

ODDBALL REGGE TRAJECTORY AND A LOW ODDERON INTERCEPT

Stephen R. Cotanch

**Department of Physics,
North Carolina State University,
Raleigh, NC 27695-8202**

- ACKNOWLEDGE**

Felipe Llanes-Estrada

Pedro Bicudo

- SUPPORTED BY DOE**

**Odderon Searches
at RHIC Workshop
BNL, Sept. 2005**

- Effective QCD gluon Hamiltonian

$$H_{eff}^g = Tr \int d\mathbf{x} [\boldsymbol{\Pi}^a(\mathbf{x}) \cdot \boldsymbol{\Pi}^a(\mathbf{x}) + \mathbf{B}_A^a(\mathbf{x}) \cdot \mathbf{B}_A^a(\mathbf{x})] - \frac{1}{2} \int d\mathbf{x} d\mathbf{y} \rho_g^a(\mathbf{x}) V(\mathbf{x}, \mathbf{y}) \rho_g^a(\mathbf{y})$$

- color charge density $a, b, c = 1, 2, \dots, 8$

$$\rho_g^a(\mathbf{x}) = f^{abc} \mathbf{A}^b(\mathbf{x}) \cdot \boldsymbol{\Pi}^c(\mathbf{x})$$

- Abelian magnetic fields

$$\mathbf{B}_A^a = \boldsymbol{\nabla} \times \mathbf{A}^a$$

- gauge fields \mathbf{A}^a , conjugate momenta $\boldsymbol{\Pi}^a = -\mathbf{E}^a$

$$\begin{aligned} \mathbf{A}^a(\mathbf{x}) &= \int \frac{d\mathbf{q}}{(2\pi)^3} \frac{1}{\sqrt{2\omega_k}} [\mathbf{a}^a(\mathbf{q}) + \mathbf{a}^{a\dagger}(-\mathbf{q})] e^{i\mathbf{q}\cdot\mathbf{x}} \\ \boldsymbol{\Pi}^a(\mathbf{x}) &= -i \int \frac{d\mathbf{q}}{(2\pi)^3} \sqrt{\frac{\omega_k}{2}} [\mathbf{a}^a(\mathbf{q}) - \mathbf{a}^{a\dagger}(-\mathbf{q})] e^{i\mathbf{q}\cdot\mathbf{x}} \end{aligned}$$

- bare gluon Fock operators

$$a_\mu^a(\mathbf{q}) \text{ momentum } \mathbf{q} , \text{ spin } \mu = 0, \pm 1$$

- Coulomb gauge transverse condition

$$\boldsymbol{\nabla} \cdot \mathbf{A}^a = 0 \implies \mathbf{q} \cdot \mathbf{a}^a(\mathbf{q}) = (-1)^\mu q_\mu a_{-\mu}^a(\mathbf{q}) = 0$$

- Bogoliubov-Valatin canonical transformation (BCS rotation) to dressed gluon operators

$$\alpha_\mu^a(\mathbf{q}) = \cosh \Theta(q) a_\mu^a(\mathbf{q}) + \sinh \Theta(q) a_\mu^{a\dagger}(-\mathbf{q})$$

- satisfy the transverse commutation relations

$$[\alpha_\mu^a(\mathbf{q}), \alpha_\nu^{b\dagger}(\mathbf{q}')] = \delta_{ab}(2\pi)^3 \delta^3(\mathbf{q} - \mathbf{q}') D_{\mu\nu}(\mathbf{q})$$

$$D_{\mu\nu}(\mathbf{q}) = \left(\delta_{\mu\nu} - (-1)^\mu \frac{q_\mu q_{-\nu}}{q^2} \right)$$

- PQCD vacuum $|0\rangle \implies$ BCS vacuum $|\Omega\rangle_{\text{BCS}}$

- minimize ground state energy

$$\frac{\delta}{\delta \Theta} \left(\frac{\langle \Omega | H_{eff} - E | \Omega \rangle}{\langle \Omega | \Omega \rangle} \right) = 0$$

- generates a gap equation

$$\omega(q)^2 = q^2 - \frac{3}{4} \int \frac{d\mathbf{k}}{(2\pi)^3} \hat{V}(|\mathbf{q}-\mathbf{k}|)(1+(\hat{\mathbf{q}}\hat{\mathbf{k}})^2) \left(\frac{w(k)^2 - w(q)^2}{w(k)} \right)$$

- gluon self-energy

$$\omega(q) = q e^{-2\Theta(q)}$$

- confinement via Cornell type potential

$$V(|\mathbf{x} - \mathbf{y}|) = \sigma |\mathbf{x} - \mathbf{y}| - \frac{\alpha_s}{|\mathbf{x} - \mathbf{y}|}$$

$$\sigma = 0.18 (\textit{from lattice}) \quad \alpha_s = 0.42$$

- solve gap eq. for gluon constituent mass

$$m_g \equiv \omega(0) \cong 0.8 GeV$$

- predicted gluon condensate (cooper pairs)

$$\langle \alpha G_{\mu\nu}^a G_a^{\mu\nu} \rangle = (433 \text{ MeV})^4$$

[agrees with lattice (441 MeV)⁴]

- TDA two gluon glueball wavefunction

$$|\Psi_{LS}^{JPC}\rangle = \sum_{am_1m_2} \int \frac{d\mathbf{k}}{(2\pi)^3} \Phi_{LSm_1m_2}^{JPC}(\mathbf{k}) \alpha_{m_1}^{a\dagger}(\mathbf{k}) \alpha_{m_2}^{a\dagger}(-\mathbf{k}) |\Omega\rangle$$

- solve for glueball mass M_{JPC}

$$H|\Psi_{LS}^{JPC}\rangle = M_{JPC}|\Psi_{LS}^{JPC}\rangle \quad [\textit{agree with lattice}]$$

- predict a Regge trajectory close to the pomeron

$$\alpha_P \approx .25t + 1$$

- variational three gluon glueball wavefunction

$$|\Psi^{JPC}\rangle = \int d\mathbf{q}_1 d\mathbf{q}_2 d\mathbf{q}_3 \delta(\mathbf{q}_1 + \mathbf{q}_2 + \mathbf{q}_3) F_{\mu_1\mu_2\mu_3}^{JPC}(\mathbf{q}_1, \mathbf{q}_2, \mathbf{q}_3) C^{abc} \alpha_{\mu_1}^{a\dagger}(\mathbf{q}_1) \alpha_{\mu_2}^{b\dagger}(\mathbf{q}_2) \alpha_{\mu_3}^{c\dagger}(\mathbf{q}_3) |\Omega\rangle_{\text{BCS}}$$

- Bose statistics $C^{abc} = f^{abc}$ ($C = 1$) or d^{abc} ($C = -1$)
- variational equation for the J^{PC} glueball

$$\frac{\langle \Psi^{JPC} | H_{eff}^g | \Psi^{JPC} \rangle}{\langle \Psi^{JPC} | \Psi^{JPC} \rangle} = M_{JPC} \quad [\text{agree with lattice}]$$

- predict leading Regge trajectory (odderon)

$$\alpha_O^{eff} = .23t - 0.88 \quad [\text{note low intercept}]$$

- compare to nonrelativistic constituent model

$$H_M = \sum_i \frac{\mathbf{q}_i^2}{2m_g} + V_0 + \sum_{i < j} [\sigma r_{ij} - \frac{\alpha}{r_{ij}} + V_{ss} \mathbf{S}_i \cdot \mathbf{S}_j]$$

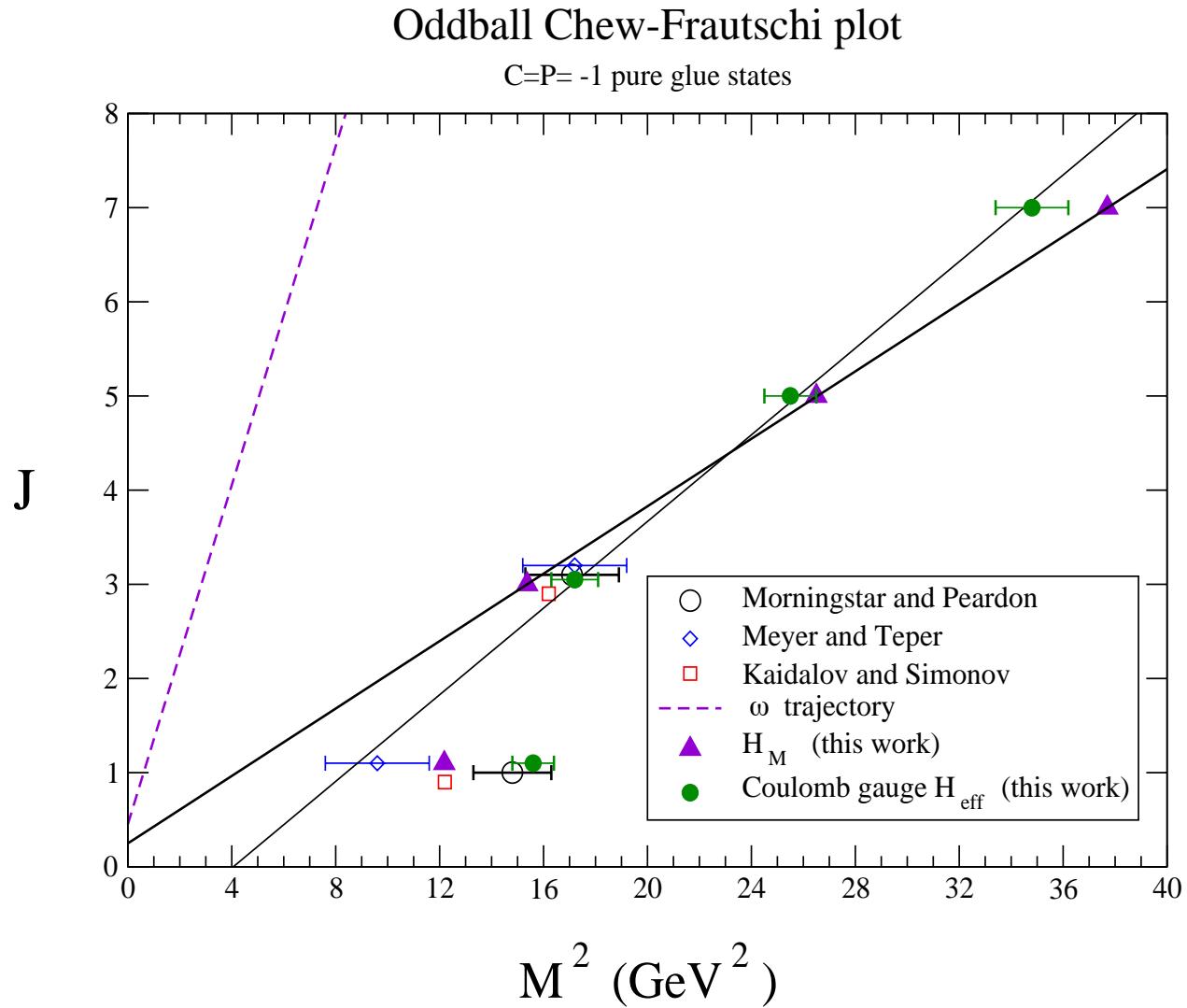
- use $q\bar{q}$ funnel potential

$V_0 = -.9 \text{ GeV}$, $\alpha = .27$, $\sigma = .25 \text{ GeV}^2$ $m_g = .8 \text{ GeV}$
and adjust $V_{ss} \rightarrow 0.085 \text{ GeV}$ to fit pomeron

- predict leading Regge trajectory (odderon)

$$\alpha_O^M = .18t + 0.25 \quad [\text{note low intercept}]$$

lattice, H_{eff}^g and H_M odderons vs. ω trajectory



SUMMARY

- lattice, H_{eff}^g and H_M models predict an odderon
- starts with 3^{--} (not 1^{--} , need lattice 5^{--})
- odderon slope is similar to pomeron
- odderon intercept is low, below 0.5
- search for odderon where ω trajectory is not dominant (examine $\frac{d\sigma}{dt}$, not $\sigma_{total} \propto s^{\alpha_\omega(0)-1}$)