

# Transverse polarization results

Hiromi Okada (BNL)

## 1. Introduction

- What is Single Spin Asymmetry (SSA) ?
- Why SSA is interesting ?
- Why SSA at RHIC

## 2. Experimental SSA results at RHIC

- STAR (Inclusive  $\pi^0$  and di-jet)
- PHENIX (Inclusive  $\pi^0$ ,  $h^\pm$ ,  $J/\psi$ )
- BRAHMS (inclusive  $\pi^\pm$ ,  $K^\pm$ ,  $p$ ,  $p\bar{p}$ )

## 3. Summary and next

# What is SSA (Single-transverse-Spin Asymmetry) ?

Definition: The ratio of the difference and the sum of the transverse **spin-dependent** differential cross-sections of **a certain interaction**.

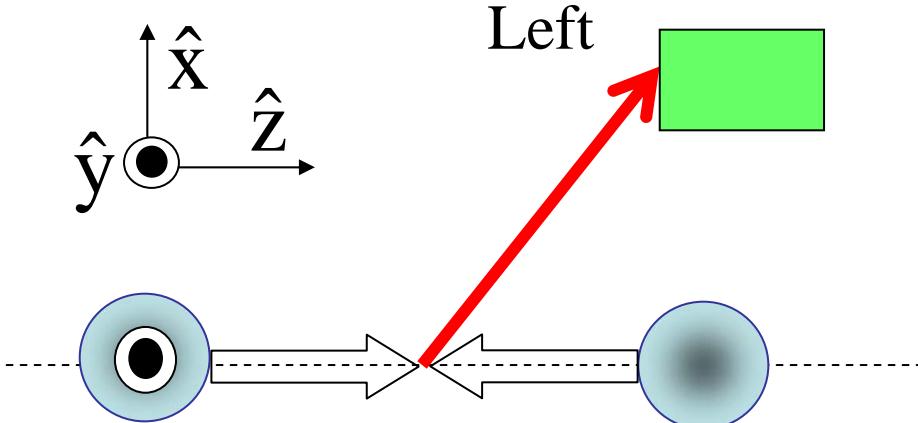
$$\text{SSA} = \frac{d\sigma_{\uparrow} - d\sigma_{\downarrow}}{d\sigma_{\uparrow} + d\sigma_{\downarrow}}$$

The diagram illustrates the calculation of Single-transverse-Spin Asymmetry (SSA). It shows two interactions between particles represented by blue circles. In the top interaction, the left particle has a spin arrow pointing up (↑), and the right particle has a spin arrow pointing down (↓). The differential cross-section for this interaction is labeled  $d\sigma_{\uparrow}$ . In the bottom interaction, the left particle has a spin arrow pointing down (↓), and the right particle has a spin arrow pointing up (↑). The differential cross-section for this interaction is labeled  $d\sigma_{\downarrow}$ . Both interactions show incoming and outgoing horizontal arrows representing the exchange of particles.

A certain interaction: elastic, hard scatter processes etc.

# How to measure SSA?

Top view



A certain interaction is detected by:

**Single arm detector (Left)**

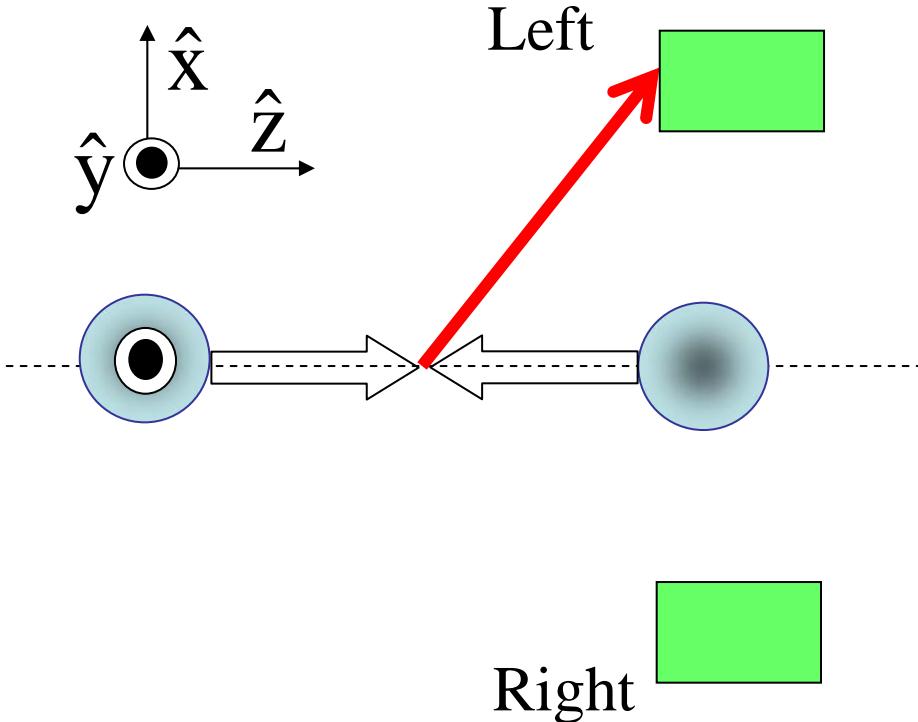
$$\text{SSA} = \frac{1}{\text{pol.}} \frac{N_L^\uparrow - RN_L^\downarrow}{N_L^\uparrow + RN_L^\downarrow}$$

R: Relative luminosity

- Normalization by relative luminosity is crucial.
- Normalization by beam polarization is crucial.

# How to measure SSA?

Top view



A certain interaction is detected by:

**Double arms detector (Left-Right)**

$$\text{SSA} = \frac{1}{\text{pol.}} \frac{\sqrt{N_{\uparrow}^L N_{\downarrow}^R} - \sqrt{N_{\uparrow}^R N_{\downarrow}^L}}{\sqrt{N_{\uparrow}^L N_{\downarrow}^R} + \sqrt{N_{\uparrow}^R N_{\downarrow}^L}}$$

Square-root-formula

- Normalization by beam polarization is crucial.

SSA of elastic process is ideal tool to measure beam pol.

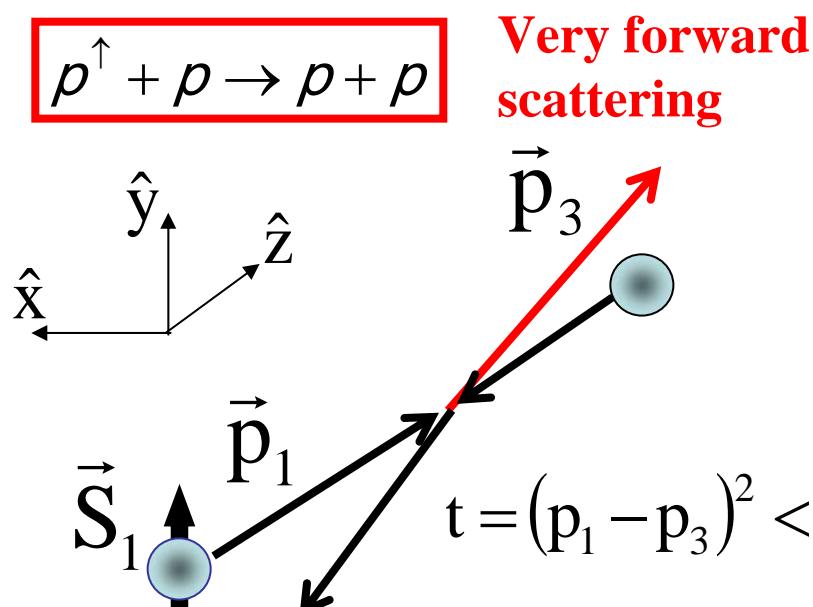
$$\text{pol.} = \frac{1}{\text{SSA}} \frac{\sqrt{N_{\uparrow}^L N_{\downarrow}^R} - \sqrt{N_{\uparrow}^R N_{\downarrow}^L}}{\sqrt{N_{\uparrow}^L N_{\downarrow}^R} + \sqrt{N_{\uparrow}^R N_{\downarrow}^L}}$$

**Non-zero and known SSA**

→ we can measure beam polarization

SSA of elastic process is ideal tool to measure beam pol.

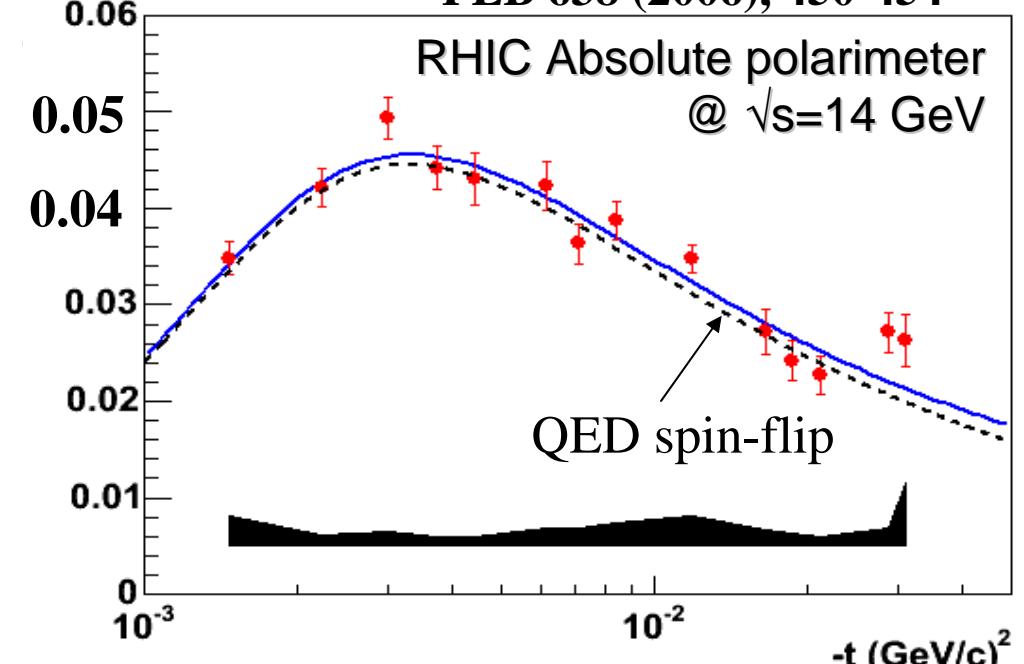
$$\text{pol.} = \frac{1}{\text{SSA}} \frac{\sqrt{N_{\uparrow}^L N_{\downarrow}^R} - \sqrt{N_{\uparrow}^R N_{\downarrow}^L}}{\sqrt{N_{\uparrow}^L N_{\downarrow}^R} + \sqrt{N_{\uparrow}^R N_{\downarrow}^L}}$$



Non-zero and known SSA

→ we can measure beam polarization  
SSA

PLB 638 (2006), 450-454



A  
p  
a  
For more details....

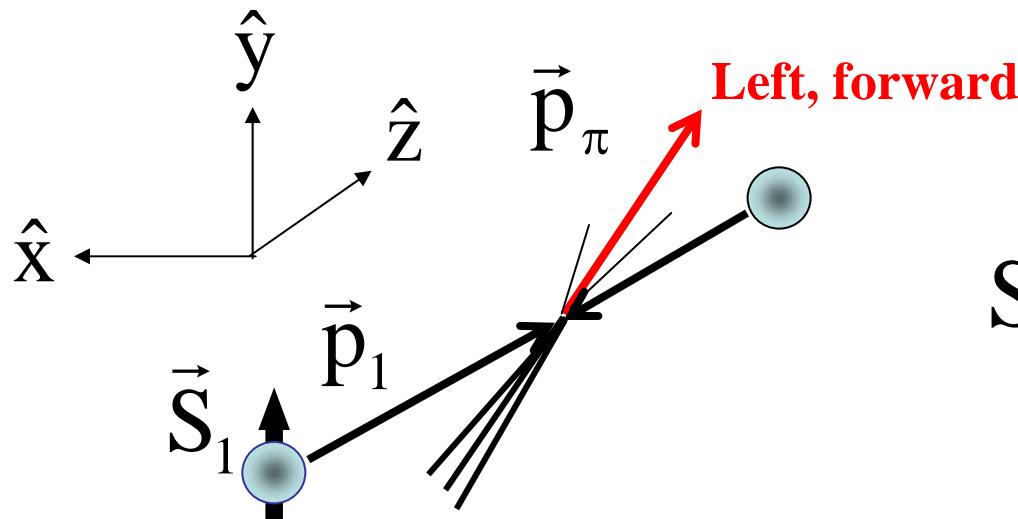
Please visit the poster session!

# How about SSA of hard scattering?

For example:

$$p^\uparrow + p \rightarrow \pi + X$$

Azimuthal asymmetry in singly polarized  $p\bar{p}$  collisions



$$\text{SSA} \propto \vec{S}_1 \cdot (\vec{p}_1 \times \vec{p}_\pi)$$

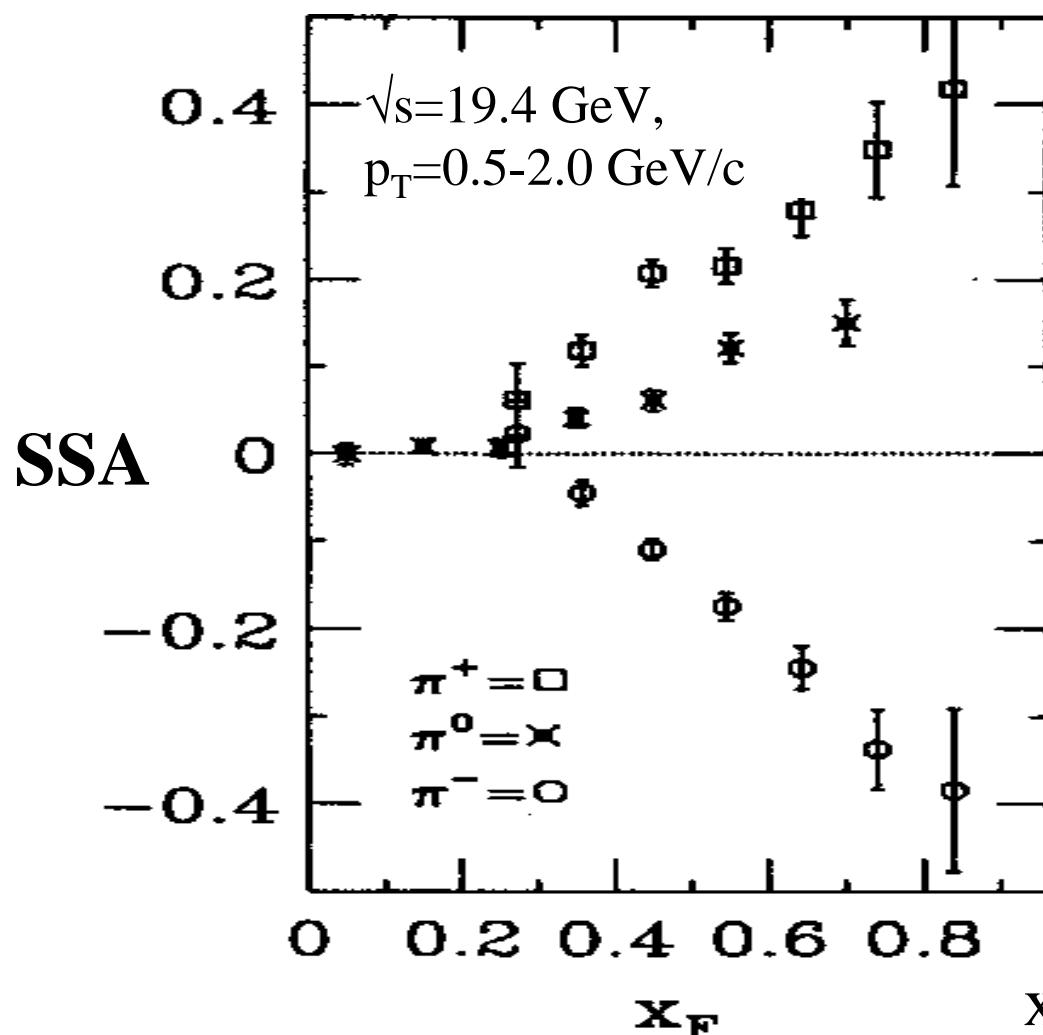
Prediction:  $\text{SSA} \sim \frac{m_q \alpha_s}{p_T} \sim 0.001$

Kane, Pumplin and Repko  
PRL 41 1689 (1978)

**SSA of hard scattering is expected to be very small**  
by the **leading twist collinear**,  $p_T \geq 4 \sim 5 \text{ GeV}/c$ .

# Huge SSAs have been measured at E704-FNAL!

$p^\uparrow + p \rightarrow \pi + X$  azimuthal asymmetry in singly polarized pp collisions



$\pi^0$  – E704, PLB261 (1991) 201.  
 $\pi^{+-}$  - E704, PLB264 (1991) 462.  
Citation amount: 258 !

- Increase linearly with Feynman x ( $x_F$ ).
- Extremely bigger than expectation!
- What is the  $p_T$  dependence?

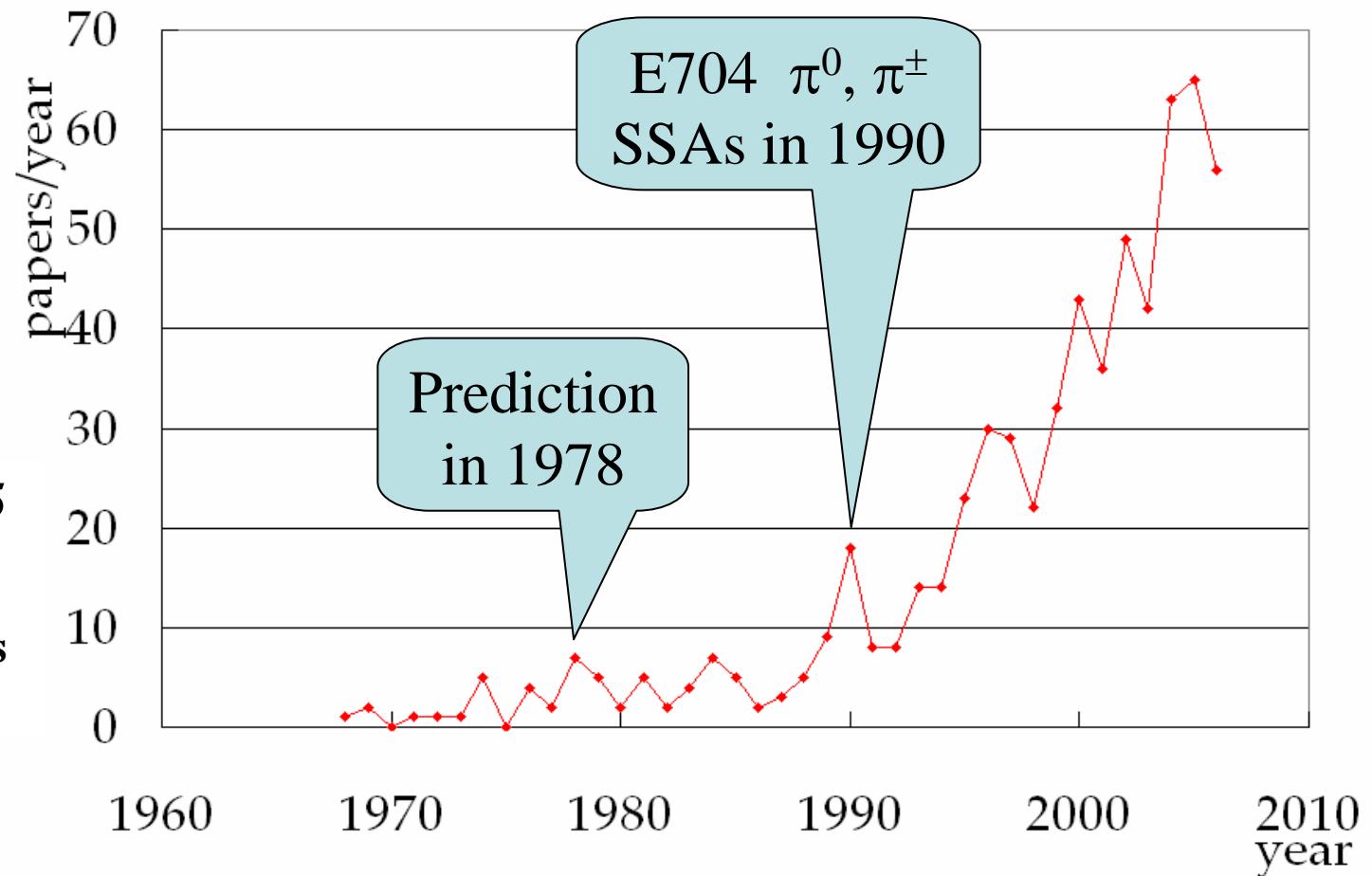
$$x_F = \frac{p_{z,\pi}}{p_{z,1}} \approx \frac{2E_\pi}{\sqrt{s}}$$

# How people excited about huge SSAs!

**SPIRES-HEP:** search title including:  
“Transverse spin, Transversity, single spin”.

**Total number: 625  
(1968~2006)**

**Experimental results  
~14%**



Theoretical approaches for huge SSA beyond the leading twist **collinear**;

### A) Higher-twist collinear QCD

→ Twist3 quark-gluon correlations for both initial and final state interactions.

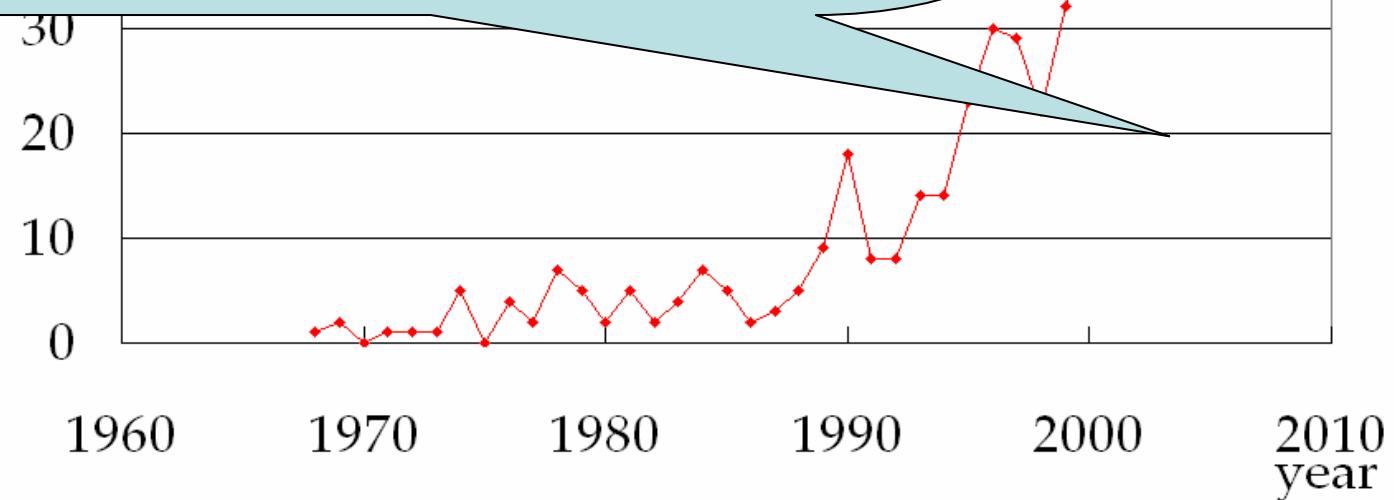
### B) Spin-correlated transverse momentum dependent (TMD) QCD

→ Initial state: Sivers effect

→ Final state: Collins effect

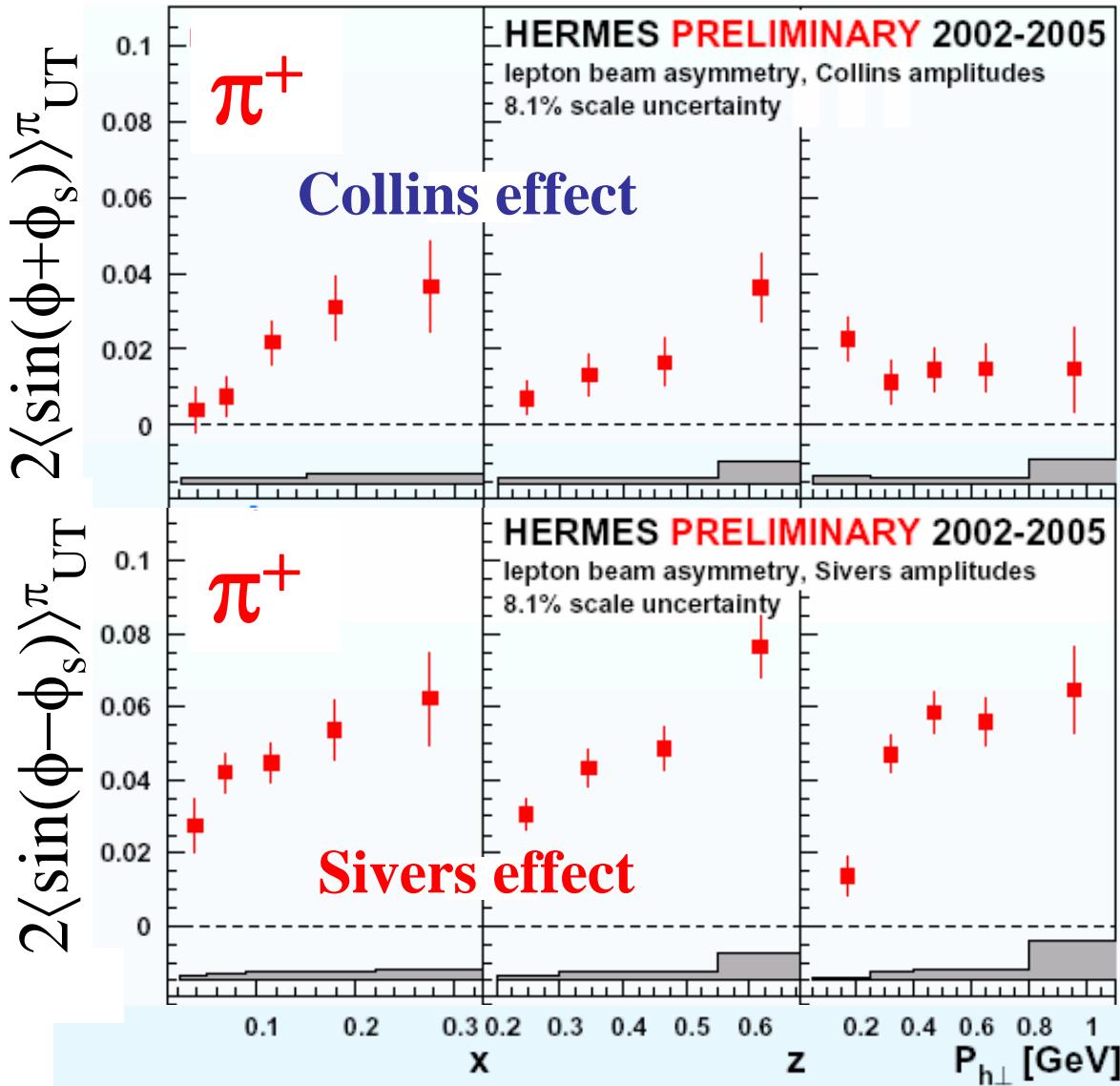
**Total number: 625  
(1968~2006)**

**Experimental results  
~14%**



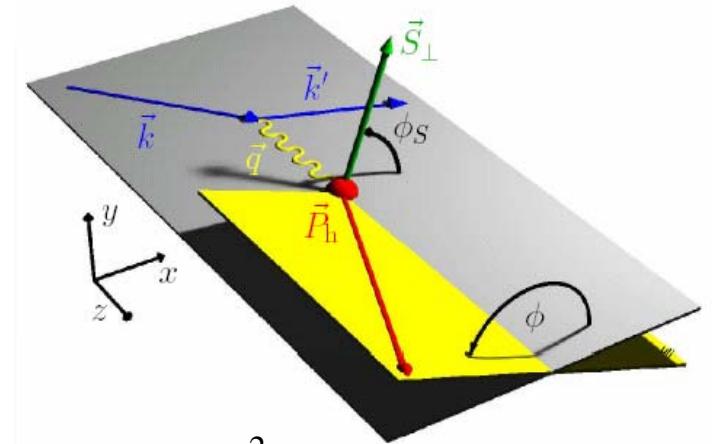
SSAs have been measured in Deep Inelastic Scattering, too.

Using lepton beam and polarized proton target.



$$\mathbf{k} = (E, \vec{k}), \mathbf{k}' = (E', \vec{k}')$$

$$-Q^2 = (\mathbf{k} - \mathbf{k}')^2 < 0$$



$$x = \frac{Q^2}{2M(E - E')}, \quad z = \frac{E_h}{E - E'}$$

$P_{h\perp}$  : Transverse momentum of  $\pi^+$

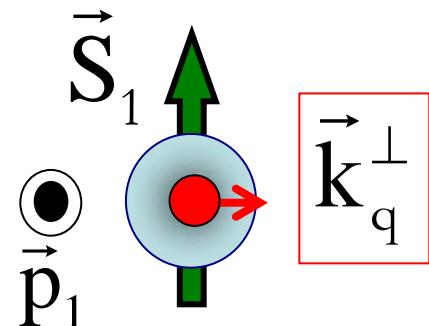
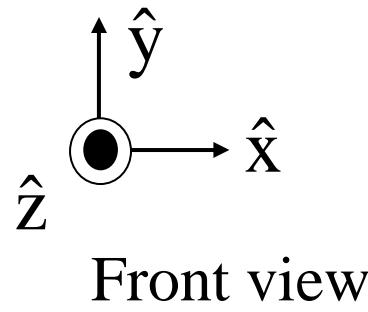
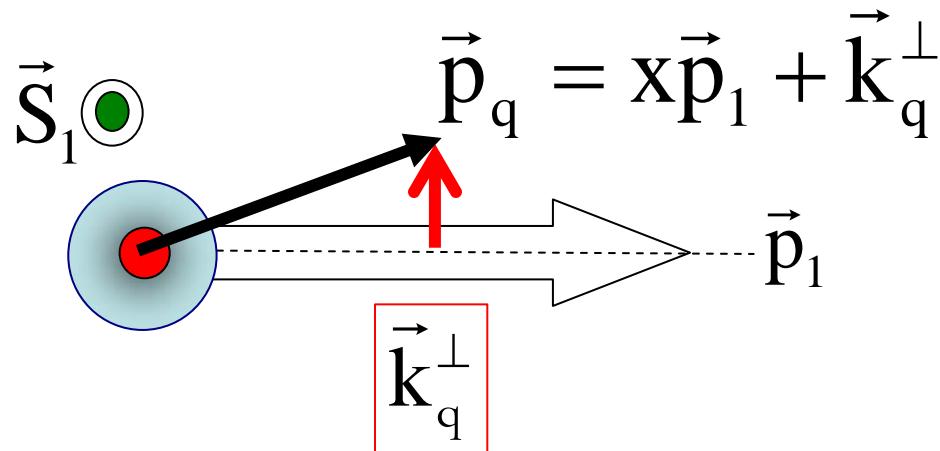
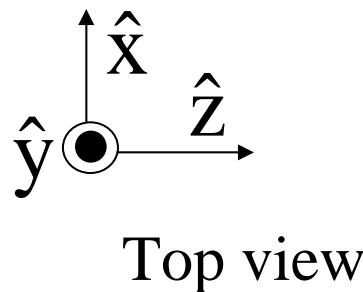
Ongoing experiments at  
**DESY, CERN etc.**

New program at **JLAB**  
12GeV upgrade.

# Sivers effect: **Initial state** of the polarized nucleon

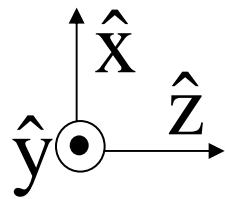
Phys Rev D41 (1990) 83; Phys Rev D43 (1991) 261

$$\text{SSA}_{\text{Sivers}} \propto \vec{S}_1 \cdot (\vec{p}_1 \times \vec{k}_q^\perp)$$



# What does Sivers effect probe?

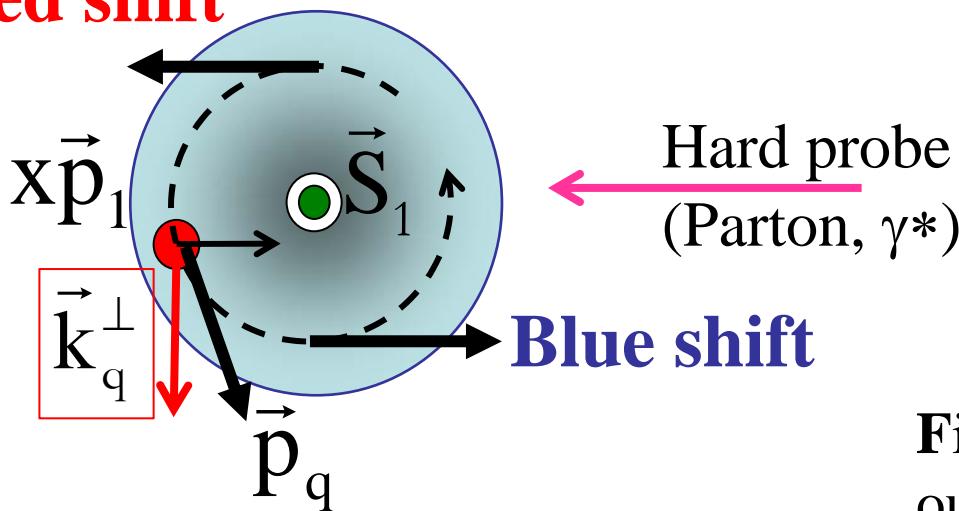
Top view, Breit frame



Quarks orbital motion adds/  
subtracts longitudinal momentum  
for negative/positive  $\hat{x}$ .

PRD66 (2002) 114005

**Red shift**



**Parton Distribution**

Functions rapidly fall in  
longitudinal momentum  
fraction x.

Final State Interaction between  
outgoing quark and target spectator.

Sivers function

$$f_{1T}^\perp(x, \vec{k}_q^\perp)$$

hep-ph/  
0703176

**Quark Orbital  
angular momentum**

Generalized Parton  
Distribution Functions

PRD59 (1999) 014013

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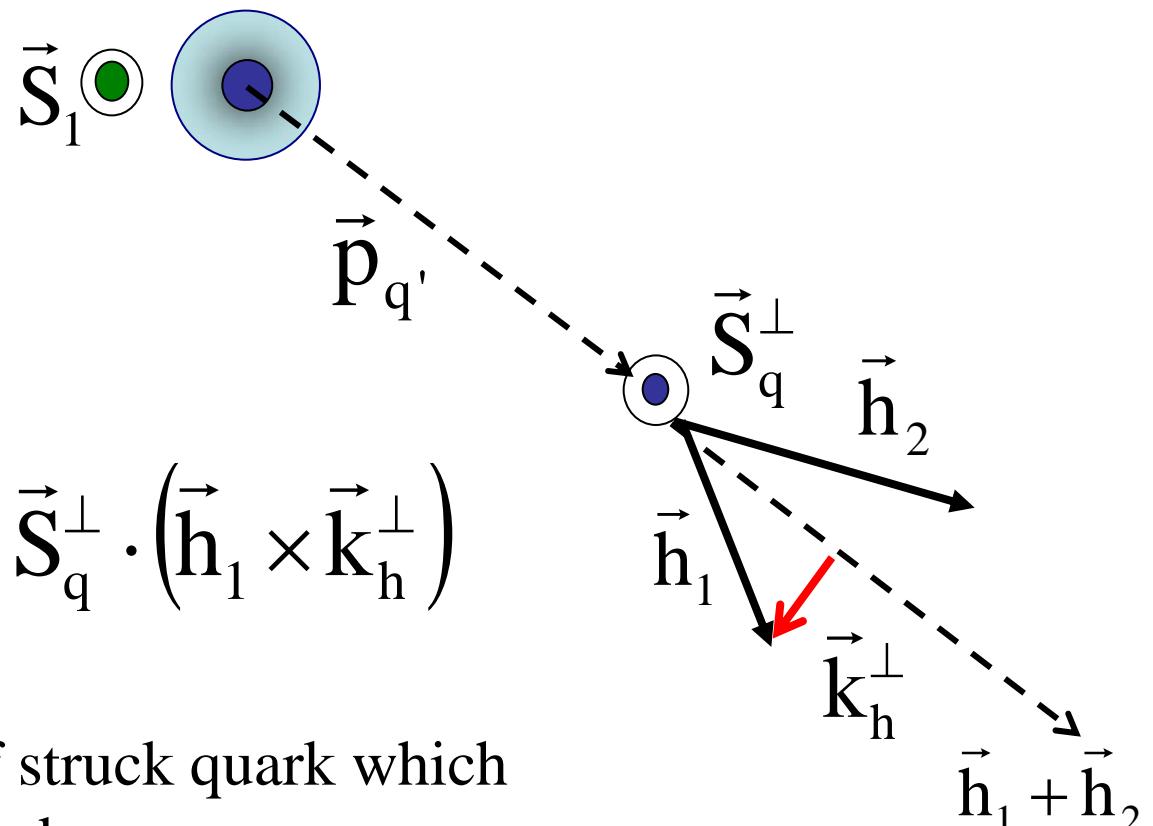
# Collins Heppelmann effect: **Final state** of fragmentation hadrons

Example:  $p^\uparrow + p \rightarrow h_1 + h_2 + X$

Nucl Phys B396 (1993) 161,  
Nucl Phys B420 (1994) 565

$$\text{SSA}_{\text{Collins}} \propto \vec{S}_q^\perp \cdot (\vec{h}_1 \times \vec{k}_h^\perp)$$

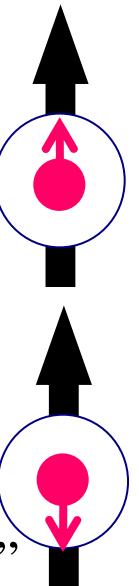
Polarization of struck quark which  
fragments to hadrons.



## Collins function: analyzer of “Transversity”

$$\text{Transversity: } \delta q(x, \mu) = q_{\uparrow}^{\uparrow}(x, \mu) - q_{\uparrow}^{\downarrow}(x, \mu)$$

$q_{\uparrow}^{(\uparrow)\downarrow}(x, \mu)$ : Probability to observe **parton** whose pol. vector is “with” or “against” the proton pol. vector with the renormalization scale  $\mu$ .



“with”

“against”

- ✓  $\delta q(x, \mu)$  has not been measured experimentally.
- ✓ Lattice QCD calculates the first moments of  $\delta q(x, \mu)$  for u,d, s quarks and the sum at  $\mu^2=2 \text{ GeV}^2$ .

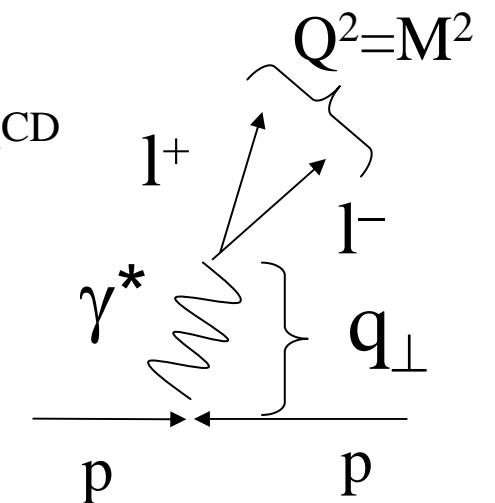
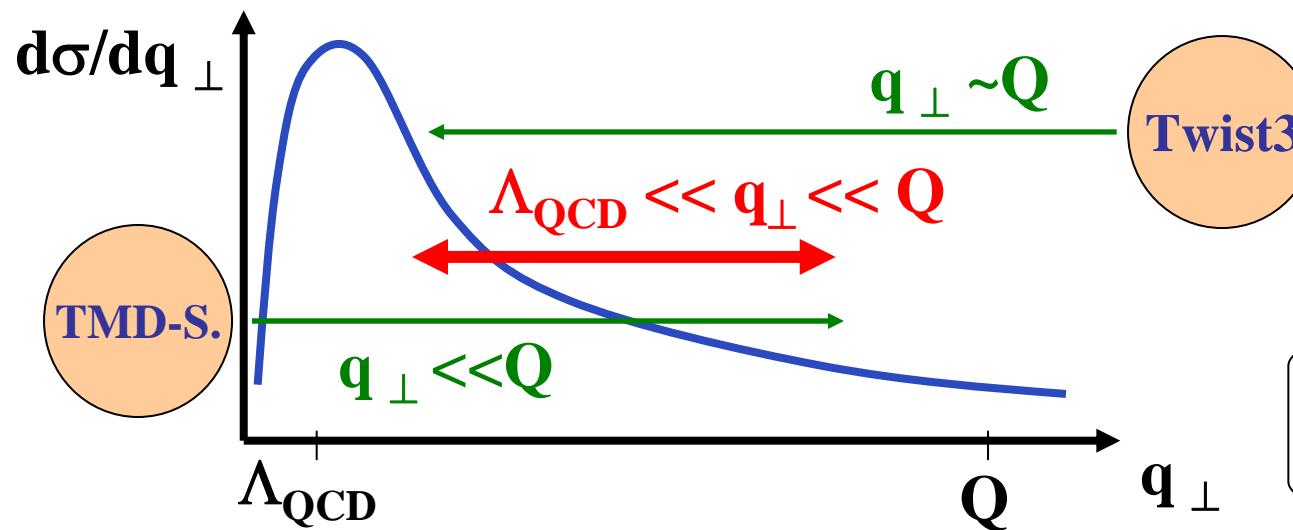
$$\delta q(\mu) = \int_0^1 dx [\delta q(x, \mu) - \delta \bar{q}(x, \mu)]$$

$$\delta \Sigma(\mu) = \delta u(\mu) + \delta d(\mu) + \delta s(\mu)$$

# Relevance of **Twist3** and **Spin-correlated TMD**

Lepton-pair (Drell-Yan) production is examined:

- Twist3 (Higher-twist Collinear QCD):  $q_\perp \sim Q \gg \Lambda_{\text{QCD}}$
- TMD-Sivers QCD:  $q_\perp \ll Q$



PRL97, 082002 (2006)  
PRD73, 094017 (2006)

**Twist3 and TMD-Sivers are related at  $\Lambda_{\text{QCD}} << q_\perp << Q$**

Towards consistency of phenomenological studies of SSA!

# Current understanding of SSA

1. SSA of hard scattering is expected to be very small and have  $p_T$  dependence by the **leading twist collinear**.
2. Huge SSAs have been measured in  **$pp$  collision** and **SIDIS** (semi-inclusive deep inelastic scattering).
3. Theoretical approaches to explain huge SSAs:
  - **Sivers** effect ( $\vec{k}_q^\perp$  is connected to quark orbital angular momentum).
  - **Collins** effect (Analyzer of transversity  $\delta q$ ).
  - **Twist3** effect which is related to both initial and final states. Relation of Twist3 to Sivers effect is introduced.

# Issues

1. Applicability of pQCD?
2. Relationship of SSAs between SIDIS and  $pp$  collision?

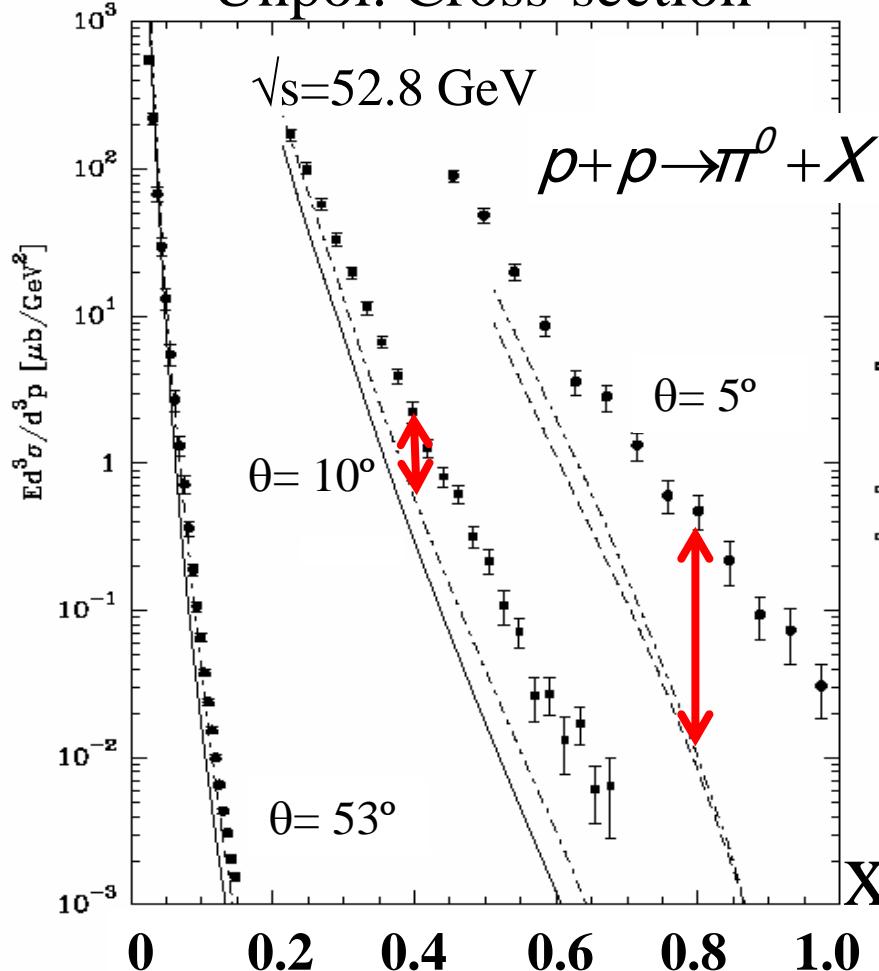
# Is pQCD applicable for $\sqrt{s} \leq 50$ GeV ?

Eur.Phys.J.C36:371-374,2004. ,

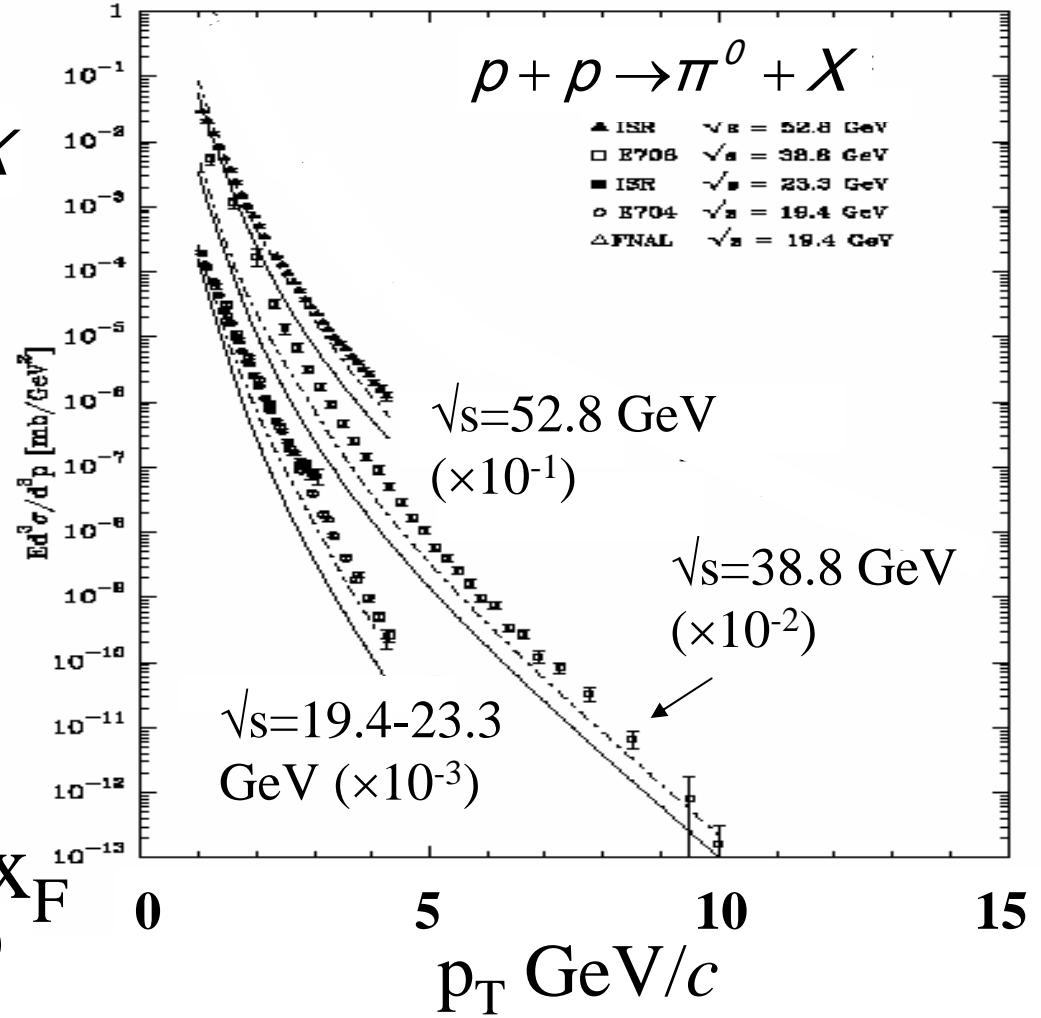
Data references therein.

2 NLO calculations with different scales  $p_T$  and  $p_T/2$

## Unpol. Cross-section



## Unpol. Cross-sections at $\theta=90^\circ$



2007/6/2

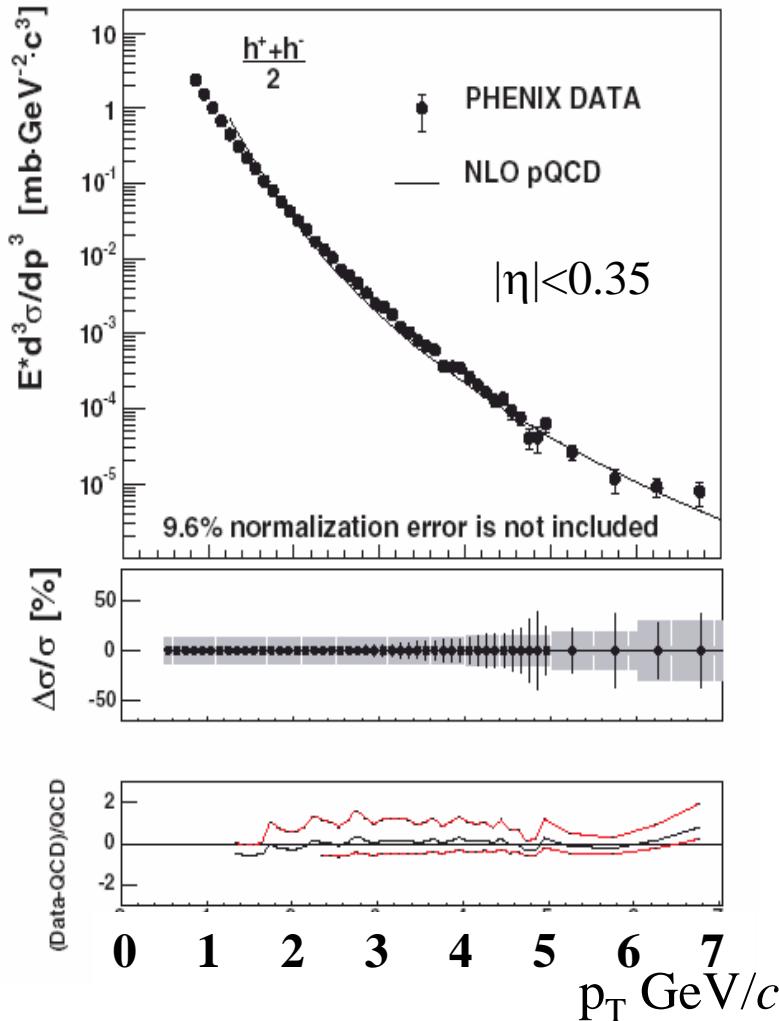
Require more than NLO pQCD...

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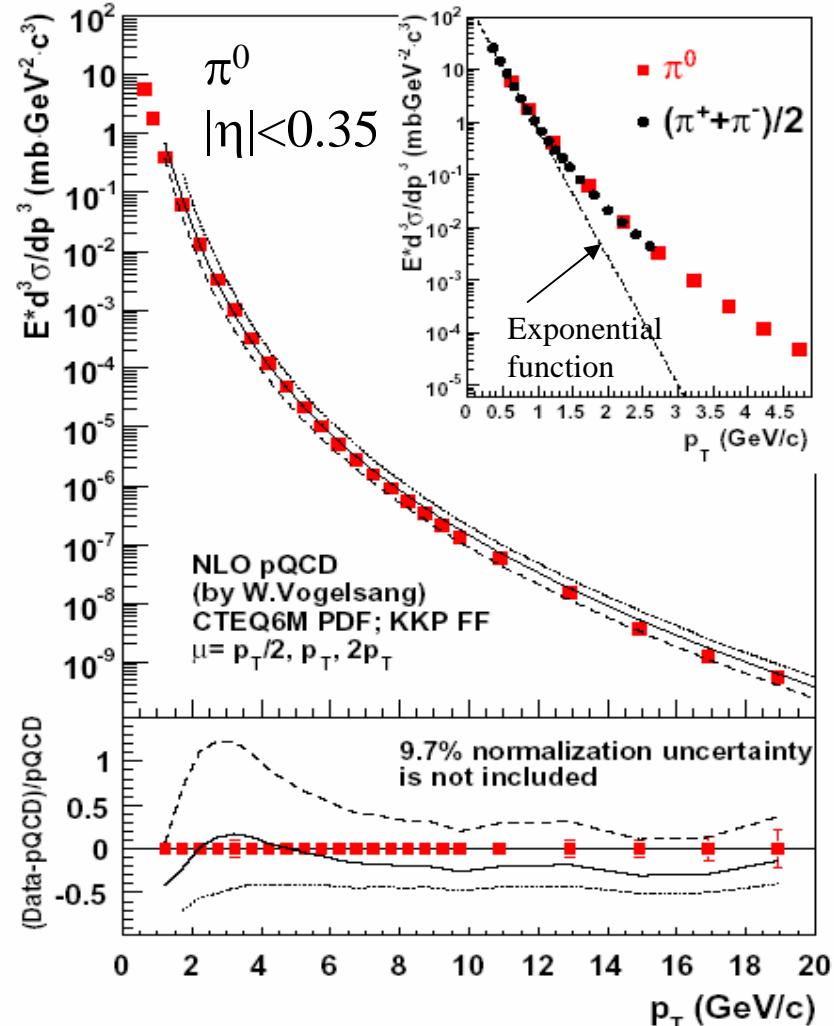
# $pp$ collision at $\sqrt{s}=200\text{GeV}$ (1)



PRL95:202001,2005.



hep-ex/0704.3599



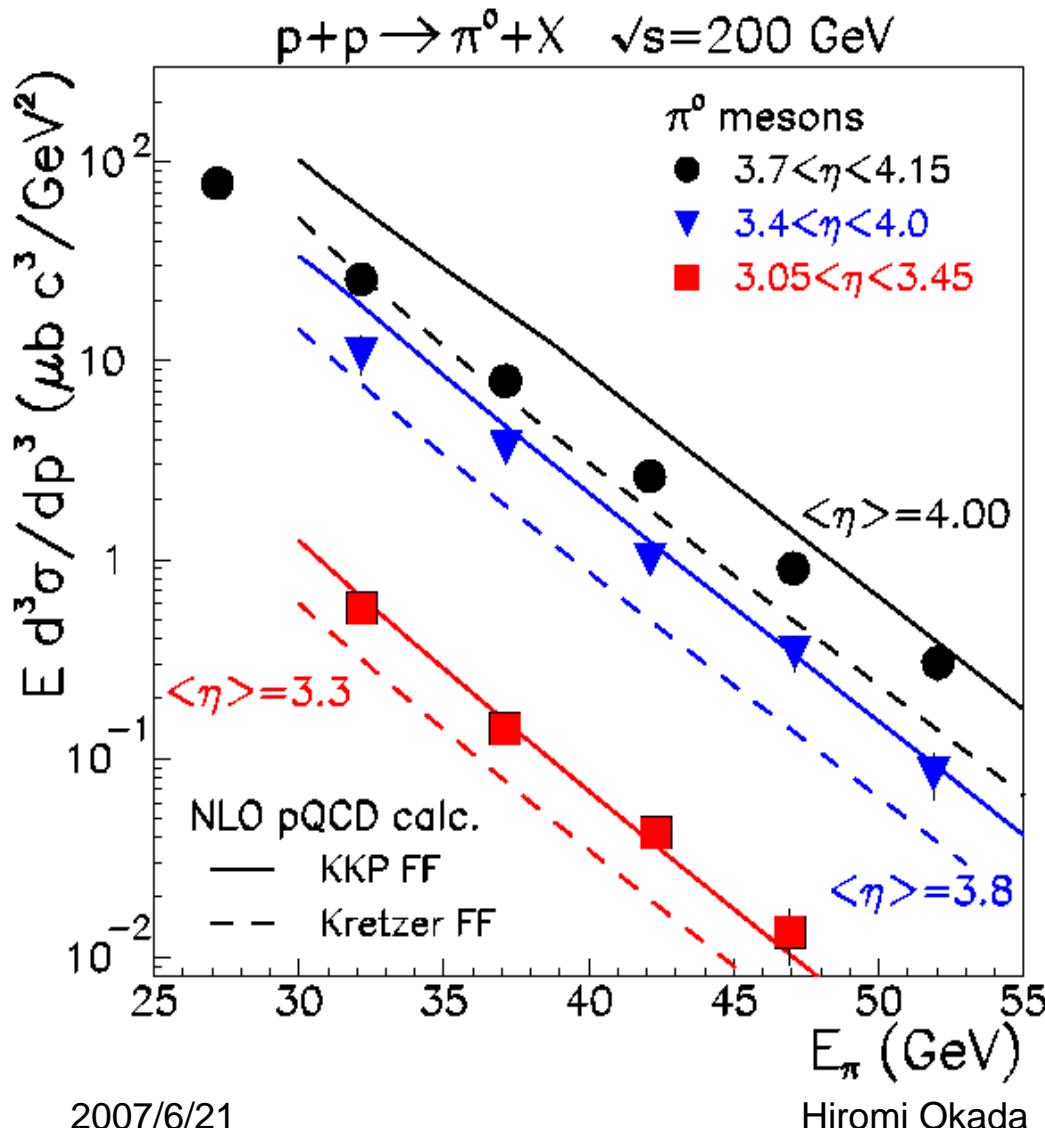
Cross-sections at  $|\eta| < 0.35$  are consistent with NLO pQCD.

W. Vogelsang, M. Stratmann *et al.*

# $pp$ collision at $\sqrt{s}=200\text{GeV}$ (2)



PRL97:152302, 2006.



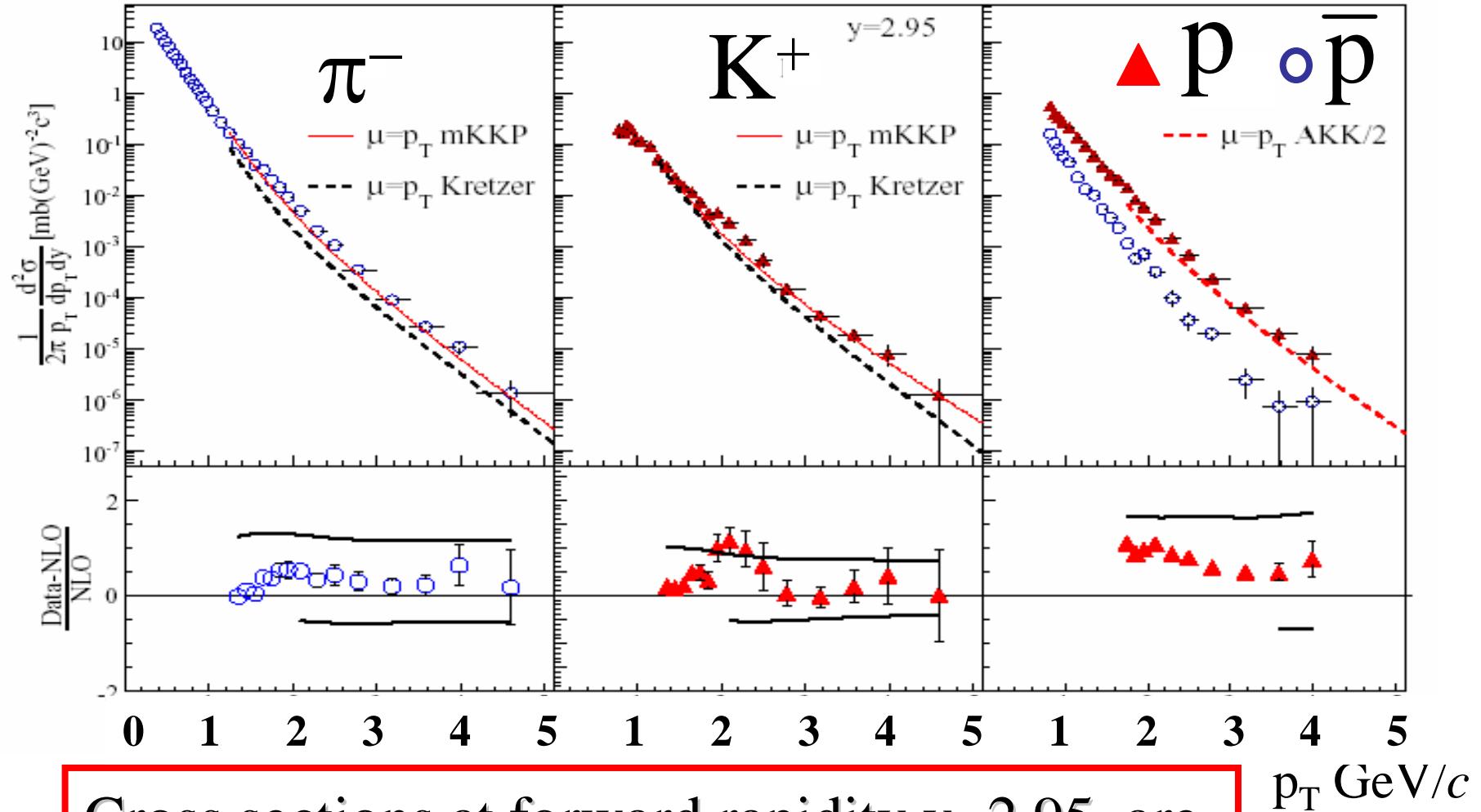
Cross-sections at the forward region,  
 $\langle \eta \rangle = 3.3, 3.8$  and  $4.0$   
are mostly consistent  
with NLO pQCD  
calculations.

KKP FF : Nucl. Phys. B597, 337 (2001)  
Kretzer FF: Phys. Rev. D62, 054001 (2000).

# $pp$ collision at $\sqrt{s}=200\text{GeV}$ (3)



Accepted for publication in Phys. Rev. Lett. hep-ex/0701041



Cross sections at forward rapidity  $y=2.95$  are consistent with NLO pQCD.

$p_T \text{ GeV}/c$

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mKKP: Nucl. Phys. B597, 337 (2001); Kretzer: Eur. Phys. J.C22, 269 (2001); AKK: Nucl. Phys. B725, 181 (2005)

- RHIC provides  $\sqrt{s} = 200$  GeV polarized  $pp$  collision.
- Hard process, pQCD is applicable.  
⇒ Advantage of SSA investigation at RHIC!

# Establish Universality of SSA

hadron-hadron high energy collision for Sivers/Collins at RHIC

Intercomparison via pQCD

Semi-inclusive deep inelastic scattering for Sivers/Collins.  
(HERMES, COMPASS, ..)

	Available Probes at RHIC	
1	$p^\uparrow + p \rightarrow h + X$	Both mix
2	$p^\uparrow + p \rightarrow \text{di-jet} + X$	Sivers?
3	$p^\uparrow + p \rightarrow h + h + X$ (far side)	Separate?
4	$p^\uparrow + p \rightarrow h + h + X$ (near side) $p^\uparrow + p \rightarrow \text{jet} + X$	Collins Sivers
5	$p^\uparrow + p \rightarrow \text{direct } \gamma + X$	Sivers
6	$p^\uparrow + p \rightarrow l^+ l^-$ (Drell-Yan)	Sivers

$e^+ e^-$  collision for Collins.  
(BELLE)

I am going to show!

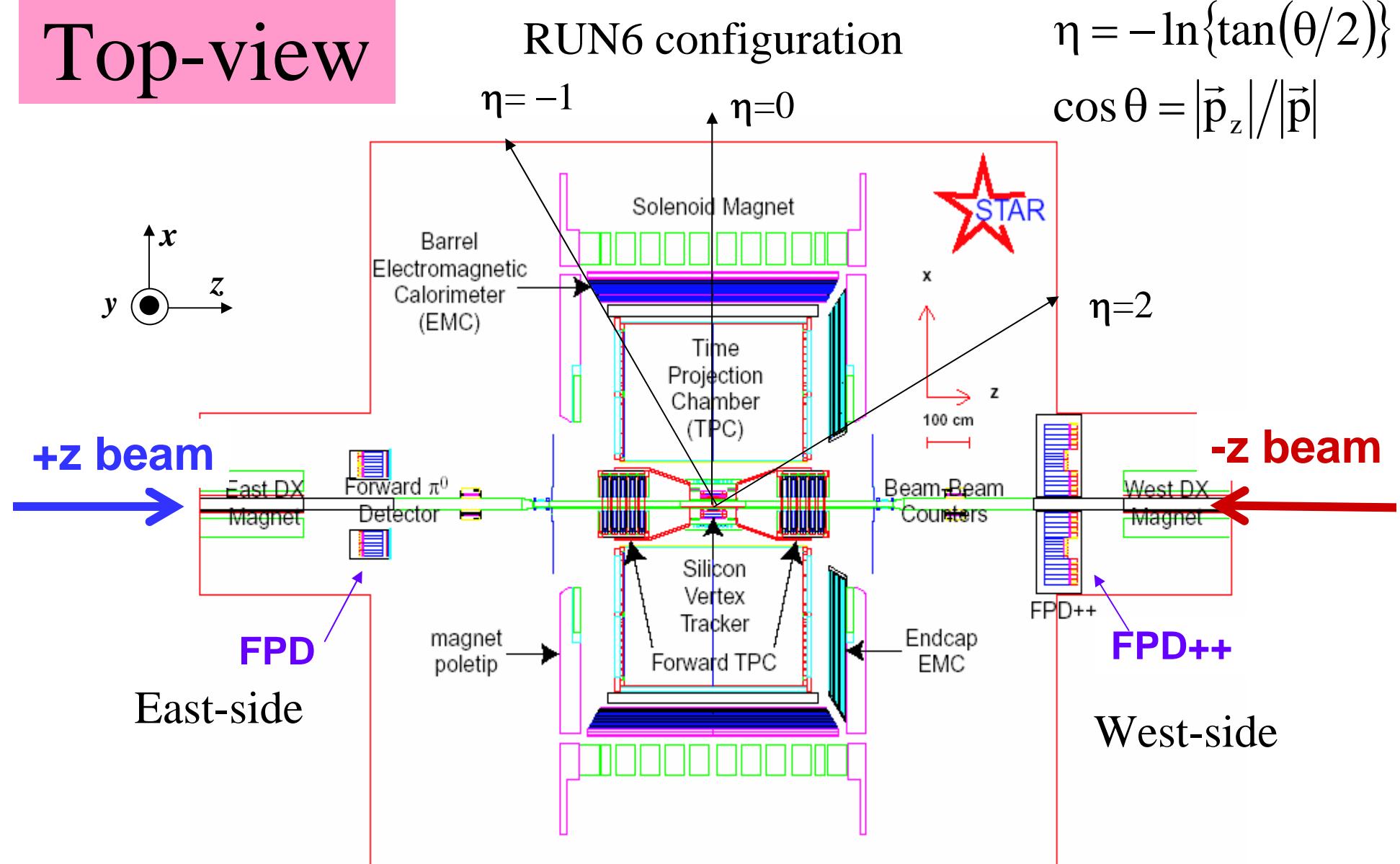
Probe “color charge” by comparing with SIDIS process. I will come back this topic, later.



Solenoidal  
Tracker  
At  
RHIC

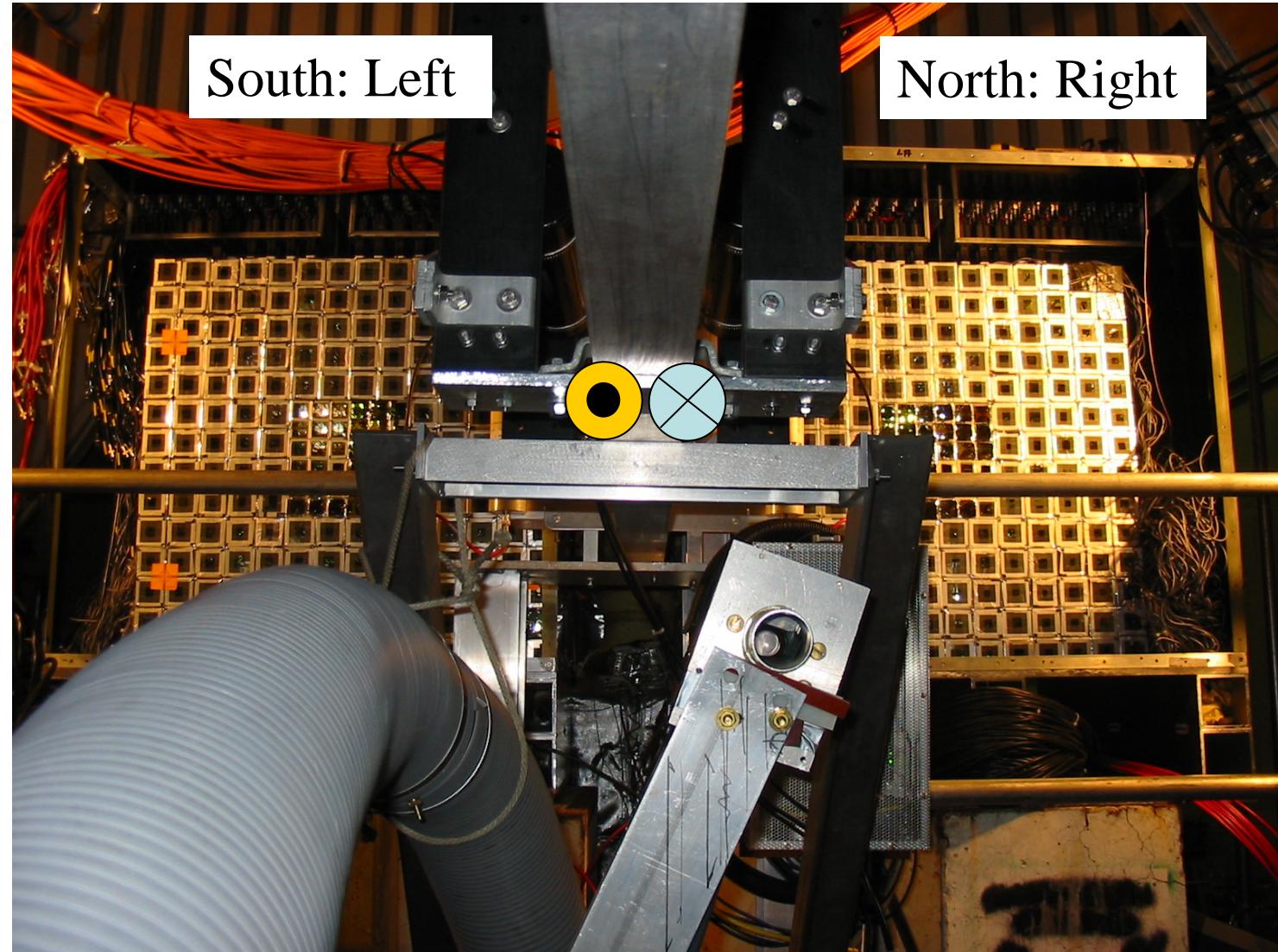
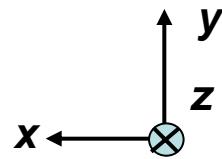
Inclusive  $\pi^0$ , Di-jet

# Top-view



- Inclusive  $\pi^0$  in forward region:  $-4 < \eta < -3$  (FPD),  $2.5 < \eta < 4$  (FPD++)
- Di-jet results:  $-1 < \eta < 2$  (Barrel EMC, Endcap EMC,  $2\pi$ )

# FPD++ (Forward Pion Detector) , West side



Inclusive  $\pi^0$



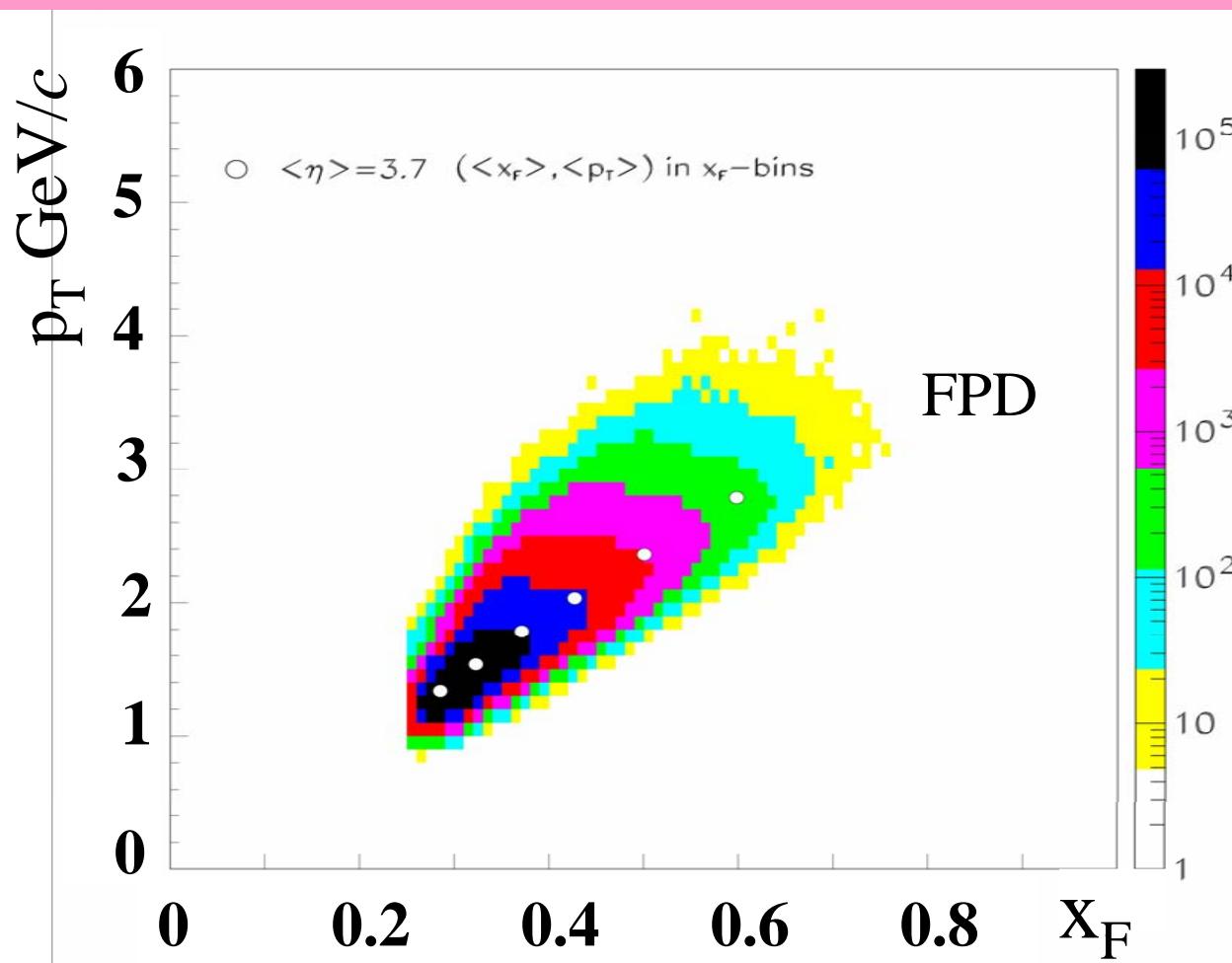
2007/6/21

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## Acceptance of FPD

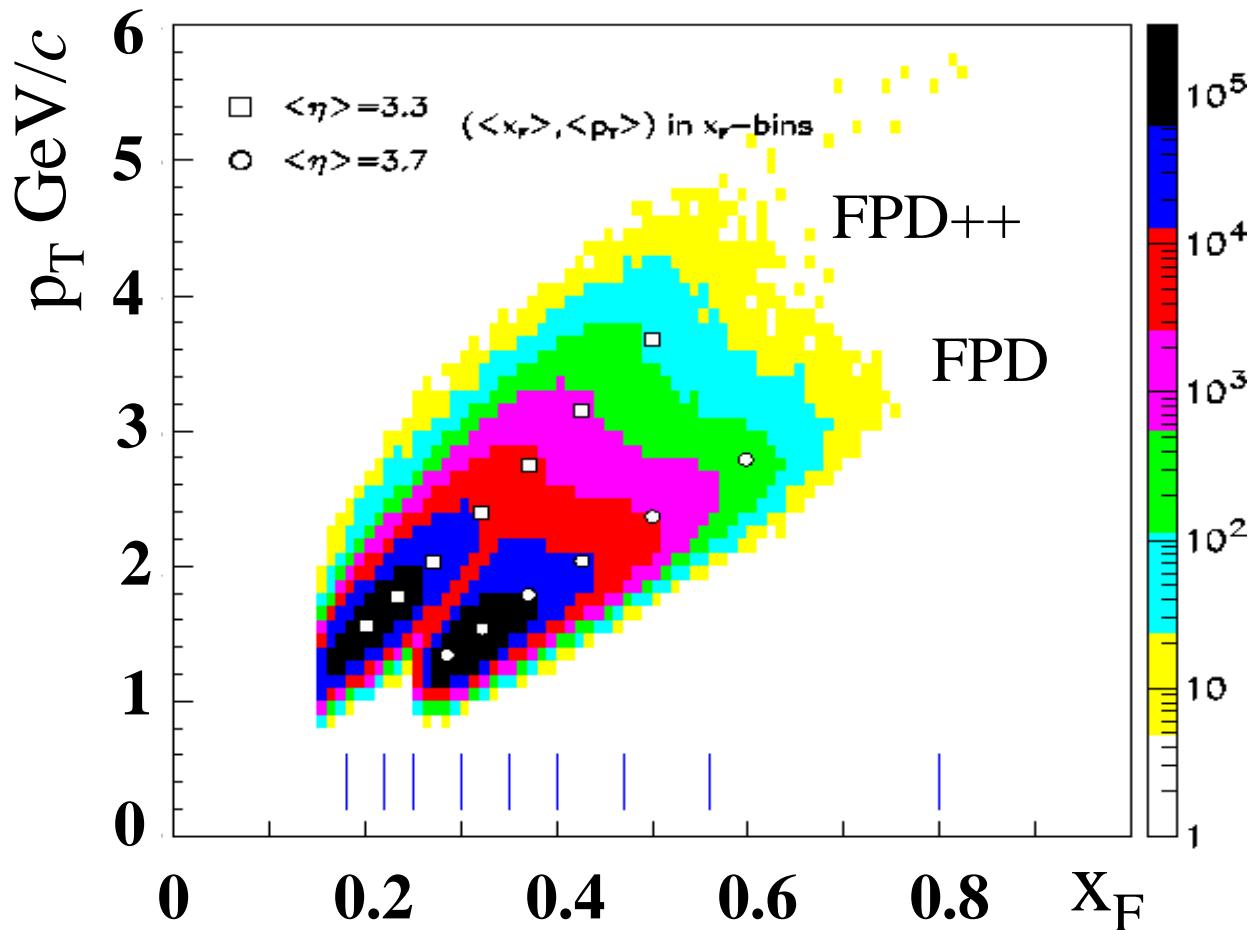
Inclusive  $\pi^0$   

Strong  $x_F$  and  $p_T$  correlation because of limited acceptance.

## Acceptance of FPD and FPD++

Inclusive  $\pi^0$   

Study of the  $p_T$  dependence needs large acceptance.

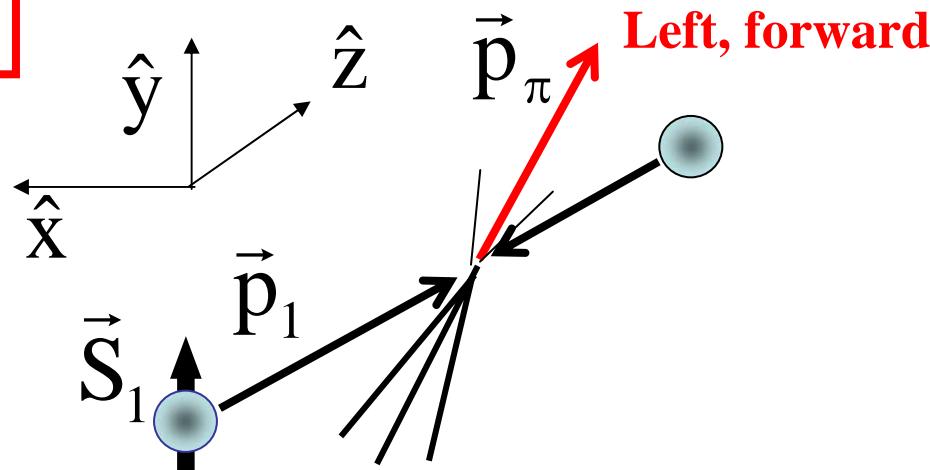
# Overview of transverse spin runs at STAR with forward calorimetry: 2001→2006

	Run2	Run3	Run5	Run6 (2006)
detector	EEMC and FPD prototypes	6 matrices of FPD	full FPD (8 matrices)	East FPD West FPD++
$P_{BEAM}^{\text{sampled}}, \%$	~15	~30	~45	~60
$\int L dt, pb^{-1}$	0.15	0.25	0.1	6.8
$< \eta >$	3.8	±3.3/±4.0	±3.7/±4.0	-3.7/3.3

FOM ( $P^2L$ ) in Run 6 is ~50 times larger than from all the previous STAR runs, and ~ 725 times larger than for Run 2

# SSA of inclusive $\pi^0$ production with FPD and FPD++

$$p^\uparrow + p \rightarrow \pi^0 + X$$

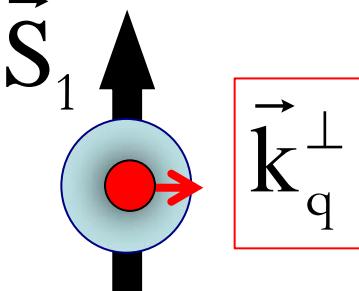


Inclusive  $\pi^0$   
**STAR**

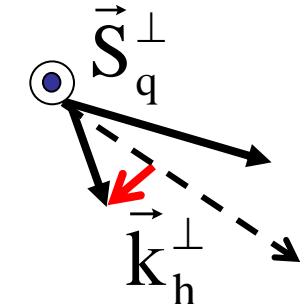
Measure;

$$\vec{S}_1 \cdot (\vec{p}_1 \times \vec{p}_\pi) \Rightarrow \frac{A \cdot [\vec{S}_1 \cdot (\vec{p}_1 \times \vec{k}_q^\perp)]}{\vec{S}_1 \cdot \vec{k}_q^\perp} + \frac{B \cdot [\vec{S}_q \cdot (\vec{h}_1 \times \vec{k}_h^\perp)]}{\vec{S}_q \cdot \vec{k}_h^\perp}$$

**Sivers**

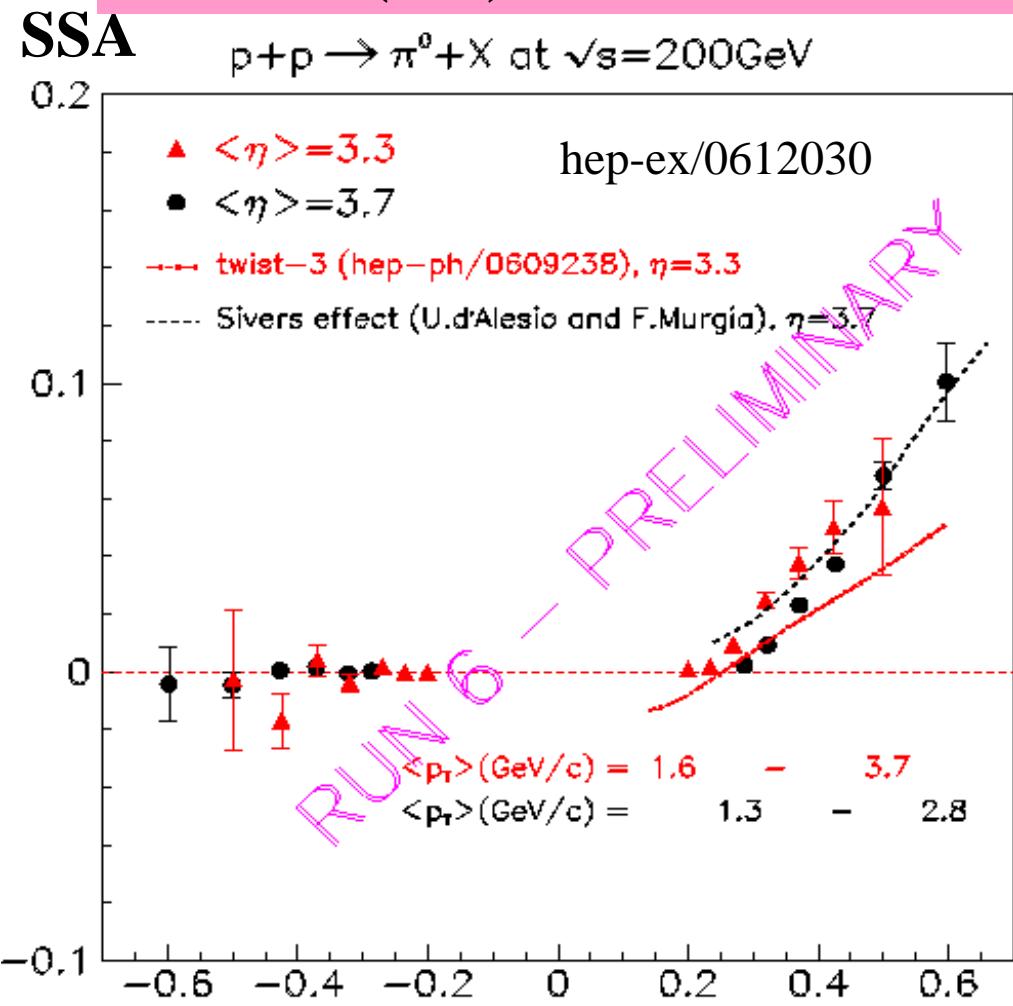


**Collins**



Does SSA survive in pQCD region?

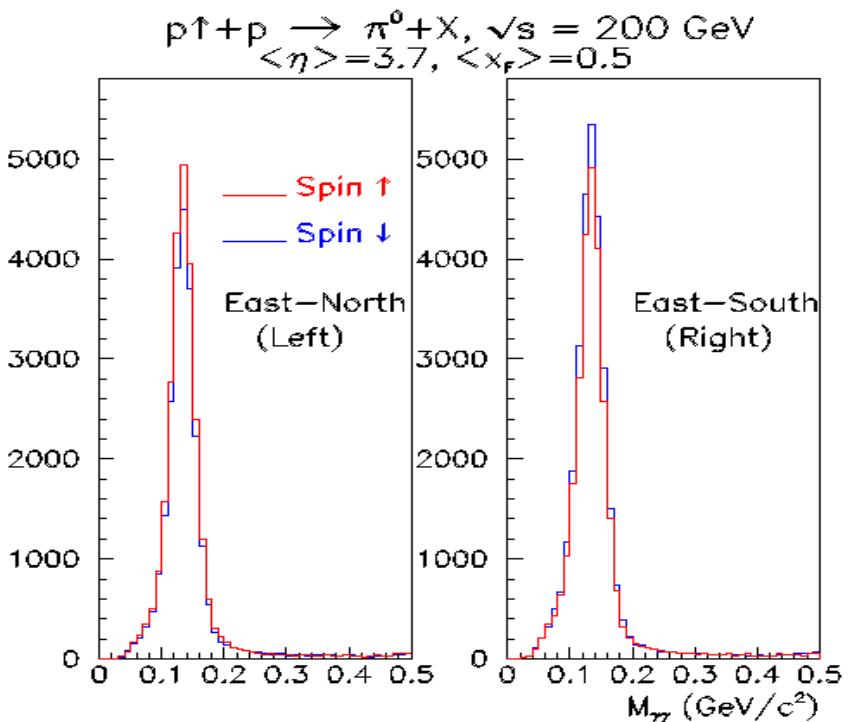
# SSA ( $\pi^0$ ) at $\sqrt{s}=200$ GeV SURVIVES !



2- $\gamma$  inv.-mass spectra

$47 \leq E_\pi < 56$  GeV

Inclusive  $\pi^0$   
**STAR**



$$x_F = \frac{p_{z,\pi}}{p_{z,1}} < 0, \quad x_F = \frac{p_{z,\pi}}{p_{z,1}} > 0$$

2007/6/21

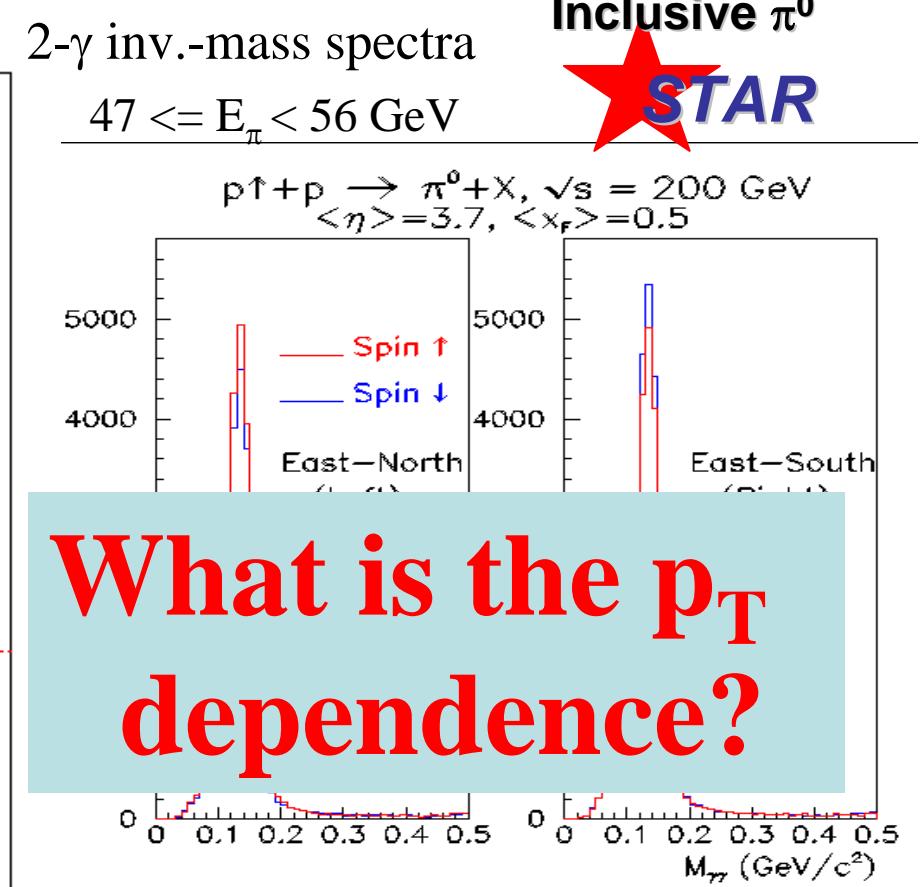
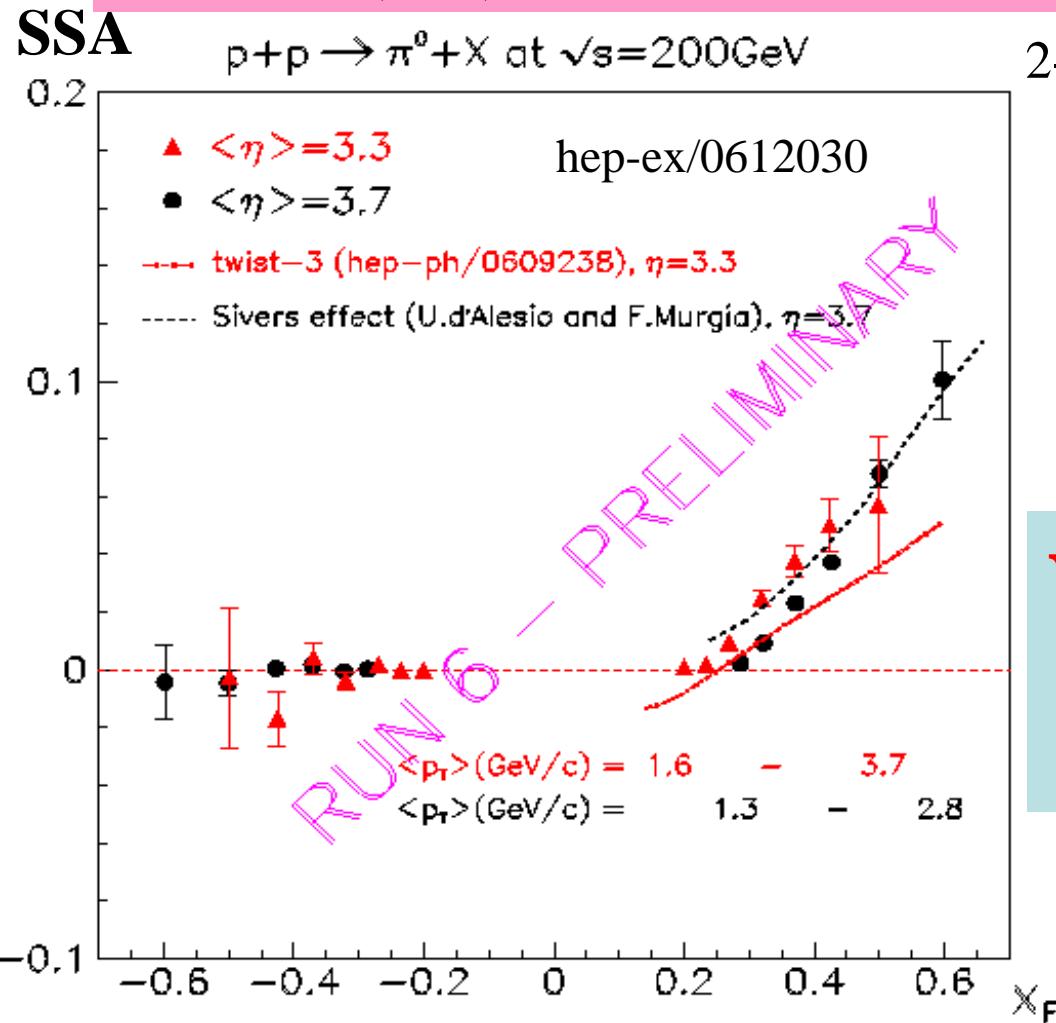
$$x_F = \frac{p_{z,\pi}}{p_{z,1}}$$

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Definition

$$\text{SSA} = \frac{1}{\text{pol.}} \frac{\sqrt{N_{\uparrow}^L N_{\downarrow}^R} - \sqrt{N_{\uparrow}^R N_{\downarrow}^L}}{\sqrt{N_{\uparrow}^L N_{\downarrow}^R} + \sqrt{N_{\uparrow}^R N_{\downarrow}^L}}$$

# SSA ( $\pi^0$ ) at $\sqrt{s}=200$ GeV SURVIVES !



$$x_F = \frac{p_{z,\pi}}{p_{z,1}} < 0, \quad x_F = \frac{p_{z,\pi}}{p_{z,1}} > 0$$

2007/6/21

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- SSA is measured by square-root formula with double-arms (left-right) detector:
- SSAs at different  $<\eta>$  do not change much.

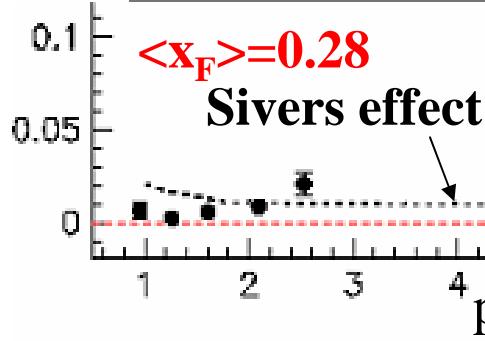
# SSA as a function of $p_T$ in $x_F$ slice

Inclusive  $\pi^0$   
**STAR**

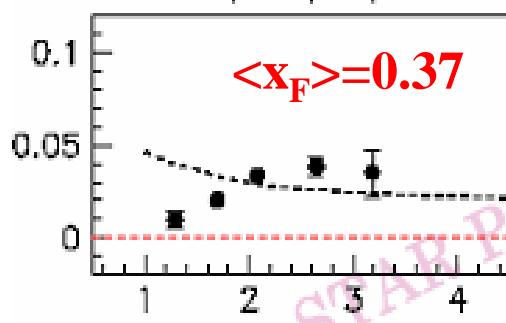
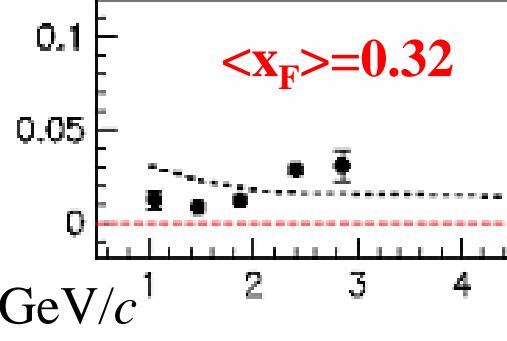
hep-ex/0612030

$p+p \rightarrow \pi^0 + X$  at  $\sqrt{s} = 200\text{GeV}$

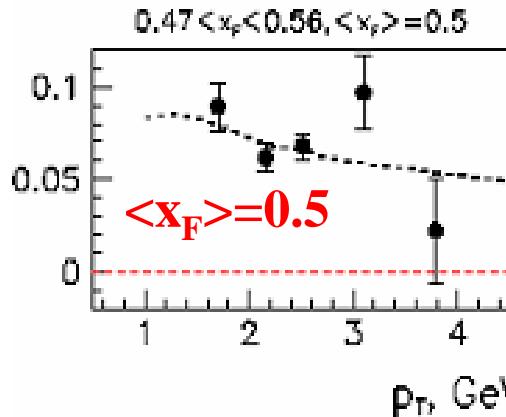
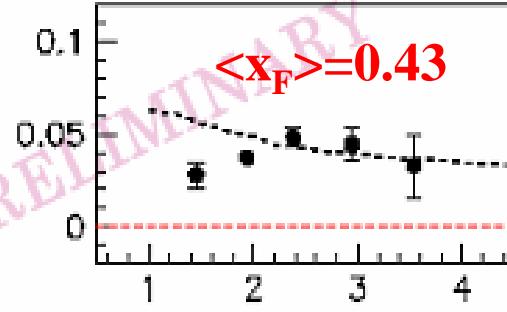
SSA  $0.25 < x_F < 0.3, \langle x_F \rangle = 0.28$



$0.3 < x_F < 0.35, \langle x_F \rangle = 0.32$

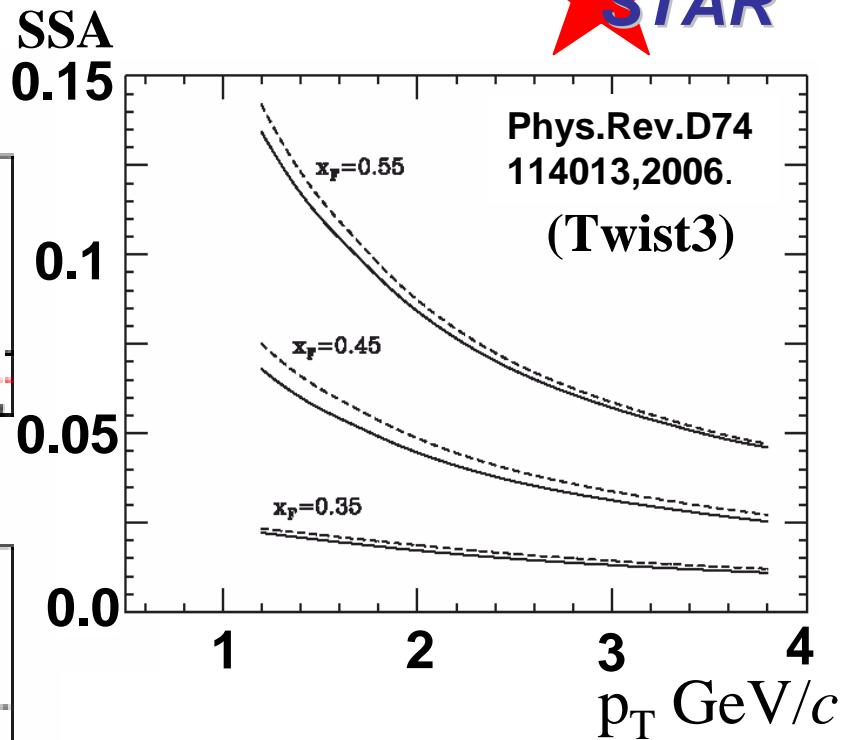


$0.4 < x_F < 0.47, \langle x_F \rangle = 0.43$



Sivers effect calculation:  
Phys. Rev. D70, 074009 (2004)

(U.d'Alesio and F.Murgia)



- Combined data from three runs at  $\langle \eta \rangle = 3.3, 3.7$  and  $4.0$ .
- Within each  $x_F$  bin,  $\langle x_F \rangle$  does not significantly change with  $p_T$ .

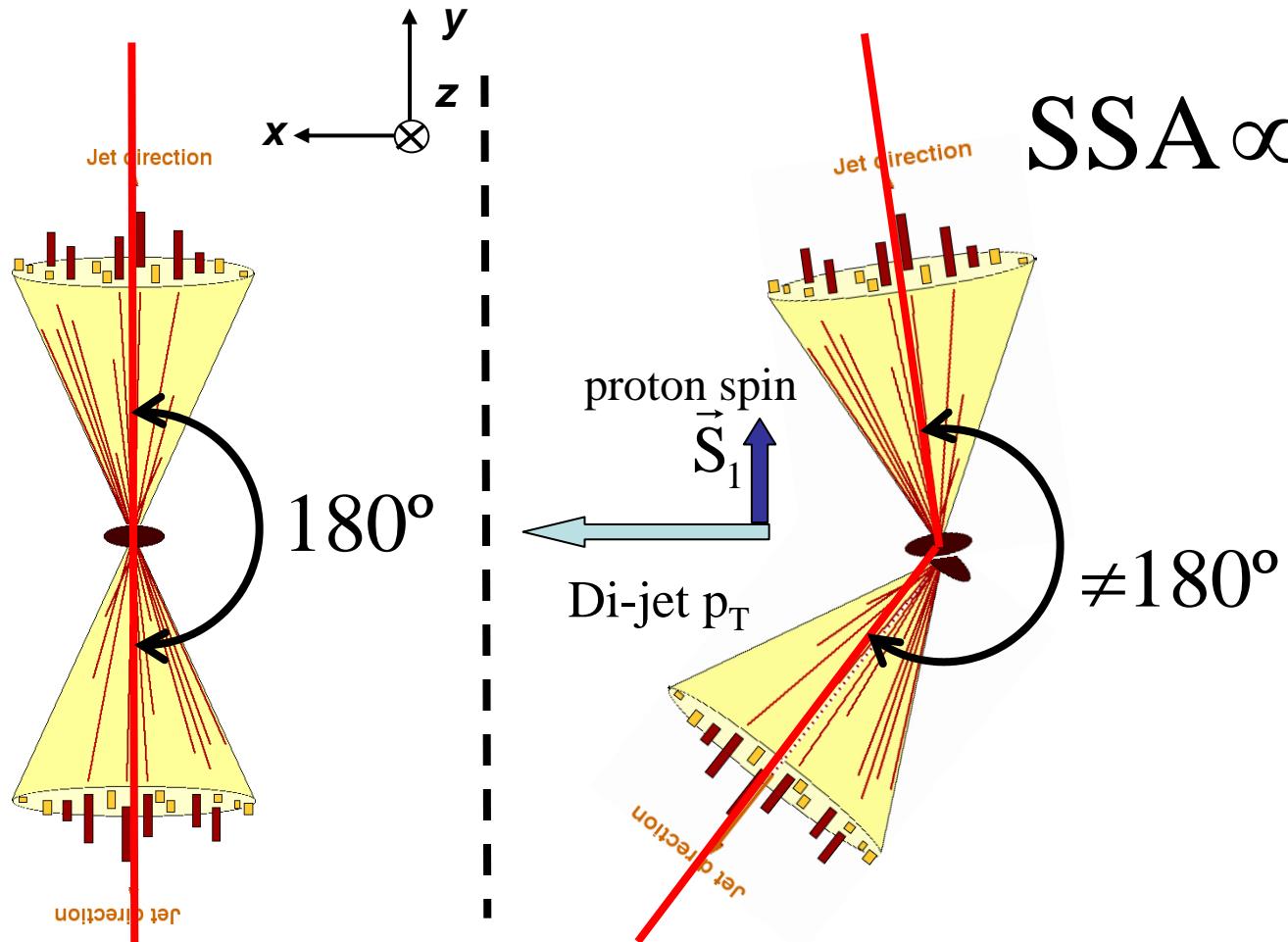
**Data do not show a simply monotonic decrease of SSA with increasing  $p_T$ .**

# SSA in di-jet production

Boer & Vogelsang, PRD 69, 094025 (2004)



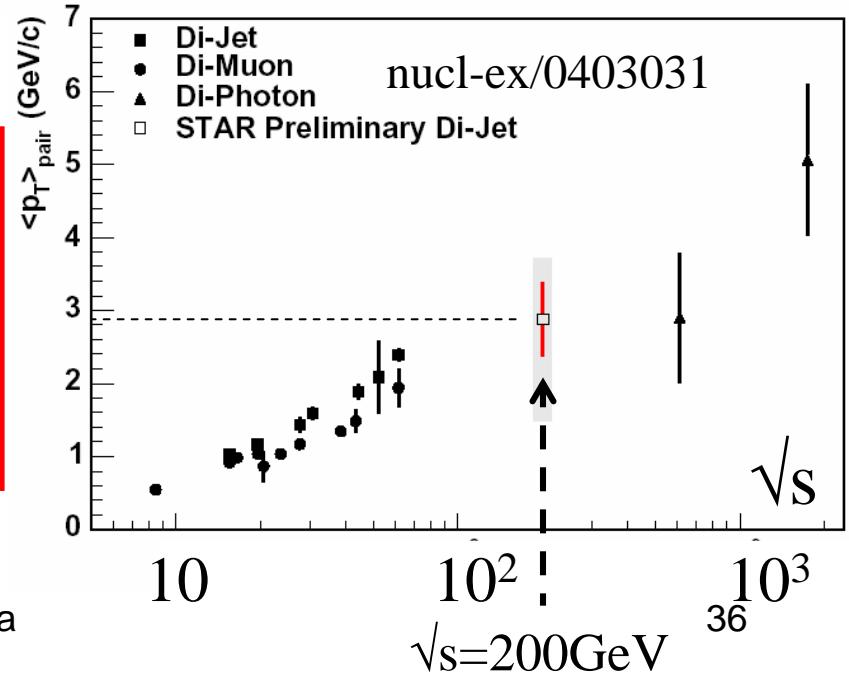
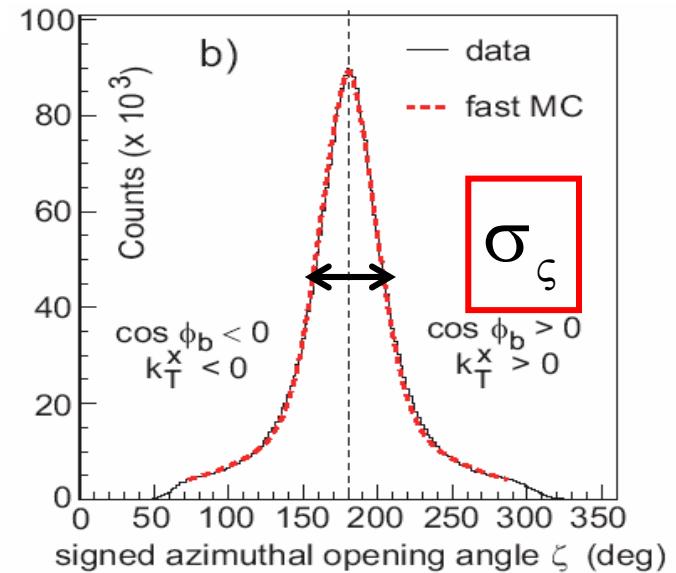
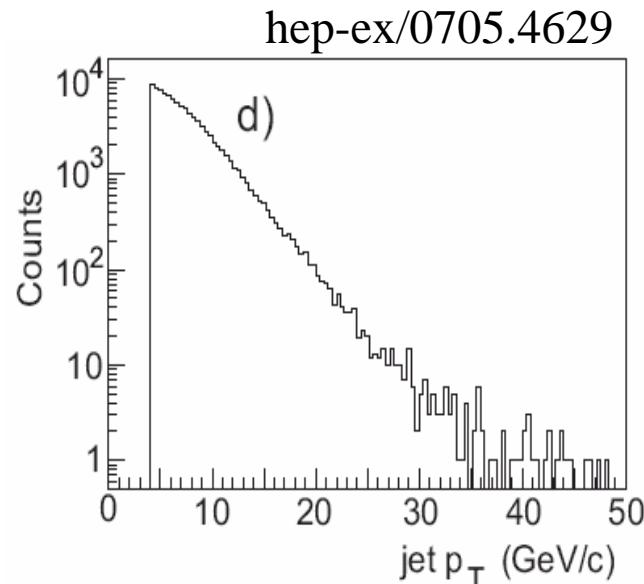
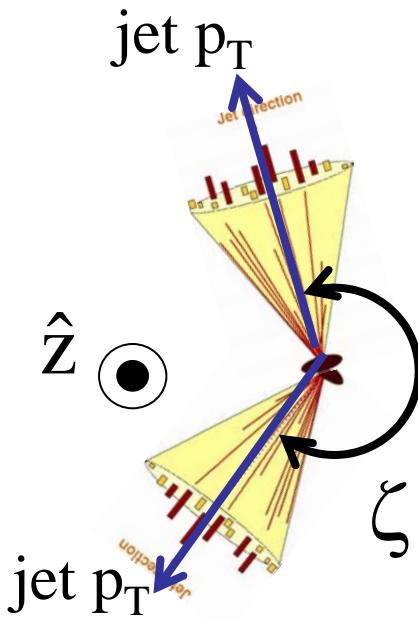
$$\text{SSA} \propto \vec{S}_1 \cdot (\vec{p}_1 \times \vec{k}^\perp)$$



Di-jet  $p_T$   
  
 $\vec{K}^\perp + \vec{k}^\perp$   
Gluon parton  
radiation

What is dominance of di-jet  $p_T$  ?

# Di-jet event measurements



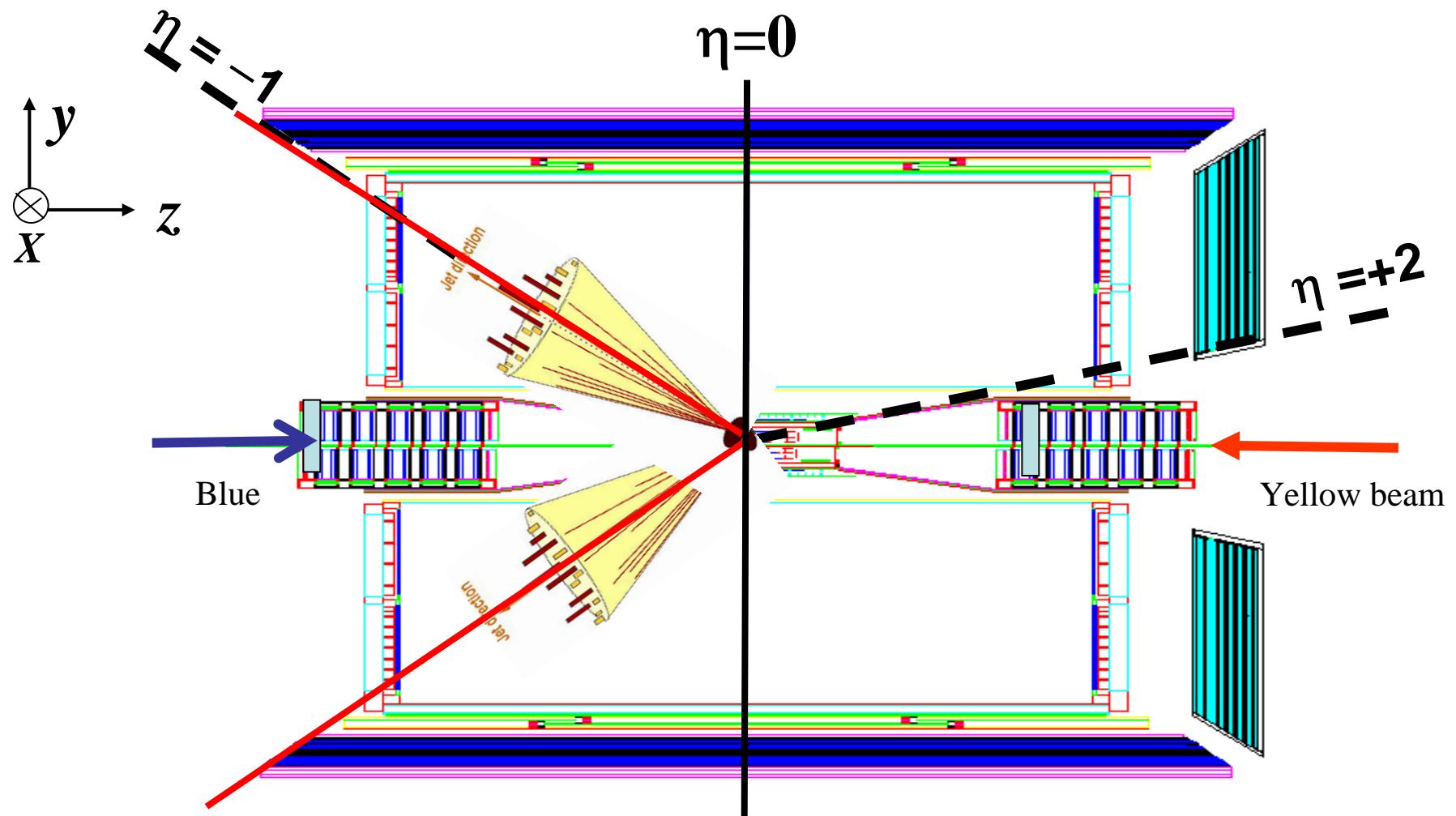
$$\langle \text{Di-jet } p_T \rangle = \frac{\pi}{2} \langle E_T \rangle \sin \sigma_\zeta \sim 2.8 \text{ GeV}/c$$

$$\text{parton } |\vec{k}^\perp| \sim \frac{\hbar c}{r_p} \sim 0.2 \text{ GeV}/c$$

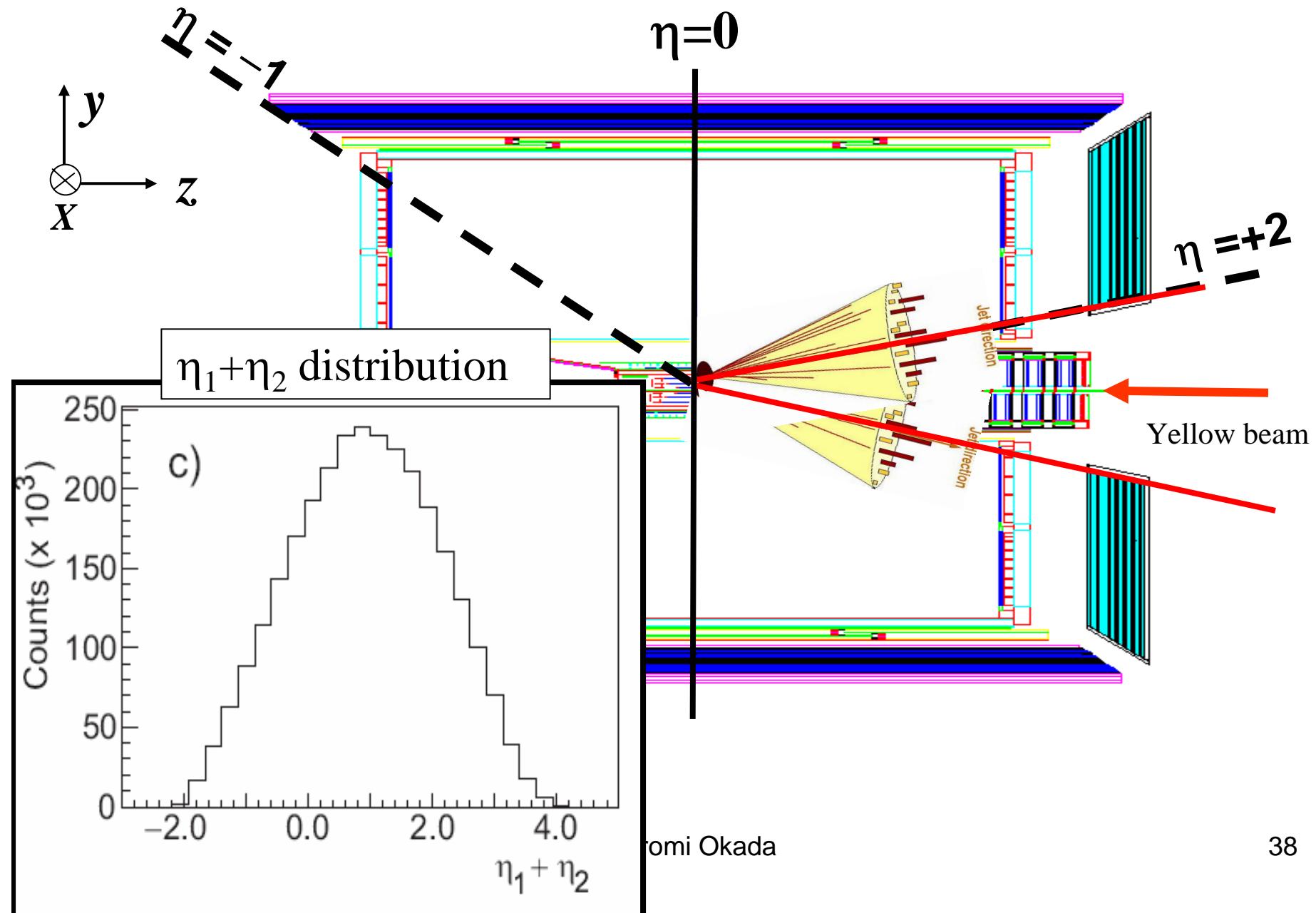
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# Di-jet event: $\eta_1 + \eta_2 = -2$



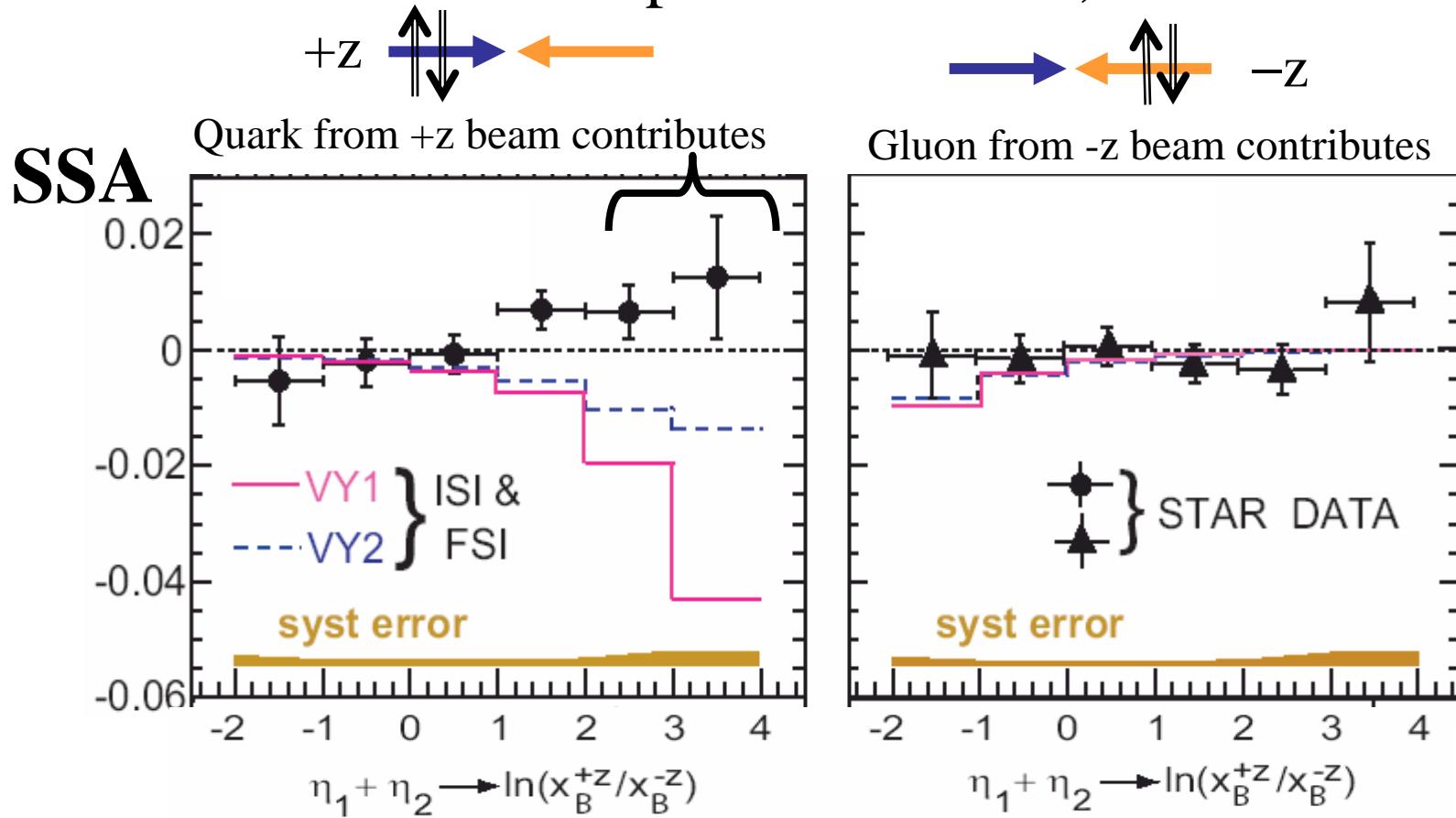
# Di-jet event: $\eta_1 + \eta_2 = +4$





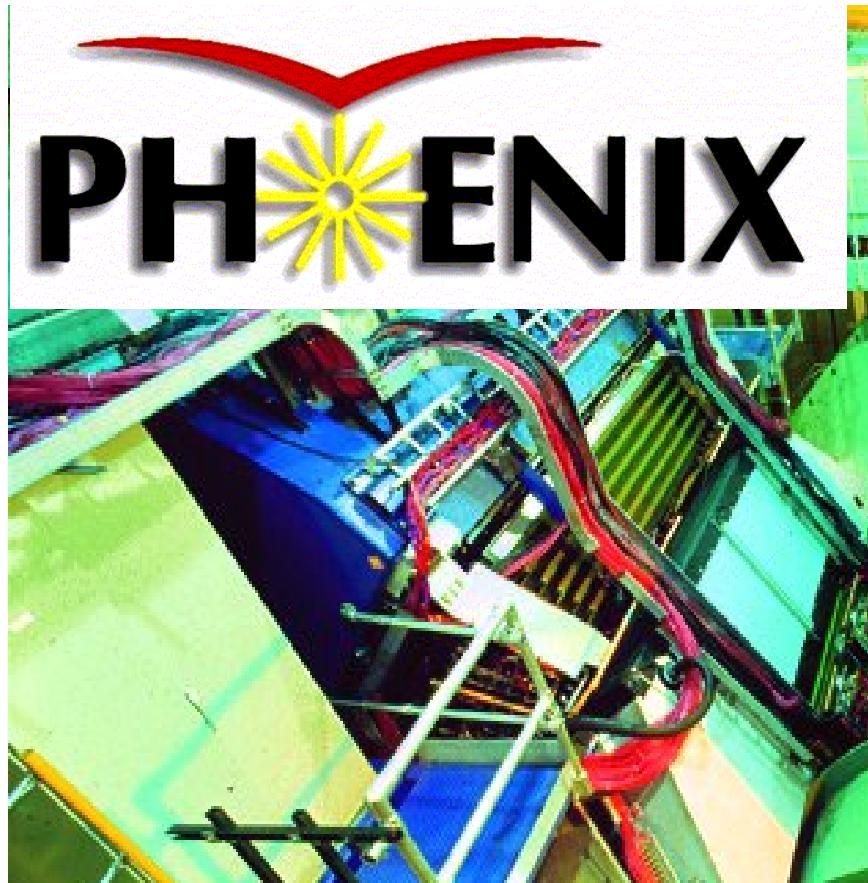
# STAR Results vs. Di-jet Pseudorapidity Sum

hep-ex/0705.4629, submitted to PRL.



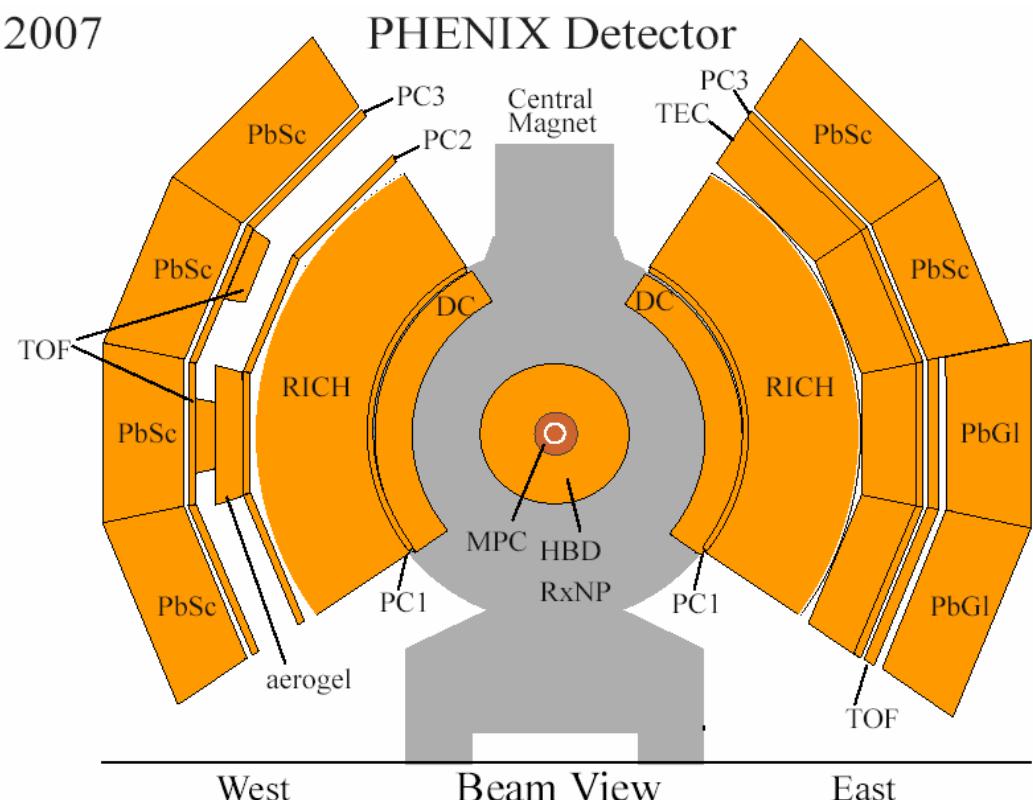
**Measured di-jet SSAs are consistent with zero.**

〔VY 1, VY 2 are calculations by Vogelsang & Yuan using  
HERMES-fitted quark Sivers function. PRD **72** (2005) 054028.〕

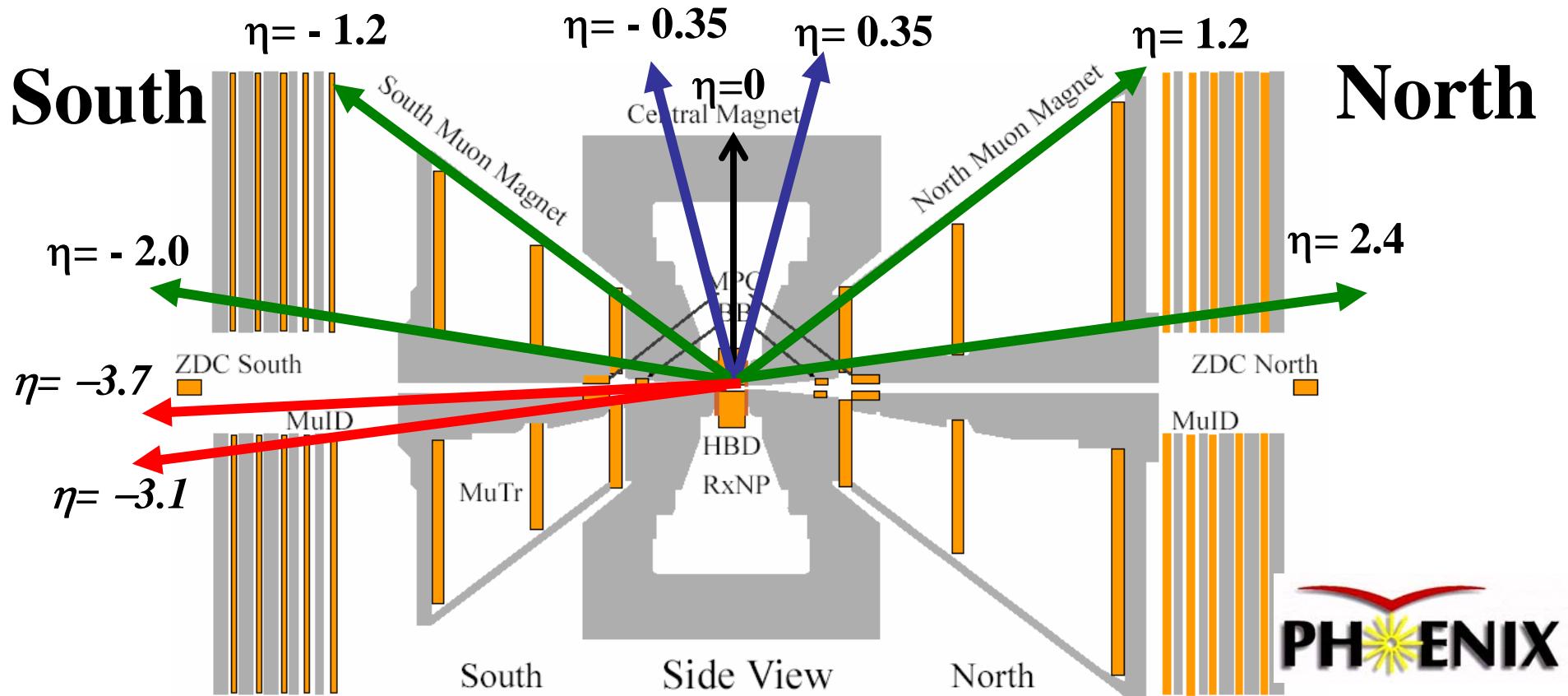


# Pioneering High Energy Nuclear Interaction eXperiment

2007



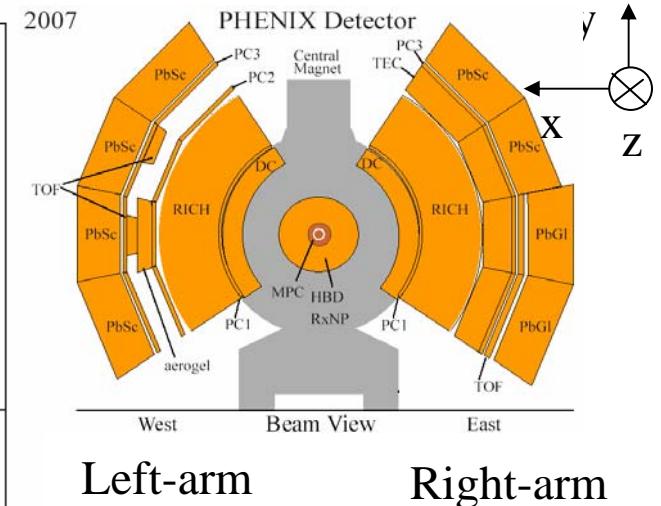
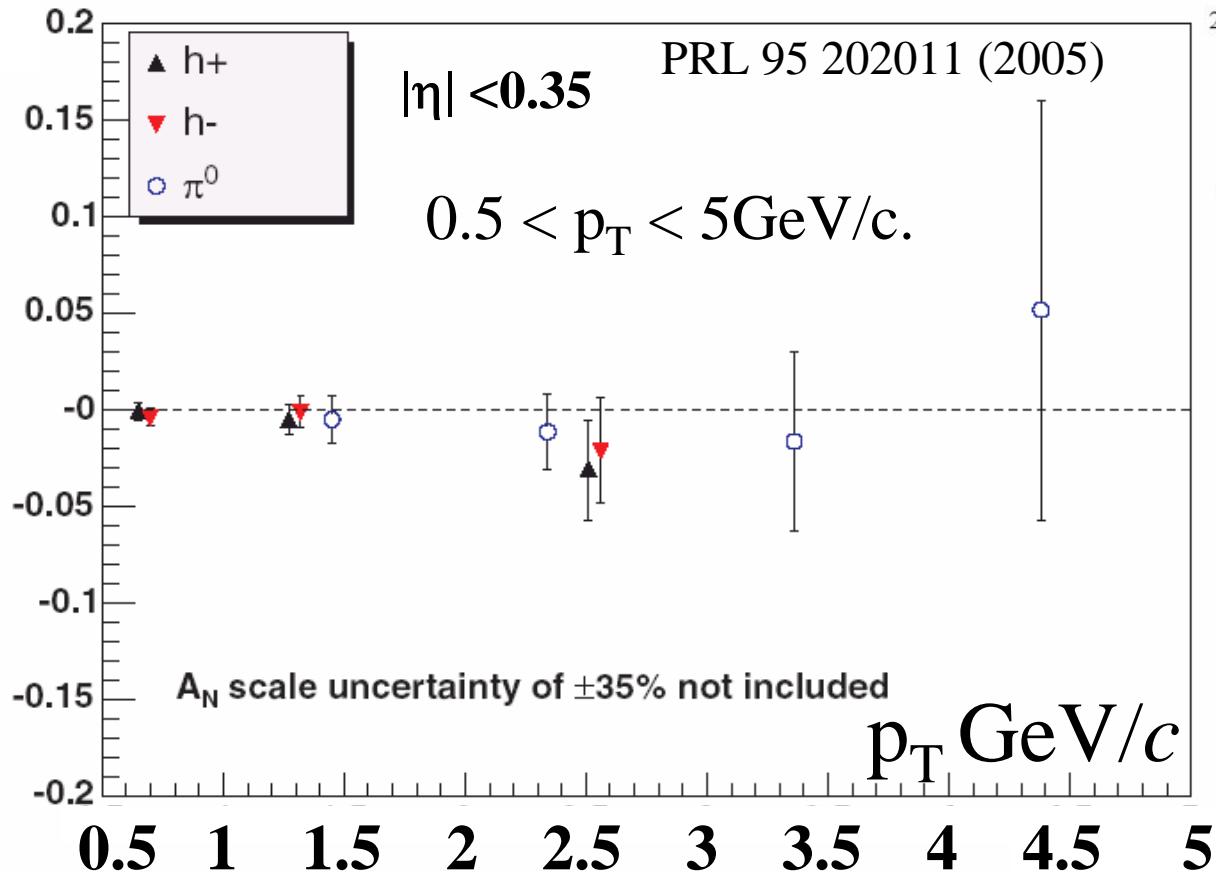
# Kinematical coverage at PHENIX (Side View)



- Central-Arms:  $\pi^0/h^\pm$  at  $\sqrt{s}=200\text{GeV}$ ,  $|\eta|<0.35$
- $\mu$ -Arms:  $J/\psi$  at  $\sqrt{s}=200\text{GeV}$ ,  $-2.0<\eta<-1.2$  &  $1.2<\eta<2.4$
- *MPC: Inclusive  $\pi^0$  at  $\sqrt{s}=62\text{GeV}$ ,  $-3.7<\eta<-3.1$*

# SSA of $\pi^0$ and $h^\pm$ from central-arms at $\sqrt{s}=200\text{GeV}$

SSA



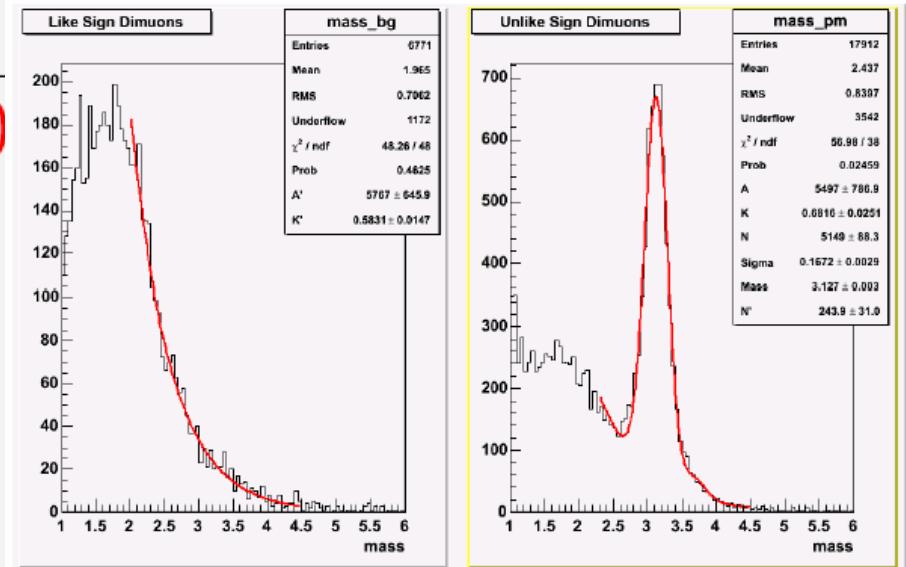
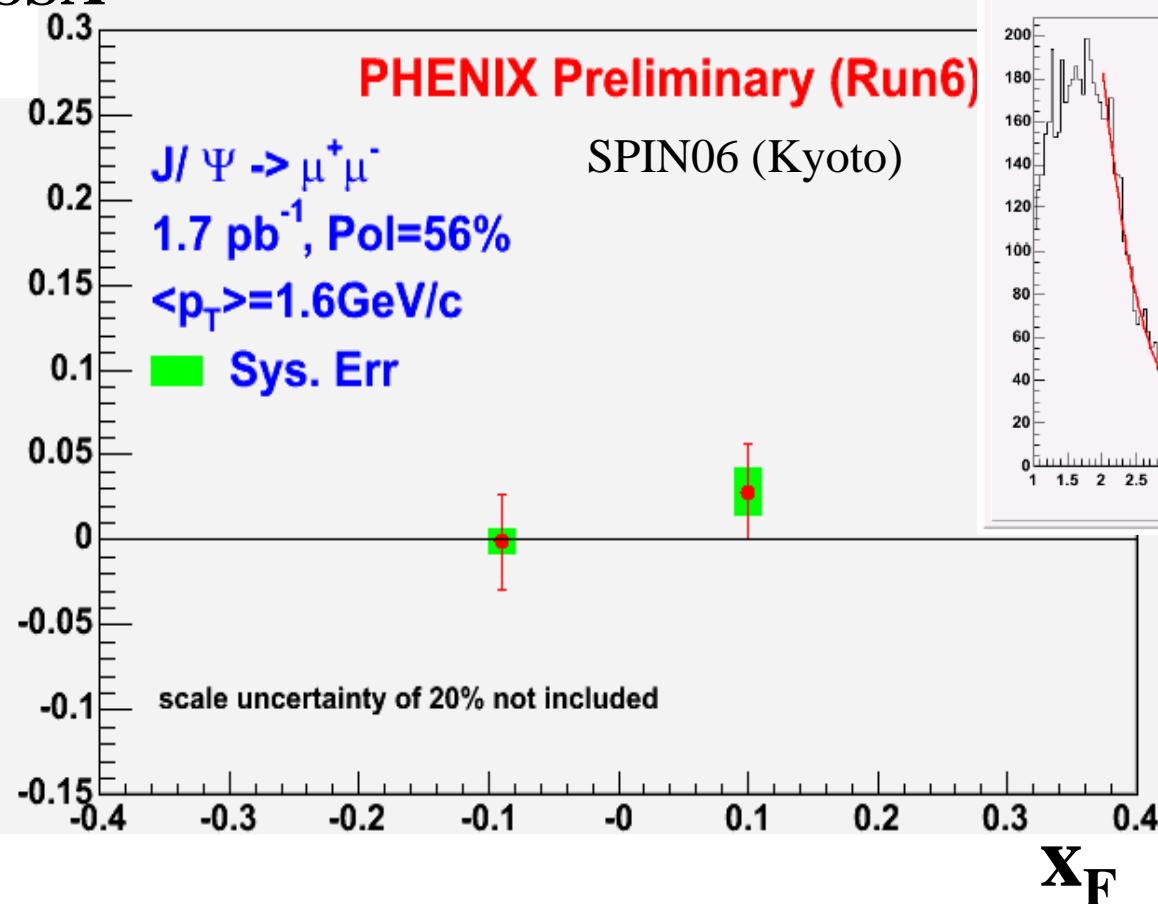
SSAs are measured by square-root-formula with two-arms (left-right) detector for either beam polarization.

- ✓ Mid-rapidity  $\langle x_F \rangle \sim 0 \rightarrow$  See “pure”  $p_T$  dependence of SSAs.
- ✓ **SSAs for mid-rapidity production of both  $\pi^0$  and  $h^\pm$  are consistent with zero.**

# SSA of J/ $\psi$ from $\mu$ -arms at $\sqrt{s}=200\text{GeV}$



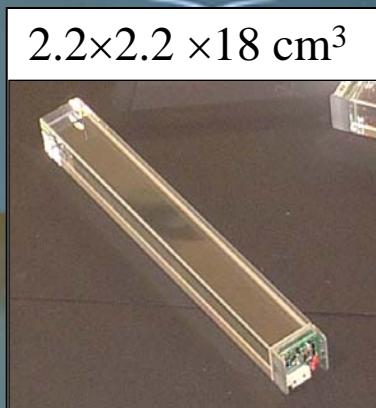
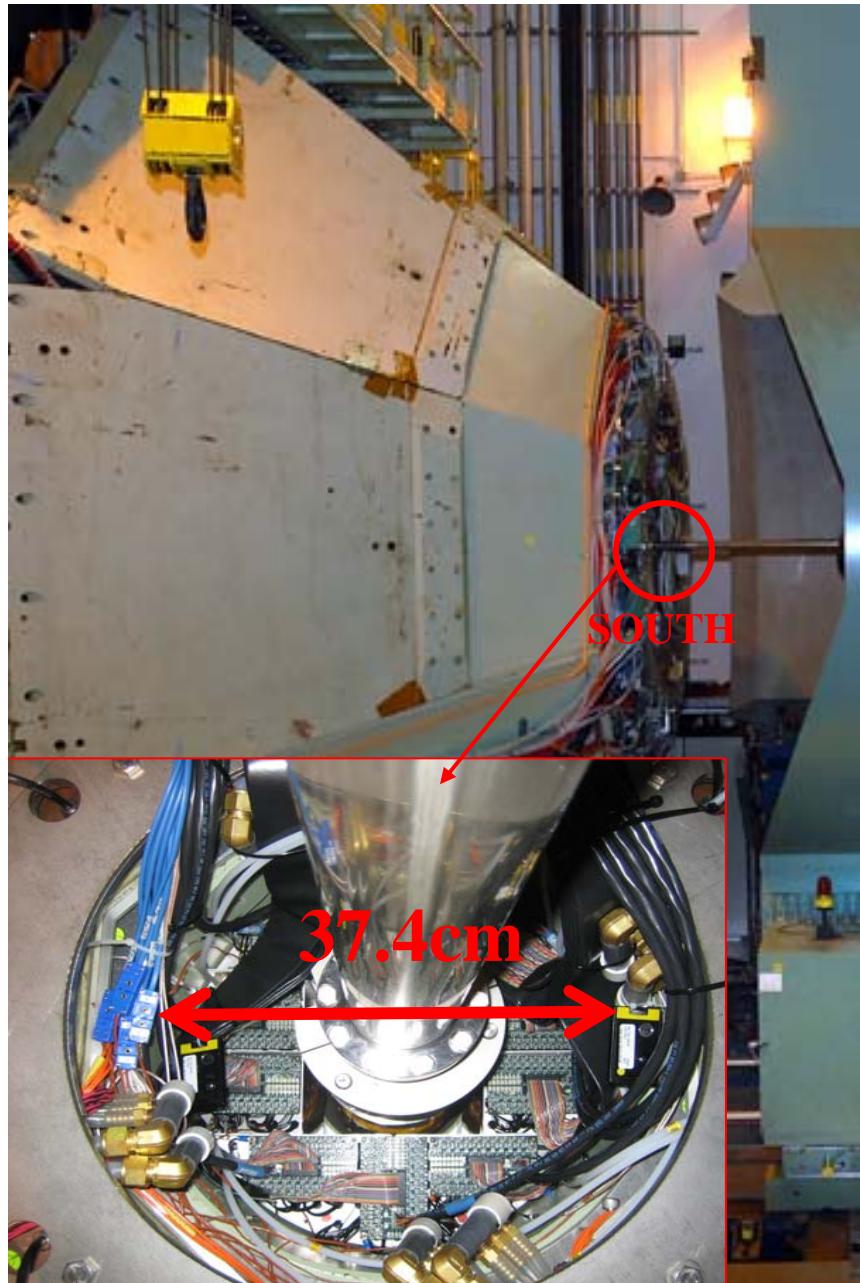
SSA



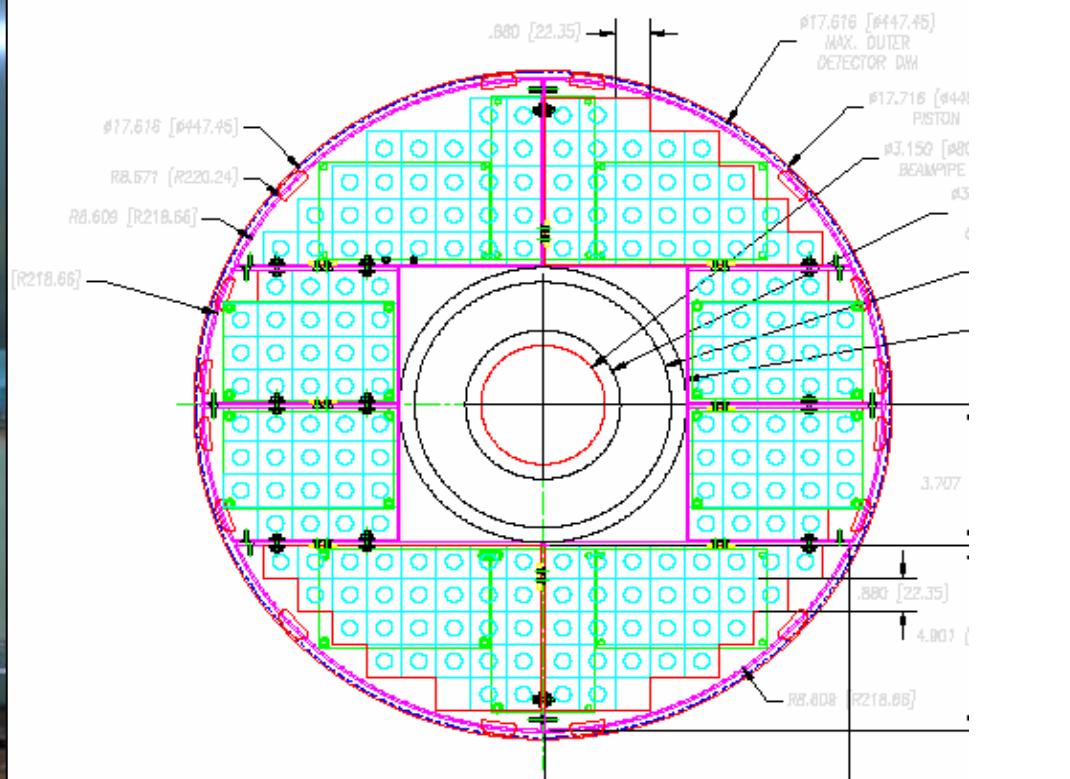
- SSAs are measured by “cross-ratio” method for either beam polarization.
- Sensitive to gluon Sivers as produced through g-g fusion

**SSAs at  $1.2 < |\eta| < 2.4$  of J/ $\psi$  production are consistent with zero.**

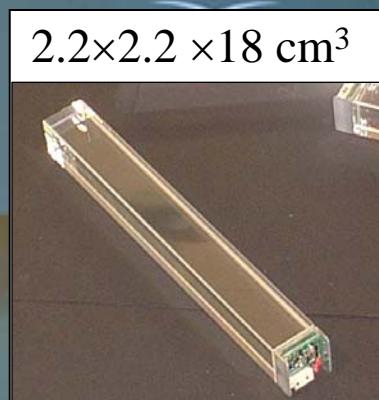
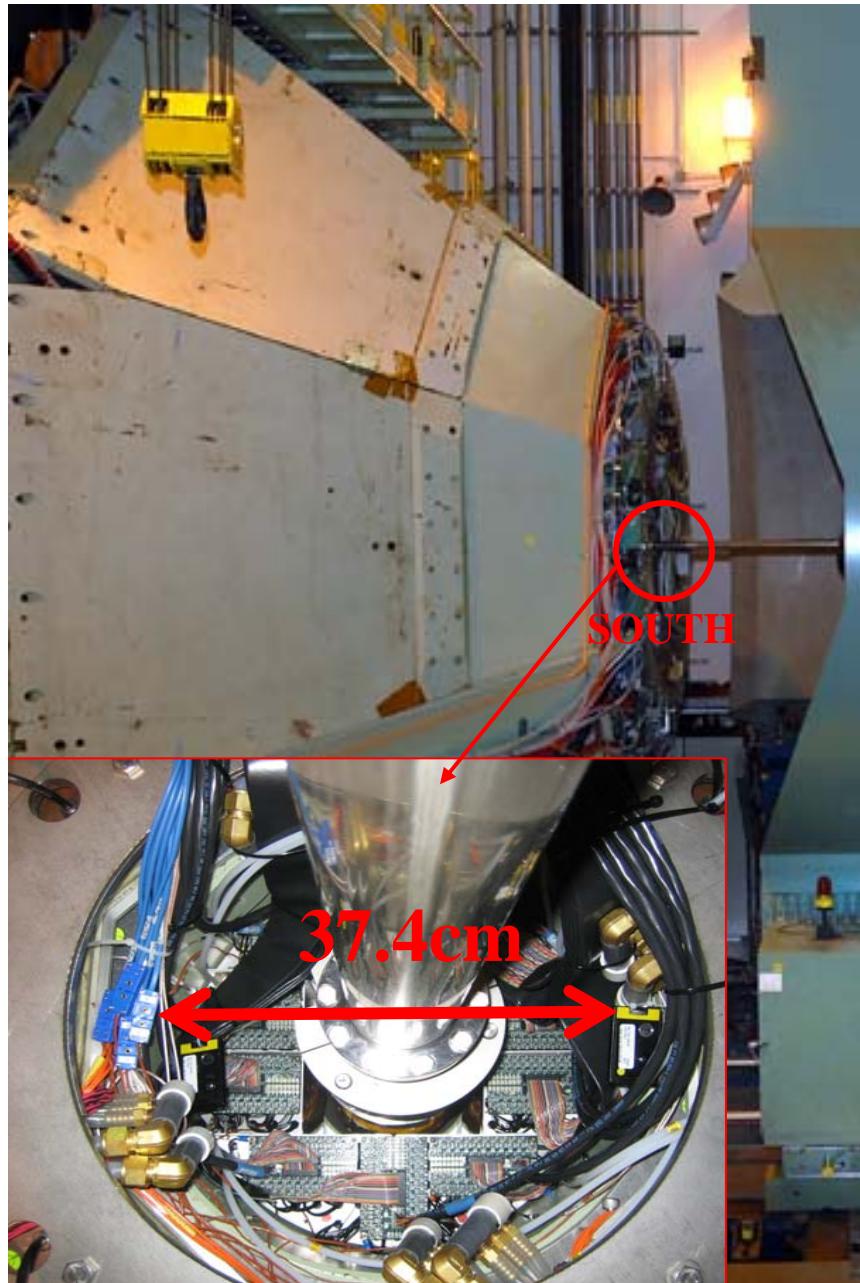
# PHENIX Muon Piston Calorimeter



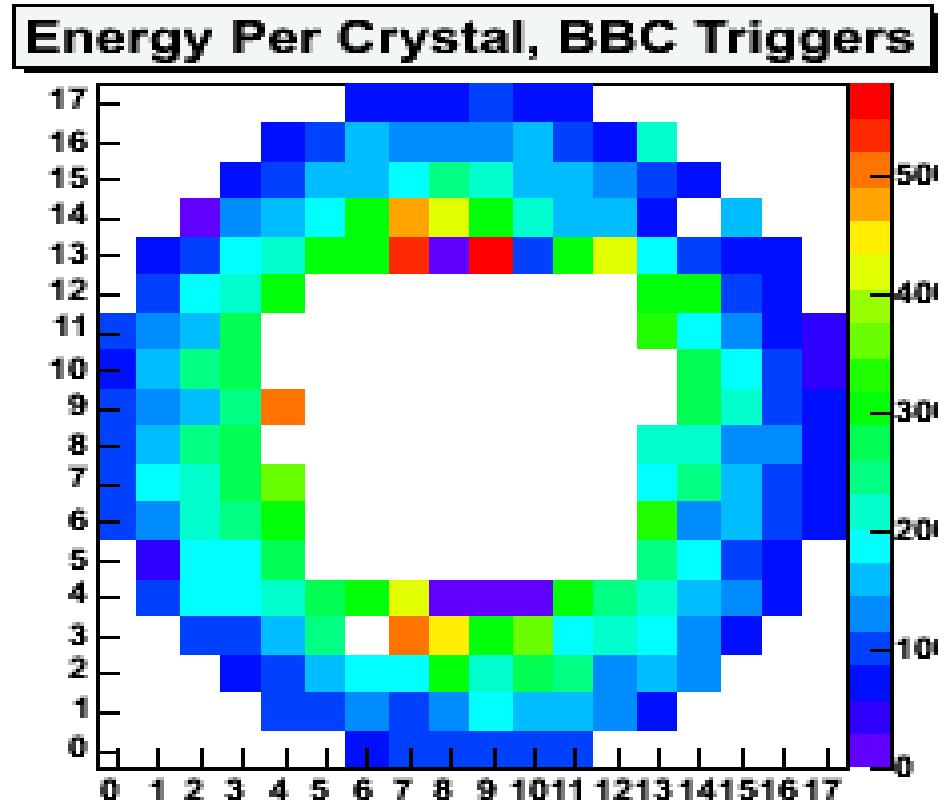
- 192 PbWO<sub>4</sub> crystals with APD readout
- Better than 80% of the acceptance is okay



# PHENIX Muon Piston Calorimeter



- 192 PbWO<sub>4</sub> crystals with APD readout
  - Better than 80% of the acceptance is okay

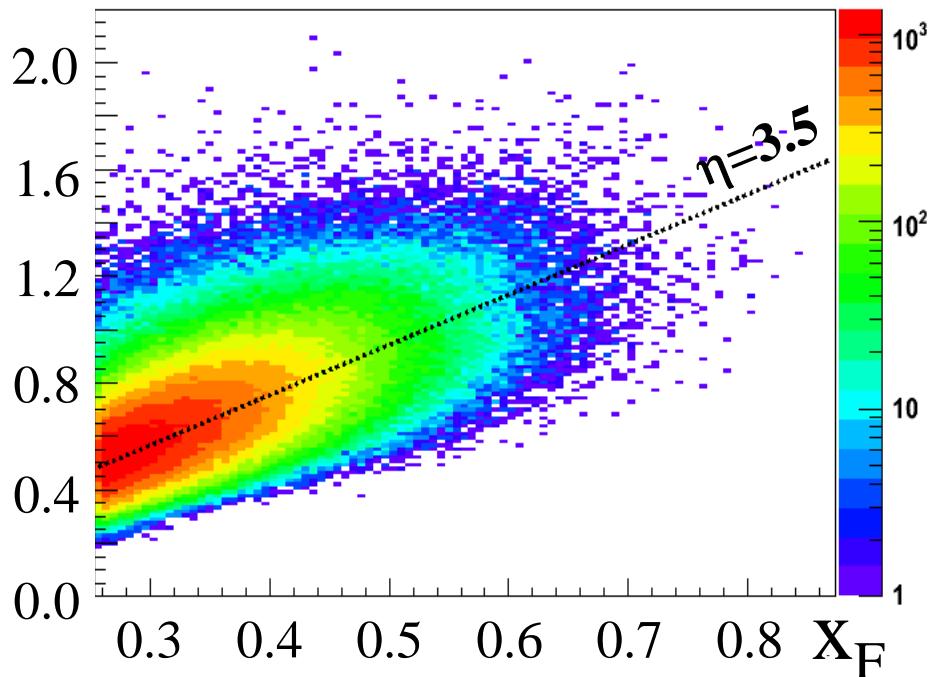


# SSA at $\sqrt{s}=62.4$ GeV from Muon Piston Calorimeter (MPC)

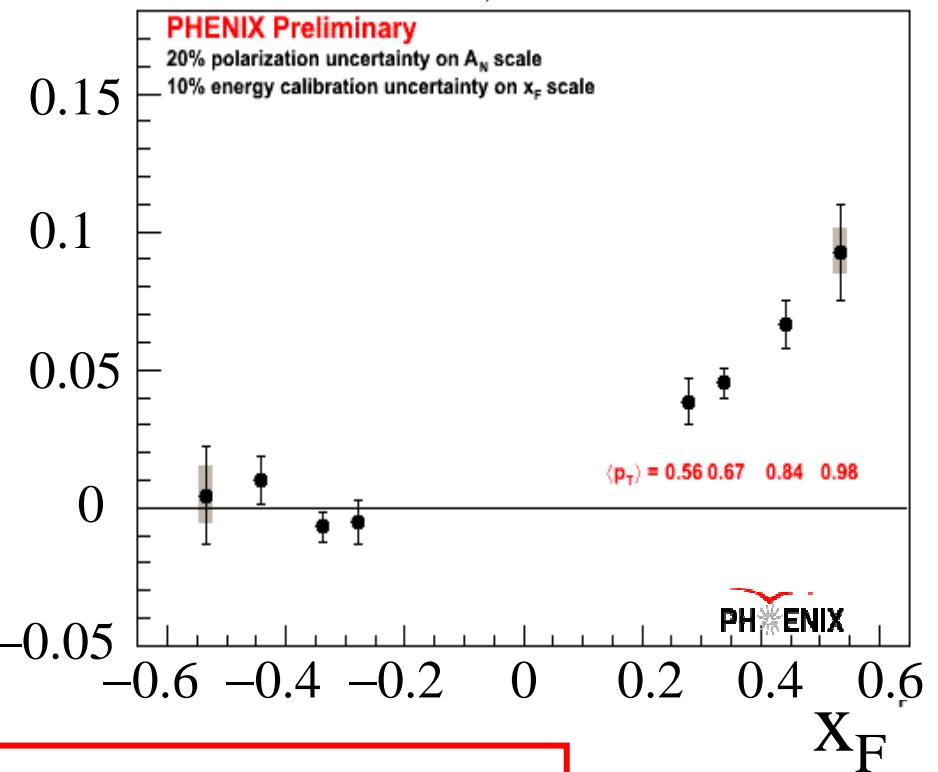
- $\pi^0$  detection in forward direction;  $3.1 < |\eta| < 3.7$
- Achieve higher  $x_F$  by decreasing  $\sqrt{s}$ .

$$x_F = \frac{p_{z,\pi}}{p_{z,1}} \approx \frac{2E_\pi}{\sqrt{s}}$$

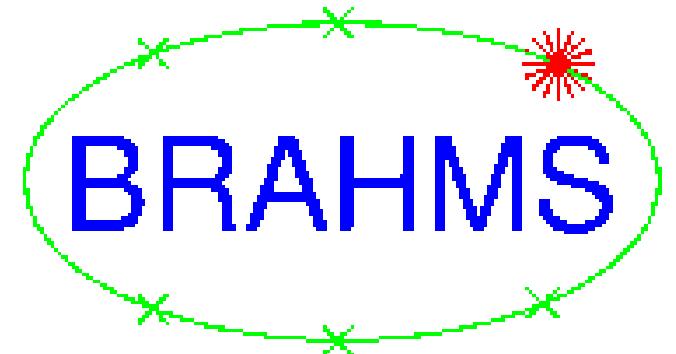
$p_T$  GeV/c



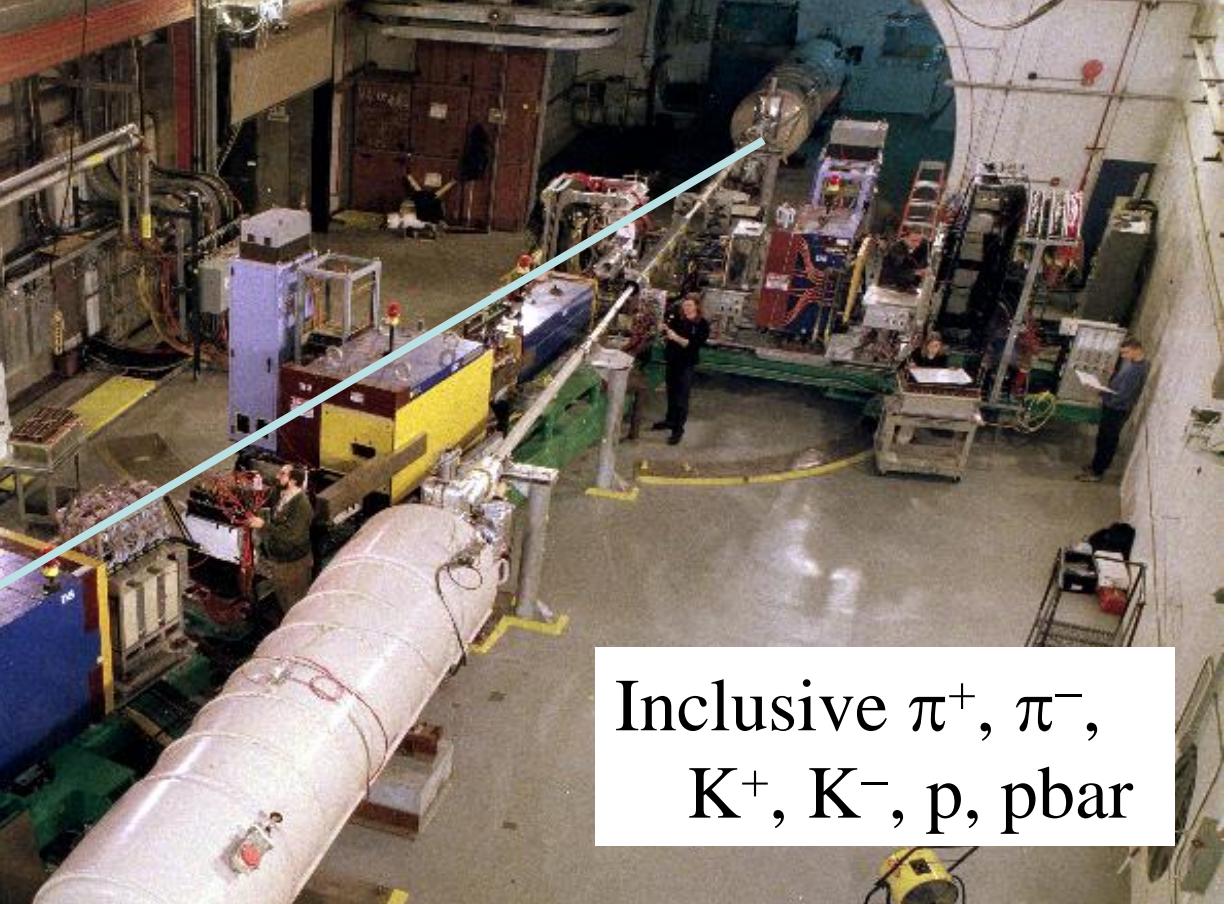
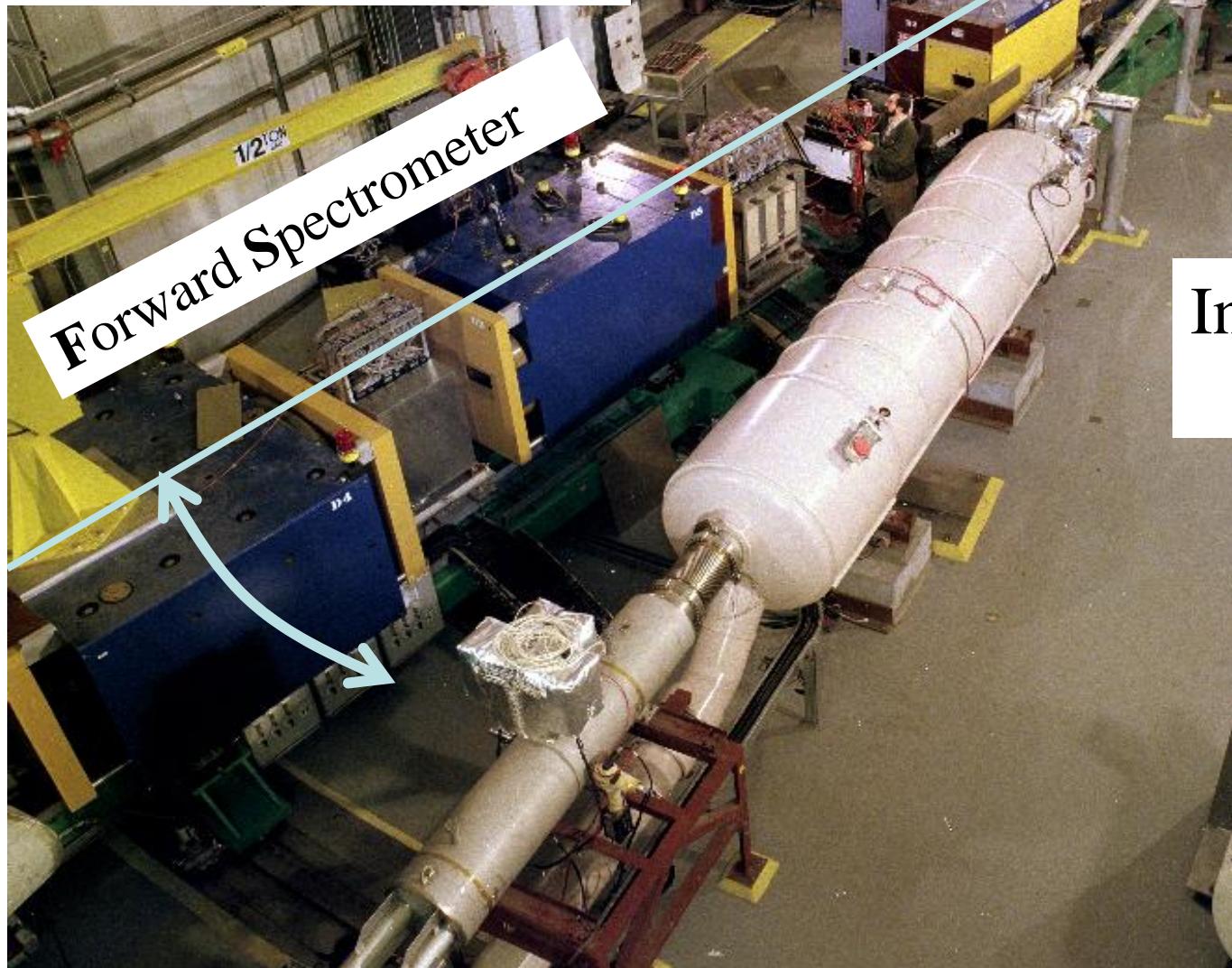
SSA



SSA survives and has similar characteristics over a broad  $\sqrt{s}$  region 20~ 200 GeV.



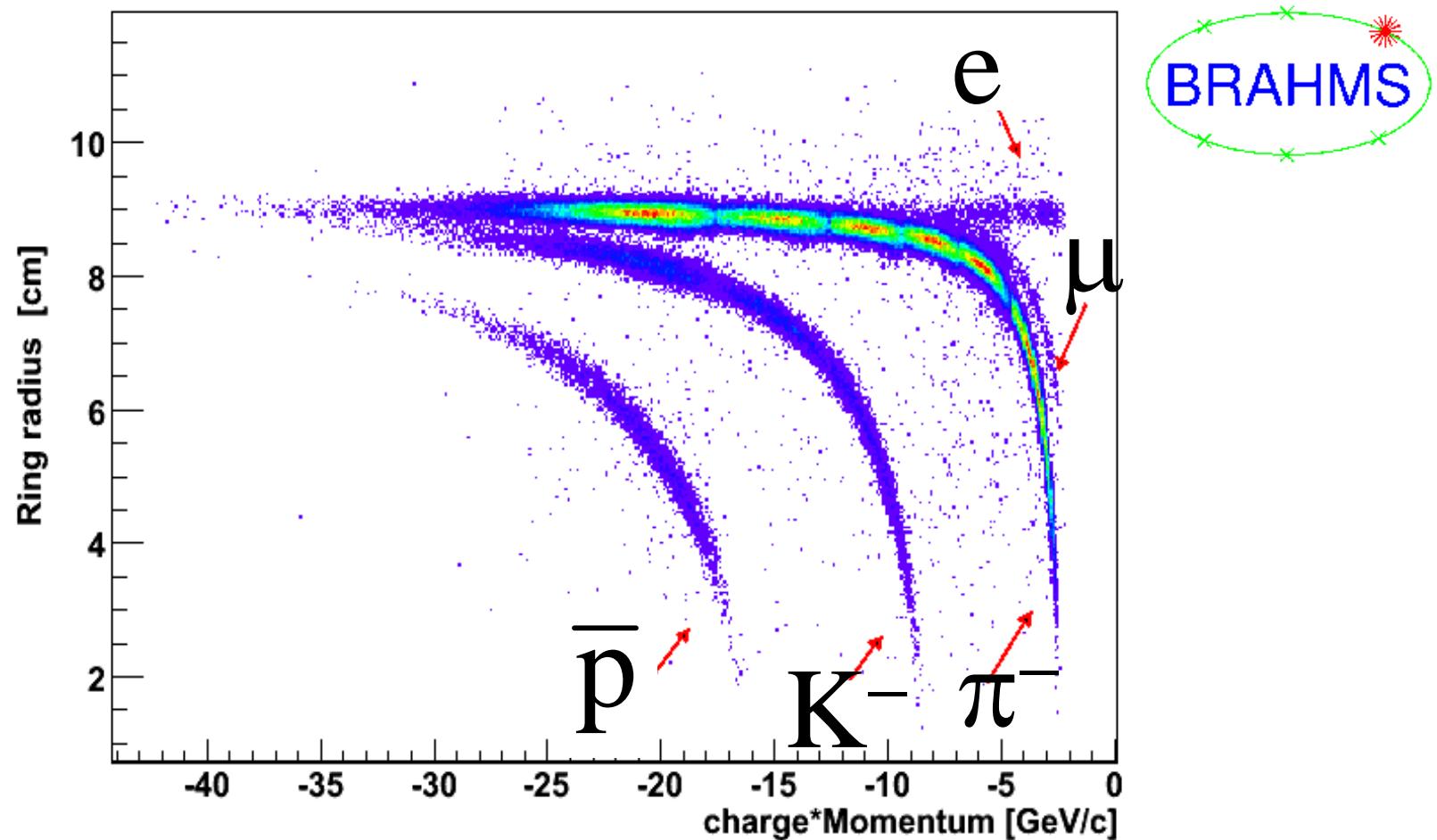
**BRAHMS**



Inclusive  $\pi^+$ ,  $\pi^-$ ,  
 $K^+$ ,  $K^-$ , p, pbar

**B**road  
**R**ange  
**A**hadron  
**M**agnetic  
**S**pectrometers

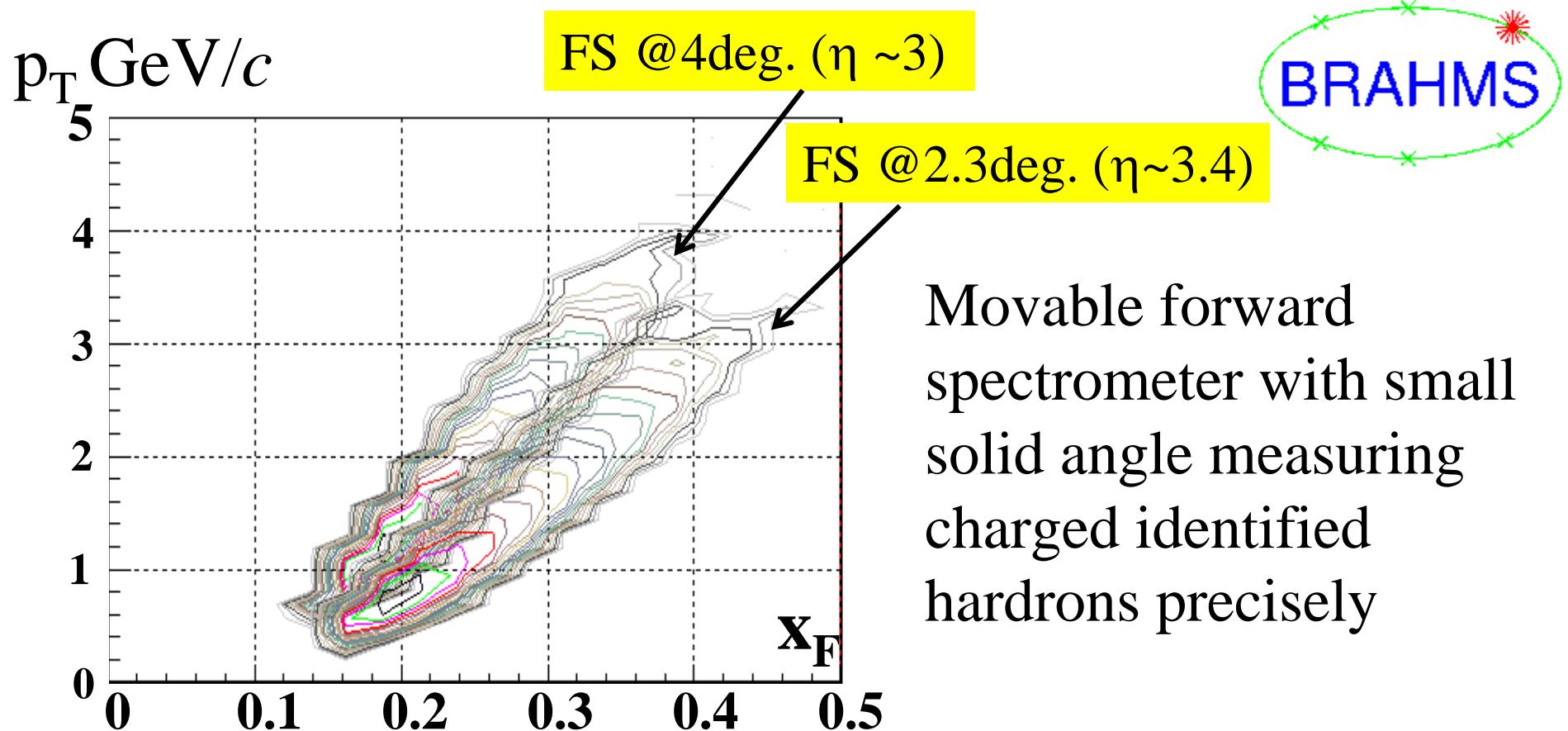
# Particle Identification using RICH



$p, K$  identification  $< 30 \text{ GeV}/c$ ,  $\bar{p} > 17 \text{ GeV}/c$  with efficiency  $\sim 97\%$

# BRAHMS FS Acceptance at 2.3 deg. and 4 deg.

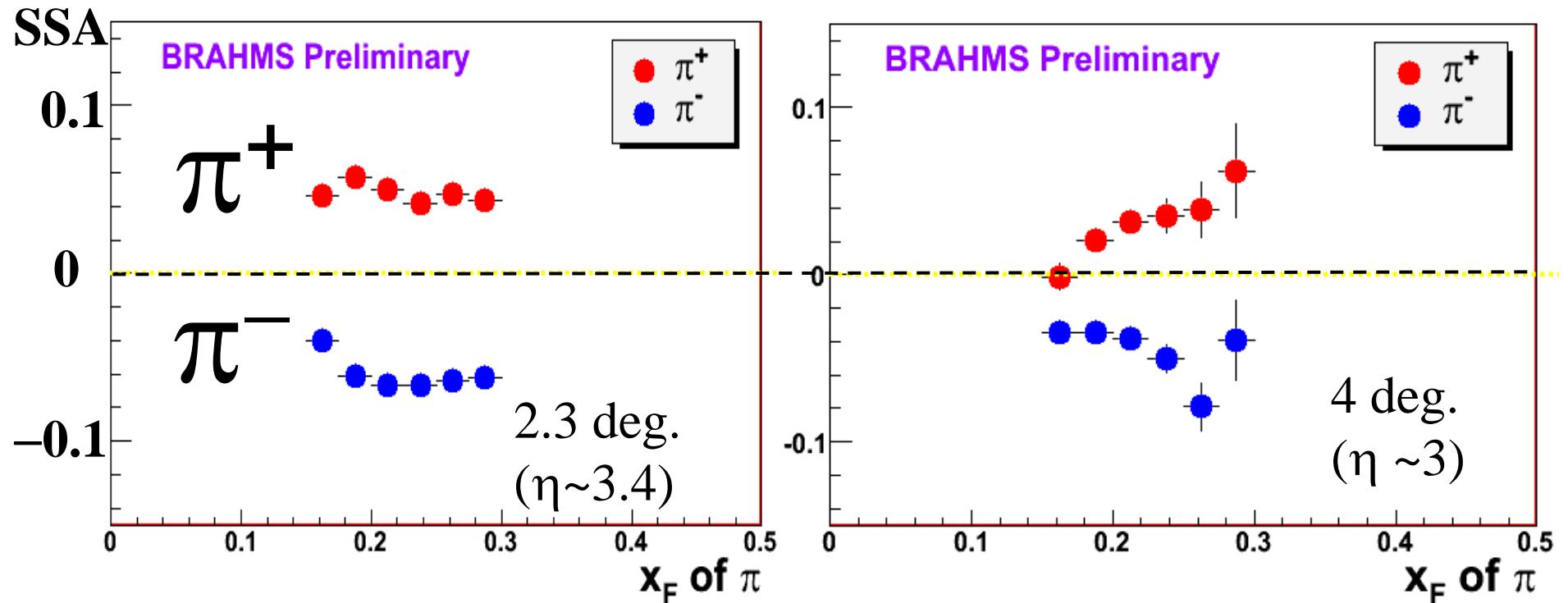
Full Field (7.2 Tm) at  $\sqrt{s} = 200$  GeV



Movable forward spectrometer with small solid angle measuring charged identified hadrons precisely

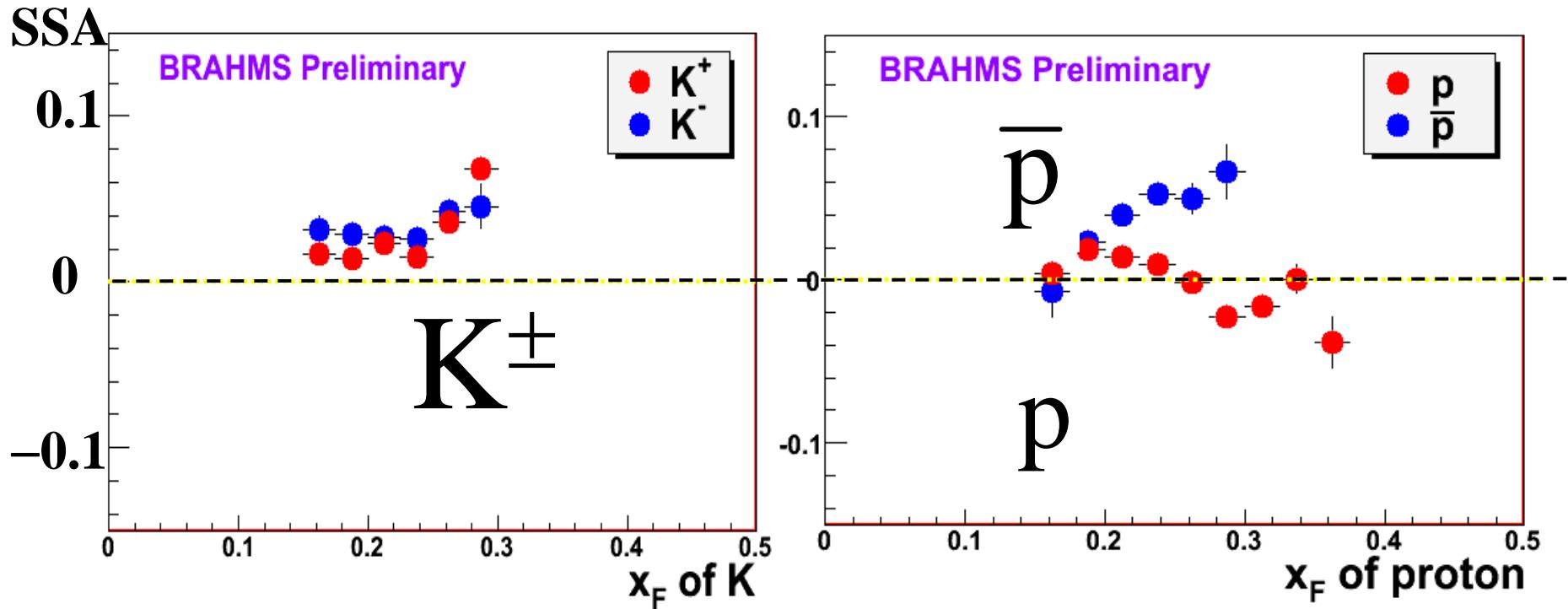
- Strong  $x_F$ - $p_T$  correlation due to limited spectrometer solid angle acceptance

# $\pi^\pm$ SSAs at 2.3 and 4 deg. at $\sqrt{s} = 200$ GeV



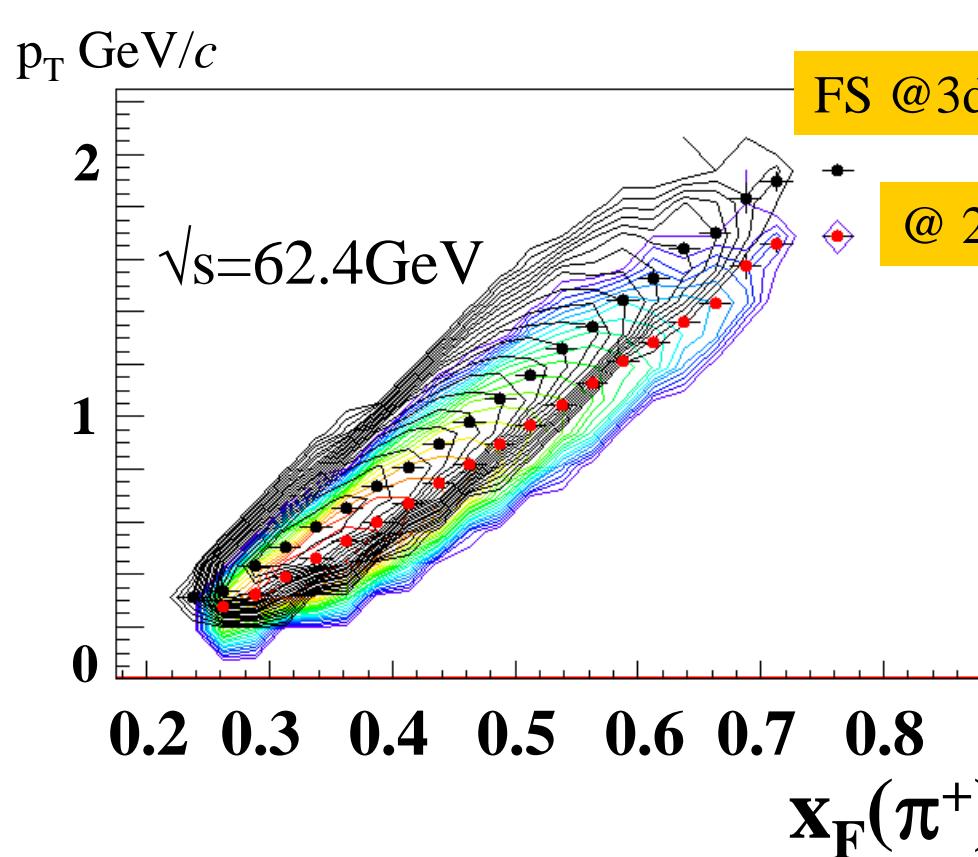
- SSA( $\pi^+$ ): positive  
SSA( $\pi^-$ ): negative
- SSA ( $\pi^\pm$ ) survive.  
 $4\text{-}6\%$  in  $0.15 < x_F < 0.3$ .

# SSAs at 2.3 deg. at $\sqrt{s} = 200$ GeV

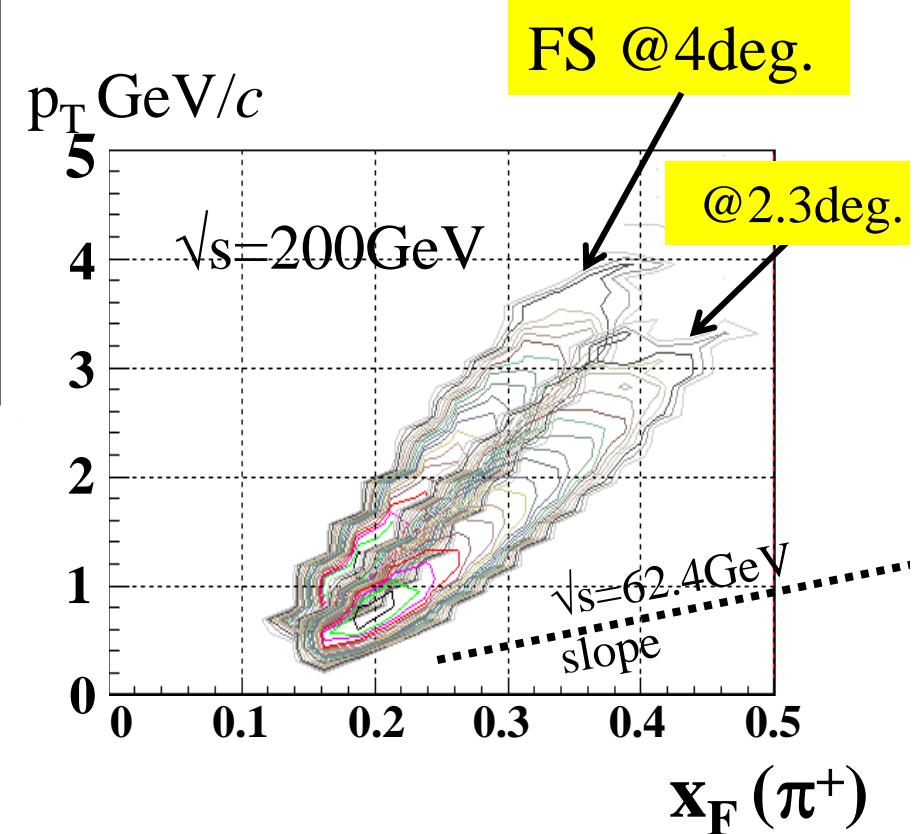


- SSA( $K^+$ ), SSA( $K^-$ ): positive 2-5% for  $0.15 < x_F < 0.3$ .
- SSA( $\bar{p}$ ), SSA( $K^-$ )  $> 0$ : Contribution from sea-quarks.  
Or Accidental?
- SSA( $p$ )  $\sim 0$ : Significant fraction of proton can be mostly from polarized beam proton, but only ones showing SSA  $\sim 0$ .

# SSA results at $\sqrt{s}=62.4$ GeV at higher $x_F$



Half Field (3.6 Tm)



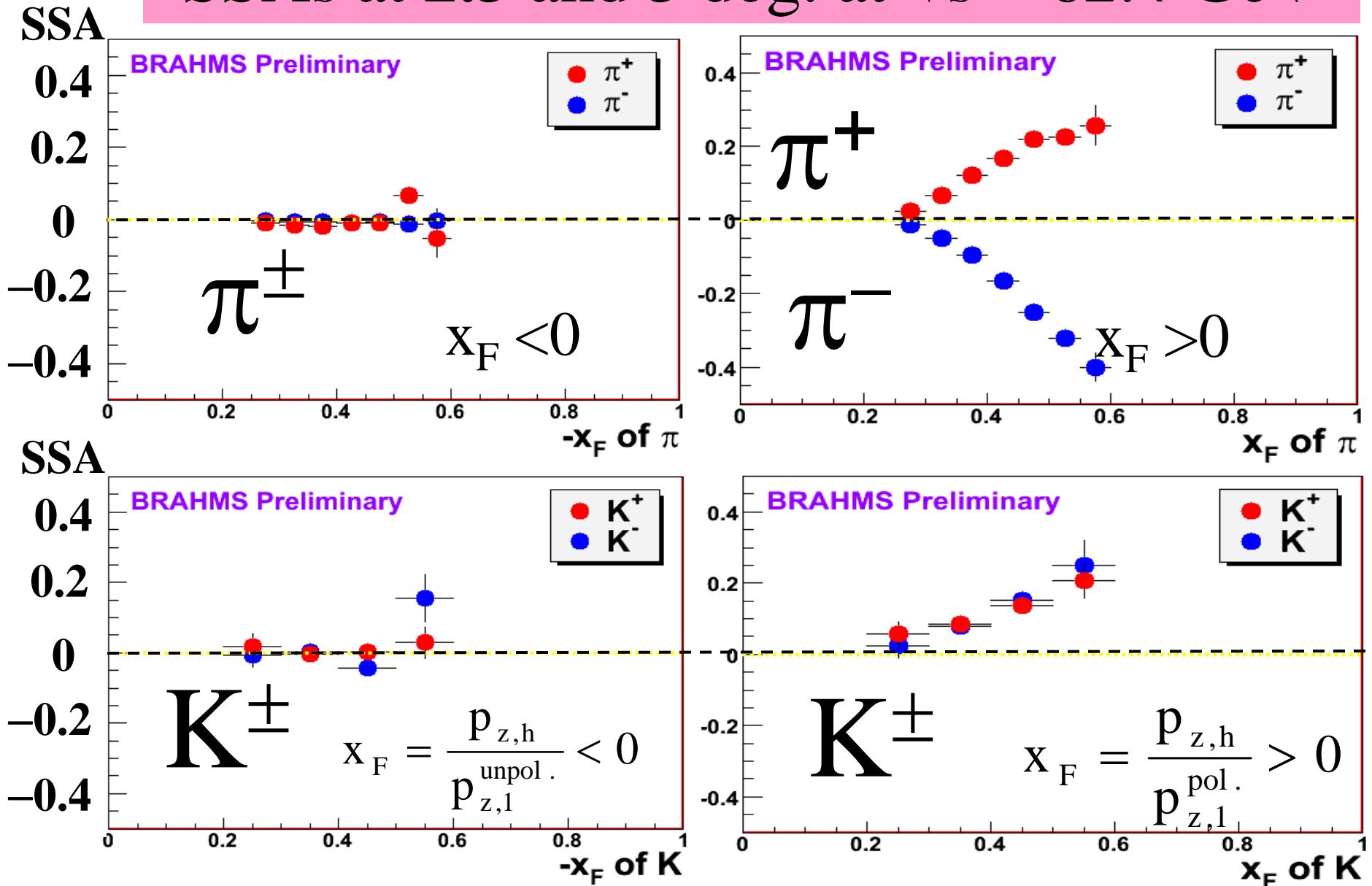
$$x_F = \frac{p_{z,\pi}}{p_{z,1}} \approx \frac{2E_\pi}{\sqrt{s}}$$

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# SSAs at 2.3 and 3 deg. at $\sqrt{s} = 62.4$ GeV



## Summary from the RHIC results

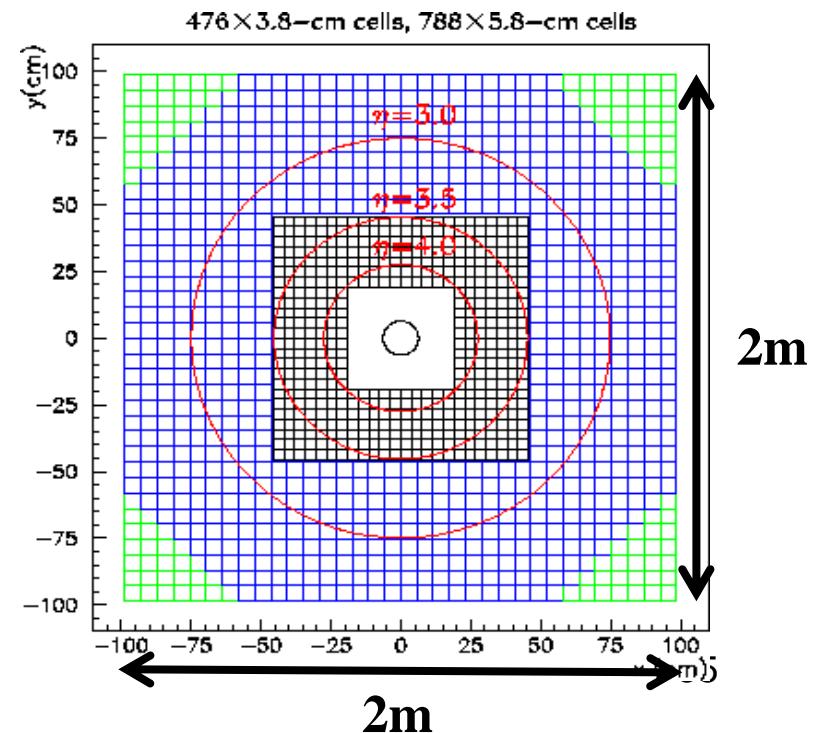
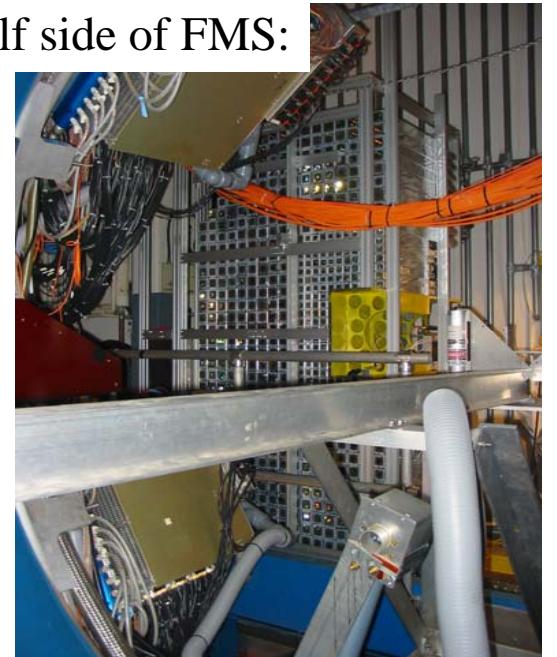
1. Study of SSA in polarized  $pp$  collisions has begun at RHIC, where cross sections are consistent with pQCD.
2. SSAs in the mid-rapidity region are consistent with zero.
3. Sizable SSAs are measured in the forward region only.
4. First observation of  $p_T$  dependence of SSA is performed in pQCD region.
  - An increase of SSA with increasing  $p_T$  is observed.
  - This is contrary to the theoretical expectation.
5. Interaction between experiment and theory is ongoing and critical.

**Forward physics is very exciting !**

# Outlook

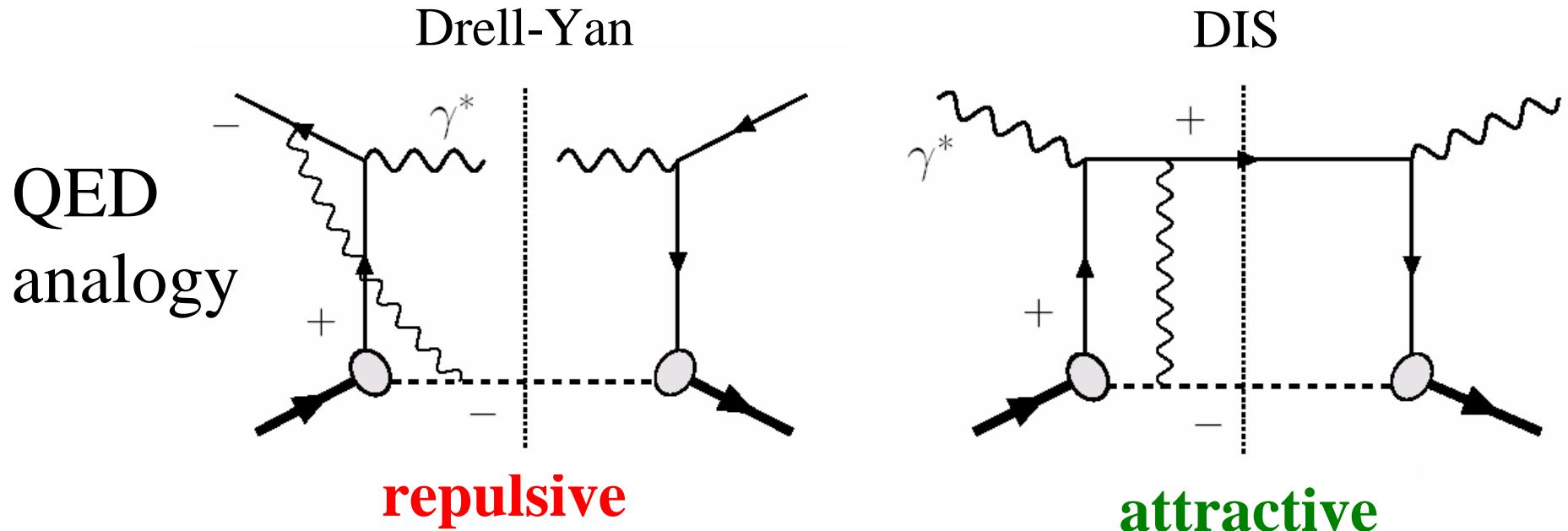
- A) Continue study of  $p_T$  dependence of SSA using **Forward Meson Spectrometer** at the STAR in RUN8. → more than factor 25 increase of acceptance for triggered events !
- B) Sivers/Collins separation by choosing proper process.
- Orbital motion via Sivers mechanism.
  - Transversity via Collins mechanism.
- C) Confirm “**color charge**” via SSA-Sivers effect in Drell-Yan process.  
(Needs tracking)

Half side of FMS:



### C) Confirm color charge via SSA-Sivers effect.

[http://spin.riken.bnl.gov/rsc/write-up/dy\\_final.pdf](http://spin.riken.bnl.gov/rsc/write-up/dy_final.pdf)



$$f_{1T}^{\perp}(x, \vec{k}_q^{\perp})_{\text{DY}} = -f_{1T}^{\perp}(x, \vec{k}_q^{\perp})_{\text{DIS}}$$

The non-universality of the Sivers functions expects opposite sign between DY and DIS processes.

**Forward physics is very very exciting !**

Thank you!!