

Shrinking the Quark-Gluon Plasma

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[arXiv: 1901.01319](https://arxiv.org/abs/1901.01319)

[arXiv: 1905.13323](https://arxiv.org/abs/1905.13323)

RHIC / AGS Annual
Users' Meeting 2019

Tues, June 4, 2019

Motivation: Turning Off the QGP in Small Systems

- **Hydrodynamics** successfully describes the bulk low-pT yields in heavy-ion collisions
- Its **basic assumptions should turn off** at lower energies and smaller collision systems
 - Does it really turn off?
 - Or do the signals of the QGP survive under the extreme conditions of a tiny droplet?
- Motivation for a **system size scan** in addition to a beam energy scan:
 - PbPb, XeXe, ArAr, OO

S.H. Lim et al., Phys. Rev. C99 (2019)

Hydrodynamics \approx Linear Response

$$\mathcal{E}_n \equiv -\frac{\langle \int r^n e^{in\phi} s(r, \phi) r dr d\phi \rangle}{\langle \int r^n s(r, \phi) r dr d\phi \rangle} \xrightarrow{V_n \approx \kappa_n \mathcal{E}_n} V_n = \frac{\langle \int d^2p dy e^{in\phi} \frac{dN_1}{d^2p dy} \rangle}{\langle N_{\text{tot}} \rangle}$$

- Because of **perfect fluidity**, the QGP is **“Nature’s Fourier Transform”**

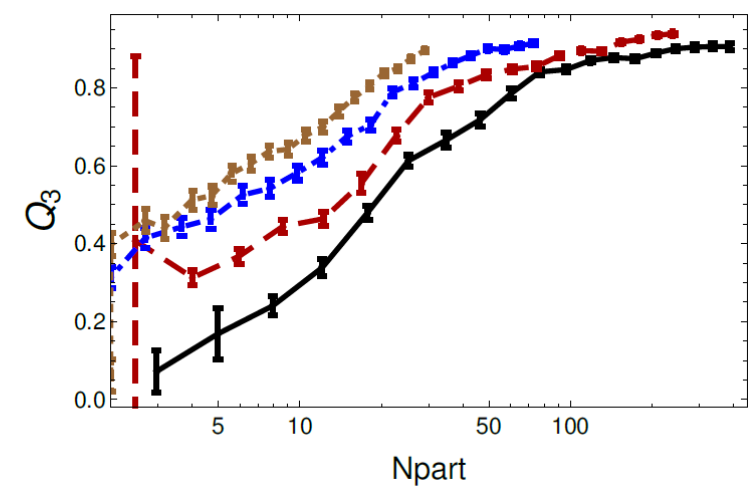
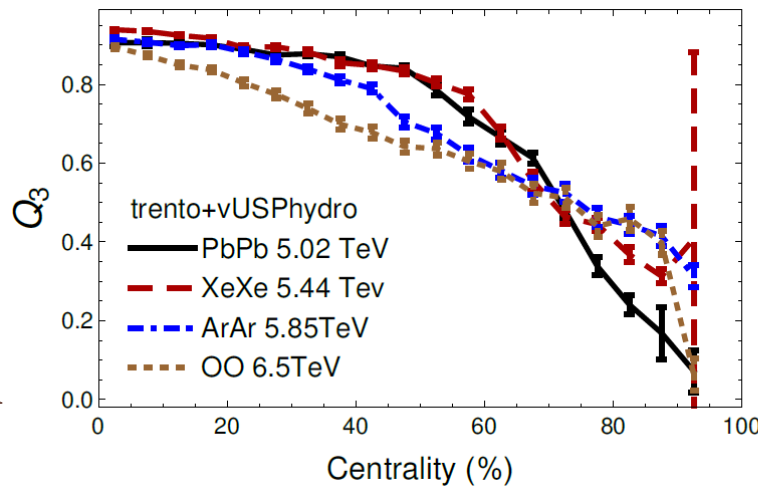
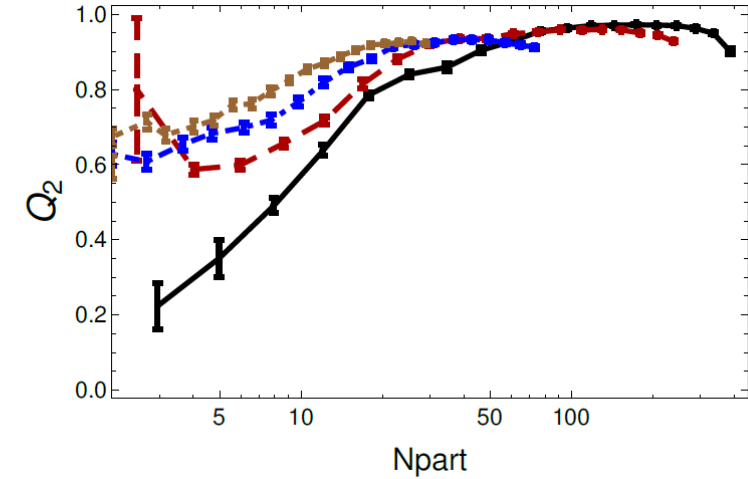
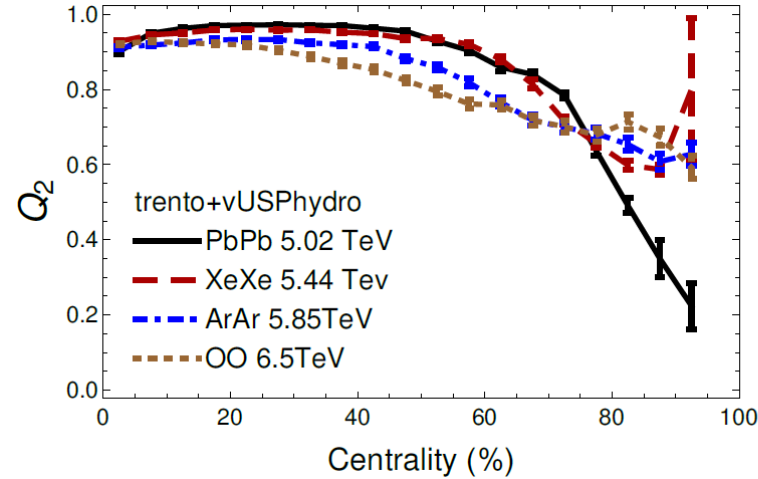
T. Dore, Rutgers University

- Quantified by the **Pearson Coefficient**

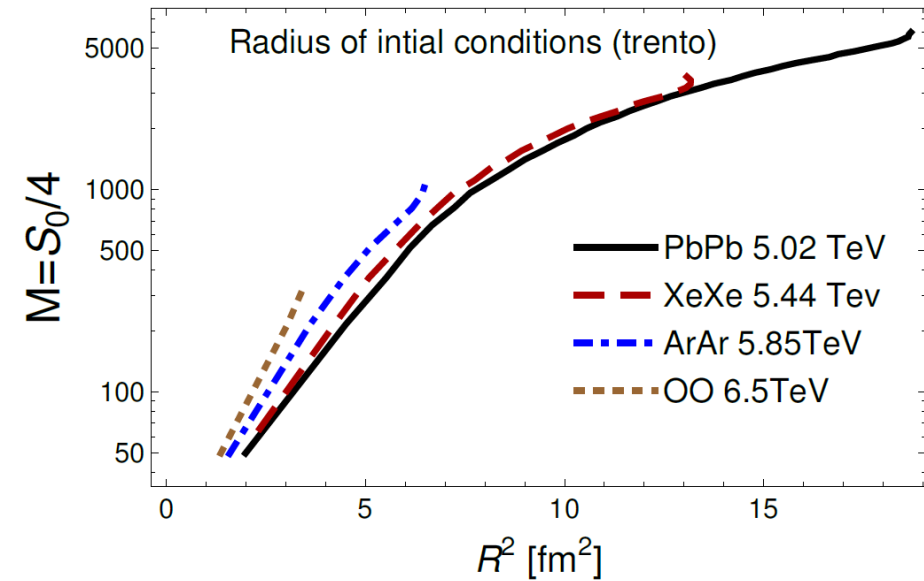
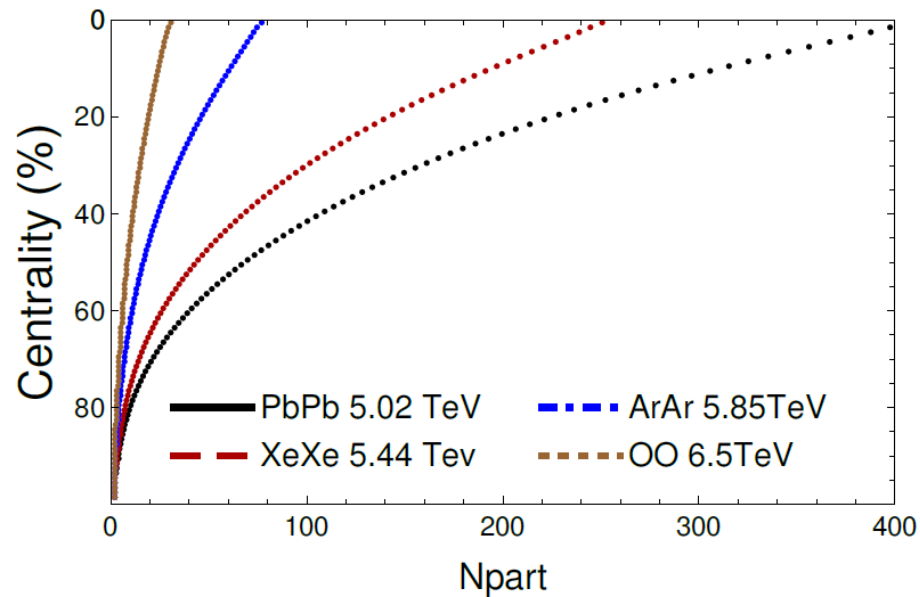
$$Q_n \equiv \frac{\langle V_n \mathcal{E}_n^* \rangle}{\sqrt{\langle |V_n|^2 \rangle} \sqrt{\langle |\mathcal{E}_n|^2 \rangle}} = \frac{\langle v_n \varepsilon_n \cos(n[\psi_n - \phi_n]) \rangle}{\sqrt{\langle |\varepsilon_n|^2 \rangle} \sqrt{\langle |v_n|^2 \rangle}}$$

Hydrodynamics \approx Linear Response

- **Linear response works well in central / mid-central collisions across system size**
- The visible trends depend significantly on **what is held fixed: centrality or Npart**
- For a **given Npart, smaller systems** are more central and have **better linear response**



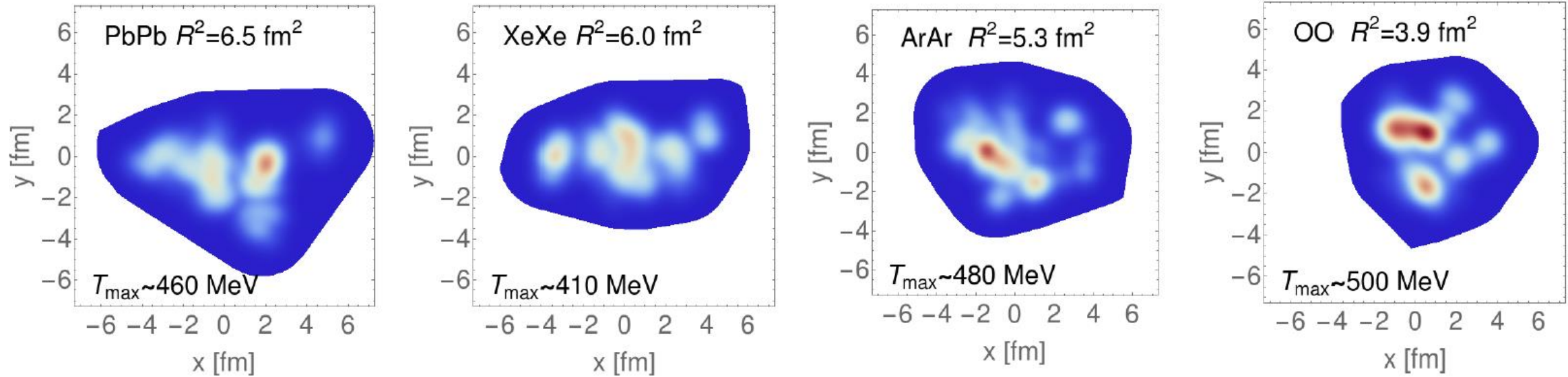
Not All Small Droplets Are Created Equal



- **A “small droplet” in PbPb is very different from one in OO**
 - **More peripheral** in PbPb, more central in OO
 - **More elliptical** in PbPb, more round in OO
 - **Lower temperature** in PbPb, higher temperature in OO

Not All Small Droplets Are Created Equal

For fixed multiplicity



Larger System



More Elliptical

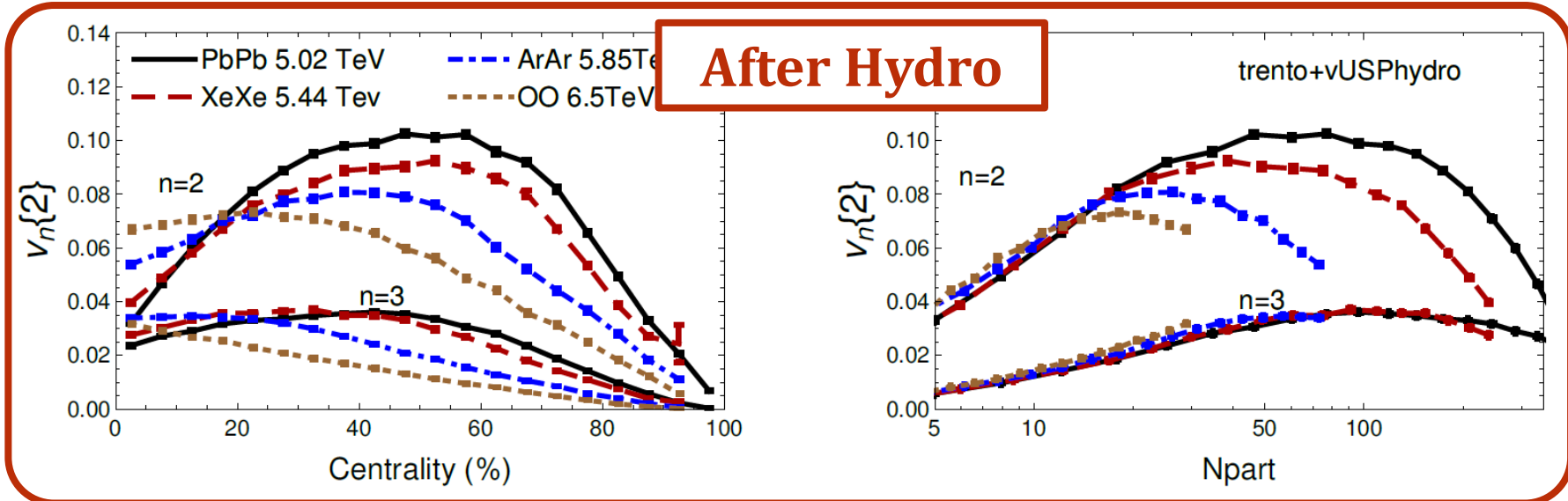
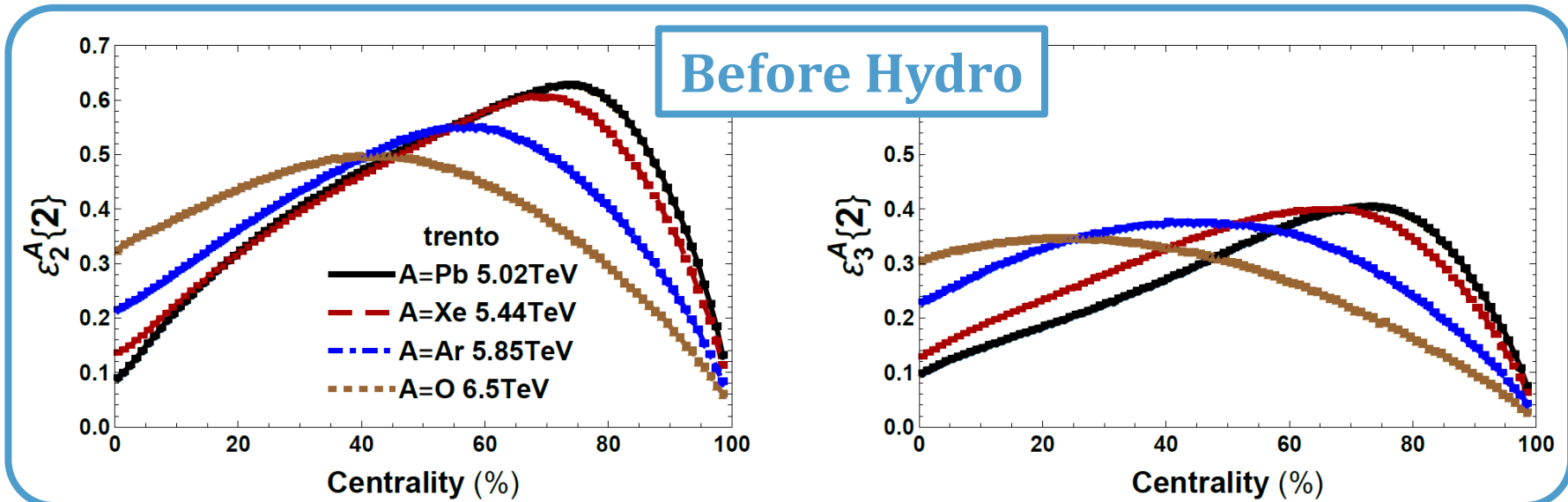


Higher Temperature



Initial Eccentricities and Final Flow Harmonics

- **System size hierarchy** exhibits a **crossing** in mid-central collisions
- **Linear response** a good qualitative predictor of v_n
- **Npart dependence** better isolates system size effects



Cumulant Ratios: Measure of Fluctuations

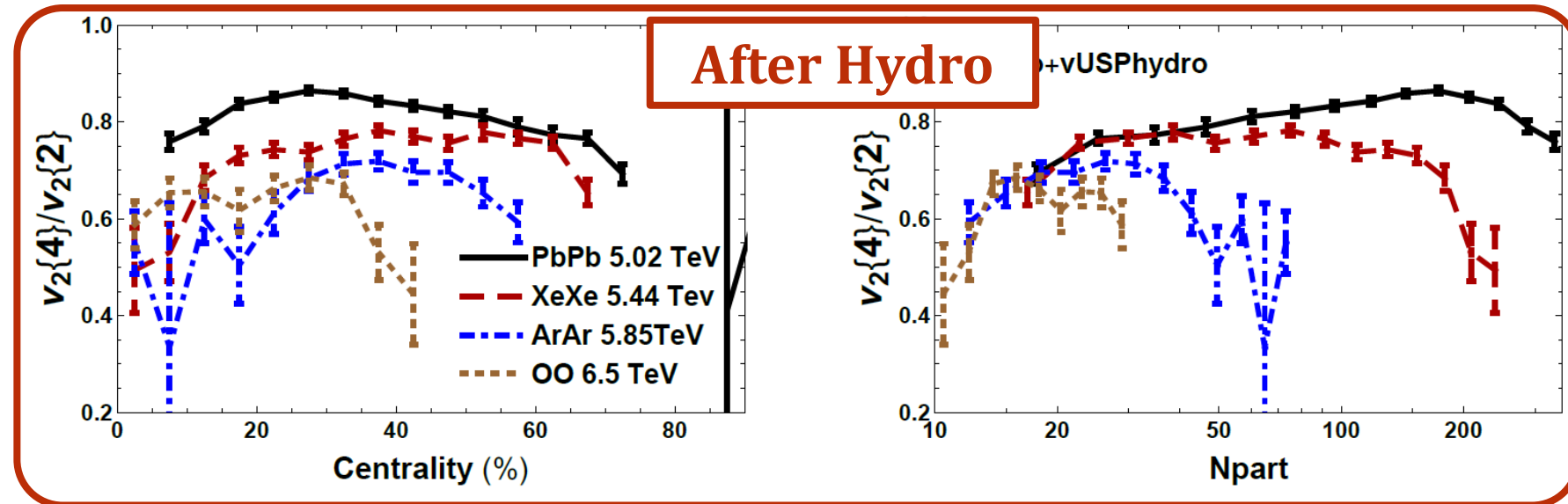
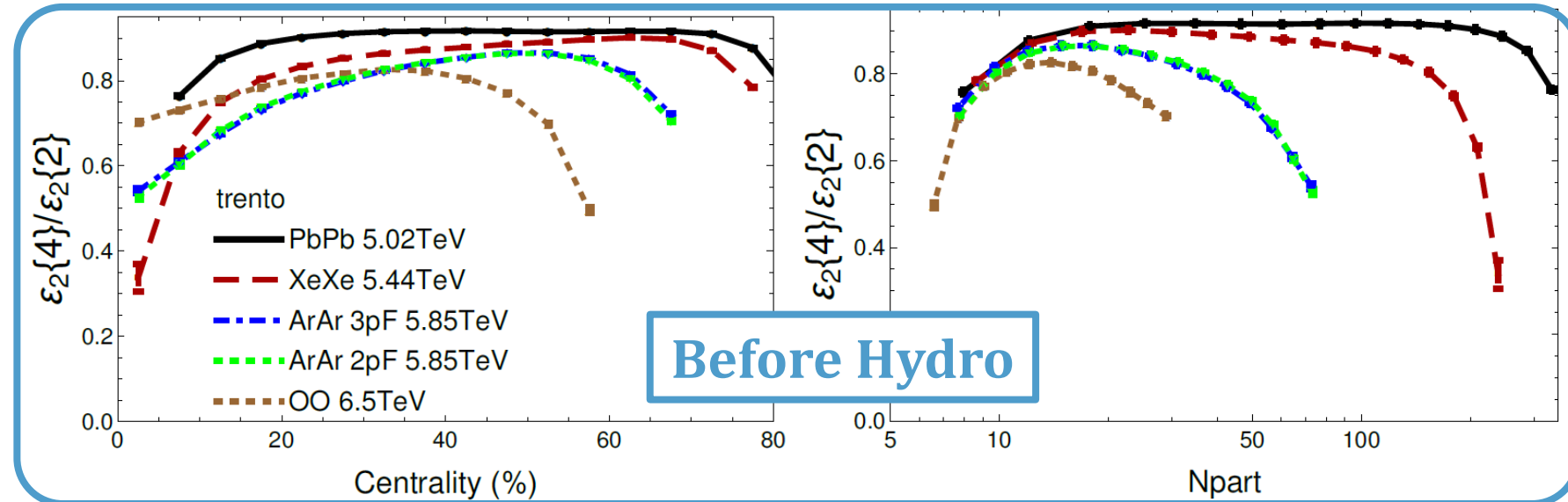
$$\begin{aligned} (v_n\{4\})^4 &\stackrel{flow}{=} 2\langle v_n^2 \rangle^2 - \langle v_n^4 \rangle \\ &= (v_n\{2\})^4 - \text{Var}(v_n^2) \end{aligned}$$

- If $\frac{v_n\{4\}}{v_n\{2\}} \approx 1$ then fluctuations are **small**
- If $\frac{v_n\{4\}}{v_n\{2\}} \ll 1$ then fluctuations are **large**

Initial state analog $\varepsilon_n\{m\}$
defined with averages of ε_n

Cumulant Ratios: Measure of Fluctuations

- **Smaller systems** tend to generate **more fluctuations**
- **Npart dependence** provides a better comparison than centrality
- **Linear response** is a reasonable estimate of fluctuations



Symmetric Cumulants

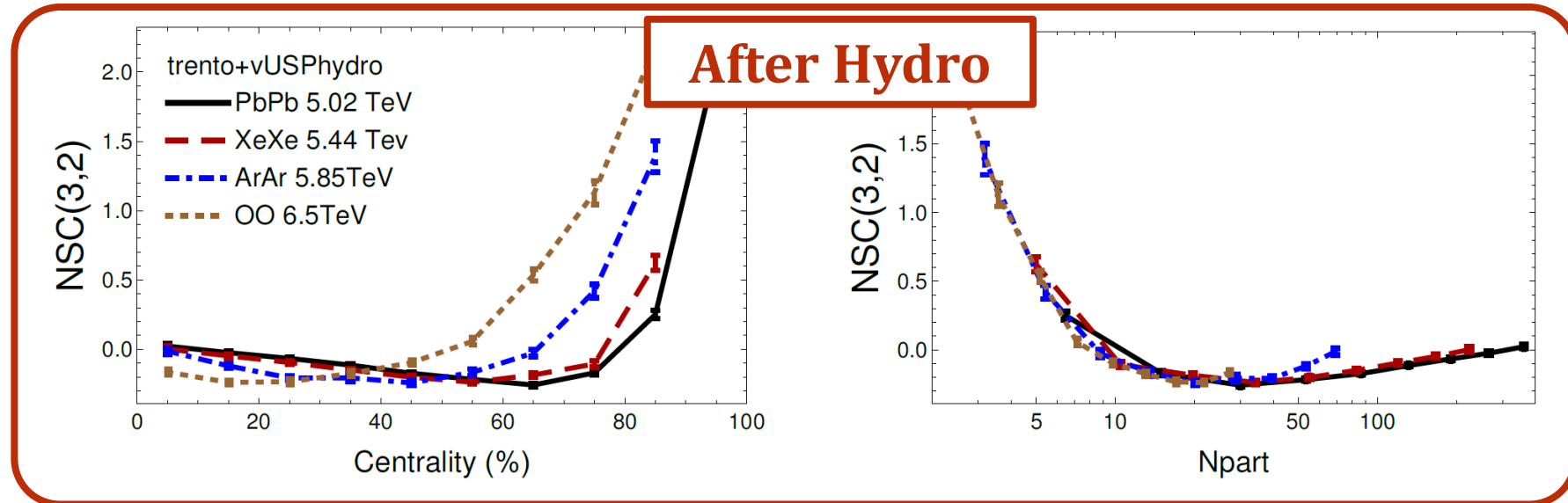
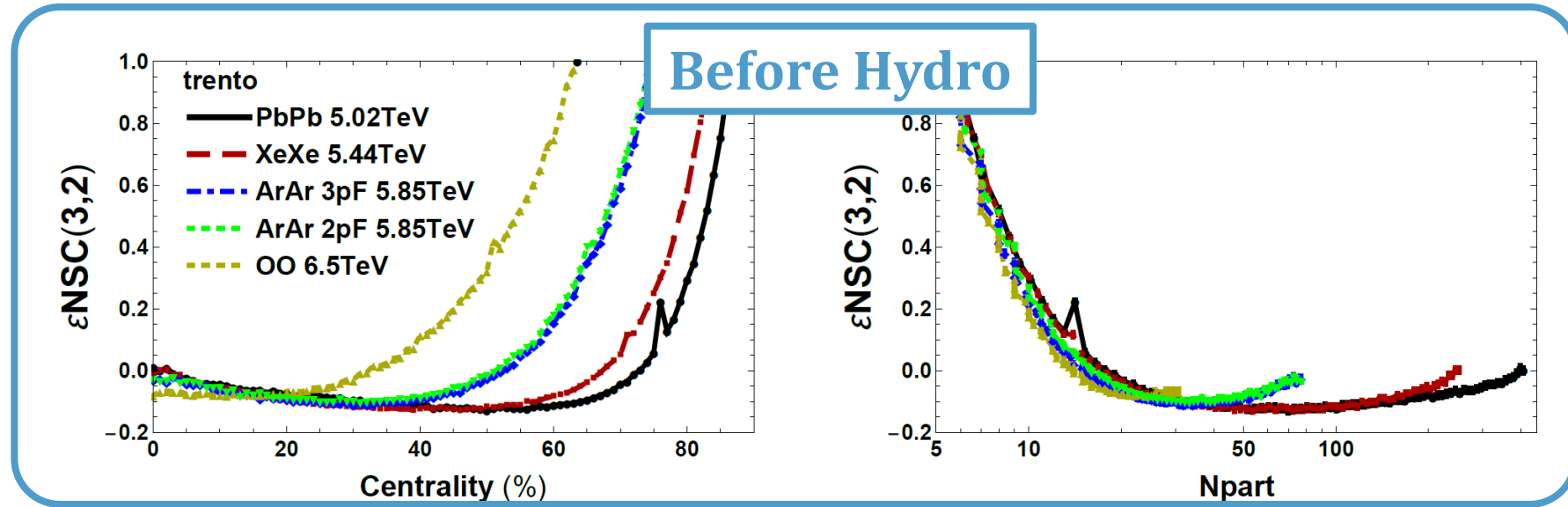
$$NSC(m, n) = \frac{\langle v_m^2 v_n^2 \rangle - \langle v_m^2 \rangle \langle v_n^2 \rangle}{\langle v_m^2 \rangle \langle v_n^2 \rangle}$$

- If $NSC(3,2) > 0$ then v_2 and v_3 are **positively correlated**
- If $NSC(3,2) < 0$ then v_2 and v_3 are **negatively correlated**
- If $NSC(3,2) = 0$ then v_2 and v_3 are **uncorrelated**

Initial state analog $\varepsilon NSC\{m\}$ defined with averages of ε_n

Symmetric Cumulants

- **Linear response** is a good qualitative predictor of correlations
- In very **central events**, NSC(3,2) is controlled by **centrality** (impact parameter)
- In **peripheral events**, NSC(3,2) is controlled by **small-number fluctuations**



Motivation: Using Deformed Nuclei as Model-Killers

- Very **central collisions** are **sensitive to deformation** of nuclei and other nuclear structure parameters

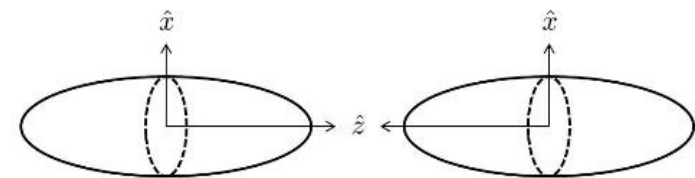
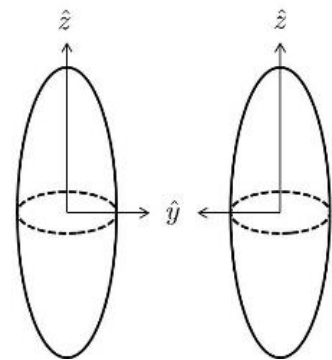
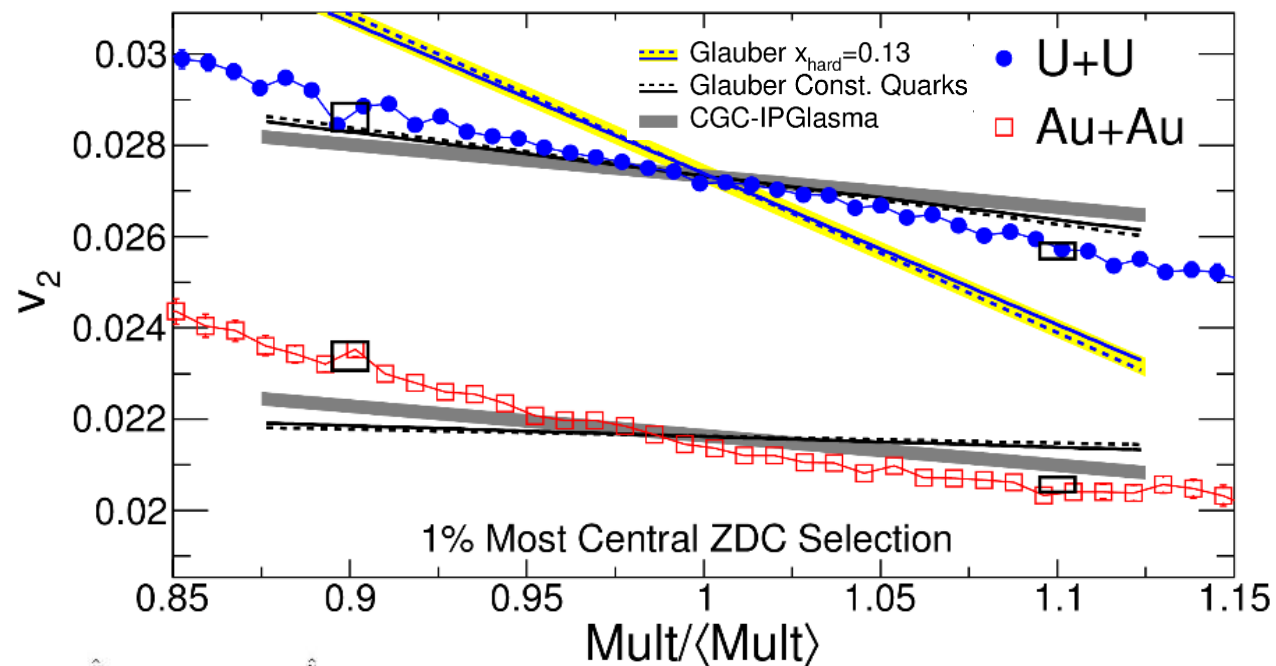
G. Giacalone et al., Phys. Rev. C97 (2018)

- The **multiplicity dependence** of ultracentral collisions of deformed nuclei can be used to **tune collision geometries**
- Different models lead to **different multiplicity dependence**, which **couple to nuclear deformations differently**
- Motivates a **deformed system size scan**
 - UU, ${}^9\text{Be}{}^9\text{Be}$ / ${}^9\text{BeAu}$, ${}^3\text{He}{}^3\text{He}$ / ${}^3\text{HeAu}$, dAu

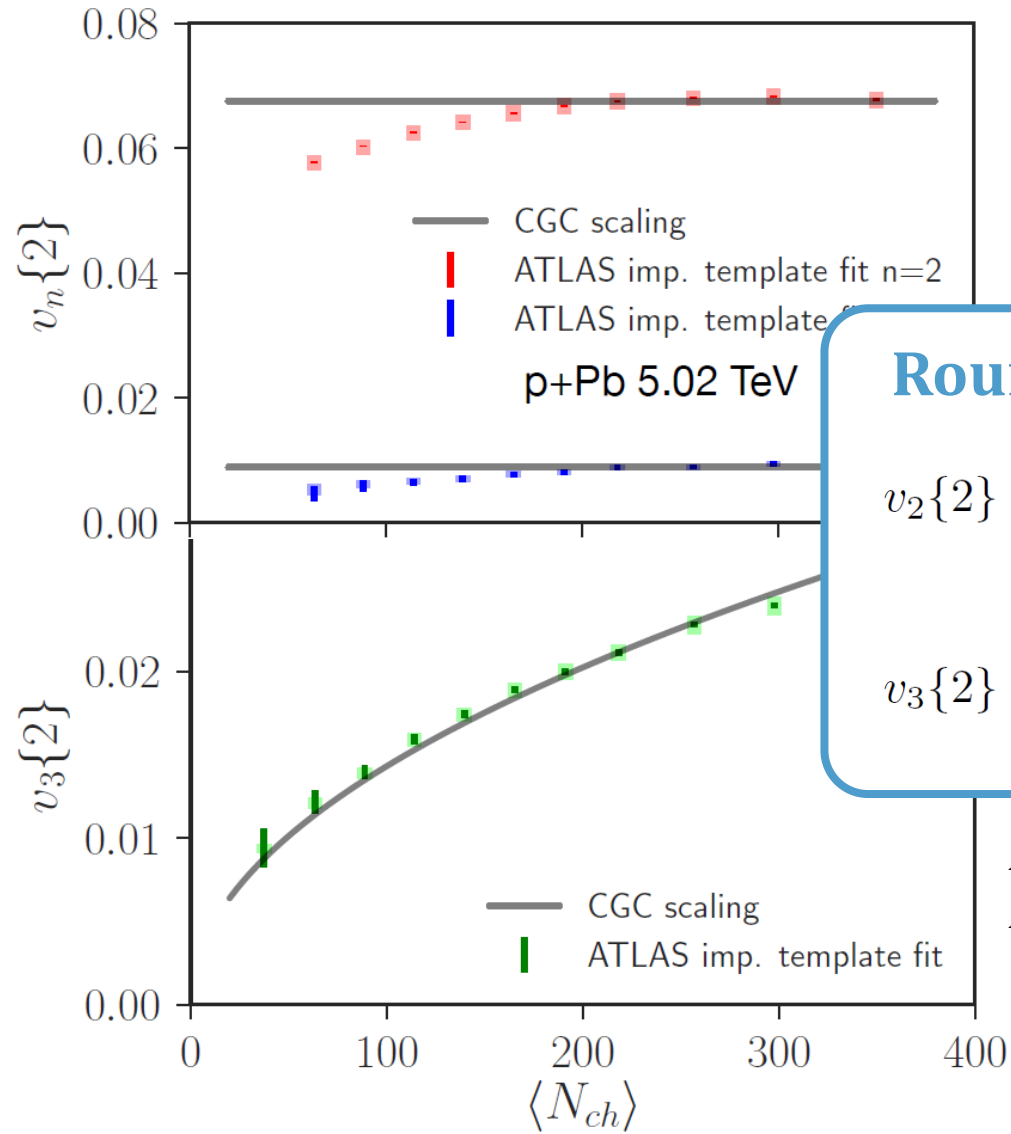
Uranium – Uranium Collisions at STAR

- STAR used a **novel centrality binning technique** for ultracentral collisions to select on **collision geometry**
- Used **ZDC cuts** to select on events with the **1% fewest spectators**
- Further **sub-binning by produced multiplicity** is sensitive to overall collision geometry
- Sensitive to the presence of **nucleonic substructure??**

L. Adamczyk et al., Phys. Rev. Lett. 115 (2015)



Multiplicity Dependence of CGC Gluon Correlations

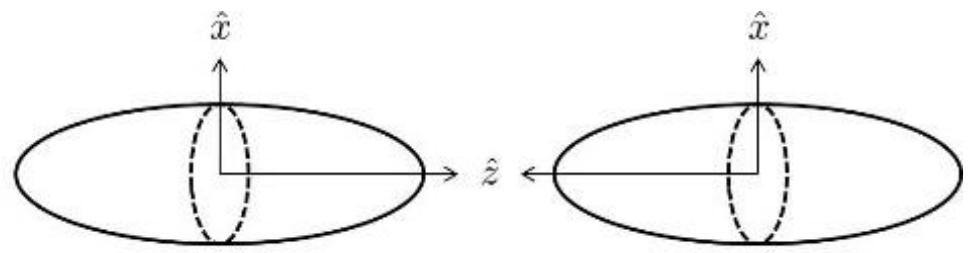


Round Nuclei

$v_2\{2\} \stackrel{T_A T_B}{\sim} \text{const}$
 $v_3\{2\} \stackrel{T_A T_B}{\sim} \sqrt{\frac{dN}{d^2x}}$

M. Mace et al.,
Phys. Lett. B788 (2019)

$$\frac{[v_2\{2\}]_{\text{tip}}}{[v_2\{2\}]_{\text{side}}} = \frac{1}{\lambda} > 1 \quad \text{Deformed Nuclei}$$



$$\rho(\vec{r}) = \rho_0 \exp \left[-\frac{x^2 + y^2}{R^2} - \lambda^2 \frac{z^2}{R^2} \right]$$

Y. Kovchegov and D. Wertepny,
Nucl. Phys. A925 (2014)

Apples to Apples: CGC Correlations vs. Hydrodynamics

- In the “dilute / dilute” limit, the **gluon correlations in the CGC** can be expressed in terms of **moments of the density profiles**

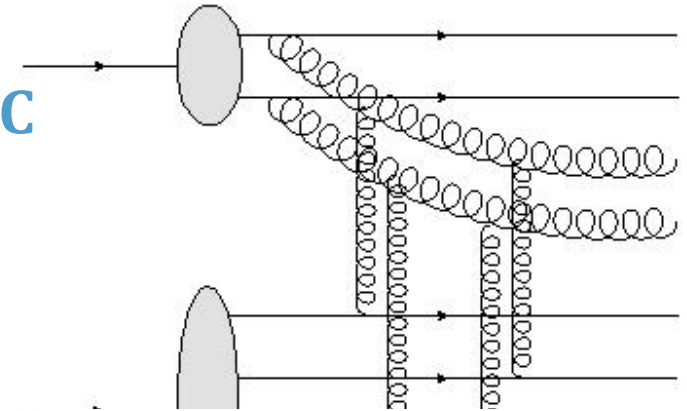
$$\frac{v_n^i\{2\}}{v_n\{2\}} = \sqrt{\frac{\langle N_{tot}^2 \rangle_{0-1\%}}{\langle N_{tot}^2 \rangle_i} \frac{\langle \mathcal{I}_{\nu_n} \rangle_i}{\langle \mathcal{I}_{\nu_n} \rangle_{0-1\%}}}$$

$$\mathcal{I}_\alpha \equiv \int d^2x_\perp T_A^\alpha(\vec{x}_\perp) T_B^\alpha(\vec{x}_\perp)$$

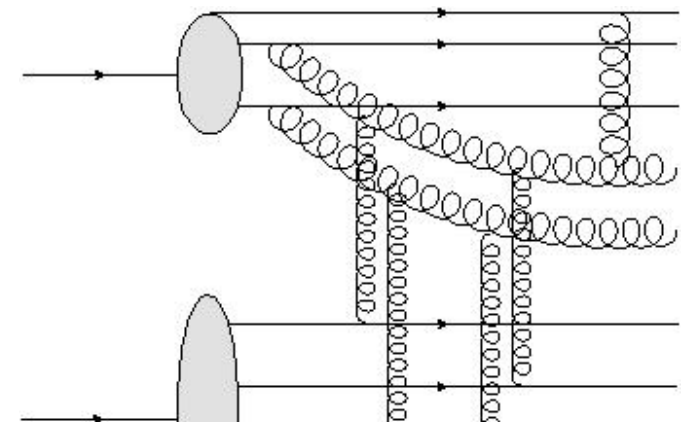
$$\nu_n = \begin{cases} 2 & \text{if } n = \text{even} \\ 3 & \text{if } n = \text{odd} \end{cases}$$

$$N_{tot} \propto T_A T_B \quad (\text{CGC})$$

$$N_{tot} \propto \sqrt{T_A T_B} \quad (\text{Trento } p=0)$$



$$\text{LO: } v_2\{2\} \sim \frac{1}{N^2} \int T_A^2 T_B^2$$

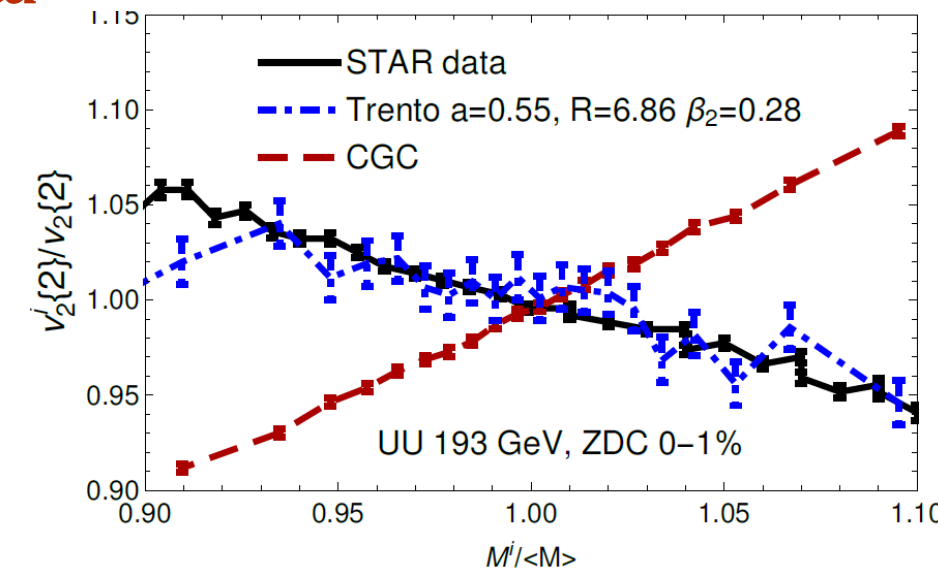
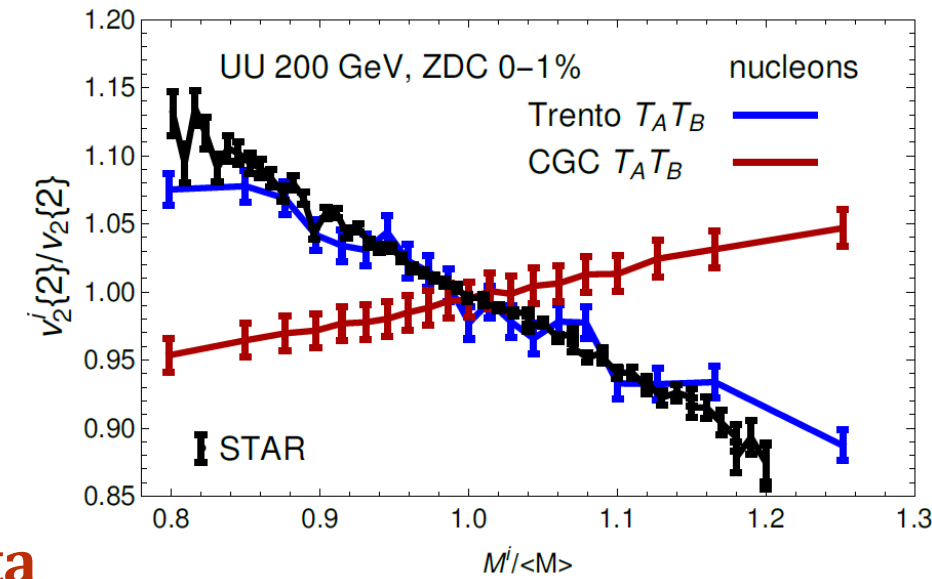


$$\text{NLO: } v_3\{2\} \sim \frac{1}{N^2} \int T_A^3 T_B^3$$

Deformed Nuclei can Discriminate Models

Using Trento + linear response
to estimate final-state hydrodynamics

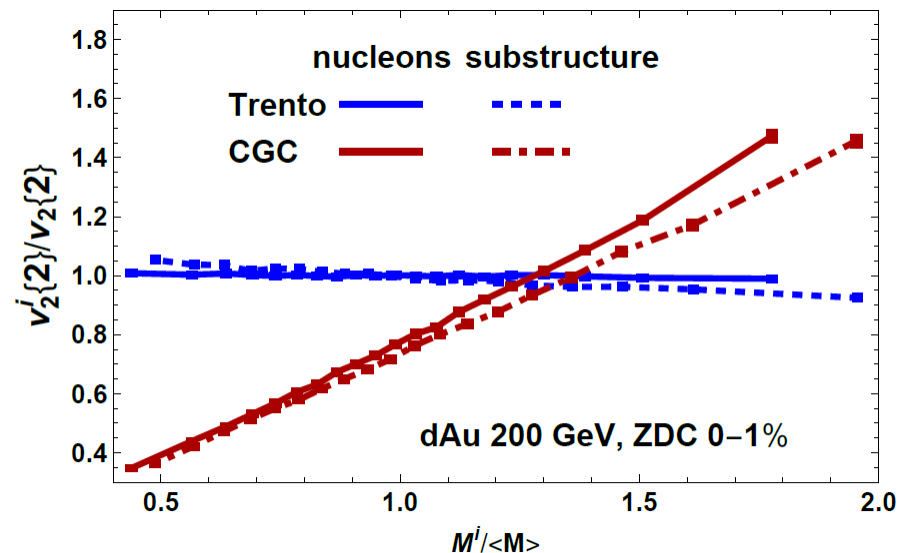
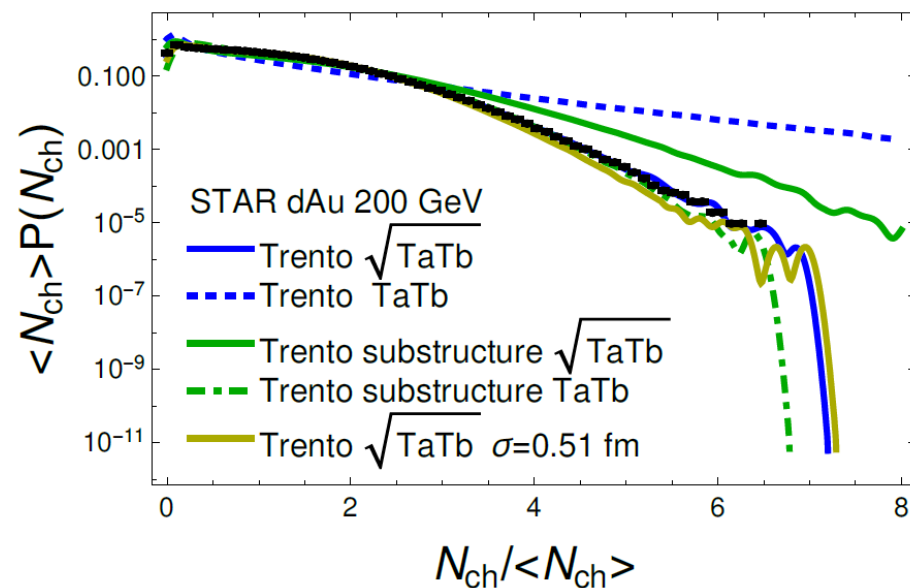
- As expected, CGC predicts **opposite multiplicity dependence** from hydro and STAR
- An initial-state only picture **cannot describe the data**
- Hydro is largely insensitive to choice of $T_A T_B$ vs $\sqrt{T_A T_B}$ scaling
- CGC slope flattened significantly by linear $T_A T_B$ scaling



Small Deformed Systems: dAu

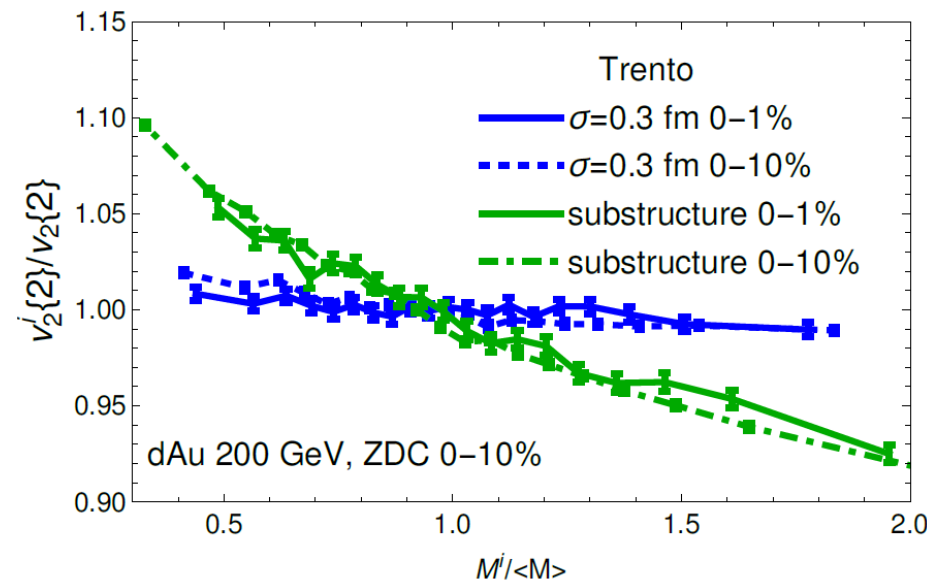
- Same qualitative trends remain in dAu

- Can study resolving power of **nucleonic substructure** in Trento 2.0



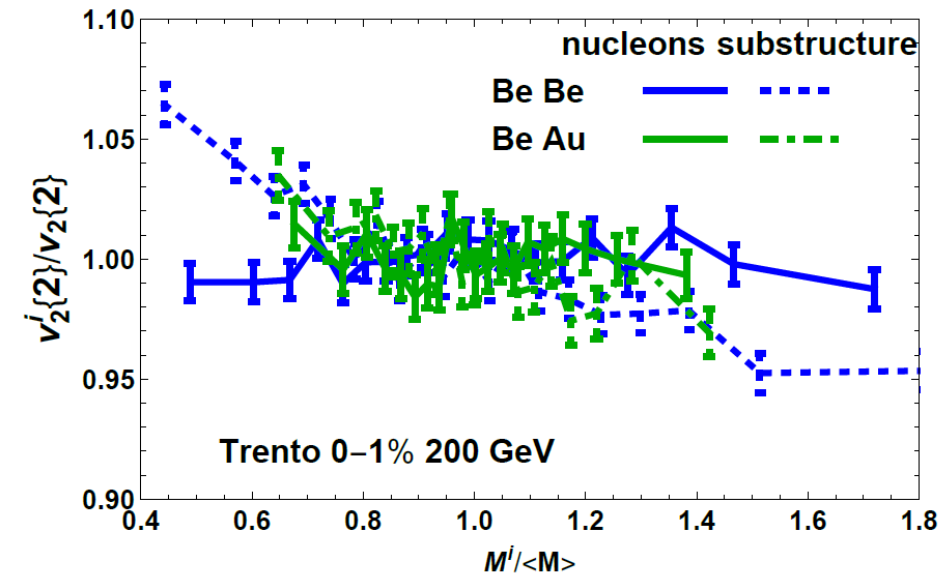
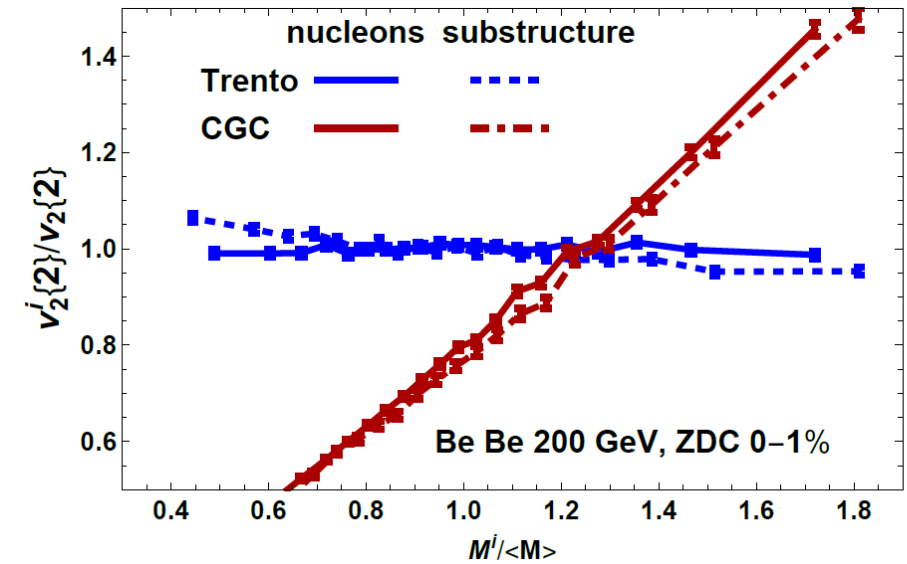
- Small **0(5%) sensitivity** to substructure (n=6)
- **Enhanced to 0(10%) sensitivity** in 0-10% centrality

- **$\sqrt{T_A T_B}$ scaling** gets multiplicity distribution right



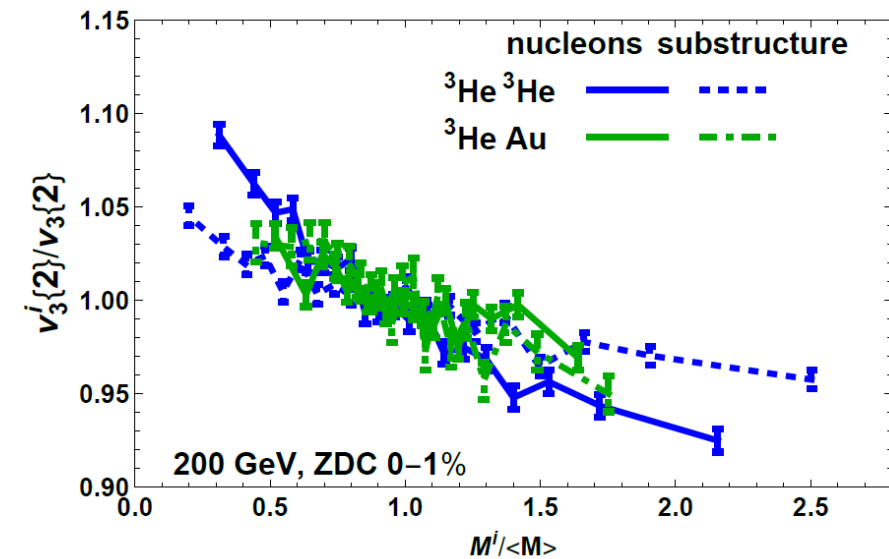
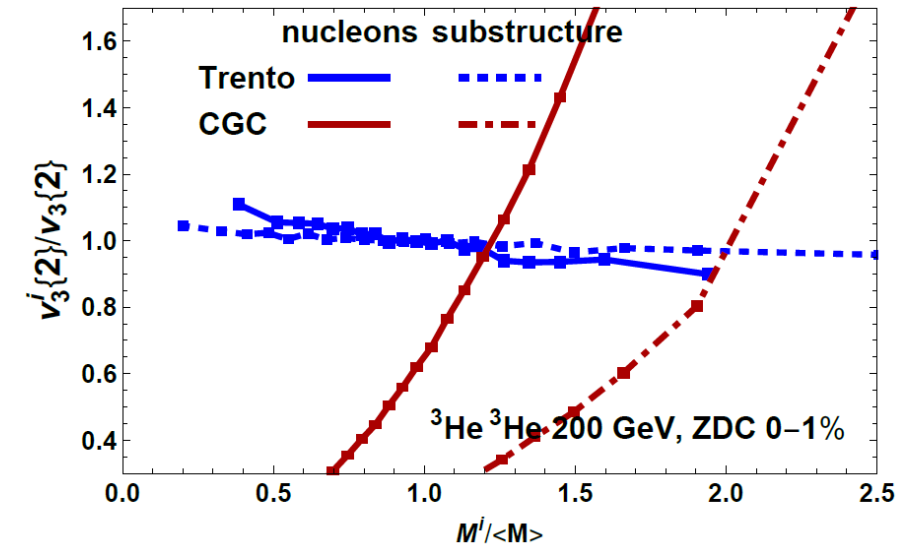
Small Deformed Systems: ${}^9\text{Be}{}^9\text{Be}$ and ${}^9\text{BeAu}$

- ${}^9\text{Be}$ is a **highly deformed ion**: its quadrupole moment is more than **twice as large as uranium**
- Intermediate collision systems with large deformation: ${}^9\text{Be}{}^9\text{Be}$ and ${}^9\text{BeAu}$ collisions
- Qualitative CGC vs hydro distinction persists
- Sensitive to nucleonic substructure at $O(5\%)$
- **${}^9\text{Be}{}^9\text{Be}$ collisions are better** than ${}^9\text{BeAu}$ collisions for identifying **sub-nucleonic fluctuations**



Small Deformed Systems: ${}^3\text{He}{}^3\text{He}$ and ${}^3\text{HeAu}$

- ${}^3\text{He}$ possesses an innate triangularity at the level of nucleons – more direct probe of v_3
- CGC multiplicity dependence is steeper than for v_2 due to entering at NLO in dilute / dense
- Still sensitive at **0(5%) to nucleonic substructure**, but this time due to a **flattening** of triangularity when turning on substructure
- ${}^3\text{He}{}^3\text{He}$ collisions are better than ${}^3\text{He Au}$ collisions for identifying **sub-nucleonic fluctuations**



Conclusions

- Drops of QGP created from smaller ions are **rounder** and **hotter** than comparable-multiplicity drops created from heavier ions.
- **Linear response** to initial eccentricities is a **good estimator for small collision systems**, with linear + cubic response improving quantitative control.
- Some features controlled by **universal small-number statistics in peripheral collisions**, others controlled by **impact parameter in more central collisions**.
- Collisions of **small, deformed systems** can **discriminate between models**, including **initial state CGC vs final state flow**, as well as nucleon vs **subnucleonic structure**.