

Spectrum Study of Nf=2 QCD with Overlap Fermions

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New JLQCD's project: Dynamical simulation with overlap fermions

- Powerful approach to low energy QCD
 - Exact chiral symmetry
→ No contamination in data
 - Any #flavor purchasable
- State of the art
 - Extremely light quarks
 - High statistics like never before (57Tflops!)
 - New techniques in gauge generation/measurement
→ Significant leap of LQCD

Neuberger, 1998

$$D_{\text{ov}} = \frac{1}{R} [1 + \gamma_5 \epsilon(H_W)]$$

Hashimoto's talk; Kaneko's talk

- Current status:

$$N_f=2, 16^3 \times 32, S_{\text{ov}} + S_{\text{exWilson}} + S_{\text{Iwasaki gauge}}$$

- ~10,000trajs for each
- 6 sea X 9 valence: $280 \text{ MeV} < m_{\text{PS}} < 740 \text{ MeV}$
- $a^{-1} \approx 1.63 \text{ GeV}$ ($r_0 = 0.50 \text{ fm}$)
- Fixed topological charge: $Q=0$

- This talk-

- Meson spectrum
- NPR for bilinear
- Kaon B-parameter

$$m_\pi, f_\pi \\ Z_A, Z_S = Z_m^{-1}$$

Yamada's talk



Chiral log



Strange mass

Improvements with low eigenmodes

DeGrand and Schaefer, 2004; Giusti et al., 2004

- Lowest 50 eigenmodes

$$D_{ov} u_i = \lambda_i u_i$$

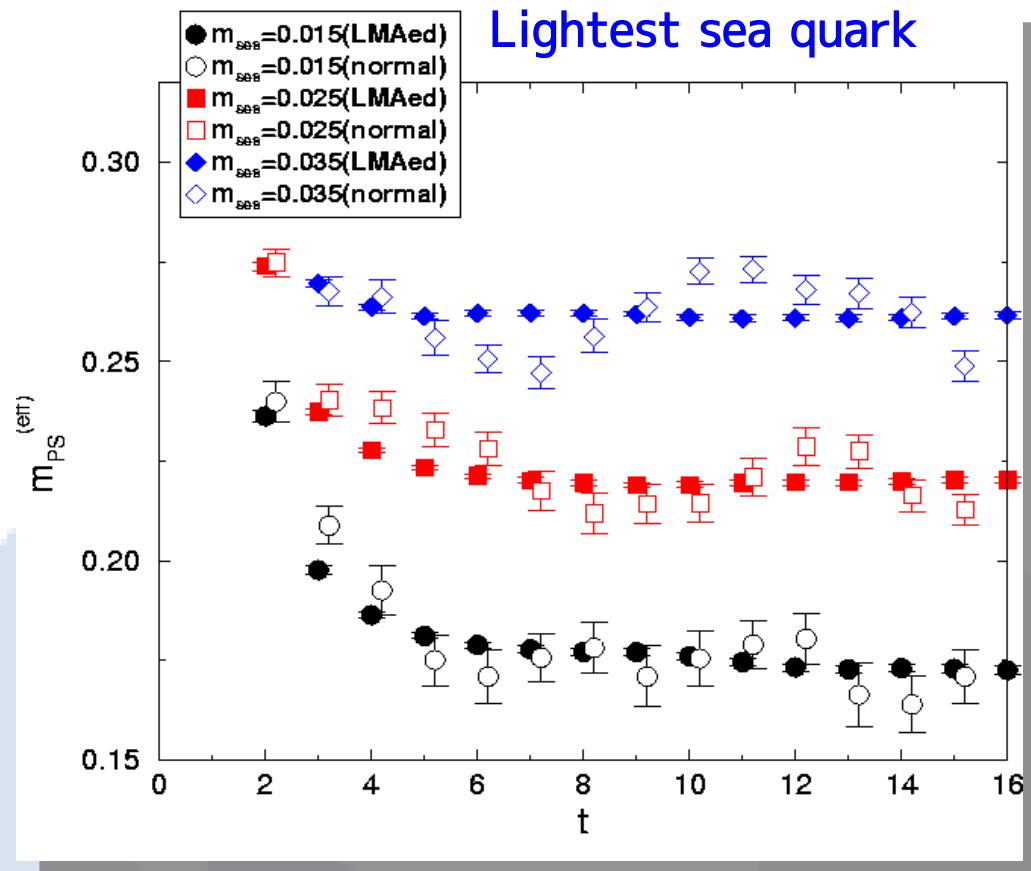
- Low mode preconditioned quark propagator

$$S_q(x, y) = \frac{\sum_{i=1}^{50} u_i(x) u_i(y)}{\lambda_i + m_q} + S_q^{\text{Higher}}(x, y)$$

- Low mode averaged meson correlator

$$C(t) = C^{\text{HH}}(t) + C^{\text{LH}}(t) + C^{\text{HL}}(t) + C^{\text{LL}}(t)$$

Average over space



Meson mass & decay constant

- Lattice data

$$\langle P(t)^{\text{pts}} P(0)^{\text{pts}} \rangle, \quad \langle P(t)^{\text{pts}} P(0)^{\text{smr}} \rangle$$

$$\rightarrow m_\pi, \quad \langle 0 | P | \pi \rangle = \frac{m_\pi}{2m_q} f_\pi^{\text{(ren)}}$$

- NLO ChPT

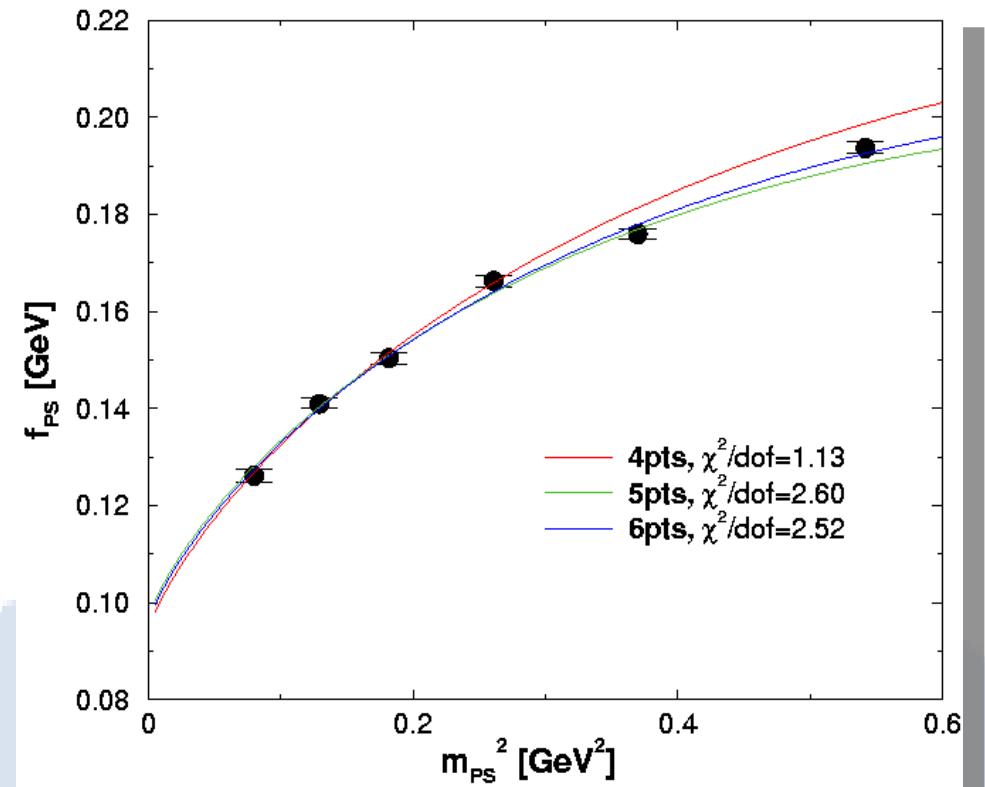
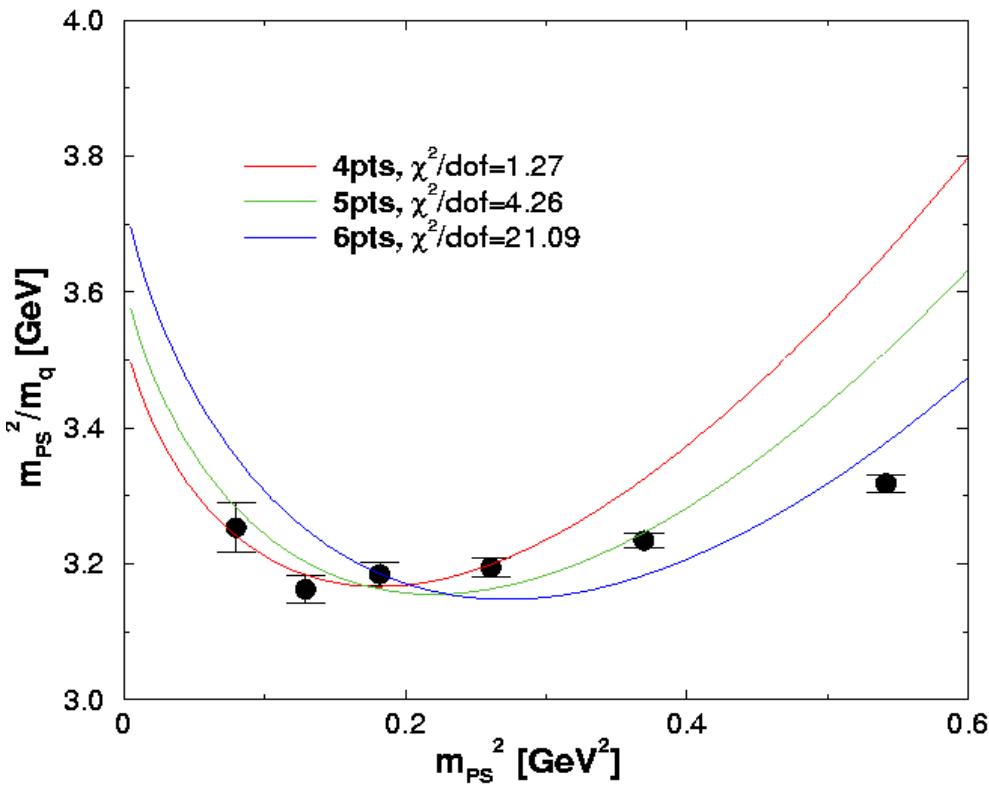
4 param. fit

$$\frac{m_\pi^2}{m_q} = 2B_0 \left[1 + \frac{2}{N_f} \frac{1}{(4\pi f)^2} m_\pi^2 \log m_\pi^2 \right] + \eta m_\pi^2$$

$$f_\pi = f \left[1 - \frac{N_f}{(4\pi f)^2} m_\pi^2 \log m_\pi^2 \right] + \zeta m_\pi^2$$

Dinamical point data

$m_{\text{valence}} = m_{\text{sea}}$

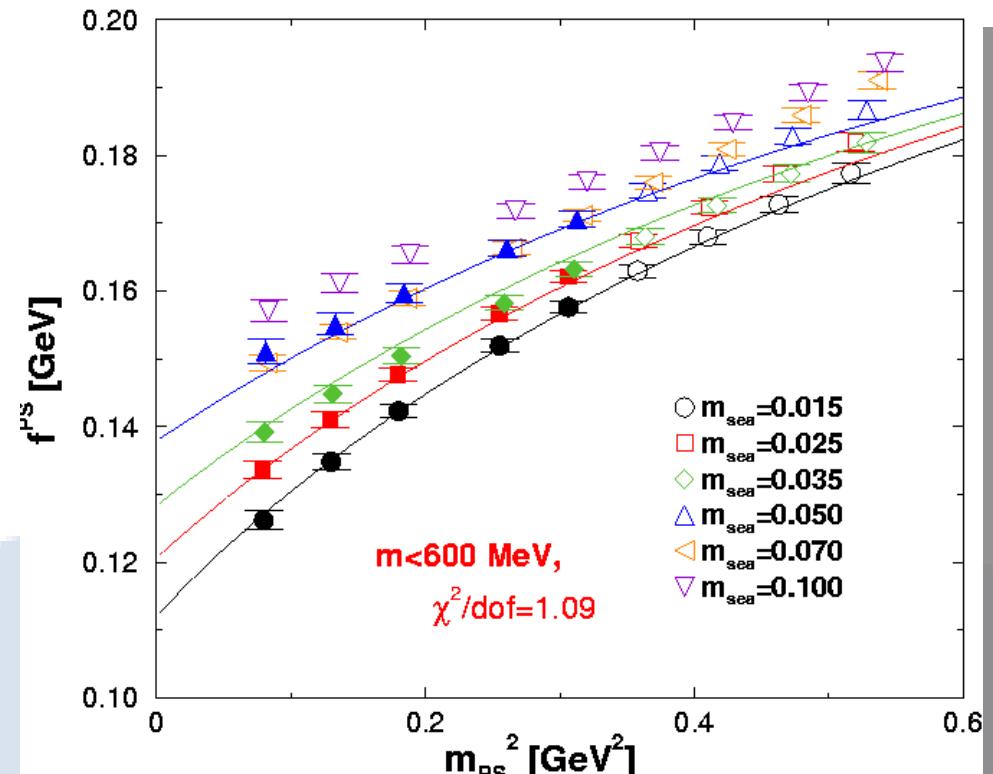
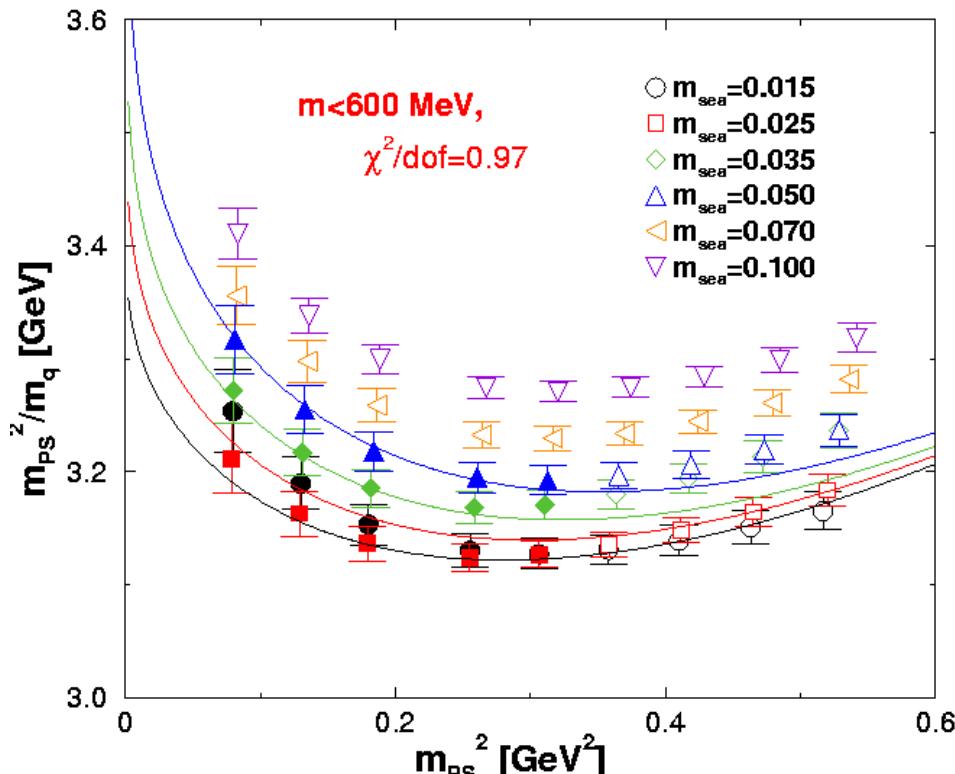


- $f = 94.4(1.4) \text{ MeV}$: 4pts fit
- FSE: $\sim 3\% / \sim +5\%$ for the lightest mass

Gasser and Leutwyler, 1987; Colangelo et al., 2005

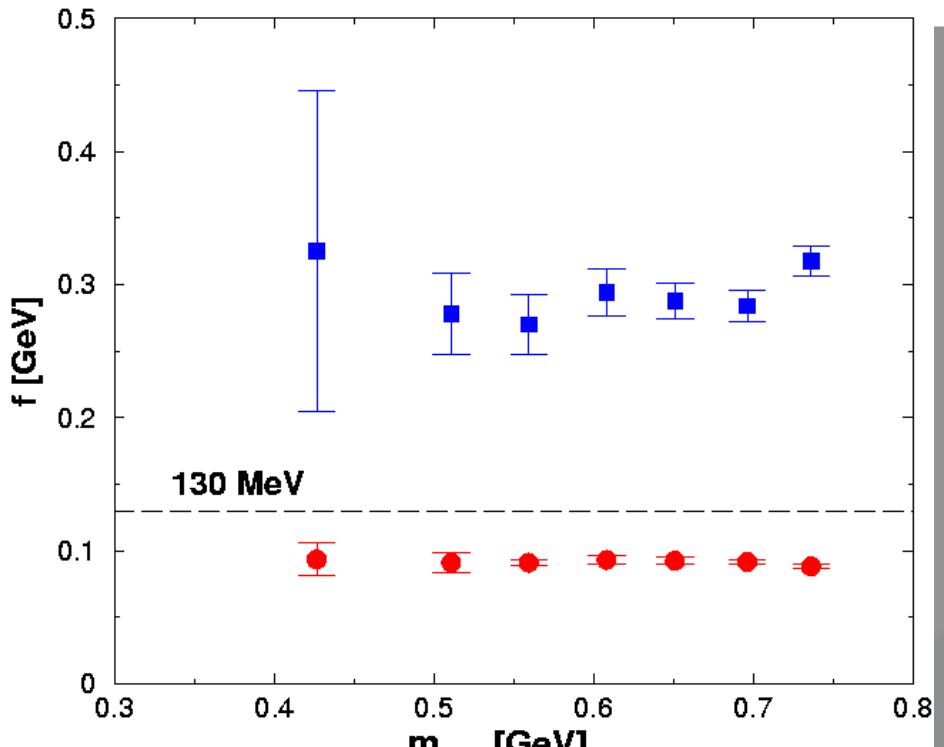
Partially quenched data

Sharpe, 1997; Golterman and Leung, 1997



- Fit with common f is impossible
- Independent fits are done

Test of (Pq)ChPT



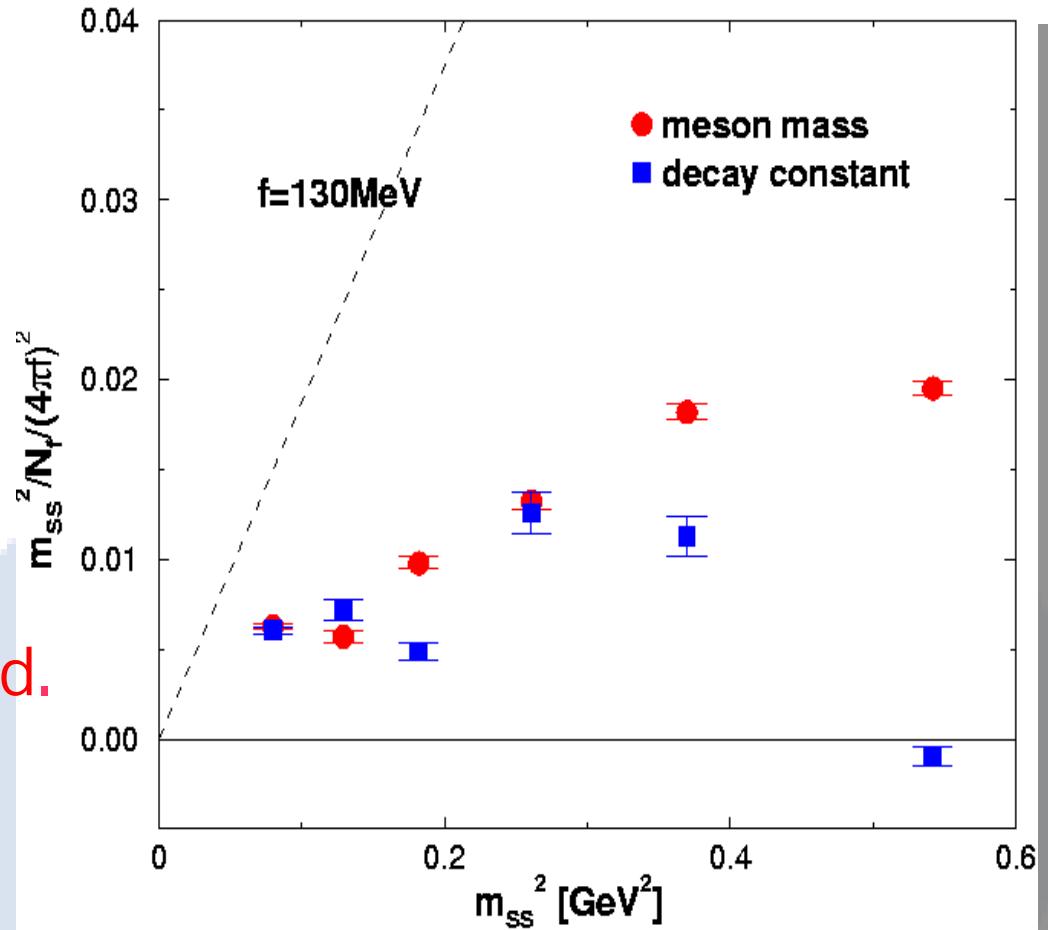
- Rather stable result
- ➡ FSE does not fill the gap.

Bernard-Golterman ratio

$$\frac{f_{VS}}{\sqrt{f_{VV}f_{SS}}} - 1 = -\frac{1}{2} \cdot \frac{m_{SS}^2}{N_f(4\pi f)^2} \cdot \tau$$

$$\frac{f_{VS}}{\sqrt{f_{VV}f_{SS}}} - 1 = -\frac{1}{2} \cdot \frac{m_{SS}^2}{N_f(4\pi f)^2} \cdot \tau$$

NNLO for meson mass is needed.
Effect of fixed topology ?



NPR for quark mass

Martinelli et al., 1995

- Renormalization condition

$$Z_q^{-1} Z_\Gamma \Lambda_\Gamma^{(\text{latt})}(p) = \Lambda_\Gamma^{(\text{tree})}$$

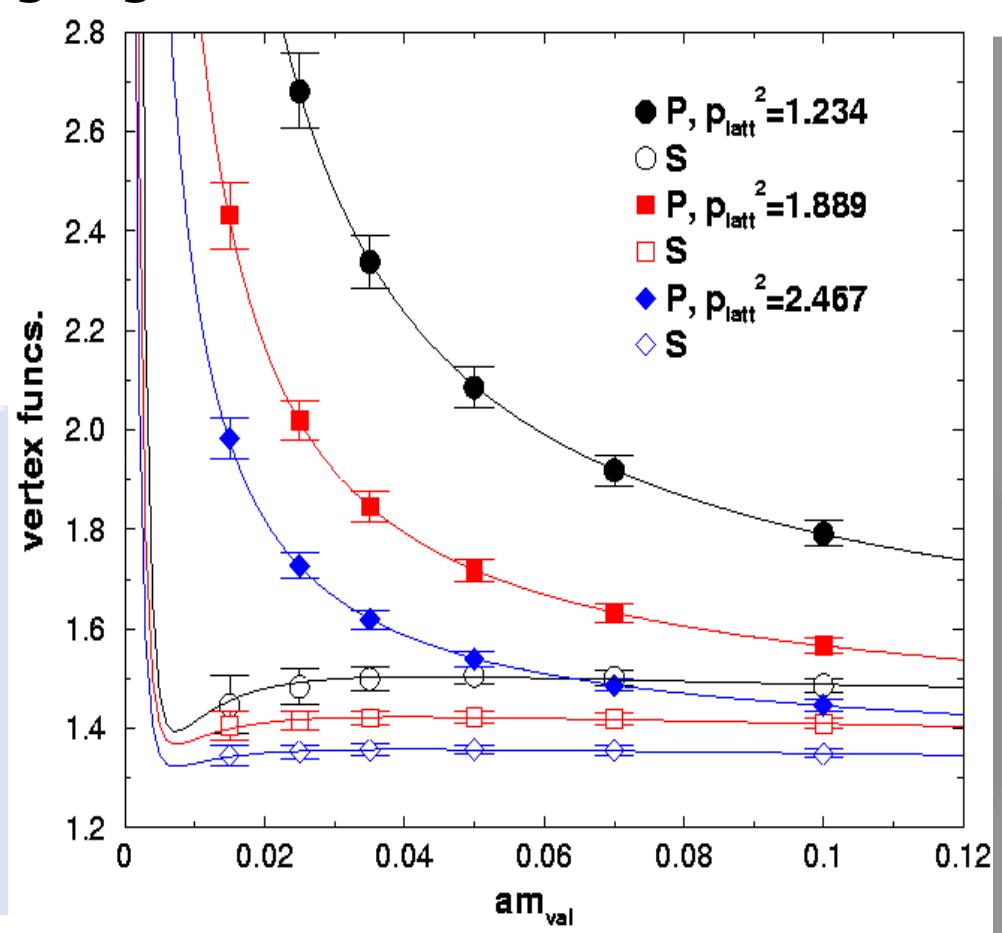
(Landau gauge)

- OPE + WTI

$$\Lambda_P(p) = A \cdot \frac{\langle \bar{\psi} \psi \rangle}{m_q} + Z_q Z_m + B_P \cdot m_q^2$$

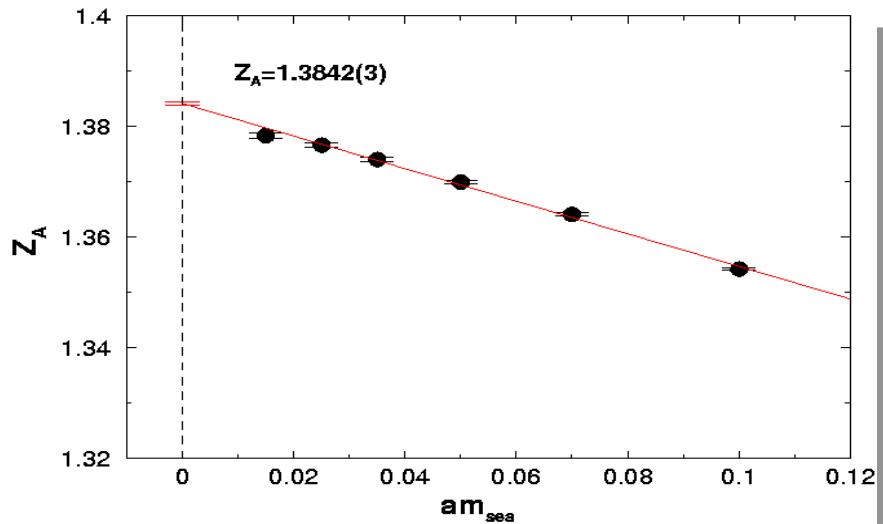
$$\Lambda_S(p) = A \cdot \frac{\partial}{\partial m_q} \langle \bar{\psi} \psi \rangle + Z_q Z_m + B_S \cdot m_q^2$$

$$\langle \bar{\psi} \psi \rangle = \left\langle \frac{1}{V_L} \sum_{i=1}^{50} \frac{2 m_q}{m_q^2 + \lambda_i^2} \right\rangle$$



- Axial vector

$$Z_A = \frac{\langle P(t)X(0) \rangle}{2m_q \Delta_t \langle A_4(t)X(0) \rangle}$$

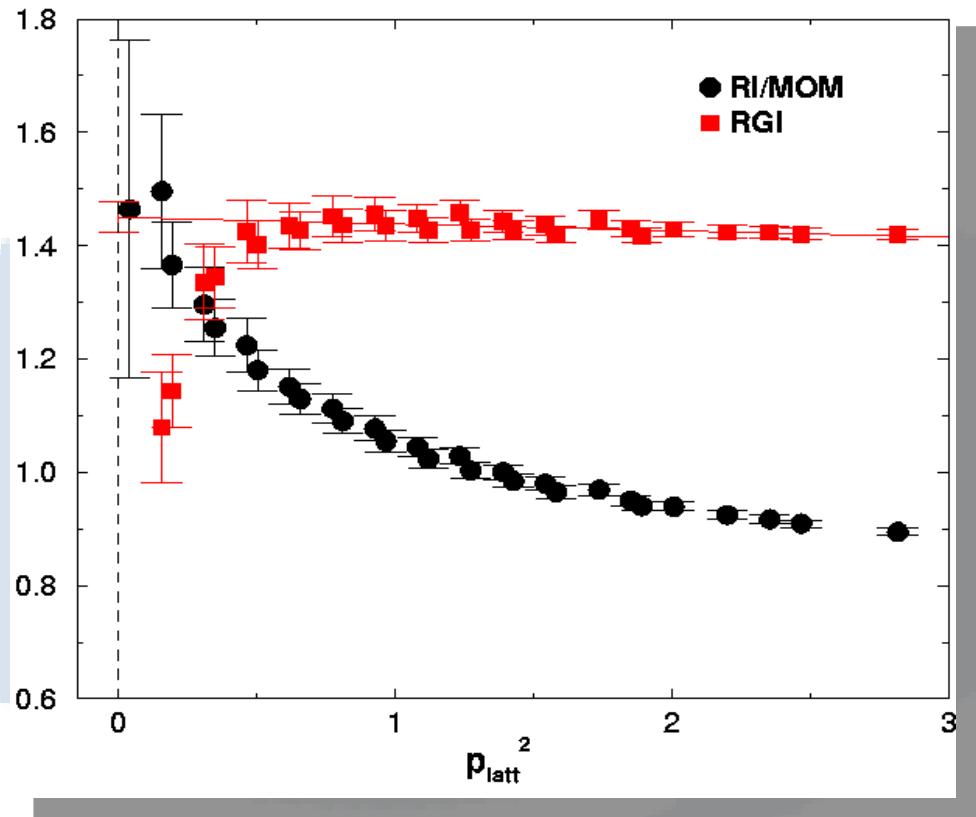


- Matching

$$Z_m^{\text{RI/MOM}} = Z_A \cdot \frac{Z_q Z_m}{\Lambda_A}$$

$$Z_m^{\text{RGI}} = w_{\text{RI/MOM}}(p) Z_m^{\text{RI/MOM}}(p)$$

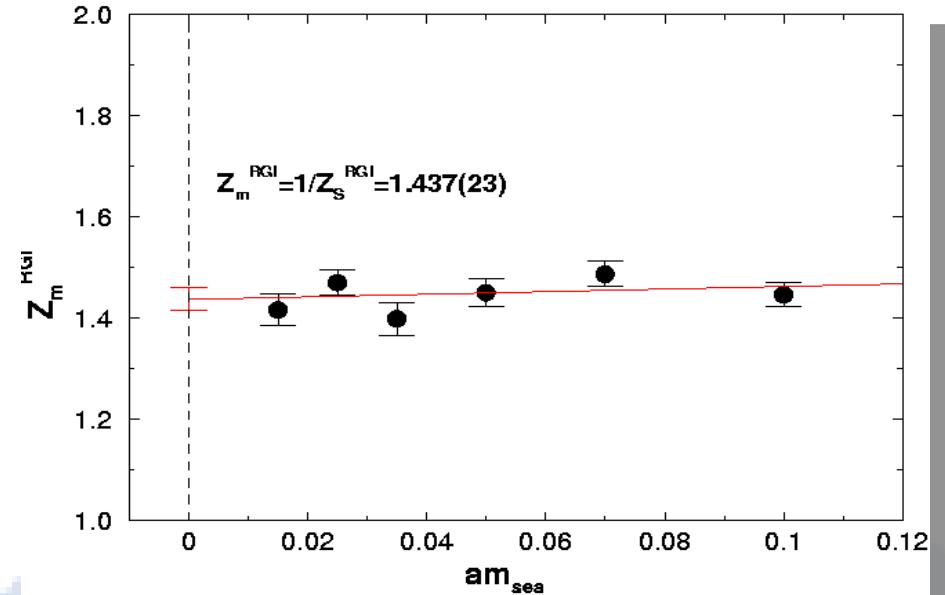
3-loop perturbation result
is used



Strange quark mass

- Dynamical limit

$$Z_m^{\overline{\text{MS}}}(2\text{GeV}) = w_{\overline{\text{MS}}}^{-1}(2\text{GeV}) \cdot Z_m^{\text{RGI}} \\ = 0.742(12)$$



- PCAC rel.

$$m_K = 0.495 \text{ MeV} \rightarrow$$

$$a \cdot m_s^{(\text{bare})} / 2 \simeq 0.050$$

$$m_s^{\overline{\text{MS}}}(2\text{GeV}) = 117(2)(\text{sys.})$$

Summary

New era has come.

