks:
can Be Inefficient
Both Theoretically
(Correlation
Times) and
Computationally
(GPU)

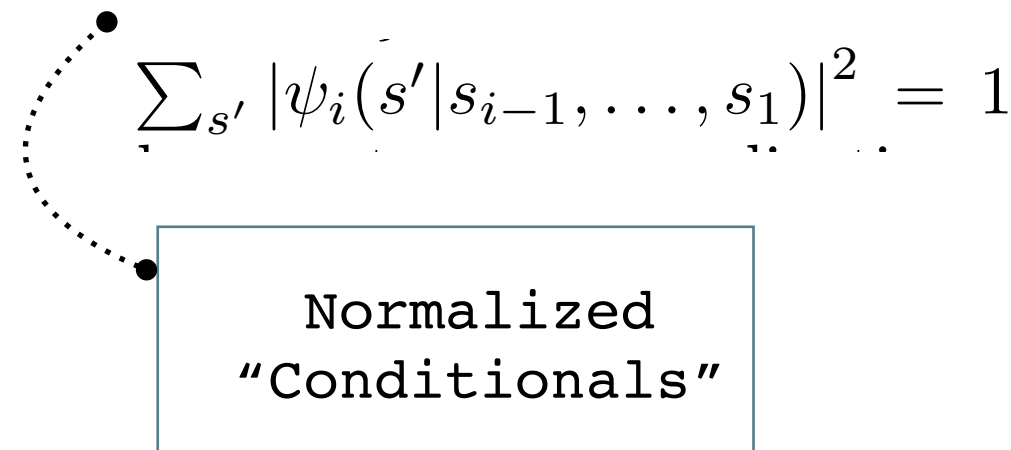
Autoregressive Quantum States

Sharir, Levine, Wies, Carleo, and Shashua
Phys. Rev. Lett. 124, 020503 (2020)

$$\Psi(s_1, \dots, s_N) = \prod_{i=1}^N \psi_i(s_i | s_{i-1}, \dots, s_1)$$

$\sum_{s'} |\psi_i(s' | s_{i-1}, \dots, s_1)|^2 = 1$

Normalized
"Conditionals"

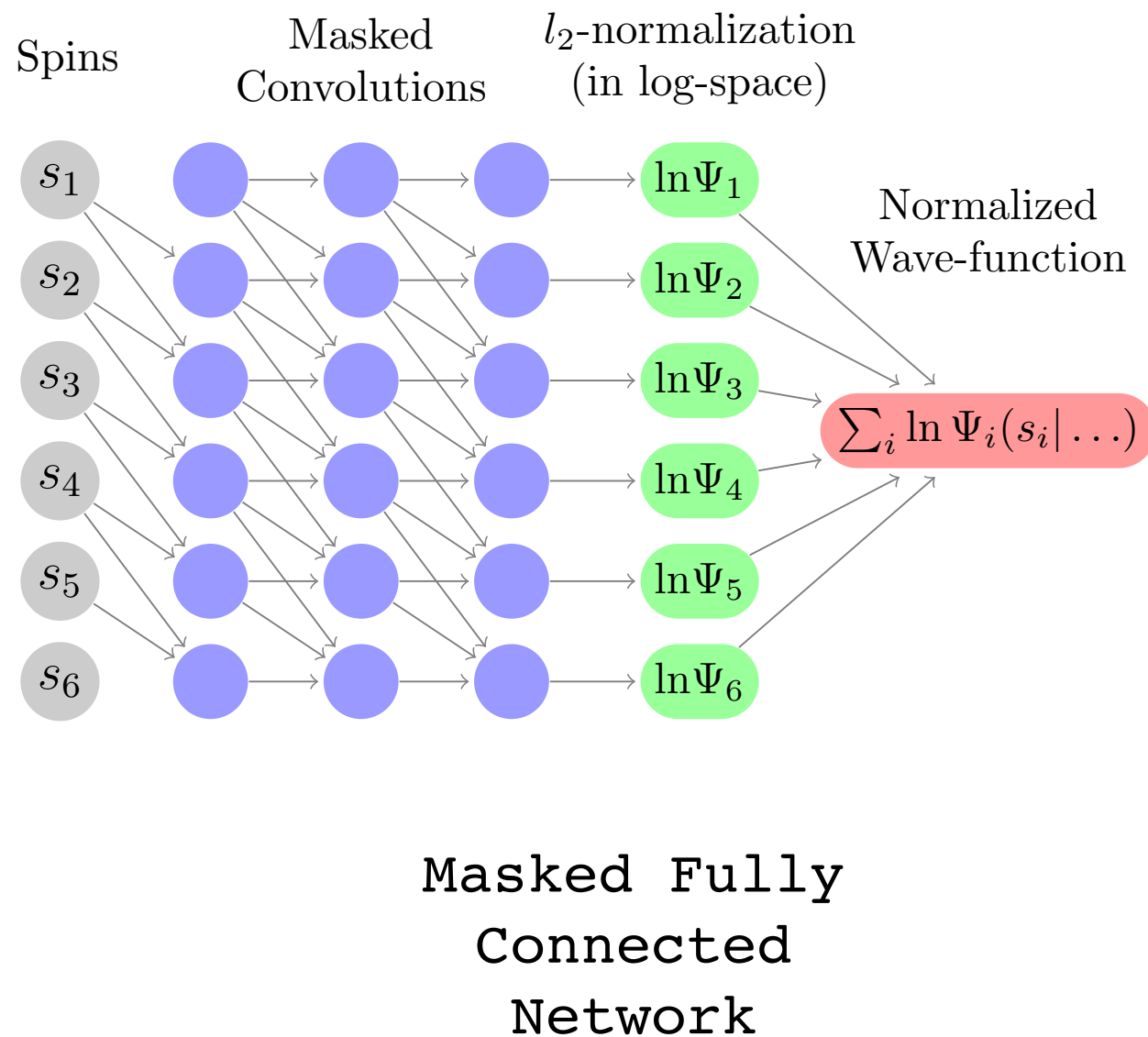


These Are Computationally Tractable

(a) Exact Sampling

(b) Computing Normalized Amplitudes is Efficient

Masked deep networks



1	1	1	1	1
1	1	1	1	1
1	1	0	0	0
0	0	0	0	0
0	0	0	0	0

Masked Convolutions

PixelCNN

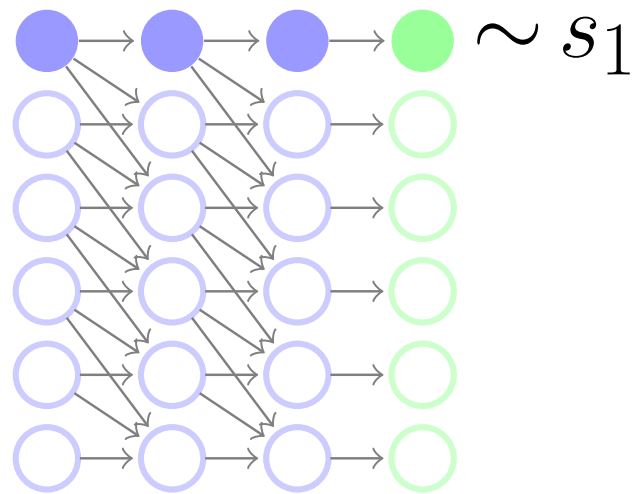
Van den Oord et al.
arXiv:1606.05328 (2016)

Salimans et al.
arXiv:1701.05517 (2017)

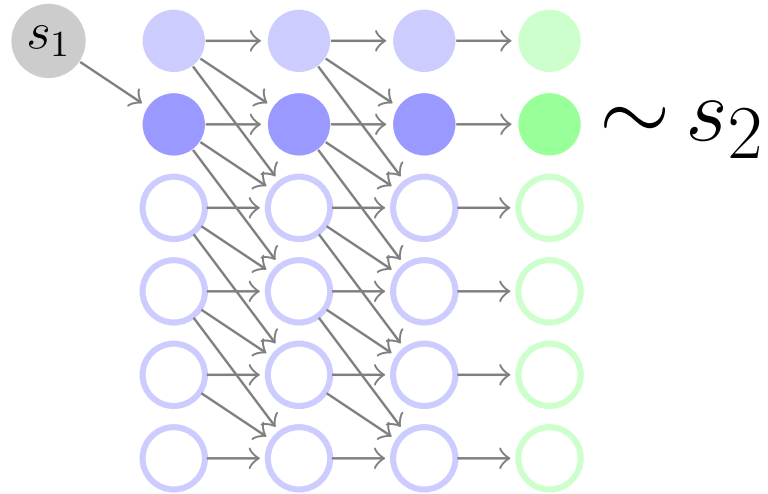
Ramachandran et al.
arXiv:1704.06001 (2017)

Exact Sampling

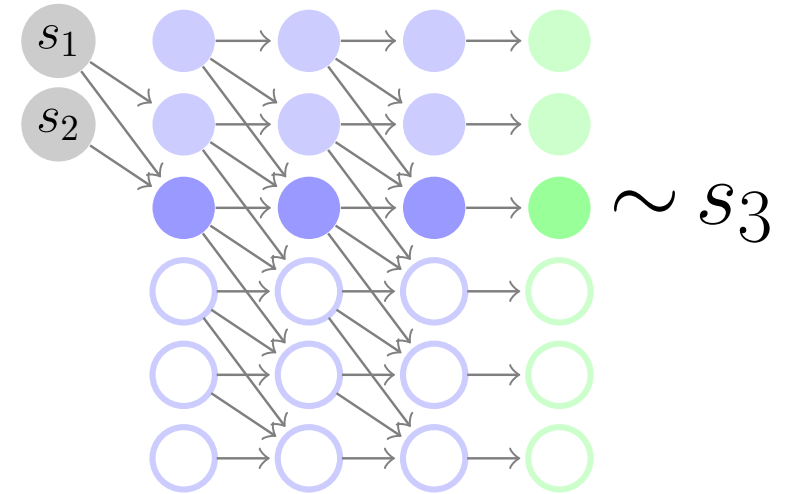
Step 1:



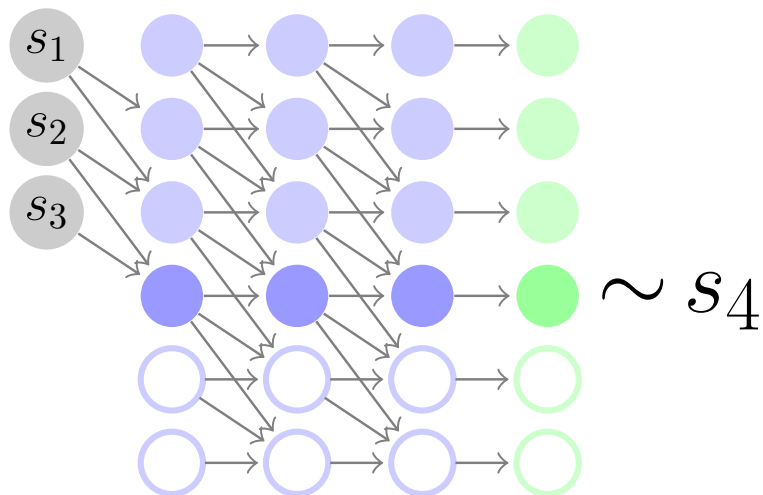
Step 2:



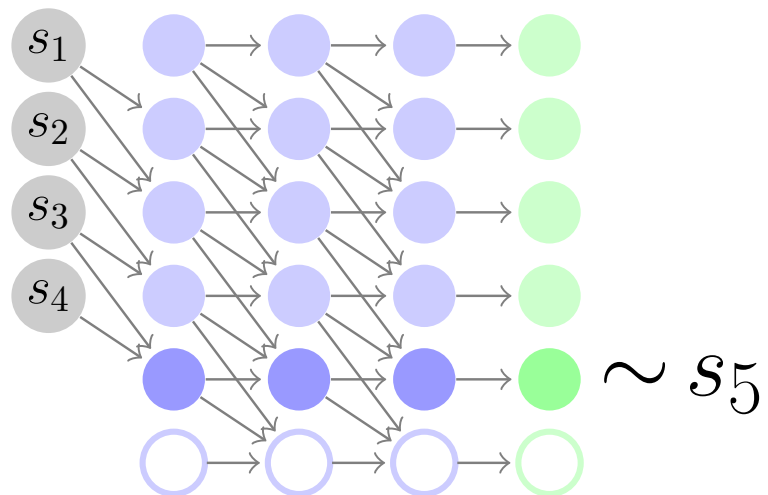
Step 3:



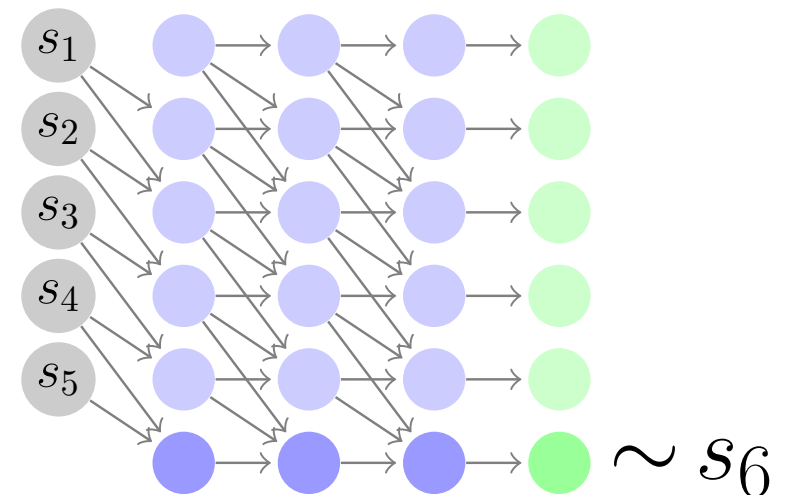
Step 4:



Step 5:

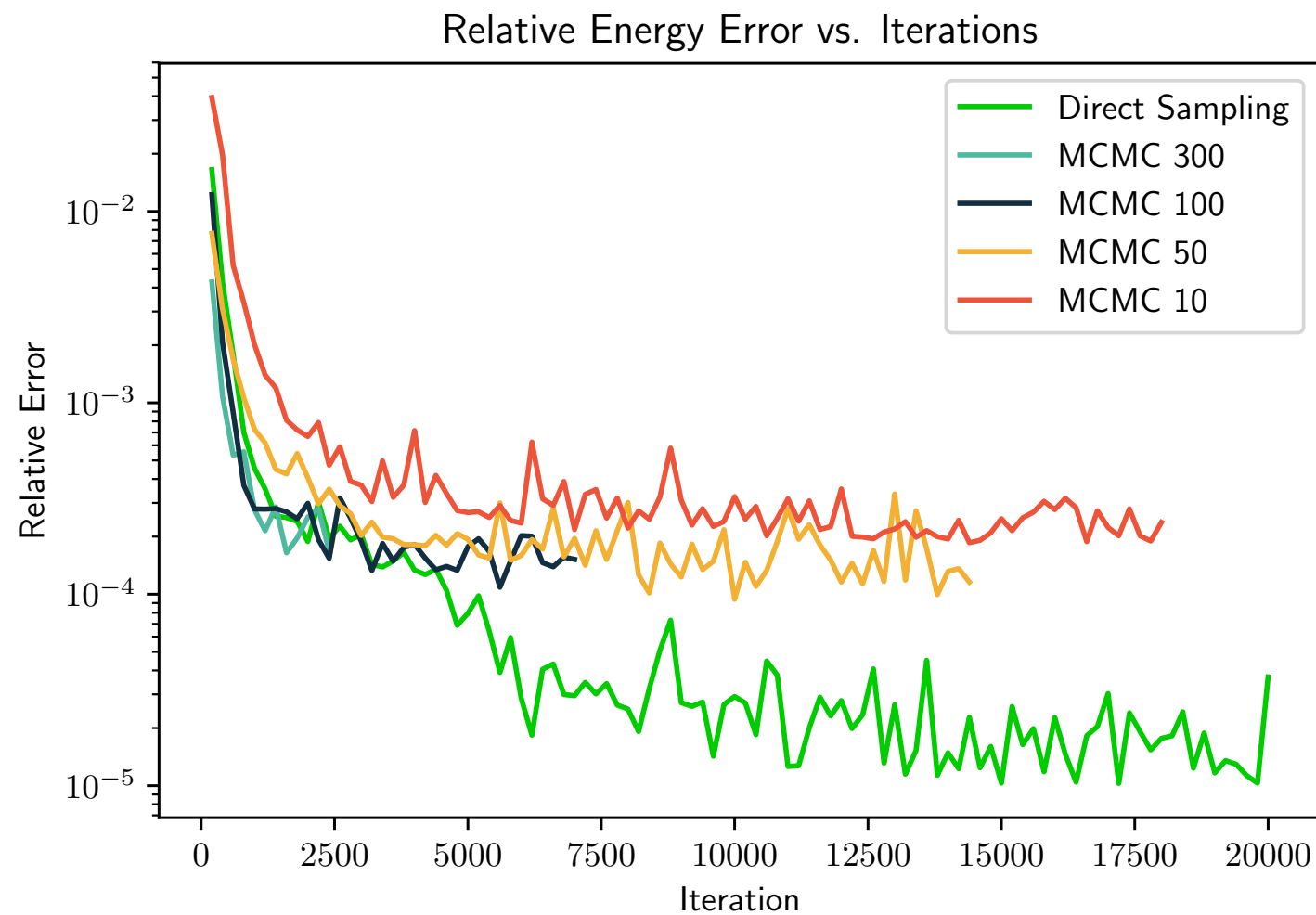


Step 6:



Removing The Sampling Bottleneck Pays Off

Sharir, Levine, Wies, Carleo, and Shashua
Phys. Rev. Lett. 124, 020503 (2020)

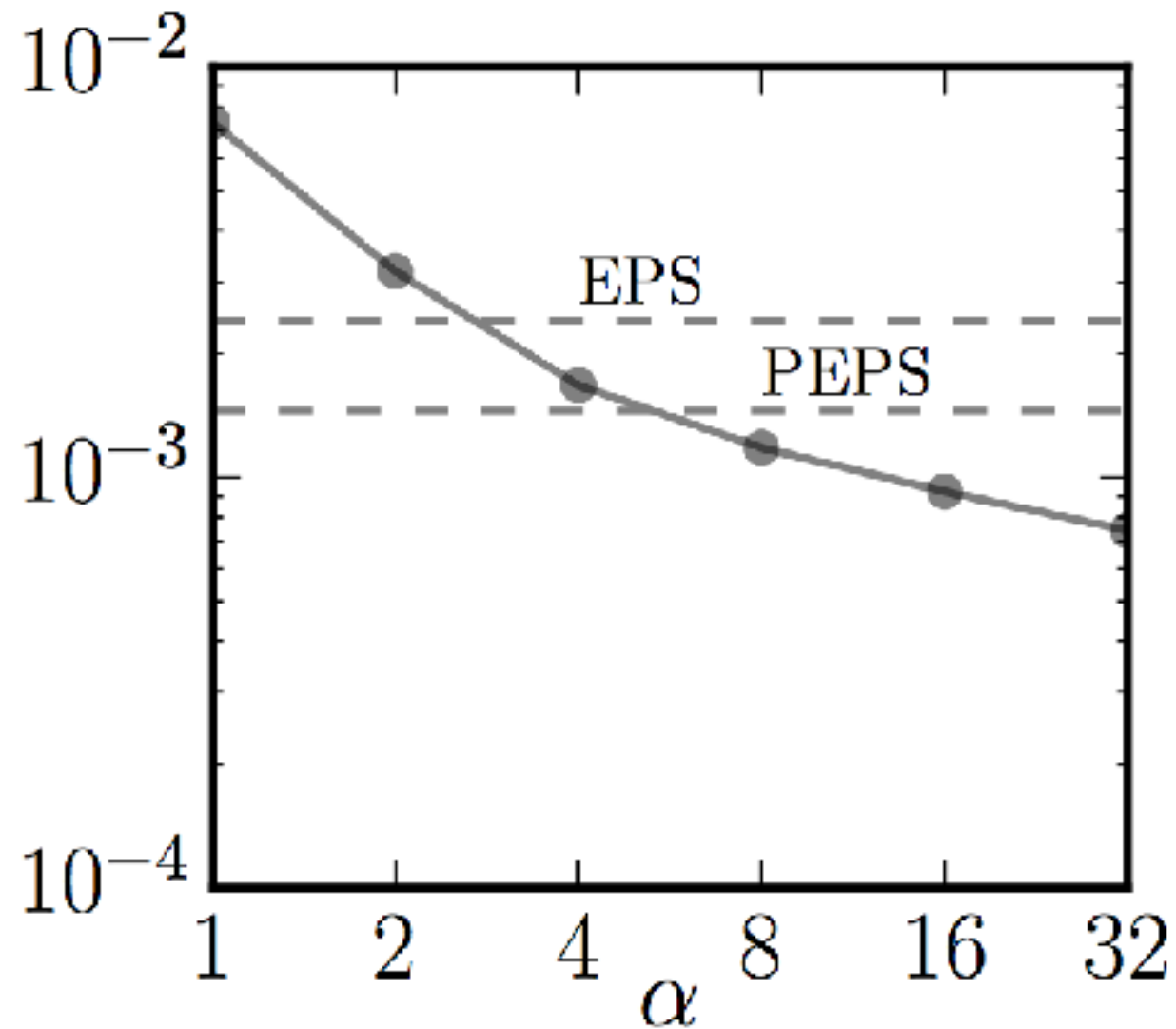


21x21
Transverse-
Field Ising
in 2d

Depth 20
About 1
Million
Parameters

Heisenberg Model

Carleo, and Troyer
Science 355, 602 (2017)

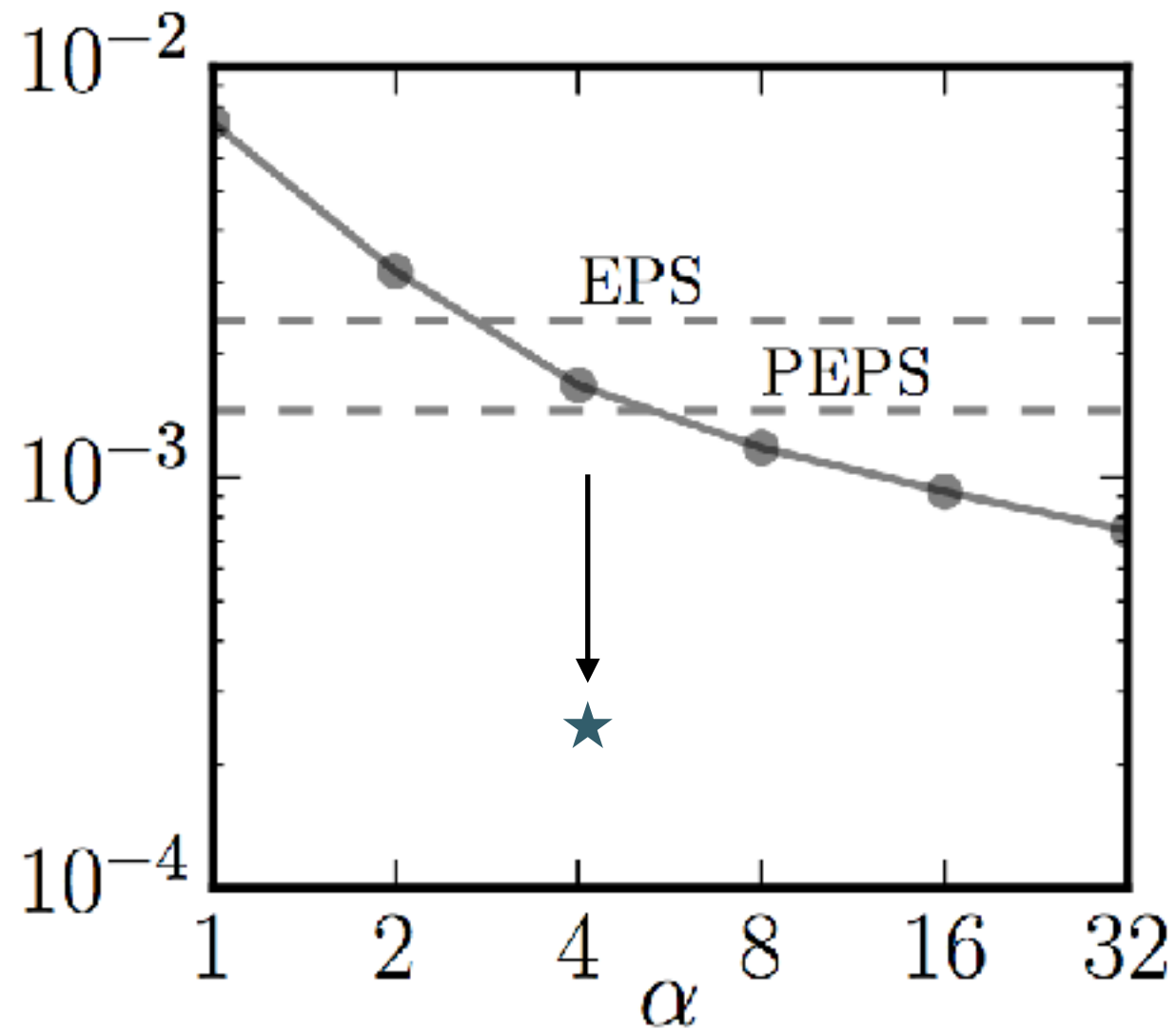


10 by 10 cluster

Early (2016) Results With
Shallow (RBM) Network

Heisenberg Model

Choo, Neupert, and Carleo
Phys. Rev. B 100, 125124 (2019)

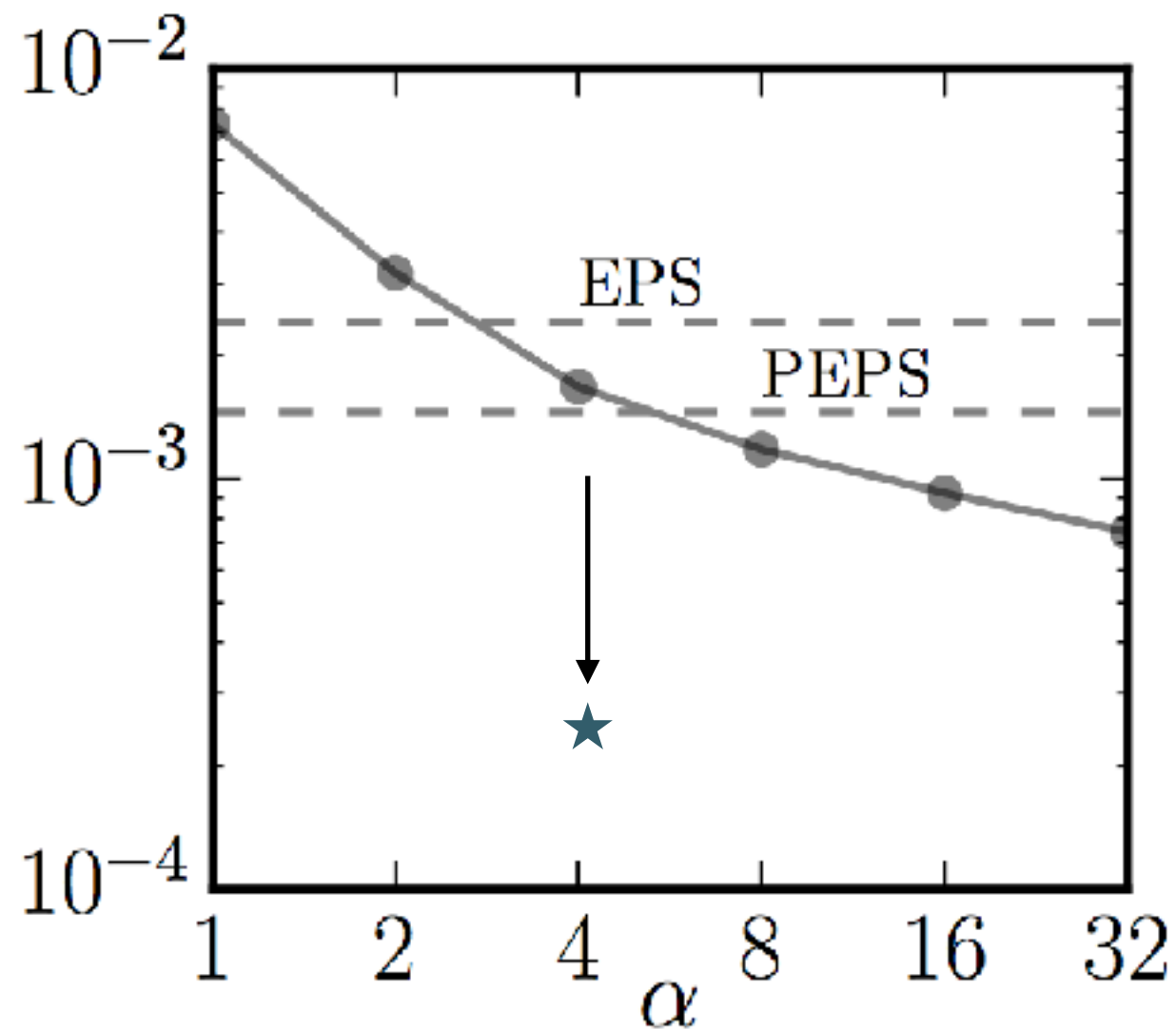


10 by 10 cluster

(Mildly) deep network further improves

Heisenberg Model

Sharir, Levine, Wies, Carleo, and Shashua
Phys. Rev. Lett. 124, 020503 (2020)



10 by 10 cluster

(Actual) deep network further improves

$\sim 4 \times 10^{-5}$

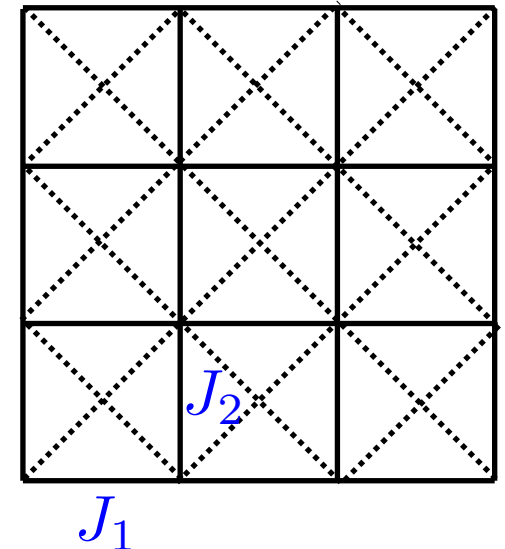


A Challenging Benchmark

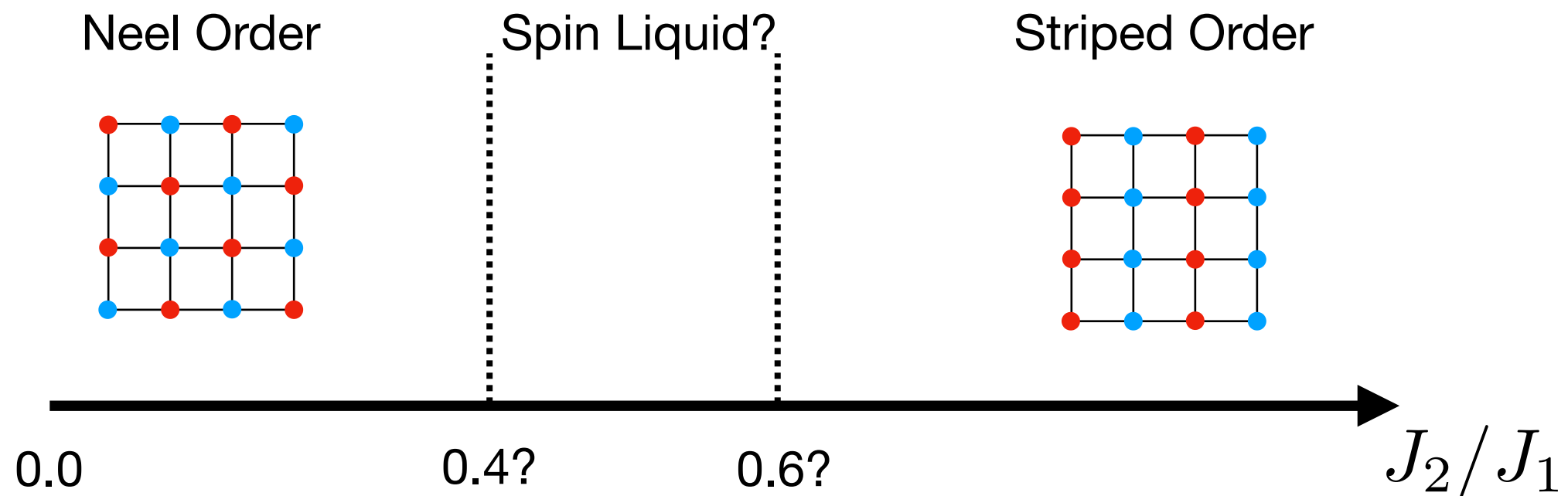
Frustrated 2D Spins

J1-J2 Model

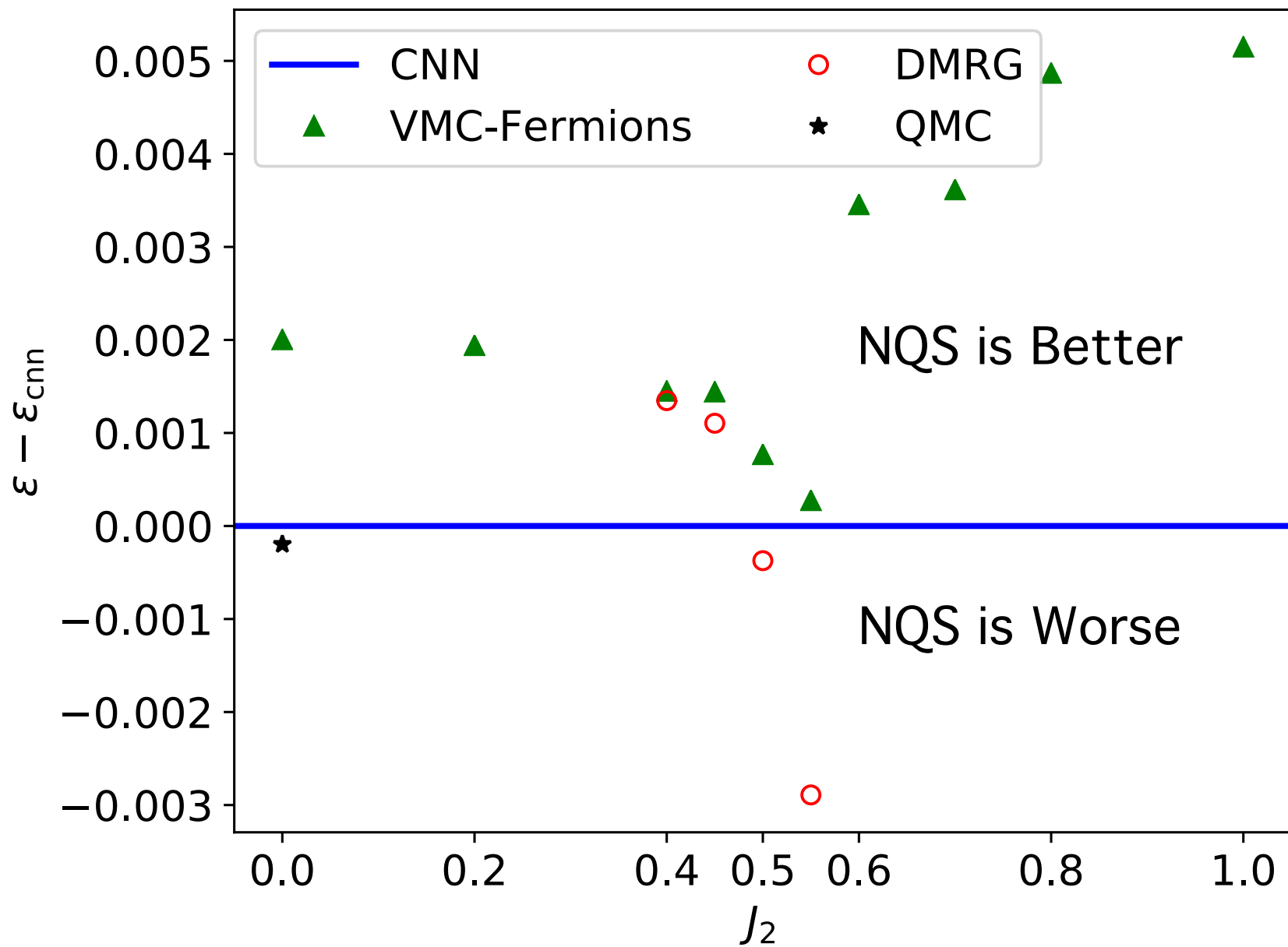
$$\hat{H} = J_1 \sum_{\langle ij \rangle} \hat{\mathbf{S}}_i \cdot \hat{\mathbf{S}}_j + J_2 \sum_{\langle\langle ij \rangle\rangle} \hat{\mathbf{S}}_i \cdot \hat{\mathbf{S}}_j$$



Phase Diagram



J1-J2 Model Square Lattice

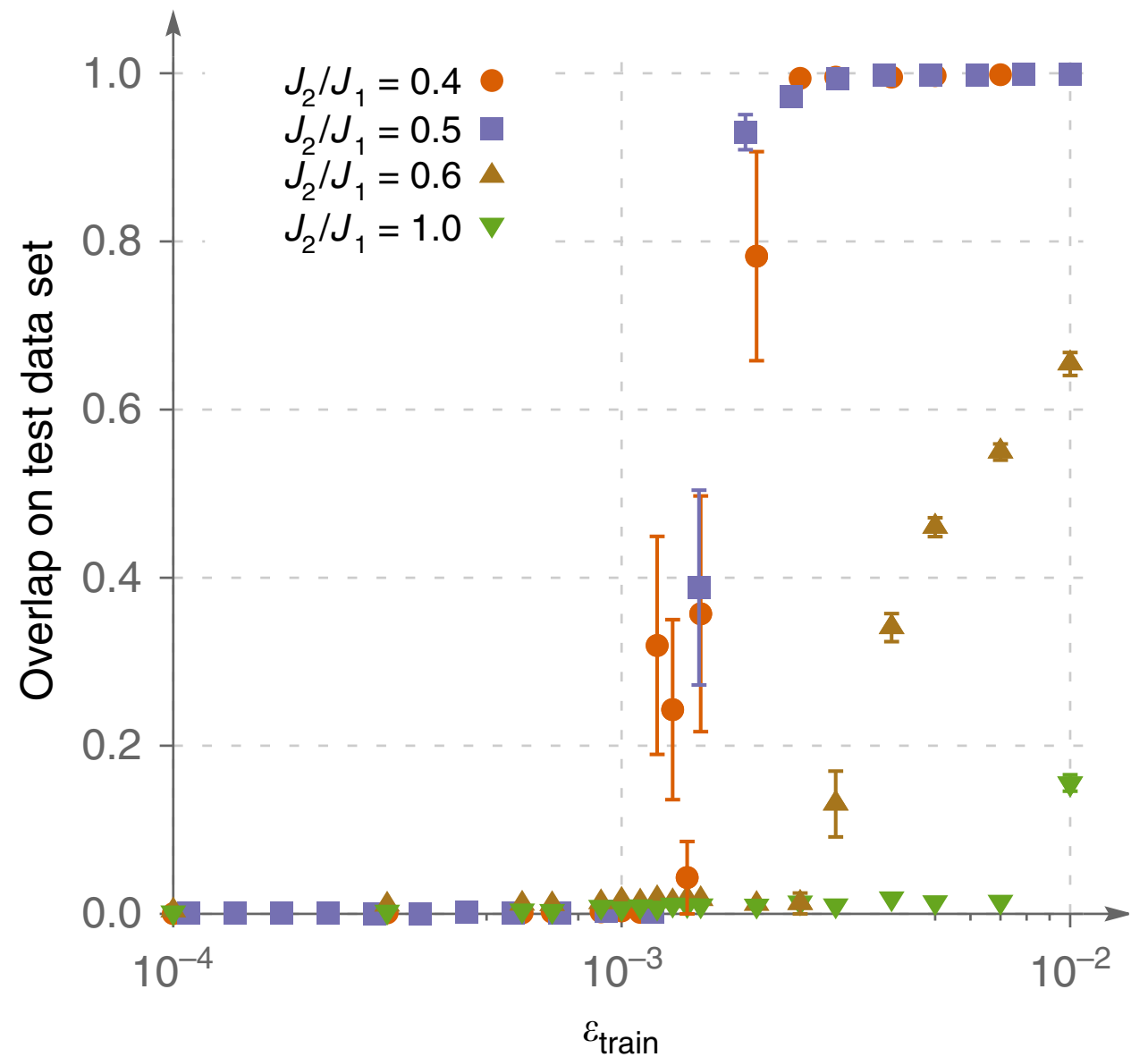
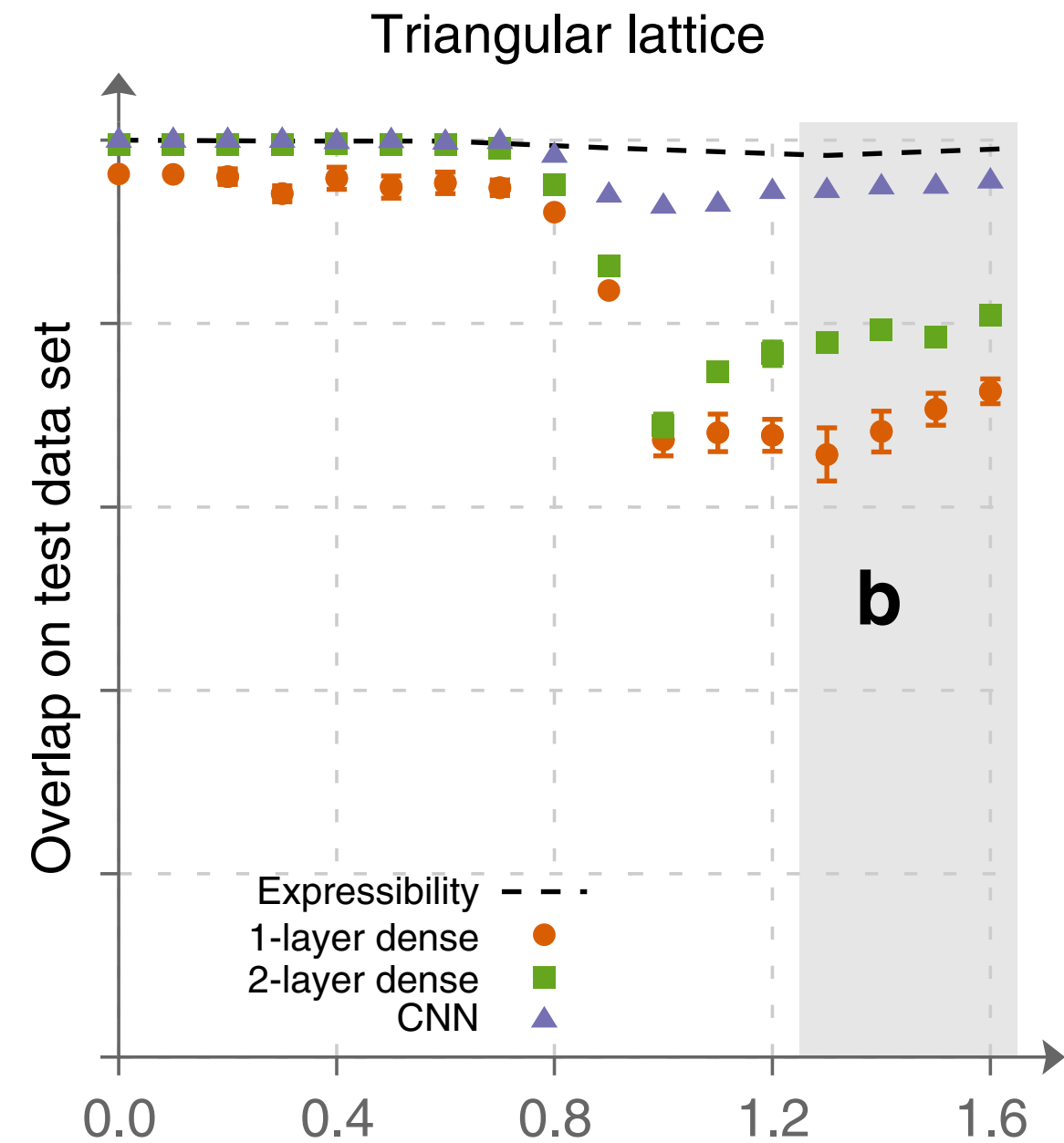


10 by 10 cluster

Choo, Neupert, and Carleo
Phys. Rev. B 100, 125124 (2019)

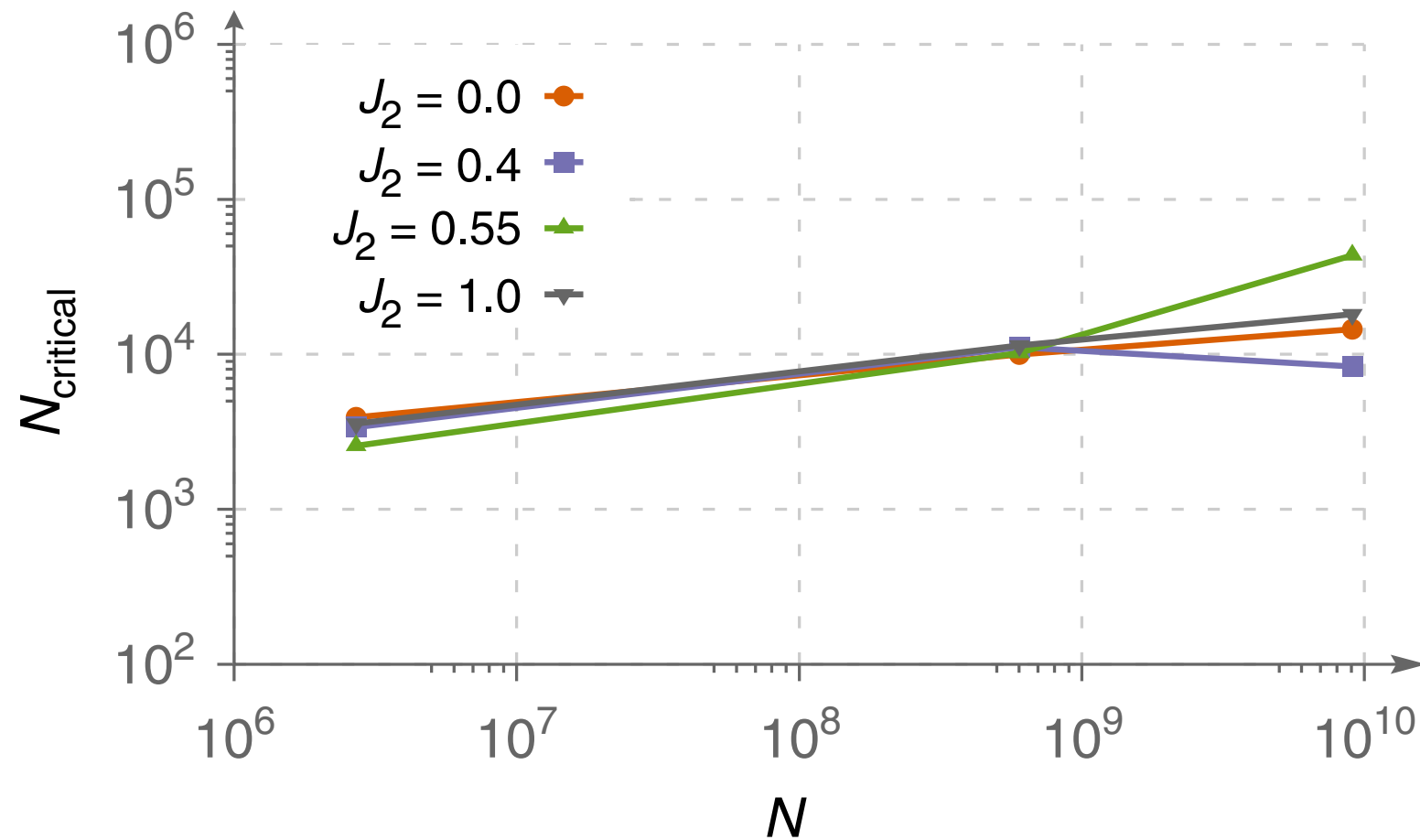
Tangible Improvements on Open
Challenging Models

Origin of the challenge: generalization for sign structure



Westerhout, Astrakhantsev, Tikhonov, Katsnelson, Bagrov
Nature Comm. 11, 1593 (2020)

Main Open Issue: Sample Complexity For Sign Structure



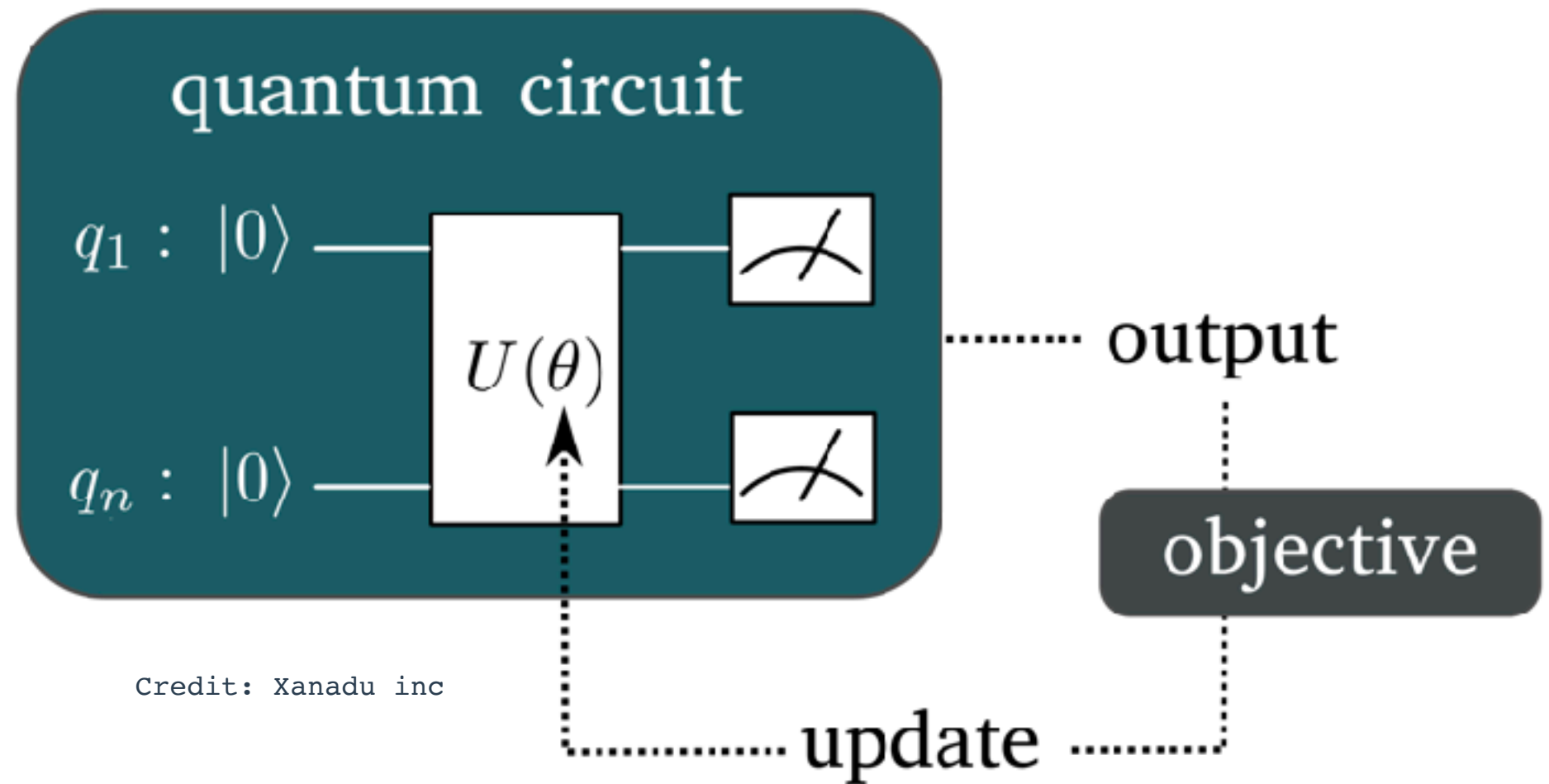
$N_s = 20$

$N_s = 36$

Westerhout, Astrakhantsev, Tikhonov, Katsnelson, Bagrov
Nature Comm. 11, 1593 (2020)

Quantum Variational States

Variational Quantum Eigensolvers



$$E(\theta) = \langle \Psi_0 | U(\theta)^\dagger \mathcal{H} U(\theta) | \Psi_0 \rangle$$

$$\mathcal{H} = \sum_k c_k \sigma_1^{(x,y,z)} \dots \sigma_N^{(x,y,z)}$$

Energy is
Estimated
Stochastically
on Samples
From
Measurements

Strong Connection With Stochastic Classical Counterparts

Stokes, Izaac, Killoran, and Carleo
Quantum 4, 269 (2020)

Higher-Order
Optimizer for
Quantum Machine
Learning

Quantum
Variational
Imaginary-Time
Evolution

Quantum
Variational
Real-Time
Evolution

Stochastic
Estimates of
the Quantum
Geometric
Tensor

Stochastic
Reconfiguration

Time-Dependent
Variational
Monte Carlo

Natural
Gradient

Stochastic Reconfiguration (Natural Gradient)

Sandro Soresella et al.
Physical Review Letters
80, 4558 (1998)

Shun-Ichi Amari
Journal Neural Computation
10, 251 (1998)

$$\sum_{k'} S_{k,k'} \Delta p_{k'} = -G_k$$

“Second-Order”
Method

$$S_{k,k'} = \langle \mathcal{O}_k^* \mathcal{O}_{k'} \rangle - \langle \mathcal{O}_k^* \rangle \langle \mathcal{O}_{k'} \rangle$$

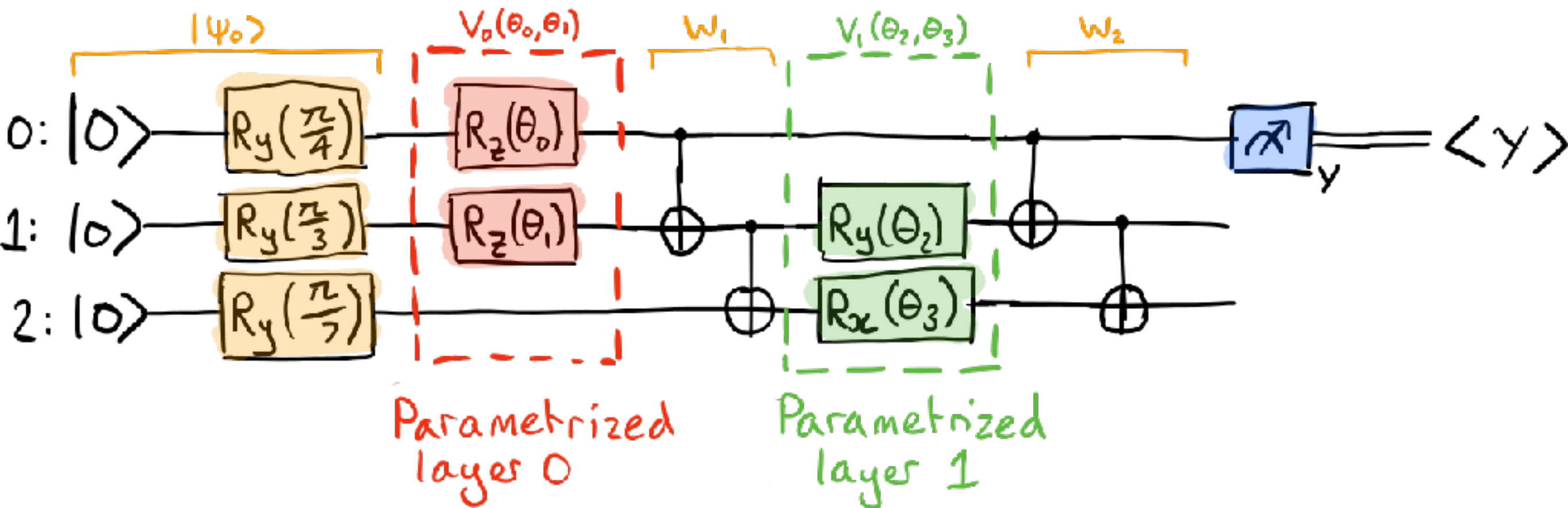
**Quantum Geometric
Tensor or Quantum
Fisher Information**

Linear System to be Solved at each
iteration of the optimizer

Sparse solvers can be used when dealing
with large number of parameters

Quantum Natural Gradient

Stokes, Izaac, Killoran, and Carleo
Quantum 4, 269 (2020)



2x2 matrix for layer 0

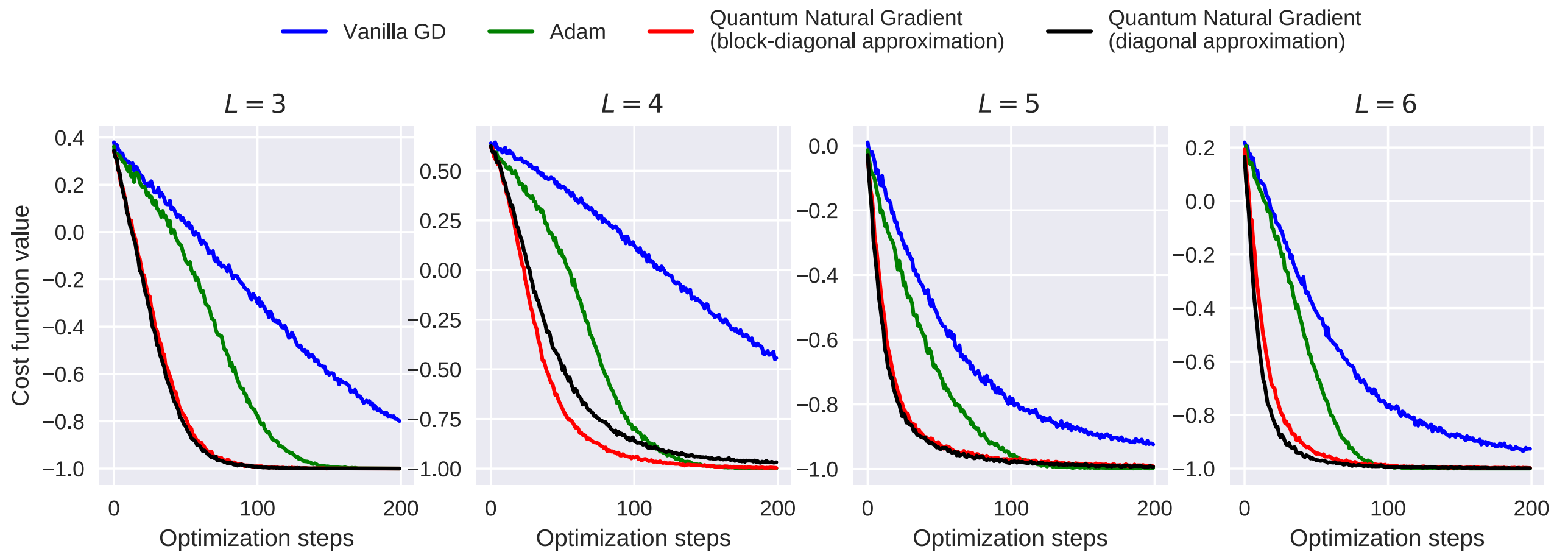
$$g = \begin{bmatrix} g^{(0)} & 0 \\ 0 & g^{(1)} \end{bmatrix}$$

2x2 matrix for layer 1

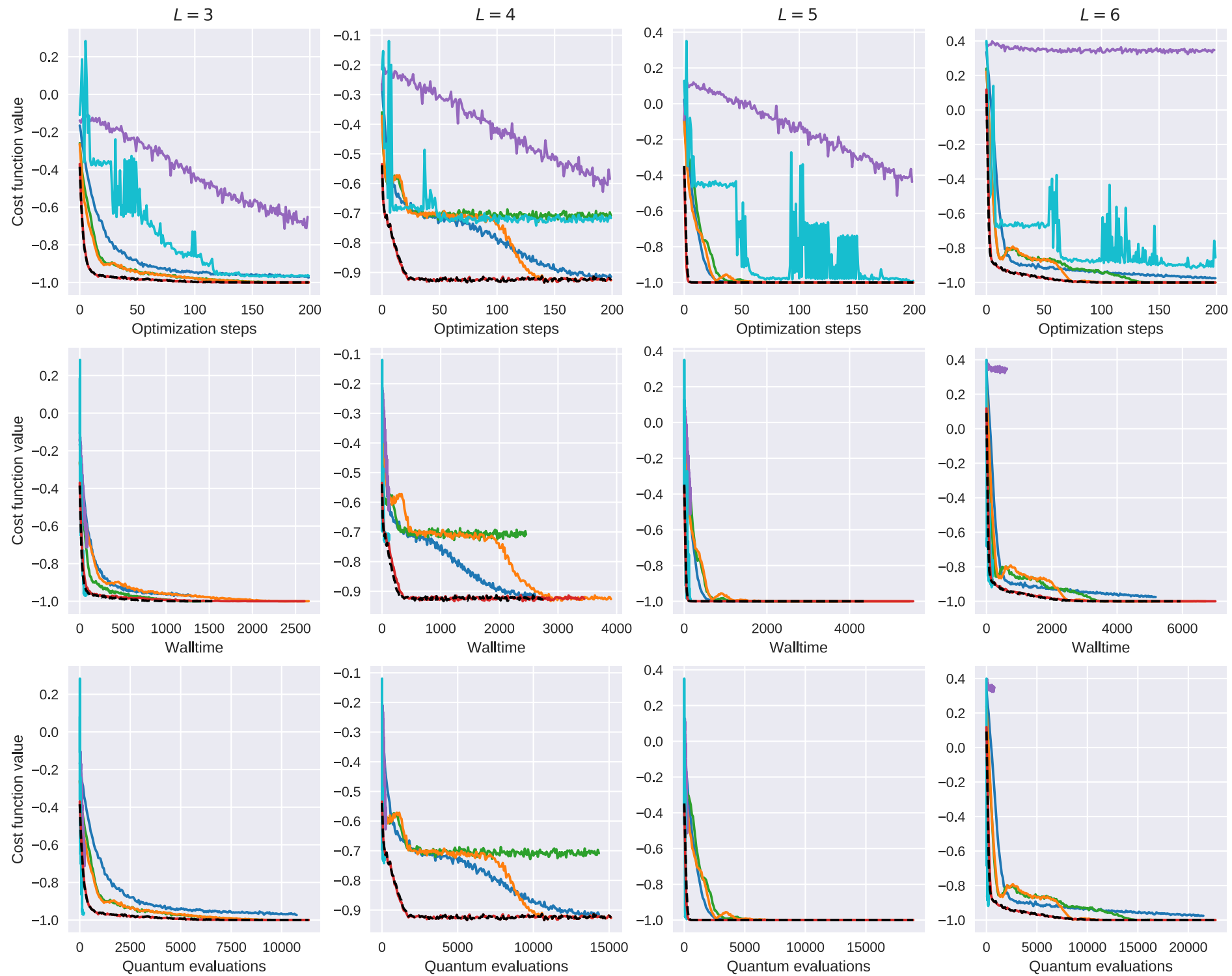
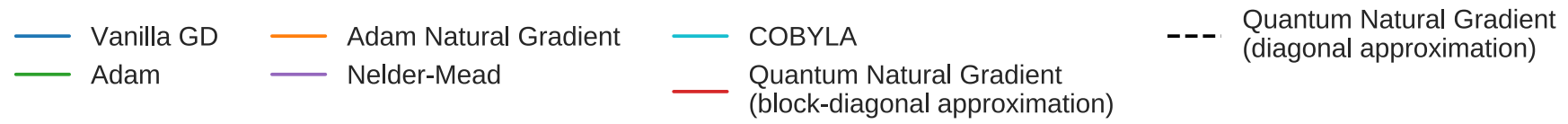
$$\theta_{t+1} = \theta_t - \eta g^+(\theta) \nabla L(\theta),$$

Faster Convergence

Stokes, Izaac, Killoran, and Carleo
Quantum 4, 269 (2020)

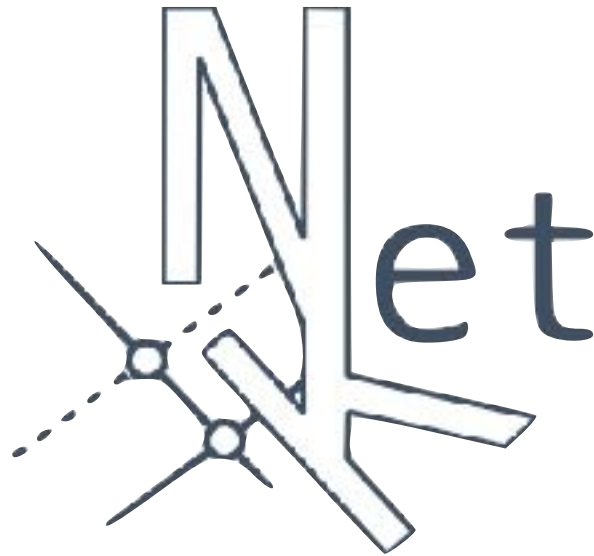


Quantum Evaluations



Software

The NetKet Project



www.netket.org

NetKet: A Machine Learning Toolkit for Many-Body Quantum Systems

Giuseppe Carleo,¹ Kenny Choo,² Damian Hofmann,³ James E. T. Smith,⁴ Tom Westerhout,⁵ Fabien Alet,⁶ Emily J. Davis,⁷ Stavros Efthymiou,⁸ Ivan Glasser,⁸ Sheng-Hsuan Lin,⁹ Marta Mauri,^{1,10} Guglielmo Mazzola,¹¹ Christian B. Mendl,¹² Evert van Nieuwenburg,¹³ Ossian O'Reilly,¹⁴ Hugo Théveniaut,⁶ Giacomo Torlai,¹ and Alexander Wietek¹

¹Center for Computational Quantum Physics, Flatiron Institute, 162 5th Avenue, NY 10010, New York, USA

²Department of Physics, University of Zurich, Winterthurerstrasse 190, 8057 Zürich, Switzerland

³Max Planck Institute for the Structure and Dynamics of Matter,
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⁵Institute for Molecules and Materials, Radboud University, NL-6525 AJ Nijmegen, The Netherlands

⁶Laboratoire de Physique Théorique, IRSAMC, Université de Toulouse, CNRS, UPS, 31062 Toulouse, France

⁷Department of Physics, Stanford University, Stanford, California 94305, USA

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¹⁰Dipartimento di Fisica, Università degli Studi di Milano, via Celoria 16, I-20133 Milano, Italy

¹¹Theoretische Physik, ETH Zürich, 8093 Zürich, Switzerland

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¹³Institute for Quantum Information and Matter,

California Institute of Technology, Pasadena, CA 91125, USA

¹⁴Southern California Earthquake Center, University of Southern California,
3651 Trousdale Pkwy, Los Angeles, CA 90089, USA

SoftwareX 10, 100311 (2019)

```
import netket as nk

# 1D Lattice
g = nk.graph.Hypercube(length=20, n_dim=1, pbc=True)

# Hilbert space of spins on the graph
hi = nk.hilbert.Spin(s=0.5, graph=g)

# Ising spin hamiltonian
ha = nk.operator.Ising(h=1.0, hilbert=hi)

# RBM Spin Machine
ma = nk.machine.RbmSpin(alpha=1, hilbert=hi)
ma.init_random_parameters(seed=1234, sigma=0.01)

# Metropolis Local Sampling
sa = nk.sampler.MetropolisLocal(machine=ma)

# Optimizer
op = nk.optimizer.Sgd(learning_rate=0.1)

# Stochastic reconfiguration
gs = nk.variational.Vmc(
    hamiltonian=ha,
    sampler=sa,
    optimizer=op,
    n_samples=1000,
    diag_shift=0.1,
    method='Sr')

gs.run(output_prefix='test', n_iter=300)
```

Release 3.0 (soon out...)

Same APIs but Almost Entire Rewriting Under the Hood

Pure Python:
Entirely Remove
C++ codebase

Numba Kernels
For Few Hotspots



Seamless
Integration with
DL Frameworks

Support for **GPU**
and **TPU**



GPU Tests on V100 Cards

Up to 100X Speedup on Gradients

Up to 30X Speedup on Sampling

Overview

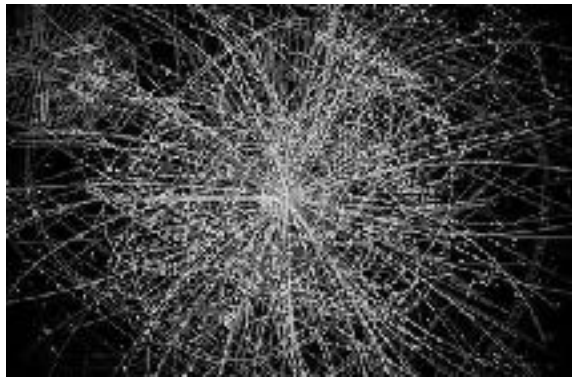
Outlook/Challenges

We need better
optimization strategies
and parameterizations
for signs/phases

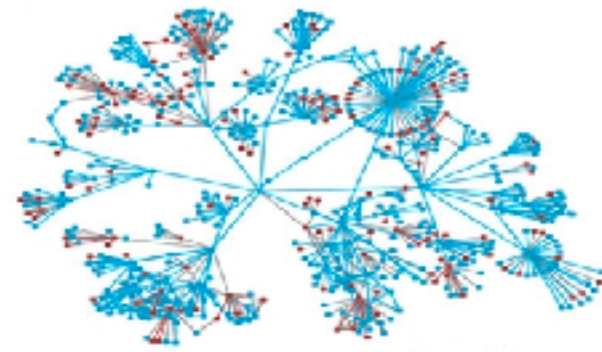
Still “Heroic”
phase for
fermions
developments

Symmetries in
networks play a
fundamental role

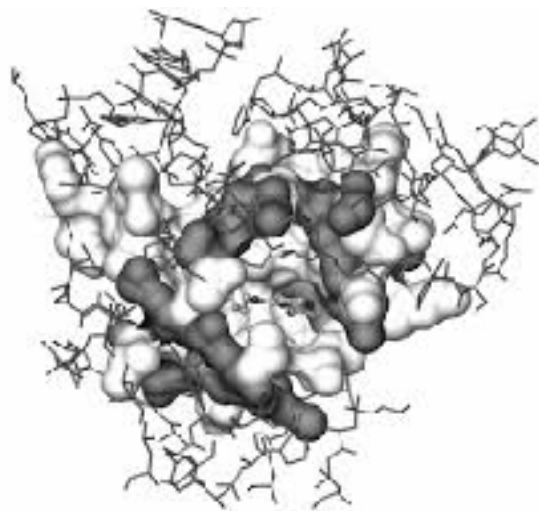
Machine Learning in Physics



**Particle
Physics**



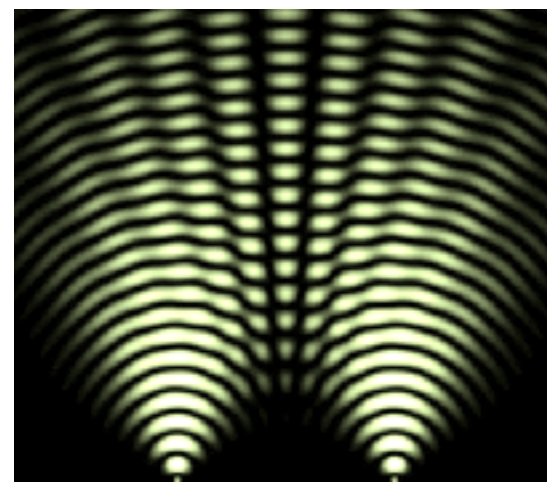
**Statistical
Physics**



**Chemistry/
Materials**



Astrophysics



**Quantum
Physics**

*Carleo, Cirac, Cranmer, Daudet,
Schuld, Tishby, Vogt, and Zdeborovà
Rev. Mod. Phys. 91, 045002 (2019)*

Issue 2: Resources

Number of Measurements

.....The Problem.....

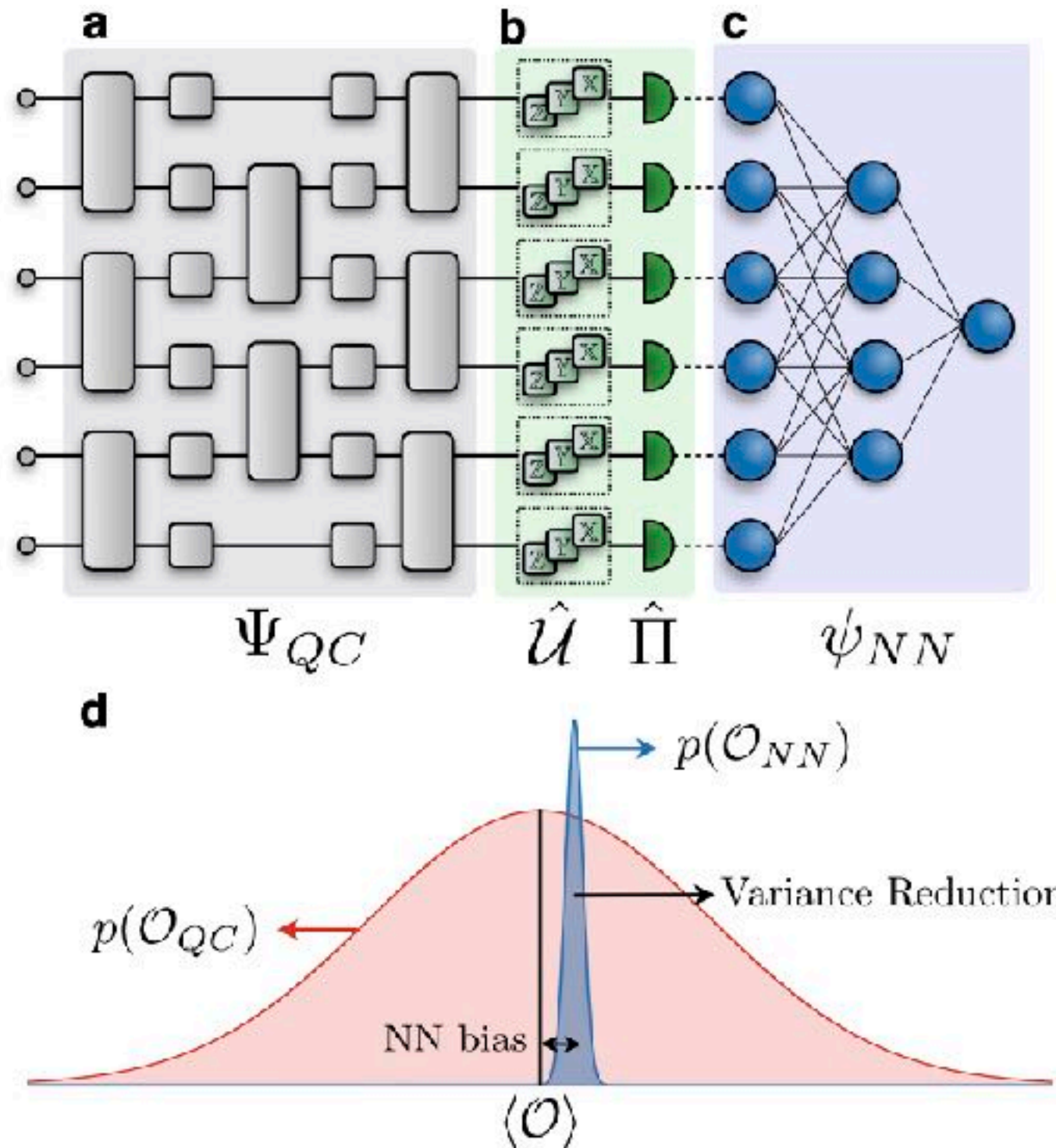
Precise Estimates of Physical Observables typically require several millions or more of measurements even on very small systems

$$\mathcal{H} = \sum_k c_k \sigma_1^{(x,y,z)} \dots \sigma_N^{(x,y,z)}$$

Example: Estimate Energy Using Measurements in Pauli Basis

Neural-Network State Parameterization

Torlai, Mazzola, Carleo, and Mezzacapo
Phys. Rev. Research 2, 022060 (2020)



Neural-Network Tomography

*Torlai, Mazzola, Carrasquilla,
Troyer, Melko, and Carleo*
Nature Physics (2018)

Neural-Network Training: find \mathbf{W} such that

$$|\Psi_b(\mathbf{s}, \mathbf{W})|^2 \simeq P_b(\mathbf{s}) \text{ in all given bases}$$

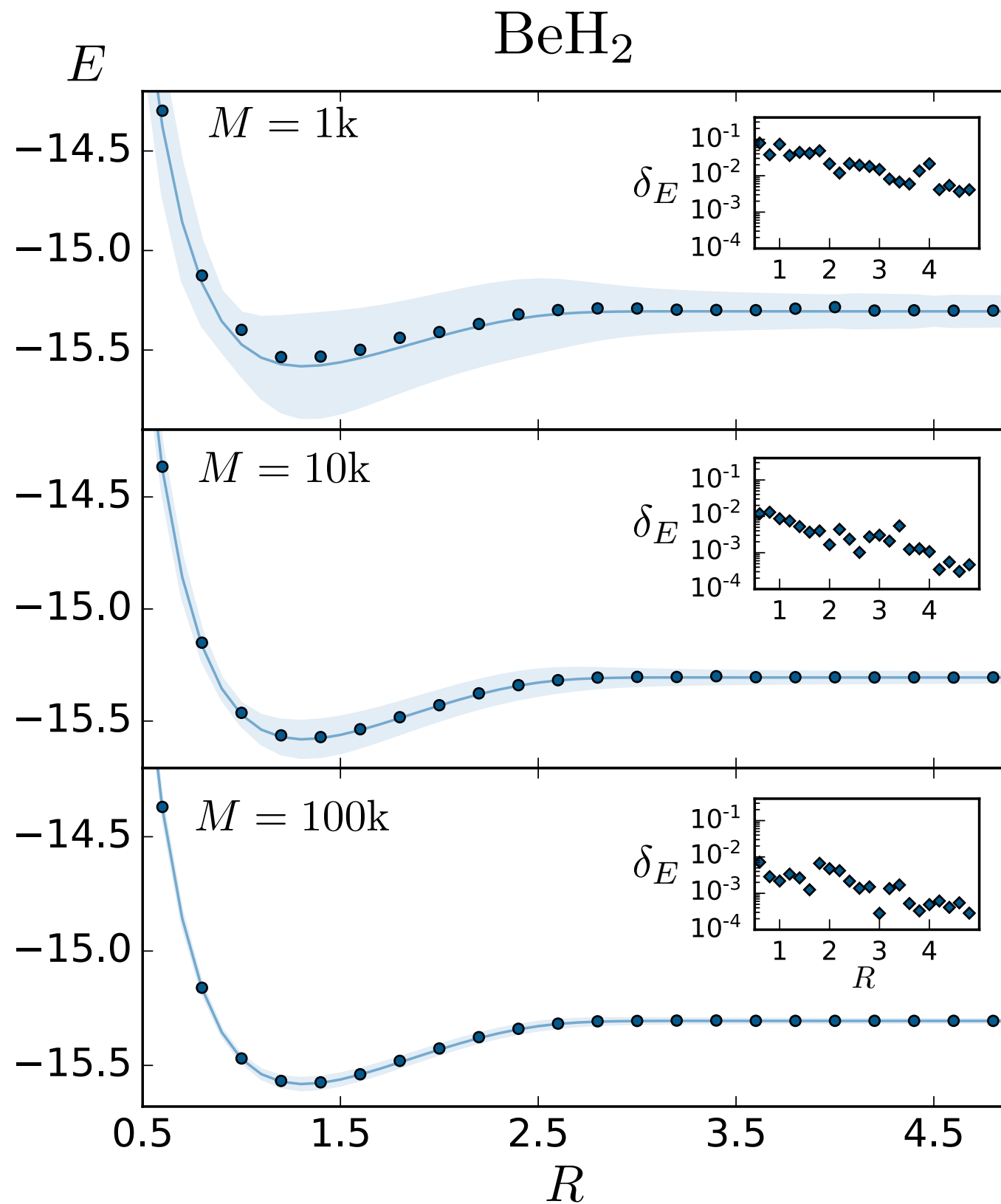
$$\Psi_b(\mathbf{s}, \mathbf{W}) = \sum_{s'} \Psi(\mathbf{s}, \mathbf{W}) U_{s,s'}^b$$

• Sparse Unitary Matrices

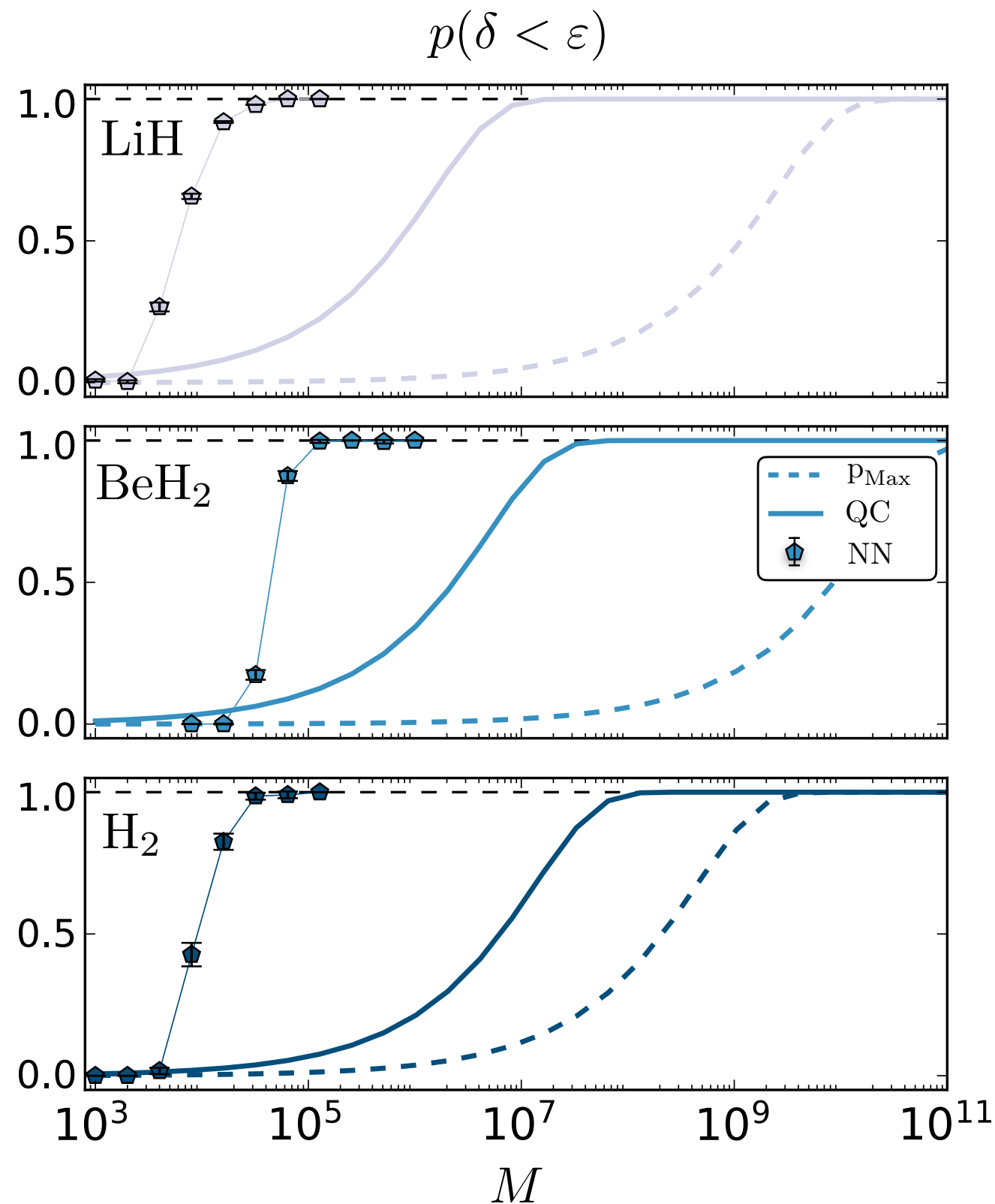
$$\mathcal{L}(\mathbf{W}) = \sum_b \sum_{\mathbf{s}} P_b(\mathbf{s}) \log \frac{P_b(\mathbf{s})}{|\Psi_b(\mathbf{s}, \mathbf{W})|^2}$$

Minimise Sum
of Kullback-
Leibler
Divergences

Quantum Chemistry Problems



Prob. Of “Chemical Accuracy”



Reduce number of measurements by few orders of magnitude

Torlai, Mazzola,
Carleo, and
Mezzacapo
Phys. Rev. Research
2, 022060 (2020)