

Progress on proton acceleration from a hydrogen gas jet

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Benefits of a gas jet as an ion beam source

- Pure (compared to solid targets which become quickly covered in impurities)
- Can employ H, He and other species difficult to make in other targets
- Allows changing target material quickly
- Can run at high repetition rate

Benefits from combining gas jet with a CO₂ laser

- Due to $\lambda = 10 \mu\text{m}$, $n_{cr} = \epsilon_0 m \omega_0^2 / e^2 = 10^{19} \text{ cm}^{-3}$ is 100 times lower than for a solid state laser.

Gas-jets easier to make at this density allowing to operate hydrogen jet (proton source) in the most efficient, near-critical regime where maximum ion acceleration has been observed (foam targets) and RPA has been predicted.

- Possibility for optical probing of overdense interactions.

For 532nm, $n_{cr} \sim 4 \times 10^{21} \text{ cm}^{-3}$ (easy transmitted through the gas jet)

BNL experiment with gas jet

- Nozzle 0.5-1mm
- Backing press.: 115 - 1100psig
- $1.3n_{cr} - 12.3n_{cr}$

- 2ω Nd:YAG, $\lambda=532\text{nm}$
- Pulse energy $\sim 200\mu\text{J}$
- Pulse length $\sim 10\text{ps}$

Parabolic mirror

Probe beam

Gas jet

Pinhole

Magnet

S

N

H^+

neutral

CCD camera
(ion spectra)

Scintillator

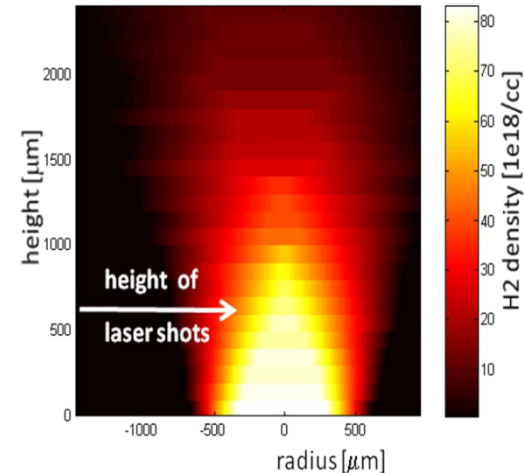
Imaging lens

CCD camera
(plasma image)

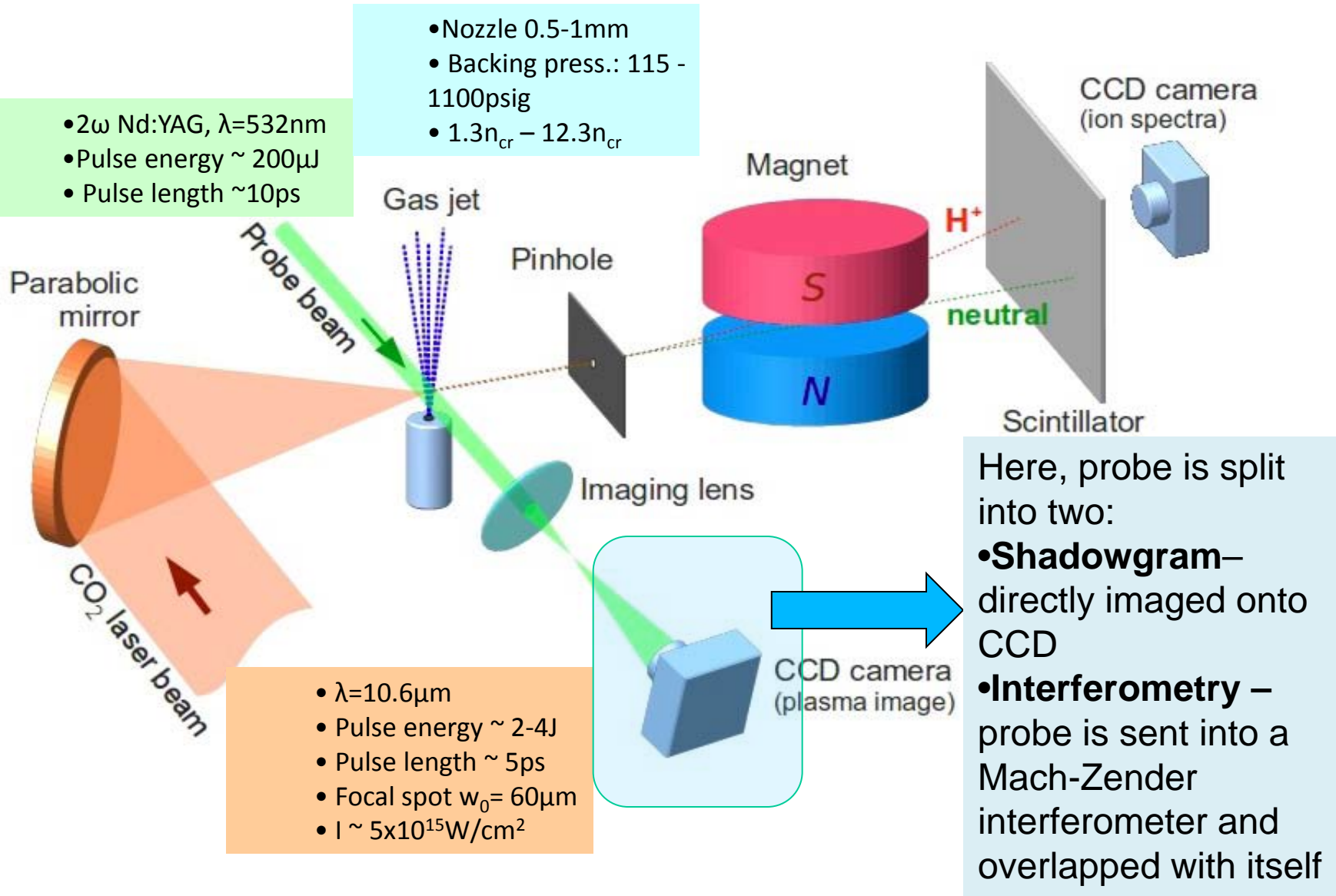
CO_2 laser beam

- $\lambda=10.6\mu\text{m}$
- Pulse energy $\sim 2-4\text{J}$
- Pulse length $\sim 5\text{ps}$
- Focal spot $w_0=60\mu\text{m}$
- $I \sim 5 \times 10^{15}\text{W}/\text{cm}^2$

600 psi, 1 mm, H_2

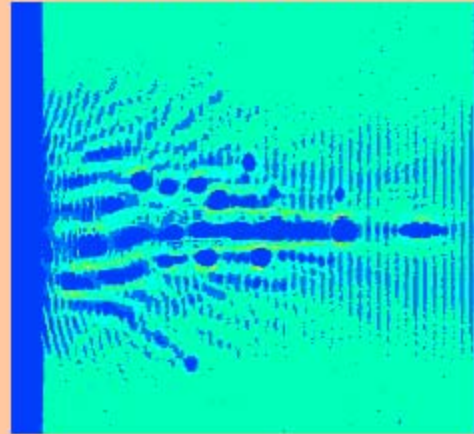


BNL experiment with gas jet



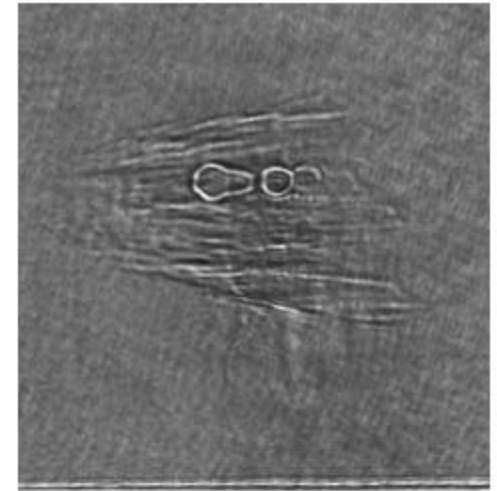
Wealth of non-linear phenomena observed

simulation

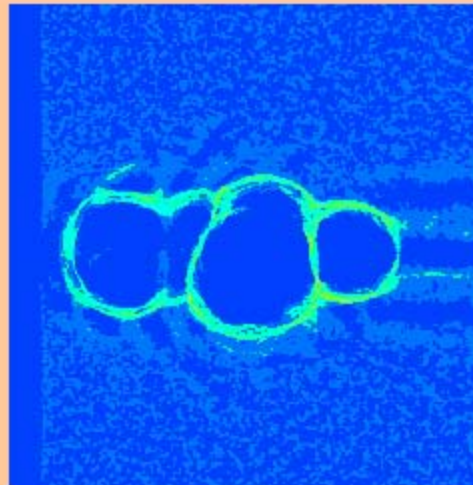


Plasma filaments
and solitons

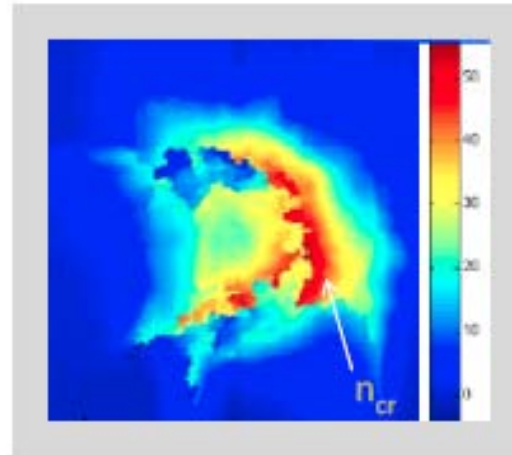
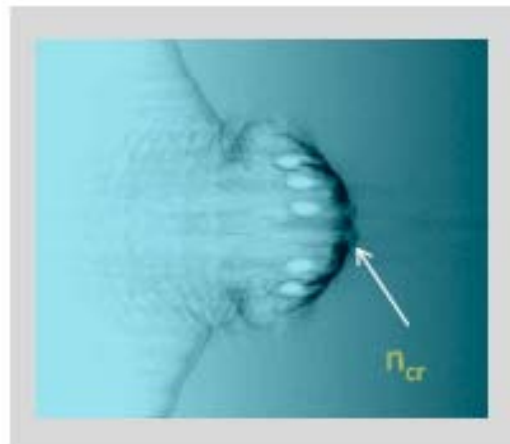
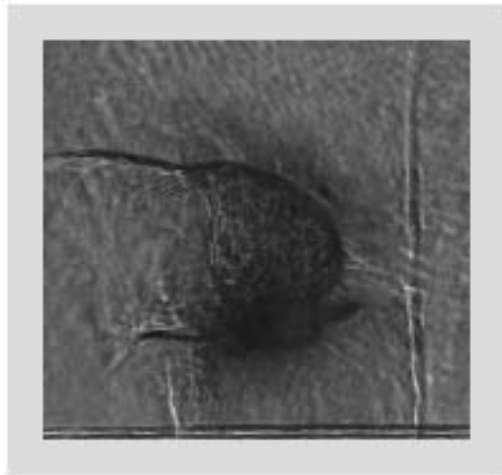
experiment



Solitons expand
and combine into
post-solitons at
the ion time scale

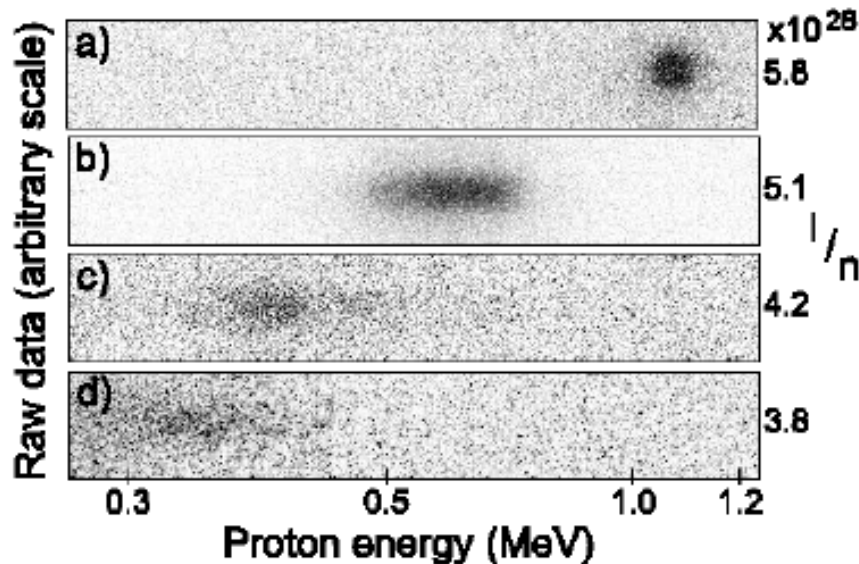
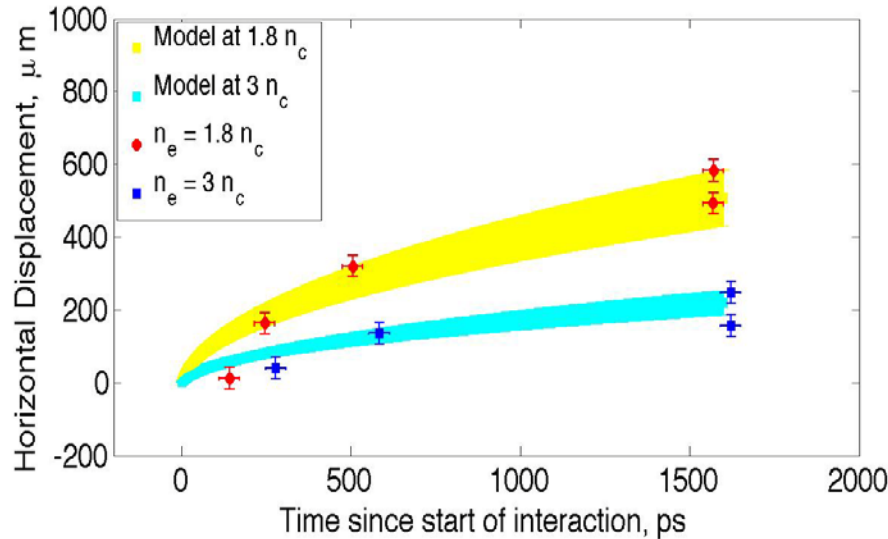


Hole boring by radiation pressure



- Plasma discontinuity represents a shock with the critical surface moving at $\sim I/n$.

Trend of increasing proton energy with I/n

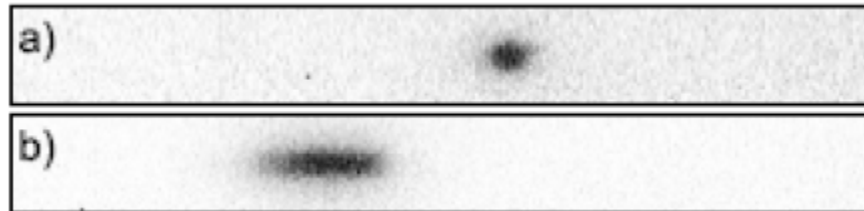
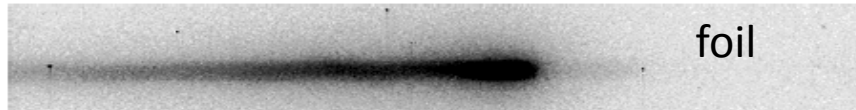


- For opaque plasma ($n > n_{cr}$), the radiation pressure, $P = 2I/c$, initially pushes plasma electrons into the target.
- Space charge field pulls along ions setting up an electrostatic shock moving at hole boring velocity $v = (2I/\rho c)^{1/2}$

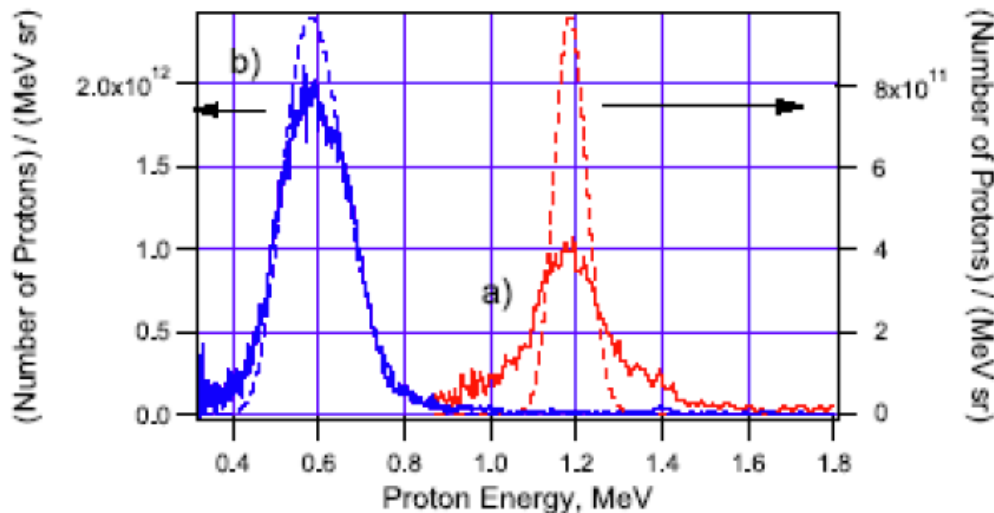
- Stationary ions in advance of the shock get accelerated by the same space charge field effectively bouncing off the shock front.*
- Associated with it proton energy $E = \frac{1}{2}m(2v)^2 = 4I/nc$

* PRL 92, 015002 (2004)
PRL 93, 155003 (2004)

Narrow energy spread

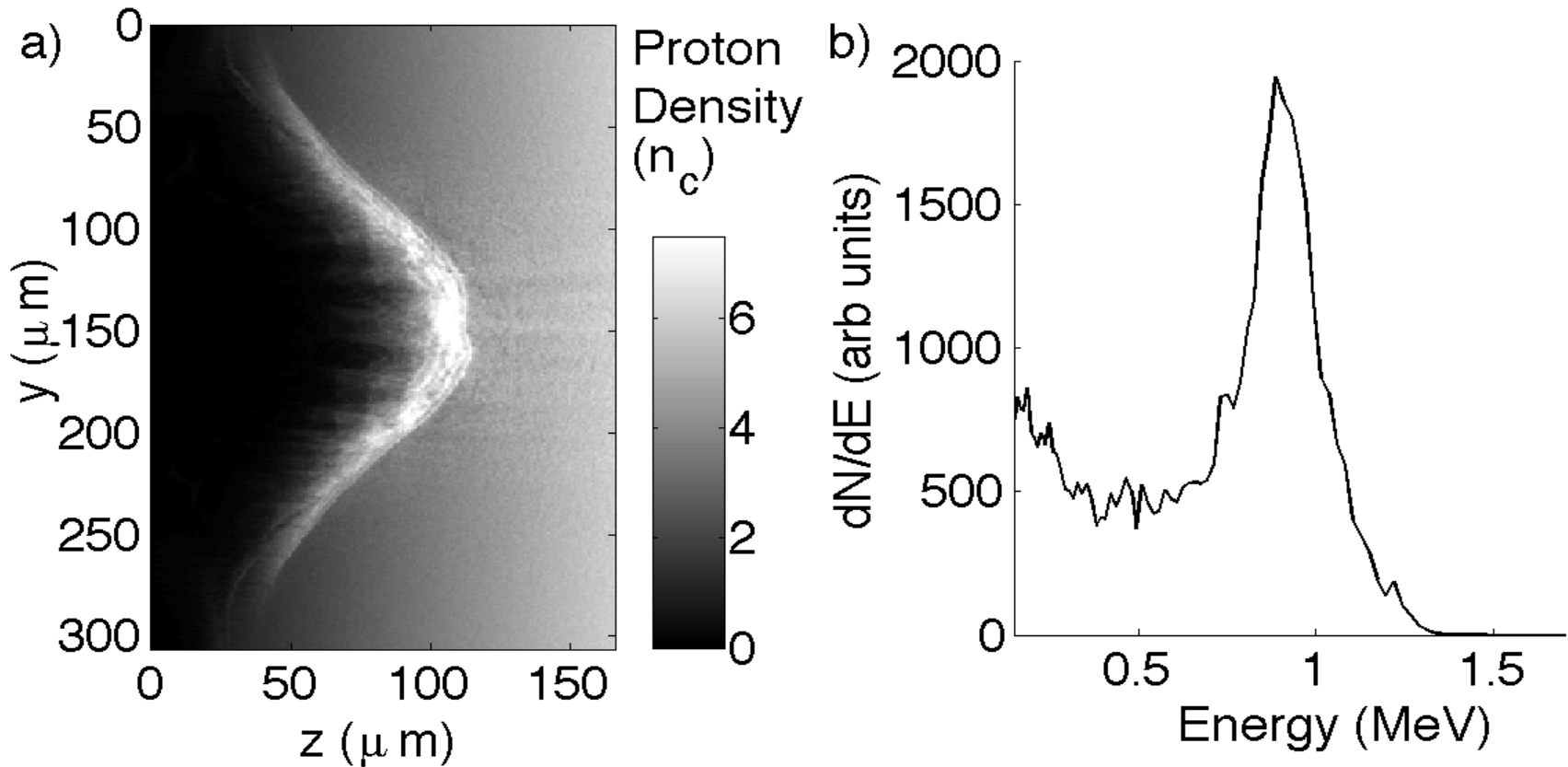


0.35 0.4 0.5 0.6 0.8 1.0 1.5 2.25 MeV



- Proton beams observed with consistently narrow energy spread.
- Deconvolution with spectrometer instrument function to determine energy spread of beams.
- Energy spread down to $\sim 4\%$ \rightarrow significantly narrower than previously observed.
- $\sim 2 \times 10^{12}$ protons/MeV/sr \rightarrow 1000 brighter than previous modulated spectra.
- Geometrical emittance $\epsilon = 0.16 \mu\text{m-rad}$
- Normalized emittance $\epsilon_n = \beta\gamma\epsilon = 8 \text{ nm-rad}$

Simulation



Simulation Conditions:

Laser:

$a_0 = 0.6$

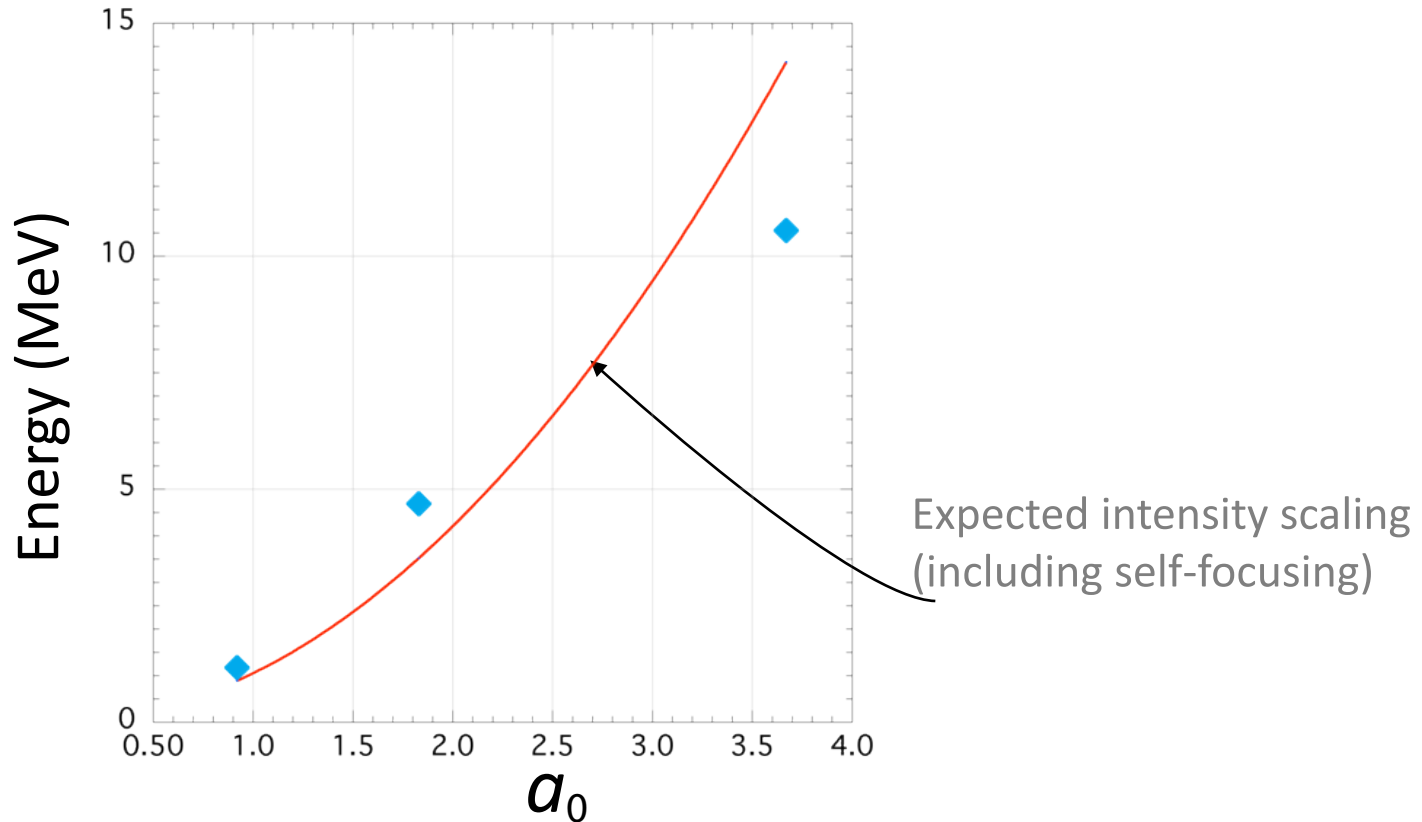
$\tau_L = 6\text{ps}$

Target:

Ionised H_2

Triangular density profile

Future prospects



- Peak energies expected to increase with increasing a_0
- $a_0 = 1.5$ easily achievable by changing to faster focusing
- higher a_0 conceivable with planned laser upgrades

Additional investigations

- As well as intensity scaling, we plan to measure:
- Angular distribution of ion beam
- Acceleration of different ion species
- Investigation of influence of double pulses
- Investigation of influence of polarization
- Investigation of micro-gas targets

